

# Leptonic (and more) Resonances at the LHC



- di-lepton resonances
  - $Z'$ -like resonances
  - $W'$
- 3<sup>rd</sup> family resonances
- Little Higgs
- leptoquarks
- excluding: Higgs, SUSY, excited fermions



## Resonances: early LHC physics

### Resonances should be relatively easy and quick to discover, but:

- what are they?, in what model do they fit? what are the couplings to fermions and bosons?
- what can we measure, and with what precision?
  - production cross section and decay branching ratios
  - width
  - interferences (and interference with Drell-Yan for Z')
  - FB asymmetry
  - rapidity distribution
  - associated with what other new physics?

**LHC:** a discovery machine

**ILC:** precision measurements

## 2-body resonances

$e/\mu$	$e/\mu$	$\tau$	$\nu$	$j$	$b$	$t$	$\gamma$	$W/Z$	$h$
$e/\mu$	$Z', G^*, \rho_{TC}, H^{++}$	LFV	$W'$	LQ, (2l-2j)	LQ, (2l-2j)	LQ, (2l-2j)	$e^*, \mu^*$	$e^*, \mu^*, \nu^*, e', \mu'$	
$\tau$		$Z', G^*, \rho_{TC}, H^{++}, A/H$	$W'$	LQ, (2l-2j)	LQ, (2l-2j)	LQ, (2l-2j)	$\tau^*$	$\tau^*, \tau', \nu^*$	
$\nu$				LQ, (2l-2j)	LQ, (2l-2j)	LQ, (2l-2j)	$\nu^*(?)$	$\nu^*, \nu'$	
$j$				$Z', G^*, \rho_{TC}, H$	$\pi_{TC}$	$\pi_{TC}$	$q^*$	$q^*, q', D$	D
$b$					$Z', G^*, \rho_{TC}, H, h$	$W'$	$b^*$	$b^*, b', T$	D
$t$						$Z', G^*, \rho_{TC}, H, g_{KK}$	$t^*$	T	T
$\gamma$							$G^*, \rho_{TC}, H$	$W'$	
$W/Z$								$Z', G^*, \rho_{TC}, H, W'$	$W', Z', H$
$h$									H/A

## Z'-like resonances

### ➔ from new U(1)

- $E_6$ , L-R symmetric model,
- little Higgs (including  $Z_H, A_H$ )
- 3-3-1 (including doubly charged)

### ➔ from extraD's

- $Z_{KK}/\gamma_{KK}$  in bulk, TeV-1-size ED: EW constraints  $\rightarrow m_{Z'} \geq \sim 4$  TeV
  - including RS bulk:  $Z'/\gamma$ :  
flavor dependent coupling [Ledroit et al., hep-ph/0703262]
  - fermiophobic  $Z', W'$  (3-site model, see talk by S. Chivukula)
- gravitons in RS
  - Higgsless models
  - measure spin-2 and observe other channels
  - RS in bulk: lepton and photon couplings negligible  $\rightarrow$  look for top or VB pairs [Agashe et al., hep-ph/0701186]
- gluon: look for decay to top pairs and polarization

### ➔ technicolor

- $\rho_{TC}, \pi_{TC}$  BESS

### ➔ Higgs ( $H/A \rightarrow \mu\mu, \tau\tau, bb, tt$ )

### ➔ Hidden Valley ....

# Experimental issues at LHC

$e^+e^-$ ,  $\mu^+\mu^-$ : clean signal, little background (Drell-Yan)

- measure  $\sigma$ ,  $\Gamma$ ,  $A_{\text{FB}}$

$\tau^+\tau^-$ : reconstruction difficult, but possible

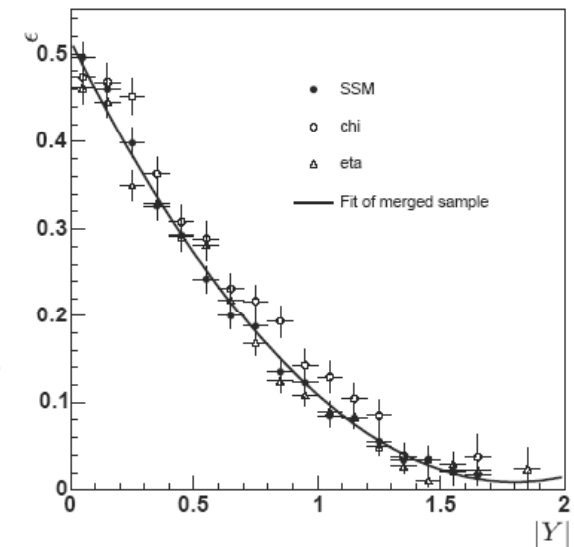
- measure  $\tau$  polarization and spin correlations

$t\bar{t}$ : very large QCD background

- can be seen if due to strong interaction resonance  $g_{KK} \rightarrow t\bar{t}$

## $A_{\text{FB}}$ is an important observable

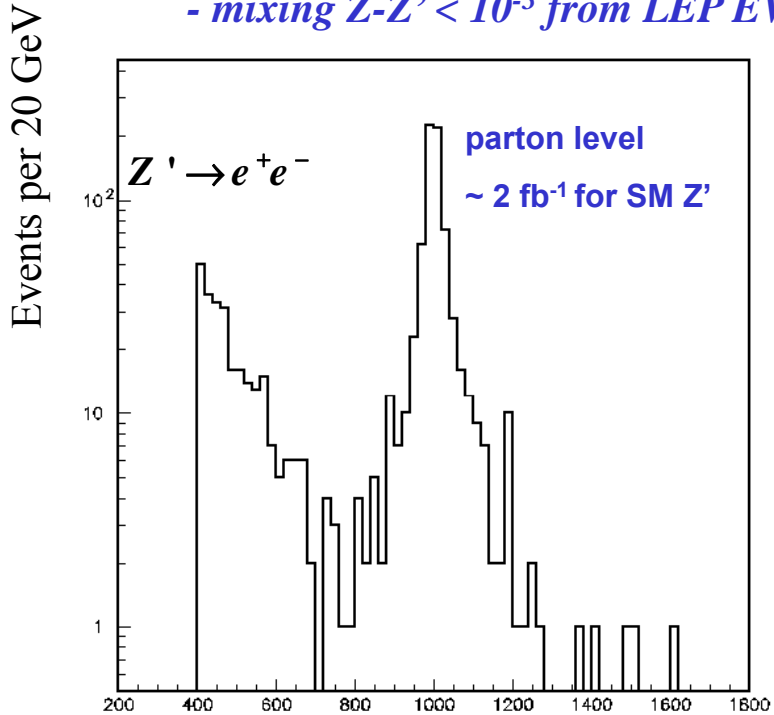
- In a pp collider, we must resort to “guess” that Z' boost is in direction of quark (by opposition to antiquark) forming it
  - reliable for large values of boost (rapidity of Z')
    - but pseudorapidity cut on leptons:  $|\eta| < 2.5$
  - can correct for dilution to a large extent
    - efficiency depends only on Y (given incoming type of quarks) because of the smmetr of the detector
    - some systematic uncertainties (PDF's...)



Probability to be wrong vs Y  
Ledroit et al., hep-ph/0703262

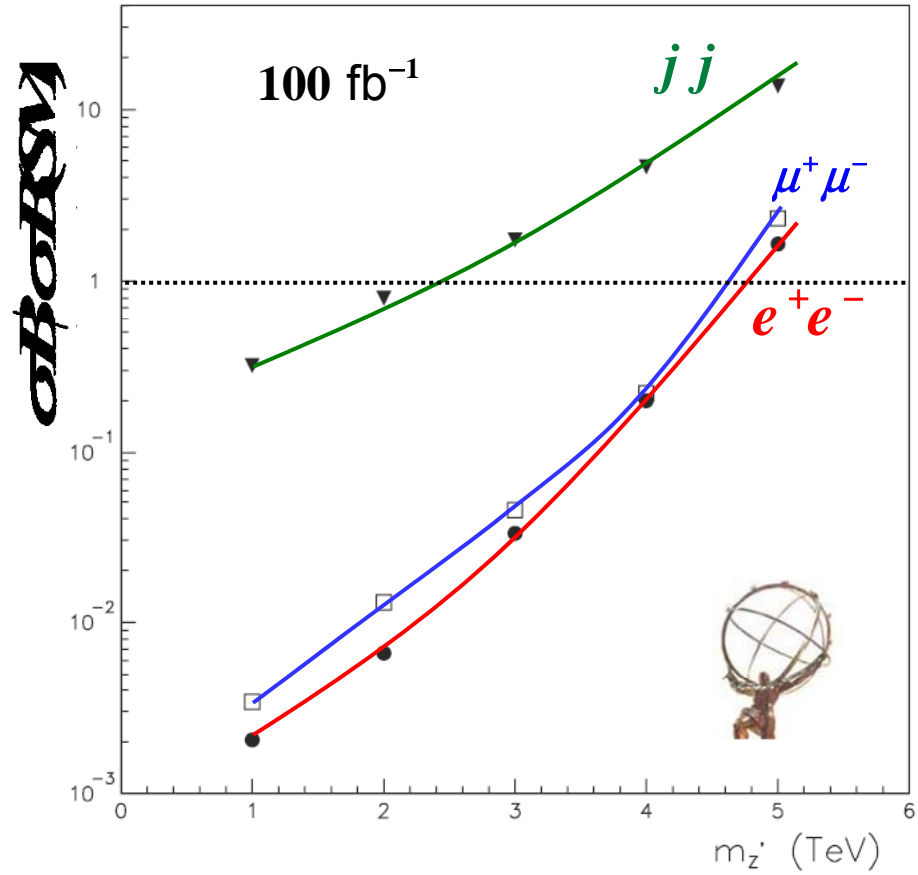
# Z' resonance discovery at the LHC

- present limits  $\sim 600$  GeV if SM couplings  $\rightarrow$  not interesting for 500 GeV ILC, unless couplings are lower than SM
- mixing  $Z$ - $Z' < 10^{-3}$  from LEP EW measurements



SM-like  $Z'$ :

- o width  $\sim 3\%$  of mass for SM  $Z'$ 
  - o could increase if  $Z' \rightarrow W W, Zh$
- o very little background
- o typical resolution:  $\frac{\sigma(E)}{E} \sim 0.75\%$



ATLAS Physics TDR

# E6 extensions of SM (superstring inspired)

$$E_6 \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_\eta \quad (\text{rank 5})$$

