

NLO Event Generation for Chargino Production at the ILC

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Tania Robens

in collaboration with W. Kilian, J. Reuter

RWTH Aachen

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 - Charginos and Neutralinos in the MSSM
 - Experimental accuracy and NLO results

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 - Photons: fixed order vs resummation
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Chargino and Neutralino sector: Reconstruction of SUSY parameters

- Charginos $\tilde{\chi}_i^\pm$ and Neutralinos $\tilde{\chi}_i^0$:
superpositions of gauge and Higgs boson superpartners
- Chargino/ Neutralino sector:

$\tan\beta$, μ (Higgs sector), M_1 , M_2 (soft breaking terms)

can be reconstructed from

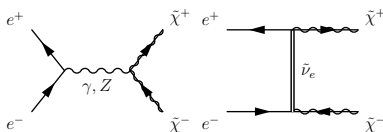
masses of $\tilde{\chi}_1^\pm$, $\tilde{\chi}_2^\pm$, $\tilde{\chi}_1^0$, 2σ in the $\tilde{\chi}^\pm$ sector

(Choi ea 98, 00, 01)

- low-scale parameters + evolution to high scales (RGEs):
 \Rightarrow hint at SUSY breaking mechanism (Blair ea, 02)
- requires high precision in ew-scale parameter determination

Chargino production at the ILC

- Charginos: (typically) light in the MSSM
 \Rightarrow easily accessible at colliders (ILC/ LHC) \Leftarrow
- LO production at the ILC:



- decays: typically long decay chains

$$\text{e.g. } e^+ e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \tilde{\tau}_1^+ \tilde{\tau}_1^- \nu_\tau \bar{\nu}_\tau (\rightarrow \tau^+ \tau^- \nu_\tau \bar{\nu}_\tau \tilde{\chi}_1^0 \tilde{\chi}_1^0)$$

Experimental accuracy and theoretical next-to-leading-order (NLO) corrections

- experimental errors: obtained from simulation studies (LHC/ ILC study, Weiglein ea, 04)
- generate “experimental data” with known SUSY input parameters
- errors: combination of statistical and systematic errors

combined **LHC + ILC**: ‰

same \mathcal{O} errors from fitting routines determining SUSY parameters

- **Theory:**

Full NLO SUSY corrections for $\sigma(ee \rightarrow \tilde{\chi} \tilde{\chi})$ at ILC:
in the ‰ regime (Fritzsche ea 04, Öller ea 04, 05)

⇒ include complete NLO contributions in analyses←

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From σ_{tot} to Monte Carlo event generators

- experiments: see final decay products
- need to compare with simulated event samples
- also: important irreducible background effects
(e.g. Hagiwara ea, 05)

⇒ include NLO results in Monte Carlo Generators ⇐

- MC Generator WHIZARD (W. Kilian, LC-TOOL-2001-039):
- so far: LO Monte Carlo Event Generator for $2 \rightarrow n$ particle processes
- includes various physical models (SM, MSSM, non-commutative geometry, little Higgs models), initial state radiation, parton shower models,...

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NLO cross section contributions

σ_{tot} contributions and dependencies:

- σ_{born}
- virtual $\mathcal{O}(\alpha)$ corrections: $\sigma_{\text{virt}}(\lambda)$
- emission of soft/ hard collinear/ hard non-collinear photons:

$$\sigma_{\text{soft}}(\Delta E_\gamma, \lambda) + \sigma_{\text{hc}}(\Delta E_\gamma, \Delta\theta_\gamma) + \sigma_{2 \rightarrow 3}(\Delta E_\gamma, \Delta\theta_\gamma)$$

- higher order initial state radiation: $\sigma_{\text{ISR}} - \sigma_{\text{ISR}}^{\mathcal{O}(\alpha)}(Q)$
 λ : photon mass, ΔE_γ : soft cut, $\Delta\theta_\gamma$: collinear angle