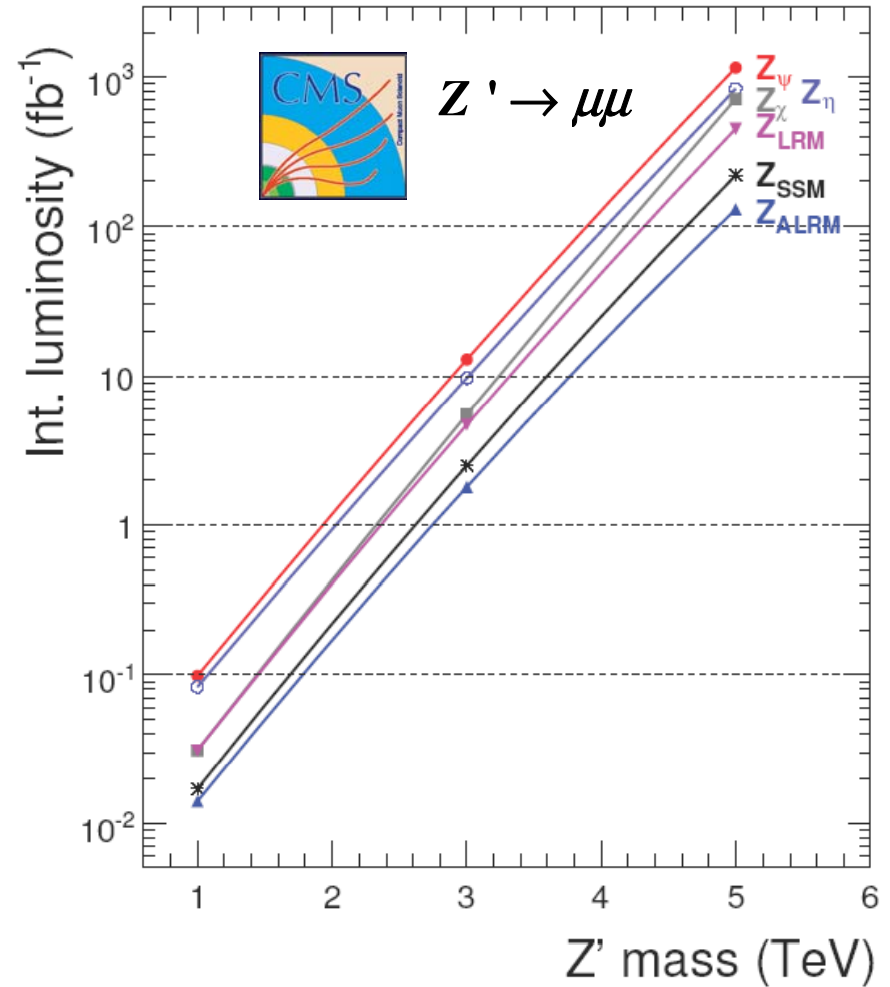
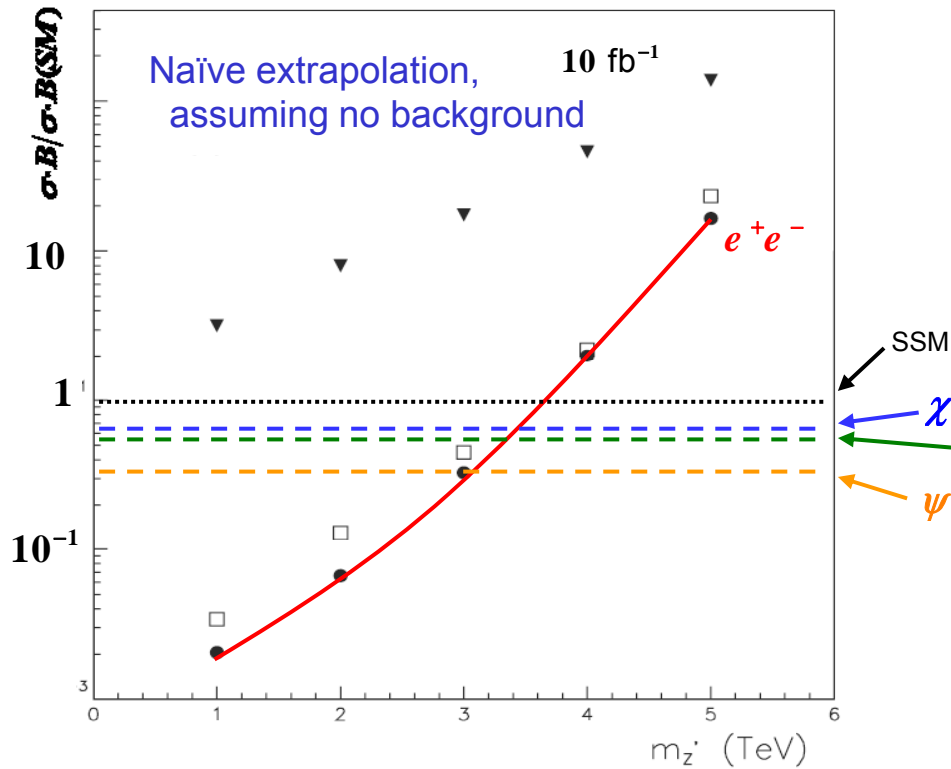
$$E_6 \rightarrow SO(10) \times U(1)_\psi$$

$$\hookrightarrow SU(5) \times U(1)_\chi$$

$$\hookrightarrow SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

J. Hewett and T. Rizzo, Phys Rep 183 (1989) 193

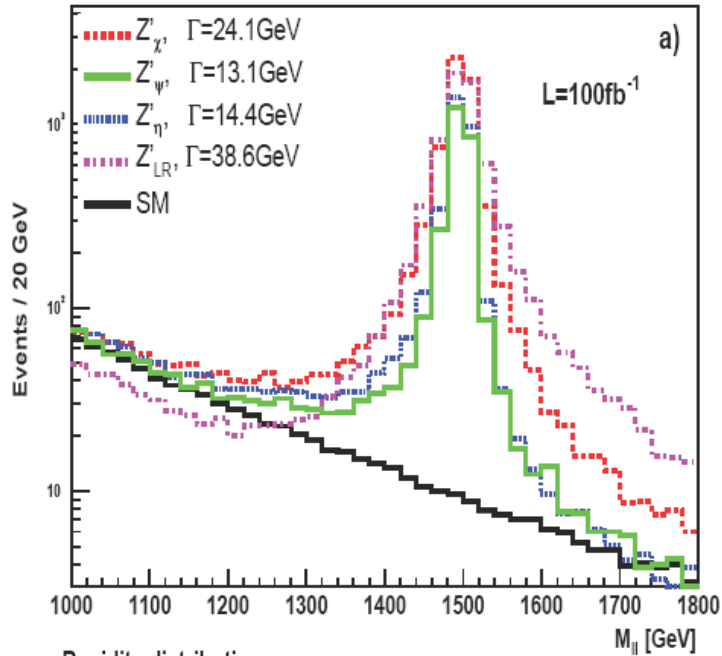
F. del Aguila, hep-ph/9404323



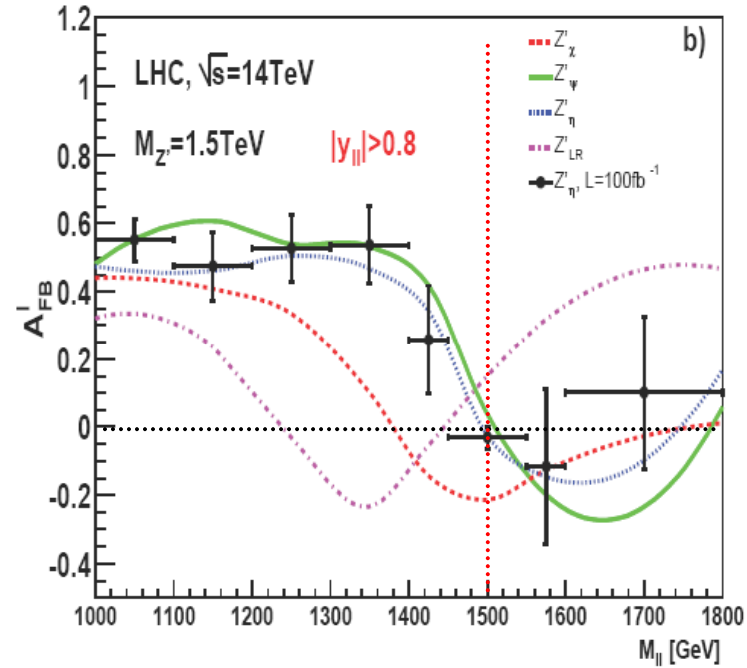
**Statistics not a problem: ILC helps with precision or investigation of some rare couplings**

# Observables

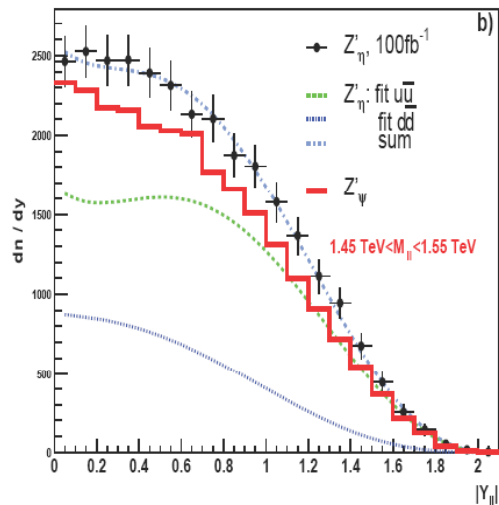
Dilepton invariant mass spectrum



Forward backward asymmetry measurement



Rapidity distribution



$$\frac{d\sigma}{d\cos\theta^*} \propto \frac{3}{8} (1 + \cos^2\theta^*) + A_{FB} \cos\theta^*$$

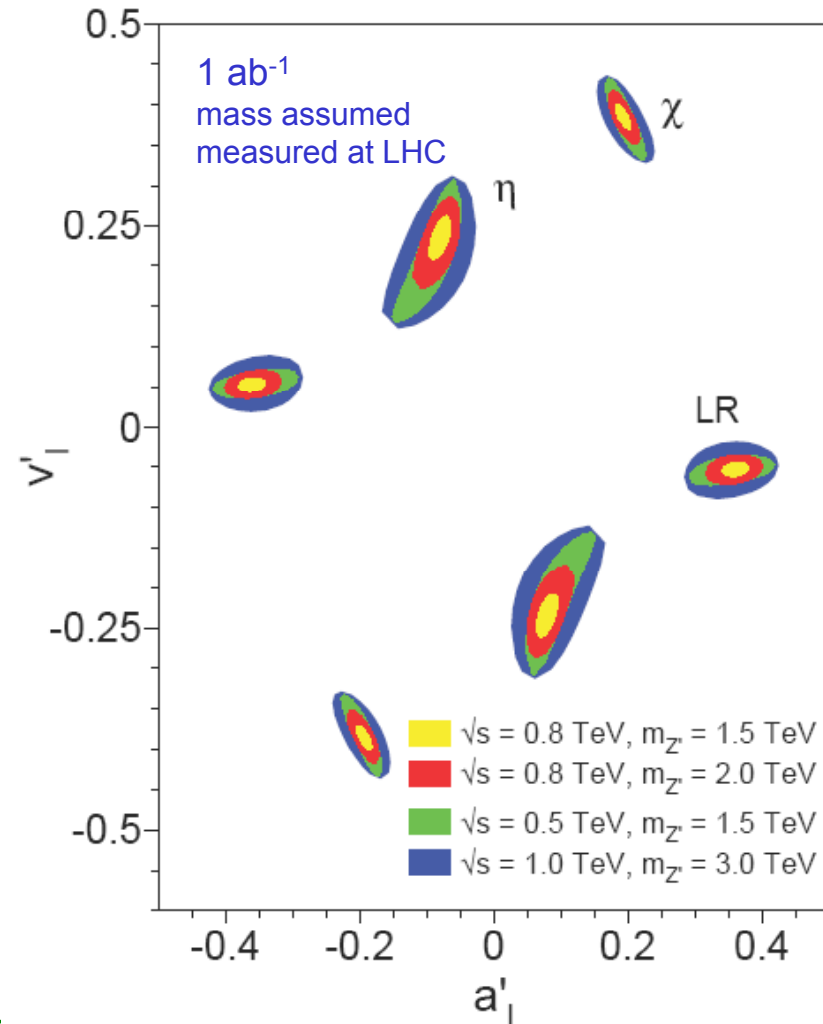
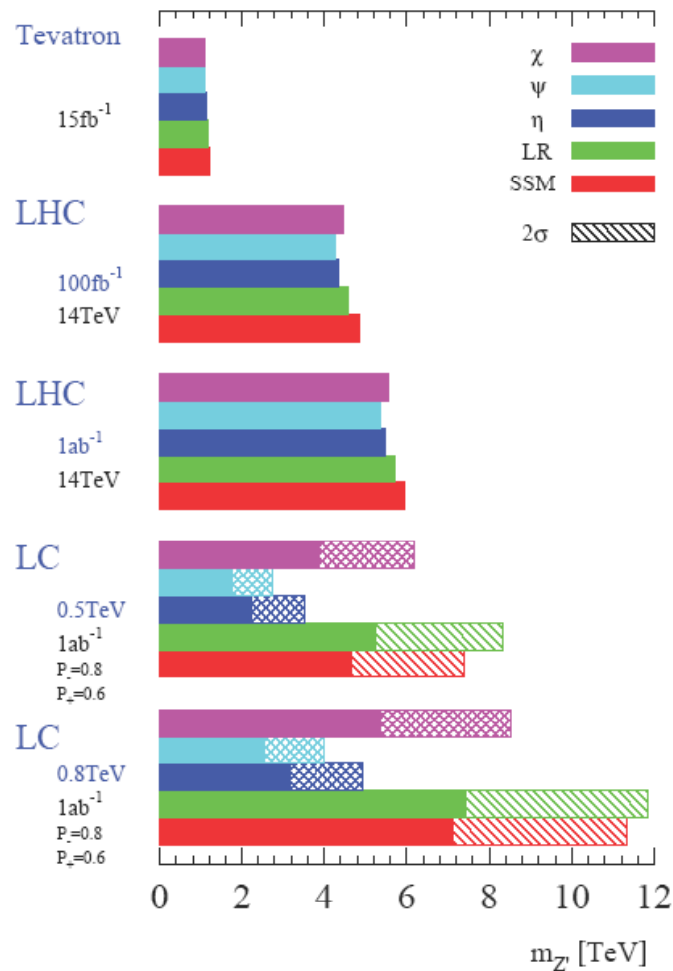
Also,  $Z'$  rapidity distribution depends on  $u/d$  couplings

M Dittmar, A Djouadi, A-S Nicollerat, Phys.Lett. B583 (2004) 111



# Z' at the ILC

At the ILC, sensitivity  $s$  from Z'-Z interference  
 - can also be sensitive to high masses (contact-like interaction)



LHC/ILC Interplay, G. Weiglein et al, hep-ph/041036

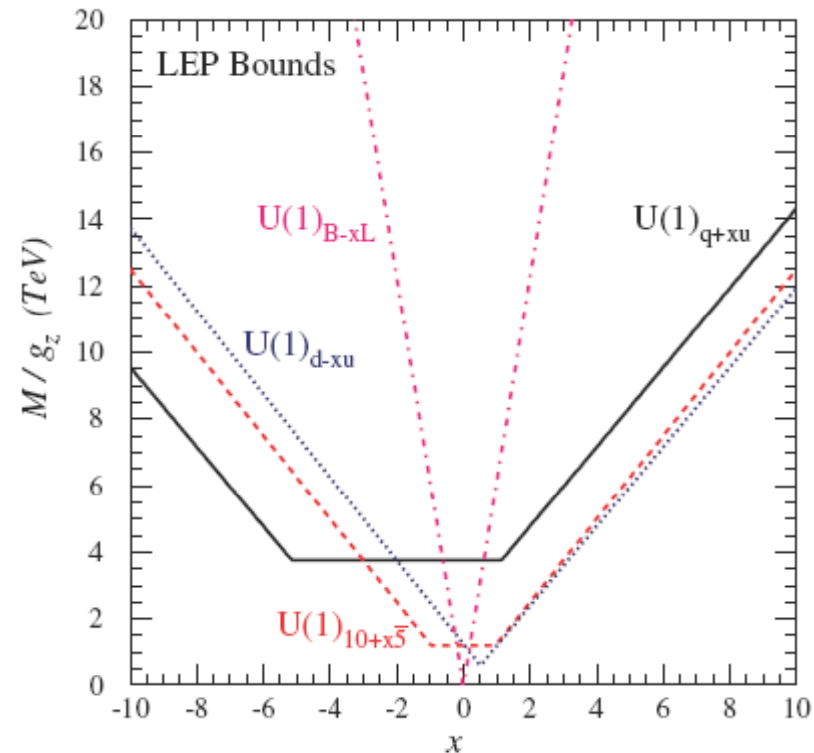
# CDDT parametrization

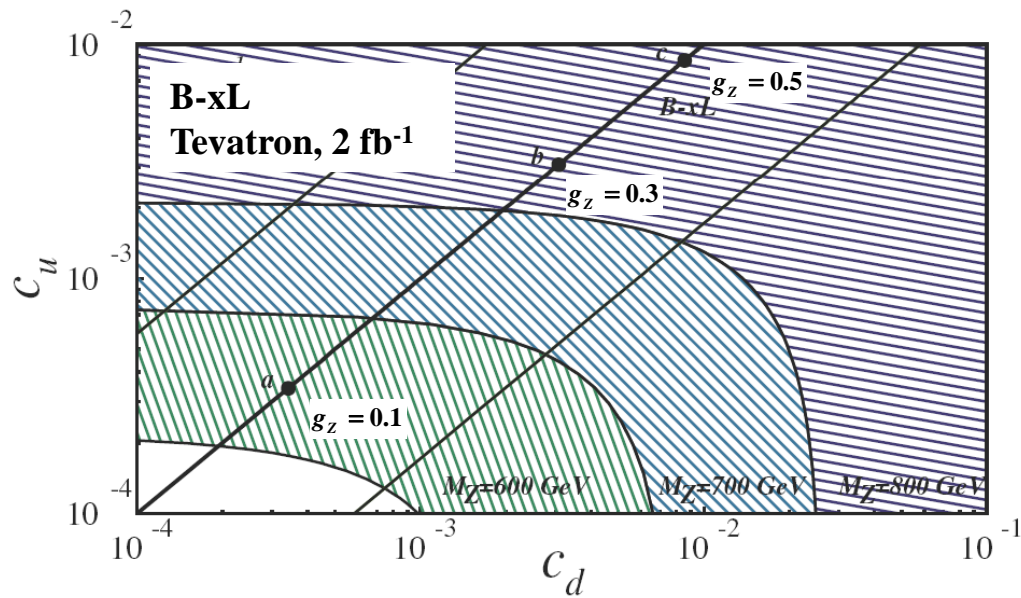
M Carena, A. Daleo, B. Dobrescu and T. Tait, PR D70 (2004) 093009