Including FormCalc $\mathcal{O}(\alpha)$ results in WHIZARD

- use FeynArts / FormCalc generated code for

$$\begin{aligned} \mathcal{M}_{\text{virt}}(\lambda) &: \text{ virtual corrections} \\ f_s(\Delta E_\gamma, \lambda) &: \text{ soft photon factor} \\ (\mathcal{M}_{\text{born}} &: \text{ born contribution}) \end{aligned}$$

- fixed order: integrate over effective matrix element:

$$|\mathcal{M}_{\text{eff}}|^2(\Delta E_\gamma) = (1 + f_s(\Delta E_\gamma, \lambda)) |\mathcal{M}_{\text{born}}|^2 + 2 \text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*(\lambda))$$

ΔE_γ : soft photon cut, λ : photon mass

- in practice: create library from FormCalc code, link this to WHIZARD

(1): Fixed $\mathcal{O}(\alpha)$ contributions

- integrate $|\mathcal{M}_{\text{eff}}|^2$ (born/ virtual/ soft photonic part)
- hard collinear photons: collinear approximation ($\mathcal{M}_{\text{born}}$)
- hard non-collinear photons: explicit $e e \rightarrow \tilde{\chi} \tilde{\chi} \gamma$ process ($\mathcal{M}_{\text{born}}^{2 \rightarrow 3}$)
- corresponds to analytic results in literature (Fritzsche ea/ Öller ea)

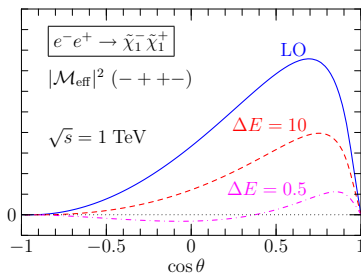
Photons: fixed order vs resummation

(1): Fixed $\mathcal{O}(\alpha)$ contributions

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problem: too low energy cuts: $|\mathcal{M}_{\text{eff}}|^2 < 0$
 \Rightarrow use negative weights
 or set $\mathcal{M}_{\text{eff}} = 0$

event generator
specific problem
 $(\sigma_{\text{tot}} \geq 0)$



\mathcal{M}^2 behaviour, different cuts [GeV]

(2): Resumming leading logs to all orders

- idea: subtract $\mathcal{O}(\alpha)$ soft + virtual collinear contributions in \mathcal{M}_{eff} :

$$|\widetilde{\mathcal{M}}_{\text{eff}}|^2 = (1 + f_s(\Delta E_\gamma)) |\mathcal{M}_{\text{born}}|^2 + 2 \text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*) - 2 f_s^{\text{ISR}, \mathcal{O}(\alpha)}(\Delta E_\gamma) |\mathcal{M}_{\text{born}}|^2$$

- fold this with ISR structure function:

$$\int d\Gamma \int_0^1 dx_1 \int_0^1 dx_2 f^{\text{ISR}}(x_1) f^{\text{ISR}}(x_2) |\widetilde{\mathcal{M}}_{\text{eff}}|^2(s, x_i)$$

- $f^{\text{ISR}}(x)$: Initial state radiation (Jadach, Skrzypek, Z.Phys. 1991)
 \Rightarrow describes collinear (real + virtual) photons in leading log accuracy \Leftarrow
- $f_s^{\text{ISR}, \mathcal{O}(\alpha)}$: soft integrated $\mathcal{O}(\alpha)$ contribution

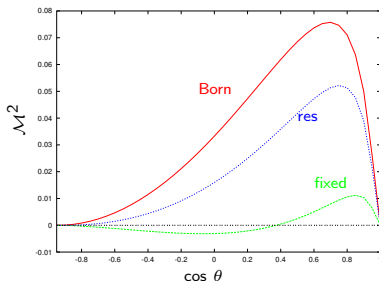
Photons: fixed order vs resummation

Resumming: What do we get ??