- very general model with 2 Higgs doublets
- 15 fermion couplings  $z_f$ :  $f = e_R^j, l_L^j, u_R^j, d_R^j, q_L^j$   $j = 1, 2, 3$
- couplings to quarks are generation-independent  $U(1)_z$  charges
  - $z_u, z_d, z_q$  and  $z_{l_j}, z_{e_j}, j = 1-3$
  - anomaly cancellations and possible new fermions must be taken into account
    - “realistic models” for Tevatron

$$U(1)_{B-xL}, \bar{U}(1)_{q-xu}, U(1)_{d-xu}, U(1)_{10-x\bar{5}}$$

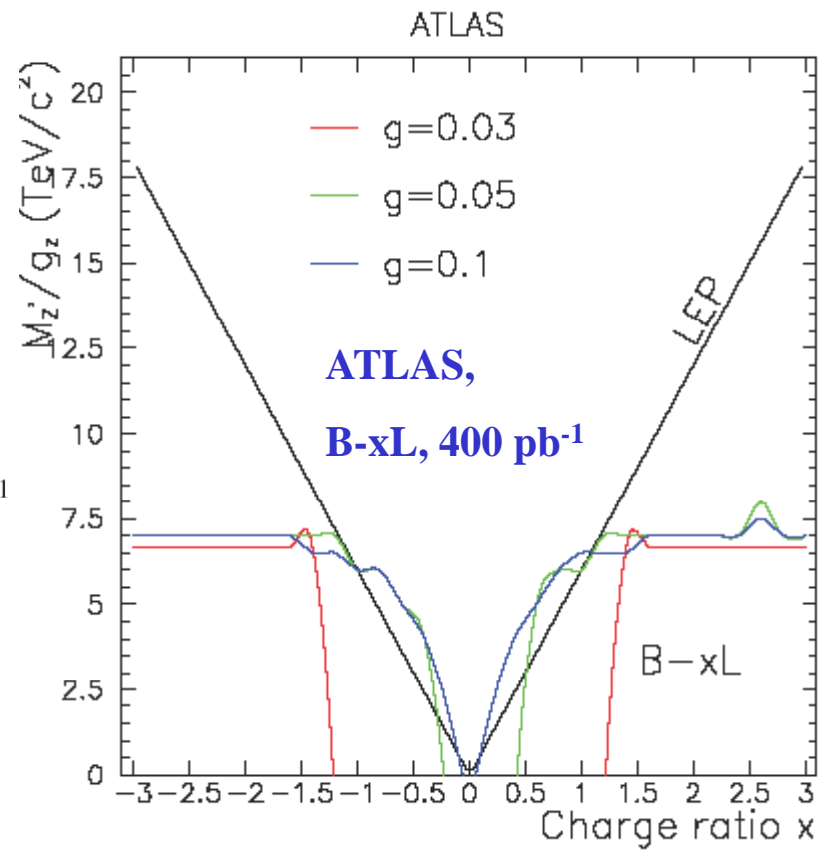
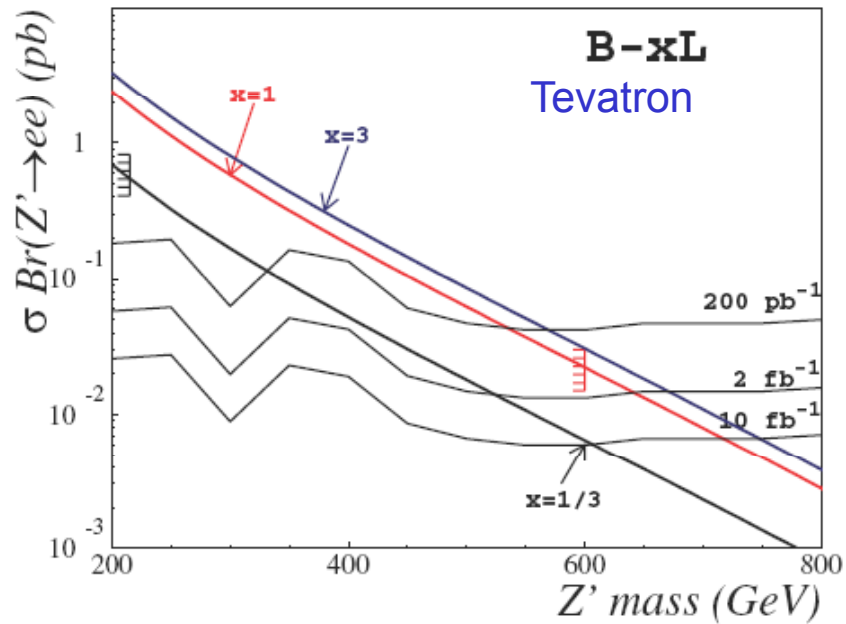
LEP bounds from contact interactions:





$$c_{u,d} = g_z^2 (z_q^2 + z_{u,d}^2) \text{Br}(Z' \rightarrow l^+ l^-)$$

$g_z = 0.1$



Ledroit et al., hep-ph/0703262

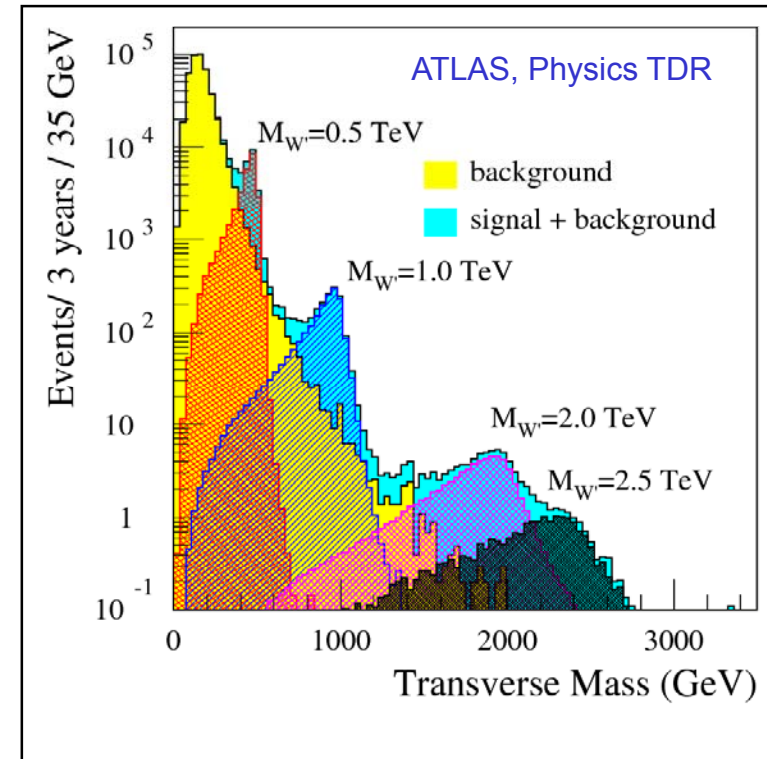
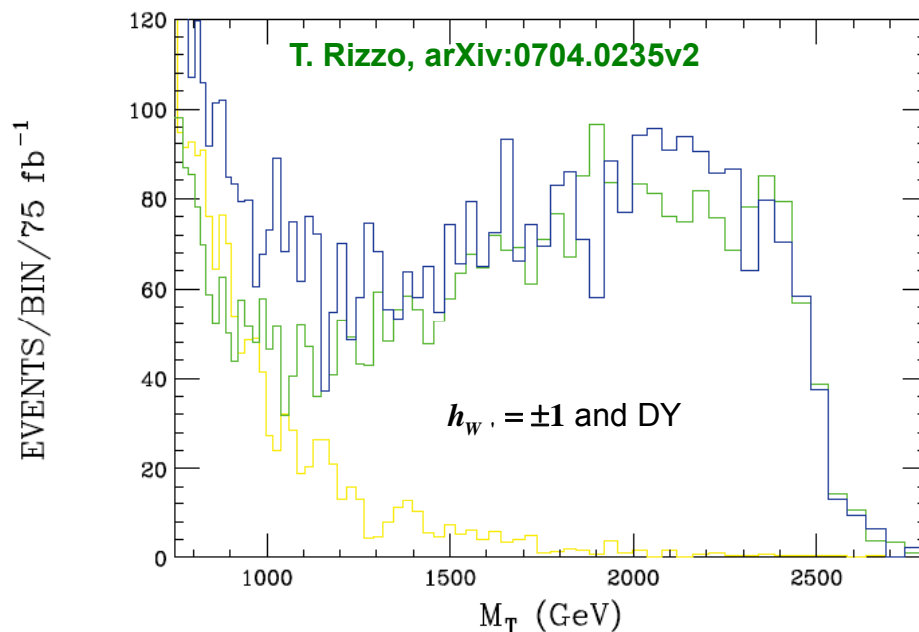
➤ **W' mass difficult to reconstruct:**

- only transverse mass observable

➤ **A<sub>FB</sub> not easily measurable**

(assume light RH neutrino)

- $\nu$  direction unknown
- $m_T$  is max when  $\cos(\theta) = 0$
- can be measured in interference region
- possibly also in  $W' \rightarrow t b$ , with  $t$  polarization



W' couplings more easily measurable at ILC, where cm is known, but is there enough mass? (produced in pairs?)

# Z', W' in LRSM

$$pp \rightarrow Z_R \rightarrow N_\ell N_\ell \rightarrow \ell jj \ell jj$$

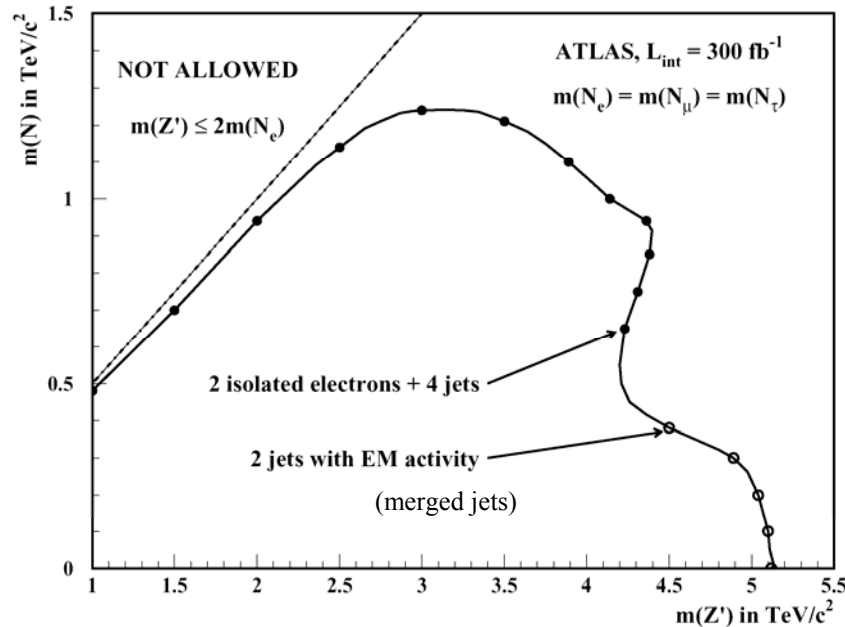
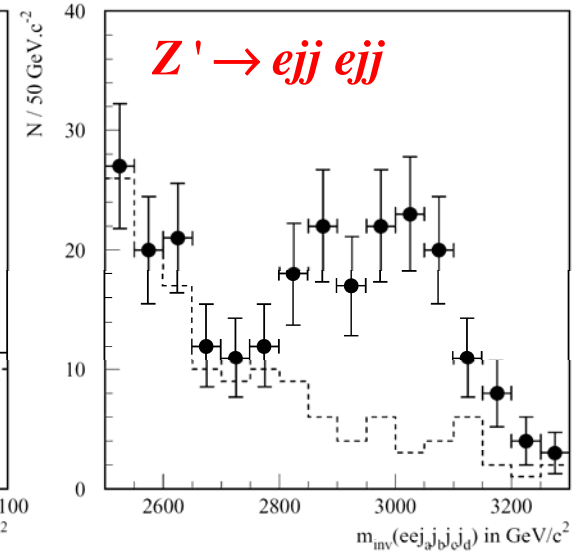
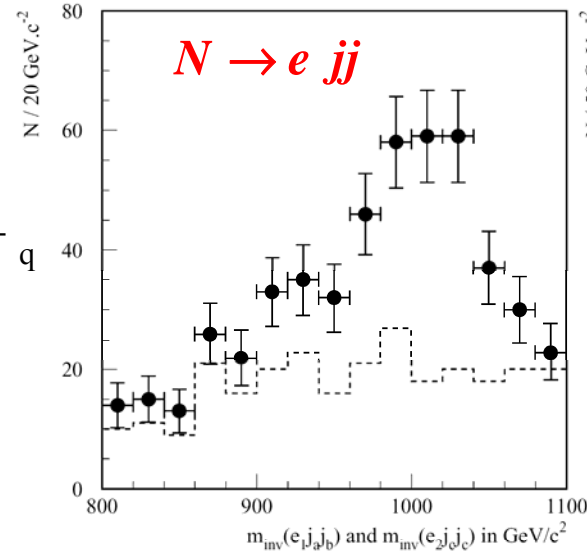
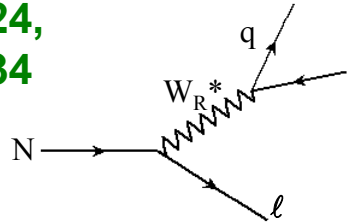
J. Collot, A. Ferrari  
ATL-PHYS-98-124,  
ATL-PHYS-99-034