- $\mathcal{O}(\alpha)$: equivalent to fixed order method

⇒ got rid of
 $|\mathcal{M}|^2 < 0$
effects !!

**no negative
weights**



(-+-),

$\Delta E_\gamma = 0.5 \text{ GeV}$

- higher orders:
higher order ISR for $|\mathcal{M}_{\text{born}}|^2$ as well as $\text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*)$!!!
⇒ new higher order effects ←

additional possibility: also fold hard noncollinear process with ISR

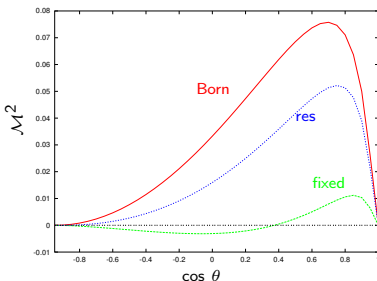
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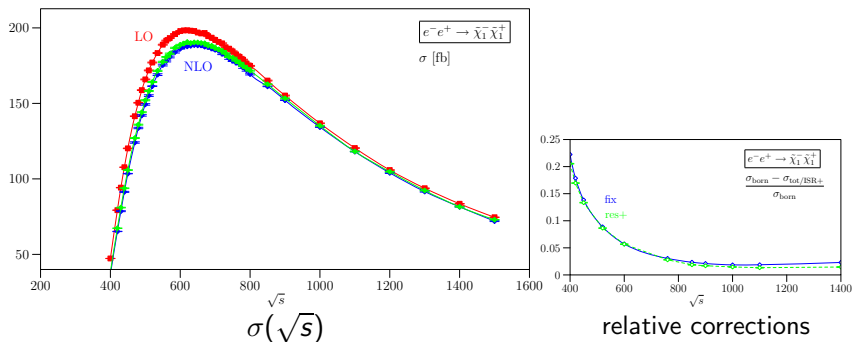


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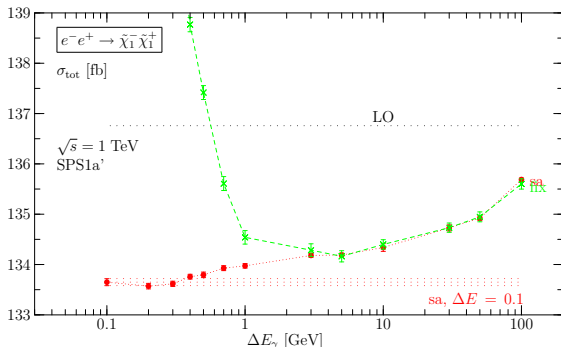
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Results: cross sections



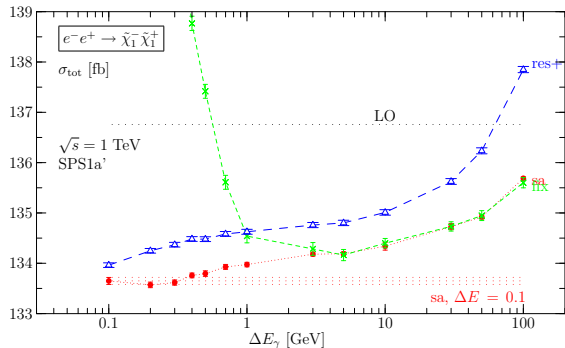
agrees with results in the literature (Fritzsche ea, Öller ea)

A closer look: ΔE_γ dependence of σ_{tot}



- **semianalytic (FormCalc)**: tests soft approximation, shifts : 2 - 5 ‰ ($\Delta E_\gamma \leq 10 \text{ GeV}$)
- **fixed order result (WHIZARD)**: same as 'sa' for $\Delta E_\gamma \geq 3 \text{ GeV}$, smaller values: $|\mathcal{M}_{\text{eff}}|^2 \leq 0$ effects