backgrounds:

*t tbar*

*DY, WW, ZW, ZZ*

*LRSM bckg: W<sub>R</sub>...*



FB asymmetry gives a  
measure of  $\kappa = g_R/g_L$

$$m_{Z'} = \sqrt{\frac{2\kappa^2 \cot^2 \theta_W}{\kappa^2 \cot^2 \theta_W - 1}} m_{W_R}$$

$= 1.7 m_{W_R}$  if  $\kappa=1$

# Z', W'

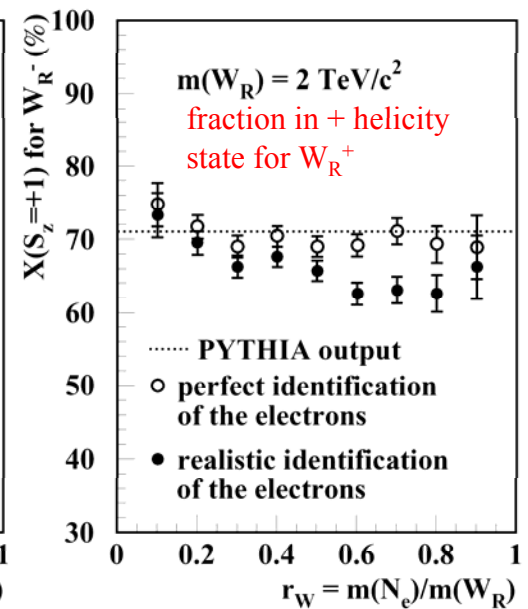
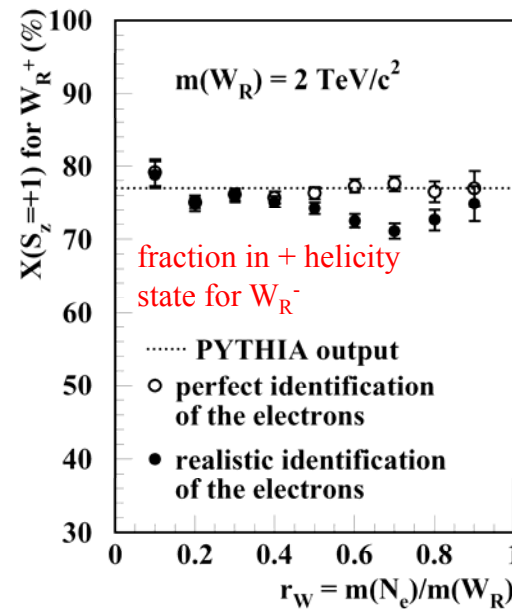
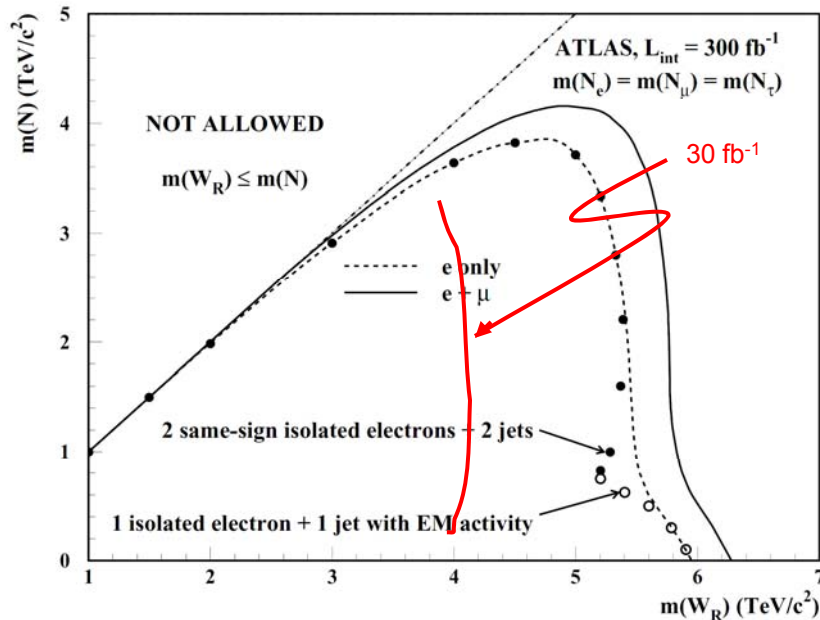
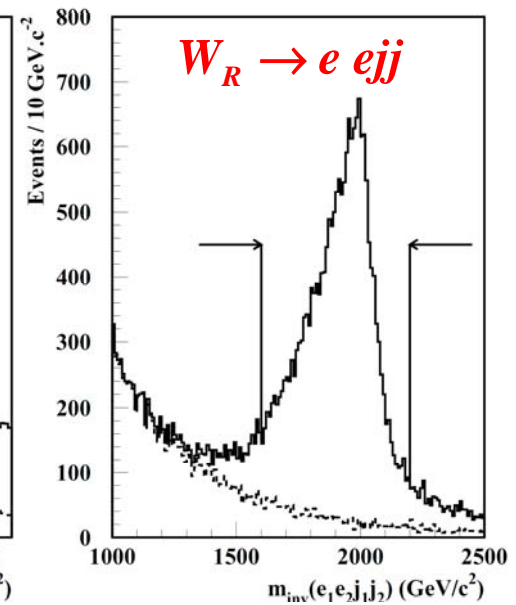
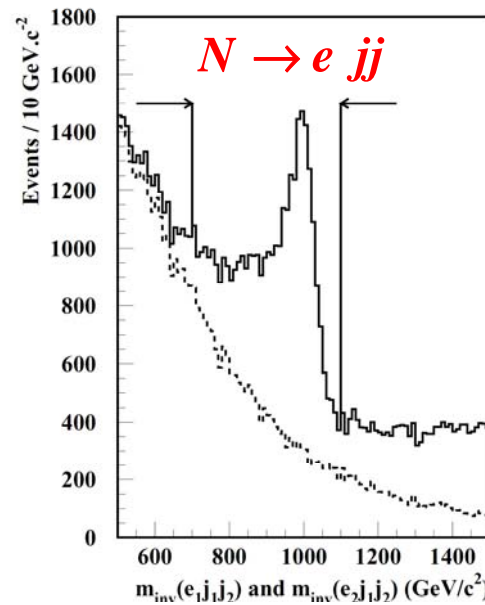
$$pp \rightarrow W_R \rightarrow \ell N_\ell \rightarrow \ell \ell jj$$

J. Collot, A. Ferrari  
ATL-PHYS-98-124,  
ATL-PHYS-99-018

backgrounds:

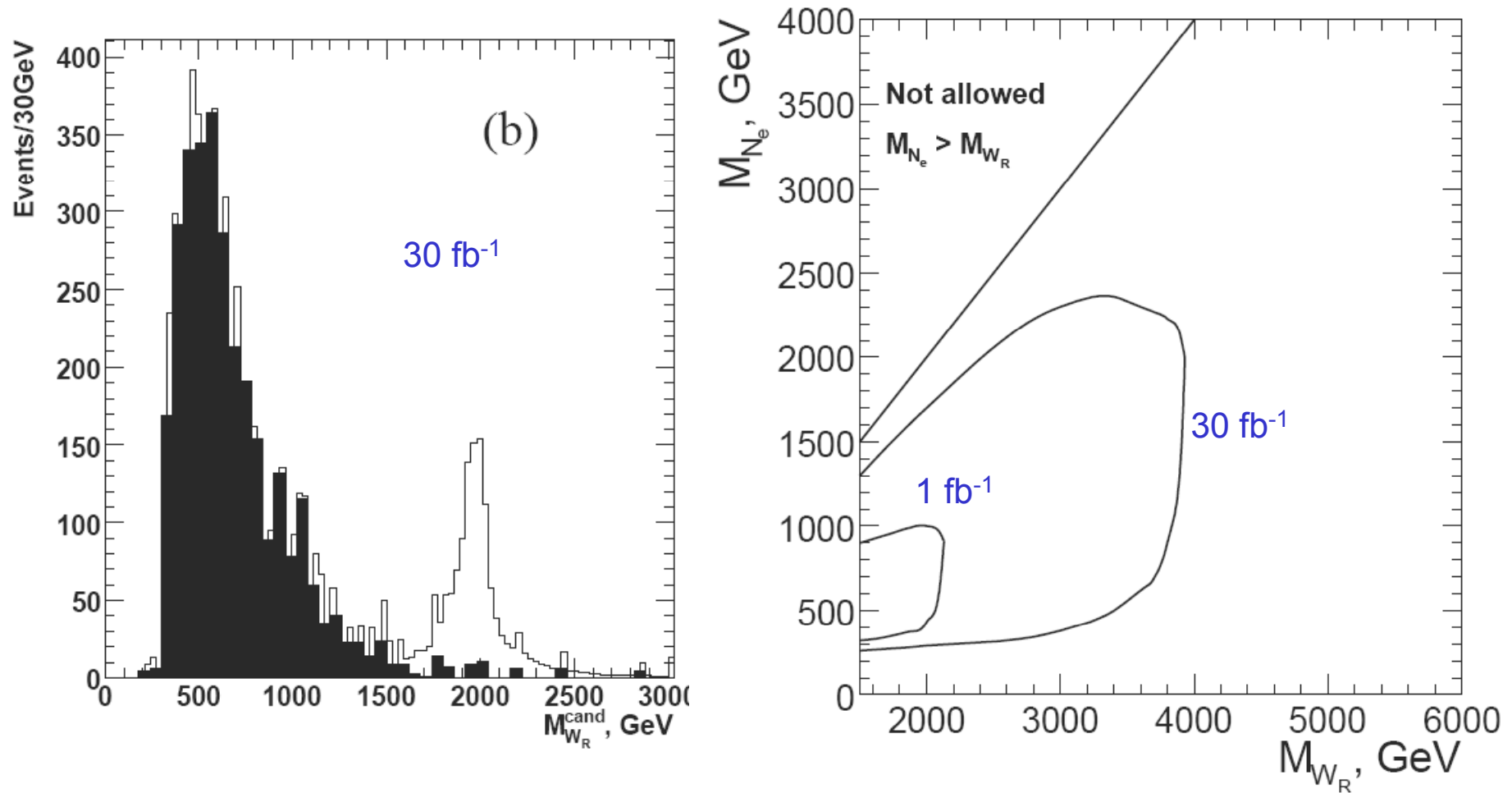
*t tbar*

*DY, WW, ZW, ZZ*



# CMS result on $W_R'$

Full GEANT detector simulation and reconstruction



CMS Physics TDR, CERN/LHCC 2006-021

# vector bi-leptons

## 3-3-1 model: $SU(3)_C \times SU(3)_L \times U(1)_X$

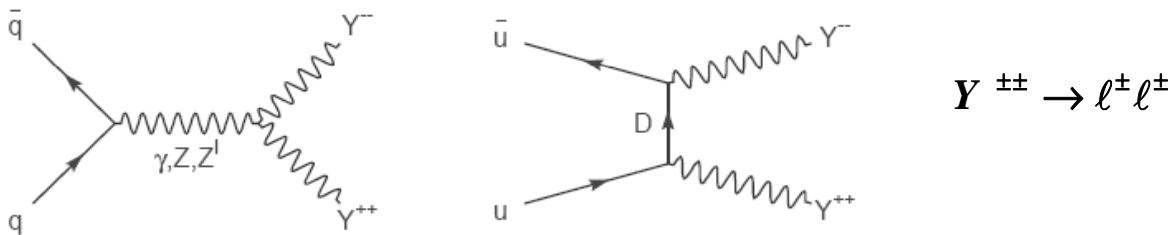
- explains 3 families: anomalies cancel, taking all 3 generations together

fermions:

$$\begin{pmatrix} e^- \\ \nu_e \\ e^c \\ u \\ d \\ D \end{pmatrix}_L \quad \begin{pmatrix} \mu^- \\ \nu_\mu \\ \mu^c \\ c \\ s \\ S \end{pmatrix}_L \quad \begin{pmatrix} \tau^- \\ \nu_\tau \\ \tau^c \\ t \\ b \\ T \end{pmatrix}_L$$

D, S, T have charges +5/3, -4/3, -4/3

new vector gauge bosons:  $Z', W'^{\pm}, Y^{\pm\pm}$



- essentially no background, and detection possible up to  $\sim 1.4$  TeV
- can measure FB asymmetry,  $Z'$  mass

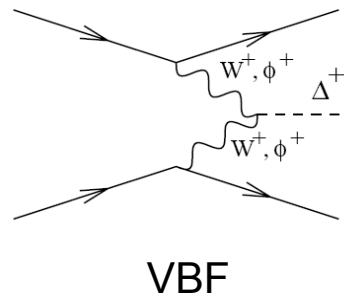
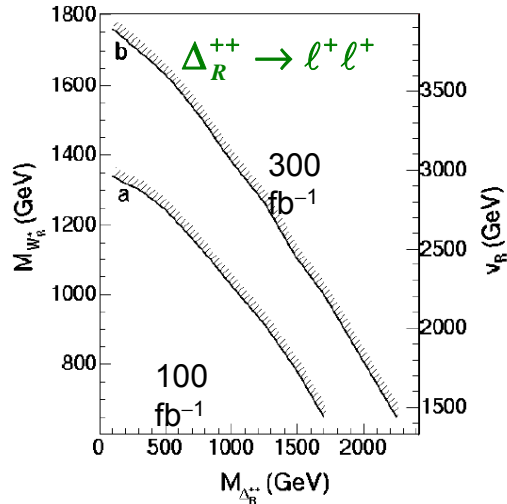
B. Dion et al., Phys.Rev. D59 (1999) 075006,  
 B. Brelier and G.A., ATLAS internal note

- at ILC, different 331 models can be distinguished (Barreto et al, hep-ph/0703099)

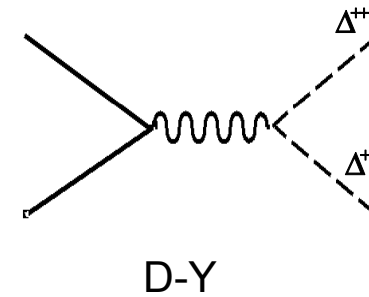


# Scalar bileptons

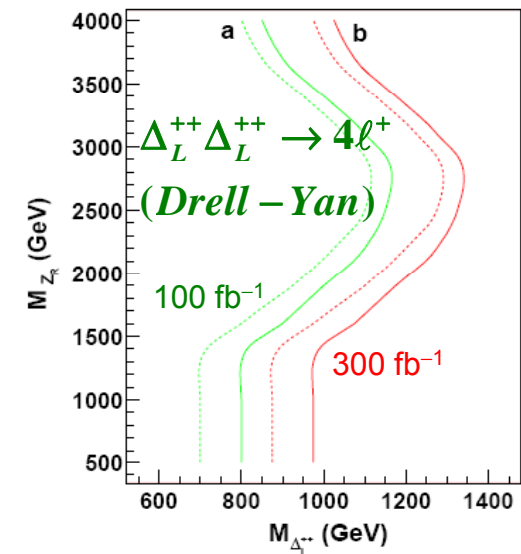
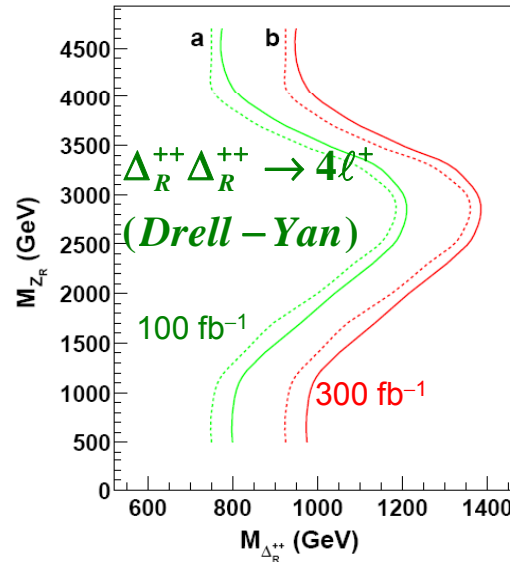
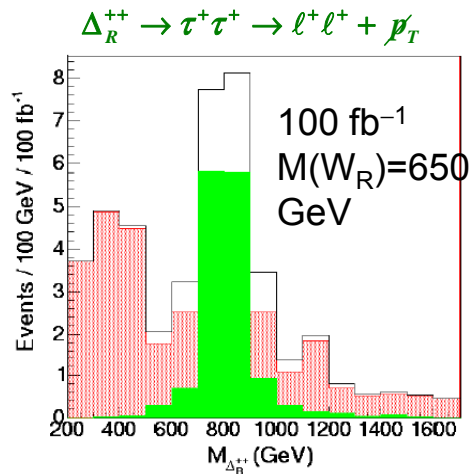
doubly (and singly) charged Higgs in Higgs triplet models  
(as in LR symmetric model)



ATL-PHYS-2004-025



$$m_{W_R}^2 = g_R^2 v_R^2 / 2$$



## The littlest Higgs Model - remarks

- the **small Higgs mass** results from non-exact symmetry  
→ **pseudoGoldstone boson**  
(pions have mass because quark masses and e.m. break chiral symmetry)
- quadratic divergences occur at two-loop level  $\sim 10$  TeV  
→ model is not complete  
UV completion required at  $\sim 10$  TeV
- Low energy EW constraints rather severe
  - FCNC's at  $\sim 100$  TeV
- New particle content