ΔE_γ dependence: resummation



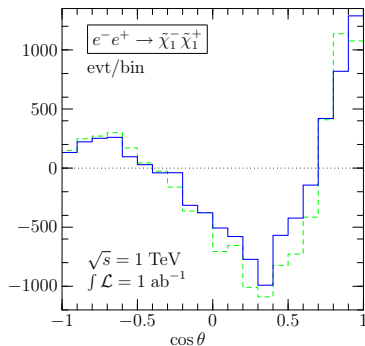
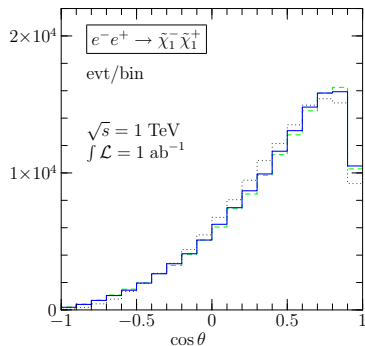
$\sigma_{\text{tot}}(\Delta E_\gamma)$:
resummation includes
 higher order effects
 5% difference to 'sa'
 for $\Delta E_\gamma \leq 10 \text{ GeV}$

In summary:

shift in ΔE_γ leads to % effects, match ILC accuracy
 \Rightarrow careful choice of ΔE_γ , method important

“best” choice: fully resummed version with low energy cut

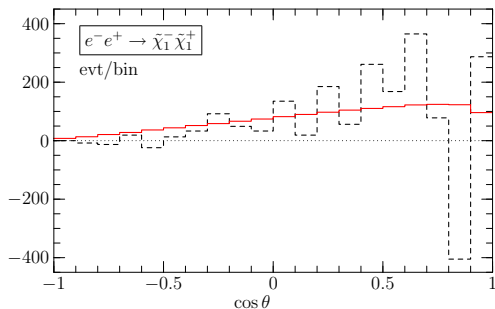
simulation results: angular distributions



Born, fixed order, resummation

!! more than 1σ deviation !! $\sqrt{n_{\max}} \approx \mathcal{O}(10^2)$; nbins = 20

Angular distributions: higher orders

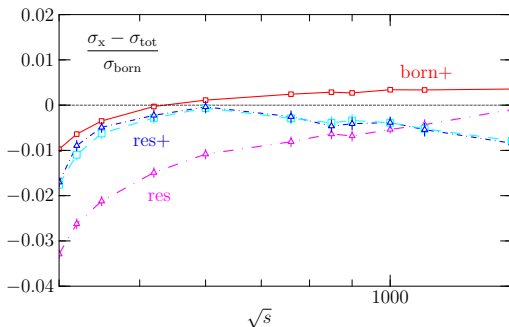


$N_{\text{res,+}} - N_{\text{ex}}$
 red: 1 standard dev
 from Born result

also higher order contributions statistically significant

Results: higher order effects

\sqrt{s} dependence of different higher order contributions



Born+: only Born folded w ISR, resummation , fully resummed result

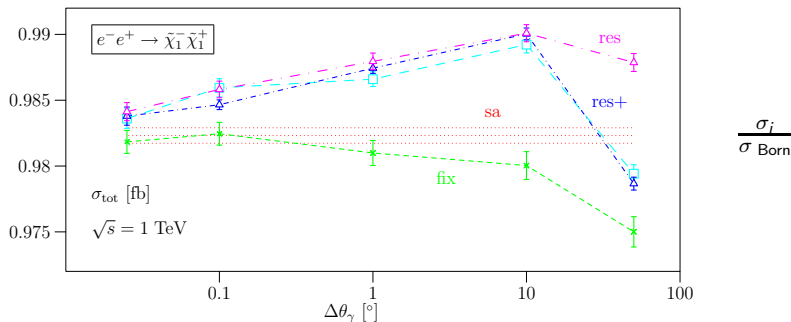
difference between **Born+** and fully resummed result: multiple photon emission from interaction term