$$W_H^\pm, Z_H, \gamma_H : \sim 1 \text{ TeV}$$

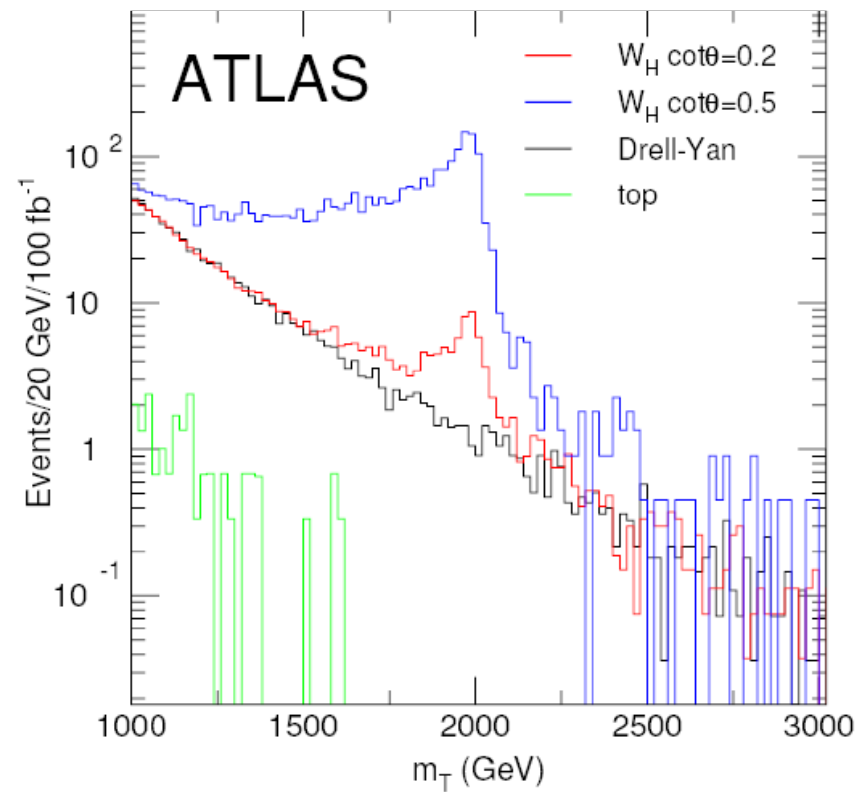
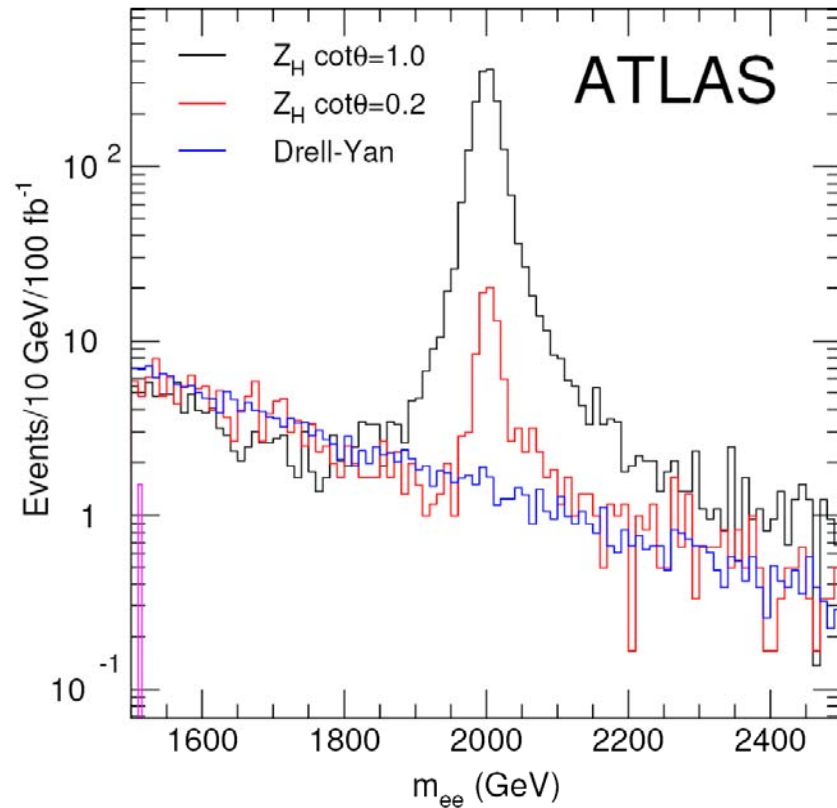
$$T : \sim 1 \text{ TeV}$$

$$\phi^{\pm\pm}, \phi^\pm, \phi^0 : \gtrsim \text{TeV}$$

# Heavy Gauge Bosons: $Z_H$

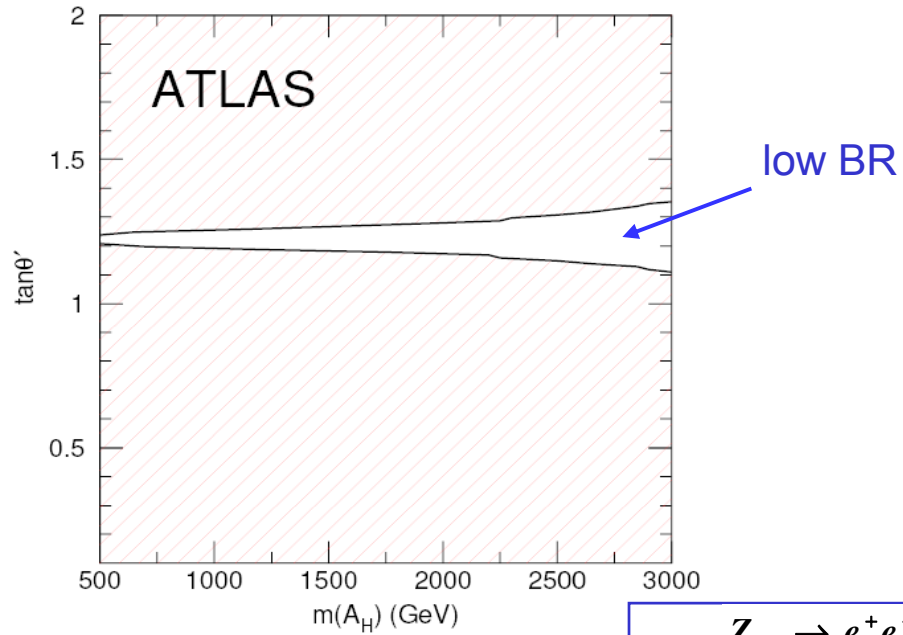
$W_H, Z_H, A_H$  arise from  $[SU(2) \otimes U(1)]^2$  symmetry

→ 2 mixing angles (like  $\theta_W$ ):  $\theta$  for  $Z_H$   
 $\theta'$  for  $A_H$

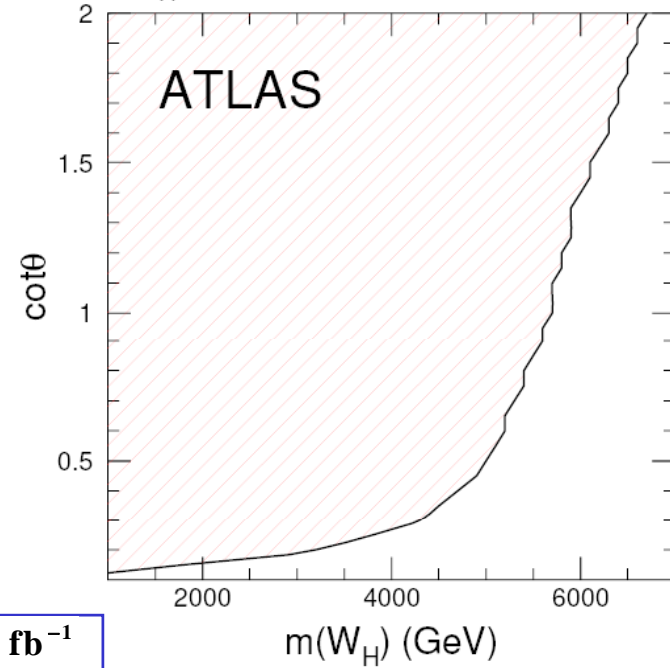


# Heavy Gauge Bosons: $Z_H, \gamma_H$

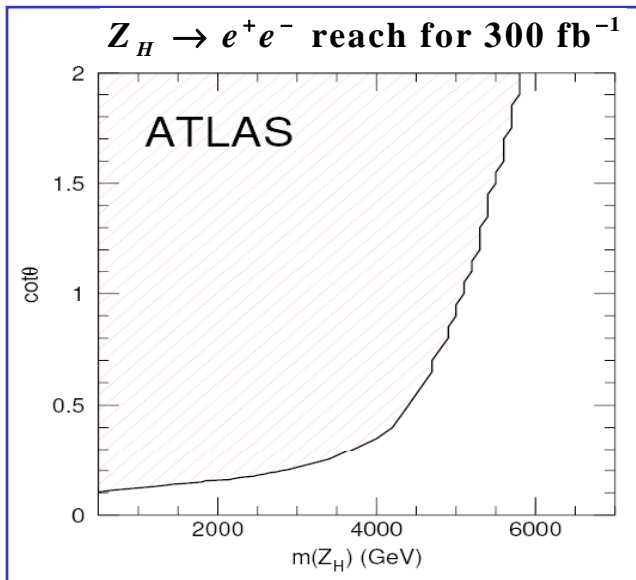
$A_H \rightarrow e^+e^-$  reach for  $300 \text{ fb}^{-1}$