Summary and Outlook

- Chargino/ neutralino sector of MSSM: high precision in SUSY parameter analysis at EW scale ($\%_0$ at ILC)
- same size/ larger NLO corrections
- ⇒ include NLO results in Monte Carlo Event generators
- resummation method for photons allows lower soft cuts/ inclusion of higher order contributions
- NLO as well as higher order contributions significant !!
- next steps: include NLO corrections to $\tilde{\chi}$ decays, non-factorizing contributions (start with photonic corrections in the double-pole approximation)
- general interface to FormCalc generated matrix elements: extendable to other processes...

cut dependencies: $\Delta\theta_\gamma$

tests: collinear photon approximation



σ_{tot} again larger for resummation method
 for higher angles: second order ISR effects between 0.05° and 0.1°
 ($\mathcal{O}(\%)$)

η, f_s , hard collinear approximation, $ISR^{\mathcal{O}(\alpha)}$

- $\eta = \frac{2\alpha}{\pi} \left(\log \left(\frac{Q^2}{m_e^2} \right) - 1 \right)$ (Q = scale of process)



$$f_s = -\frac{\alpha}{2\pi} \sum_{i,j=e^\pm} \int_{|\mathbf{k}| \leq \Delta E} \frac{d^3 k}{2\omega_k} \frac{(\pm) p_i p_j Q_i Q_j}{p_i k p_j k},$$

(Denner 1992)

$\omega_k = \sqrt{\mathbf{k}^2 + \lambda^2}$, p_i initial/ final state momenta, k : γ momentum

- hard collinear factor (\pm helicity conserving/ flipping):

$$f^+(x) = \frac{\alpha}{2\pi} \frac{1+x^2}{(1-x)} \left(\ln \left(\frac{s(\Delta\theta)^2}{4m^2} \right) - 1 \right), \quad f^-(x) = \frac{\alpha}{2\pi} x.$$

(Dittmaier 1993)



$$f_s^{ISR, \mathcal{O}(\alpha)} = \left[\int_{x_0}^1 f_{ISR}(x) dx \right]_{\mathcal{O}(\alpha)} = \frac{\eta}{4} \left(2 \ln(1-x_0) + x_0 + \frac{1}{2} x_0^2 \right)$$

ISR in its full beauty (Skrzypek ea, 91)

$$\begin{aligned}
 \Gamma_{ee}^{LL}(x, Q^2) = & \frac{\exp(-\frac{1}{2}\eta\gamma_E + \frac{3}{8}\eta)}{\Gamma(1 + \frac{\eta}{2})} \frac{\eta}{2} (1-x)^{(\frac{\eta}{2}-1)} \\
 & - \frac{\eta}{4} (1+x) + \frac{\eta^2}{16} \left(-2(1-x) \log(1-x) - \frac{2 \log x}{1-x} + \frac{3}{2} (1+x) \log x - \frac{x}{2} \right. \\
 & - \left. \frac{5}{2} \right) + \left(\frac{\eta}{2} \right)^3 \left[-\frac{1}{2} (1+x) \left(\frac{9}{32} - \frac{\pi^2}{12} + \frac{3}{4} \log(1-x) + \frac{1}{2} \log^2(1-x) \right. \right. \\
 & - \left. \left. \frac{1}{4} \log x \log(1-x) + \frac{1}{16} \log^2 x - \frac{1}{4} \text{Li}_2(1-x) \right) \right. \\
 & + \left. \frac{1}{2} \frac{1+x^2}{1-x} \left(-\frac{3}{8} \log x + \frac{1}{12} \log^2 x - \frac{1}{2} \log x \log(1-x) \right) \right. \\
 & - \left. \frac{1}{4} (1-x) \left(\log(10x) + \frac{1}{4} \right) + \frac{1}{32} (5-3x) \log x \right] ; \eta = \frac{2\alpha}{\pi} \left(\log \left(\frac{Q^2}{m_e^2} \right) - 1 \right)
 \end{aligned}$$