$W_H \rightarrow e\nu$   $5\sigma$  reach for  $300 \text{ fb}^{-1}$

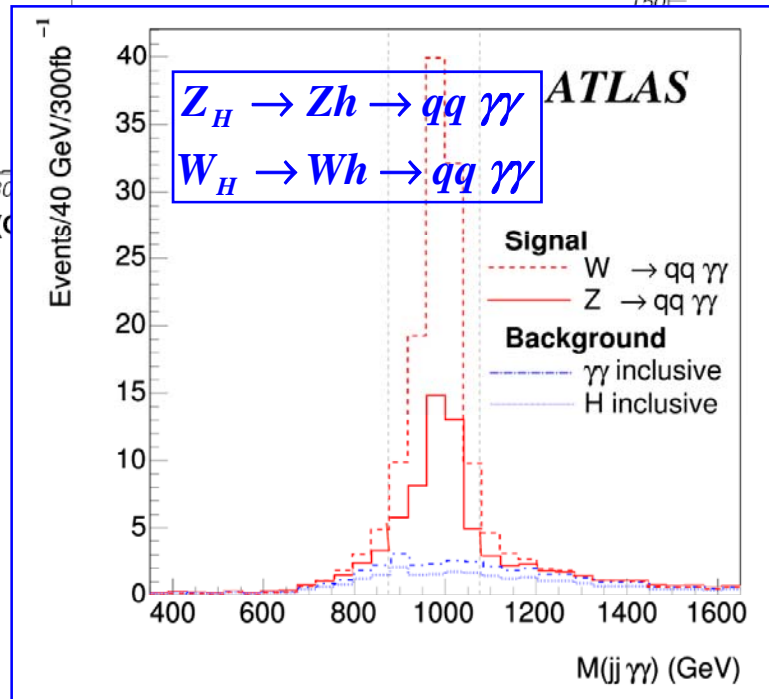
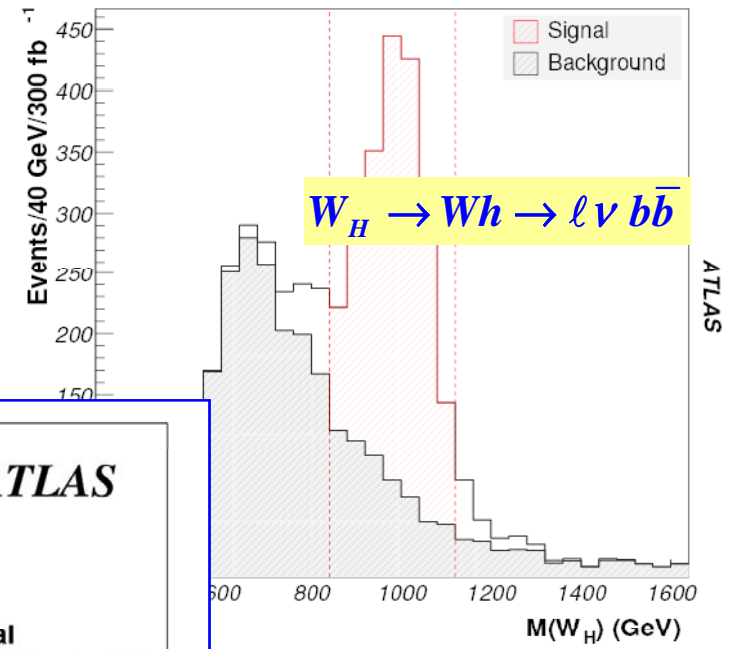
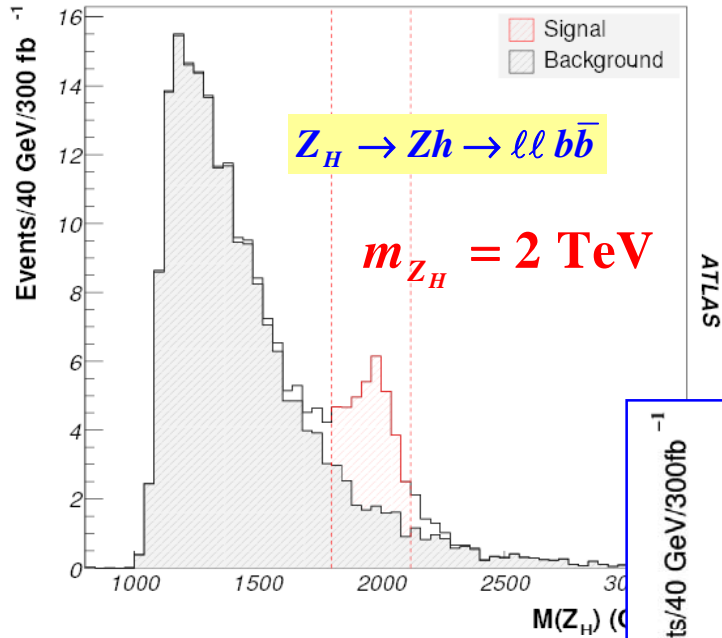


$Z_H \rightarrow e^+e^-$  reach for  $300 \text{ fb}^{-1}$



# Higgs-Gauge boson couplings

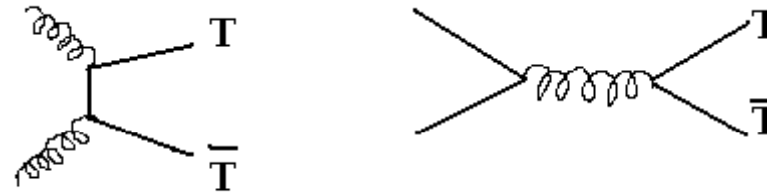
## Measurement of $Z_H Z h$ and $W_H W h$ couplings needed to test model



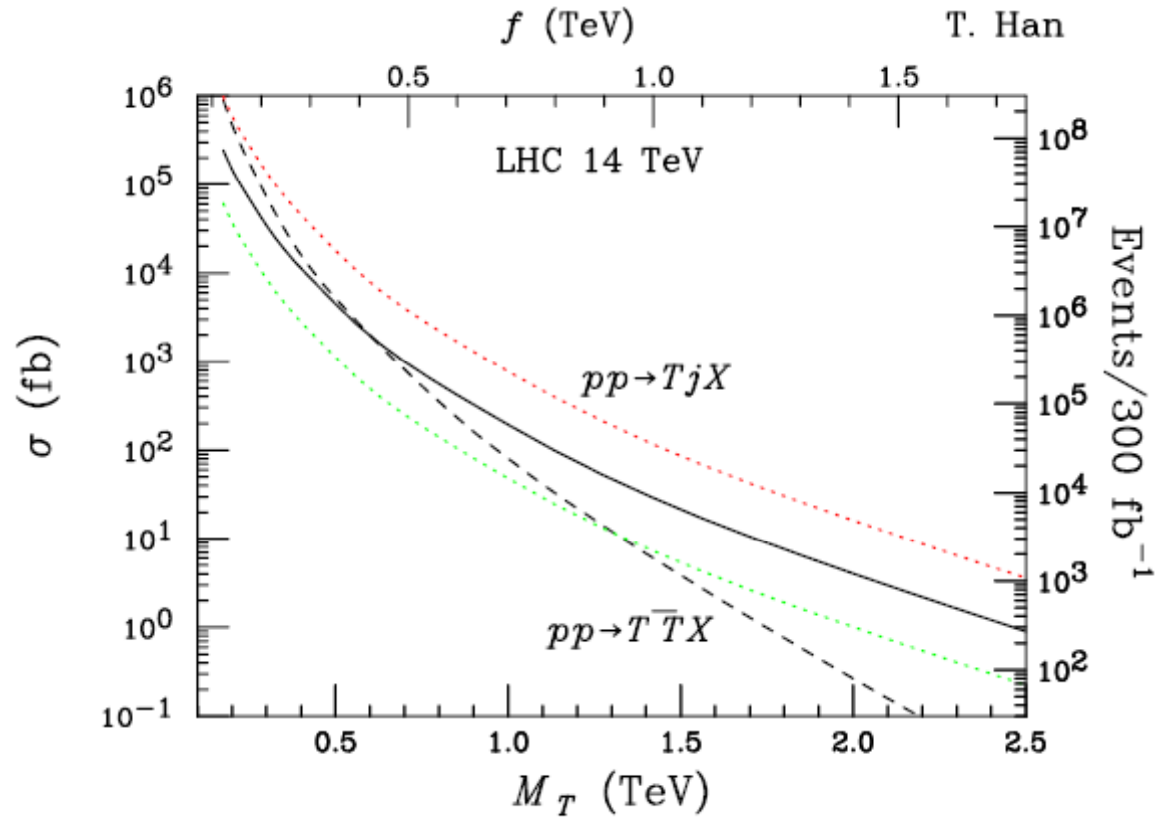
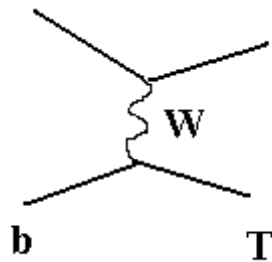
b-tagging  
at high energy

# Search for the heavy T quark

Pair production



Single production:



# Search for the heavy T quark

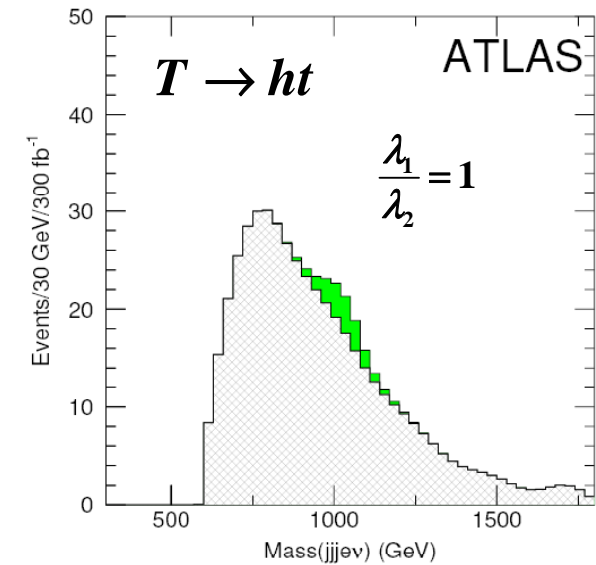
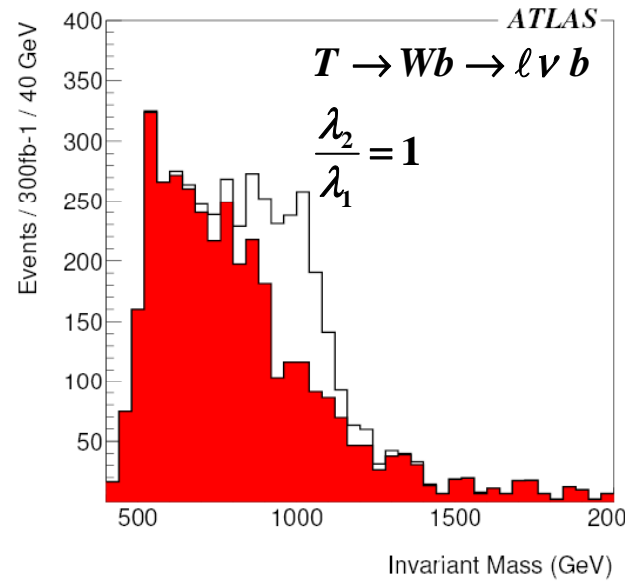
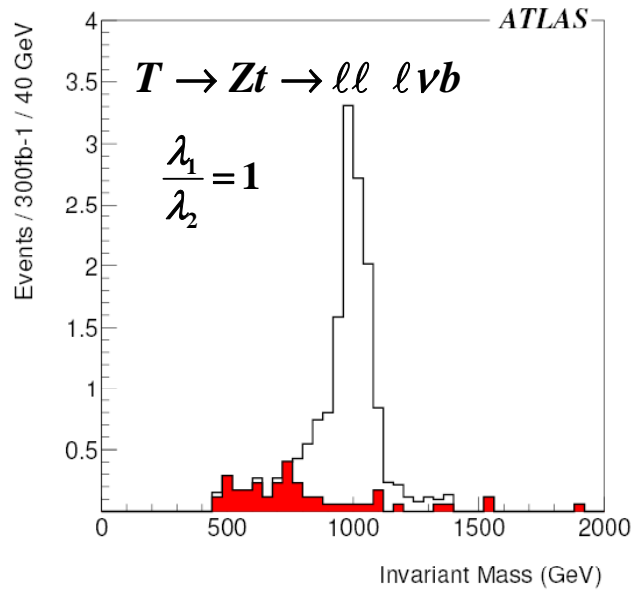
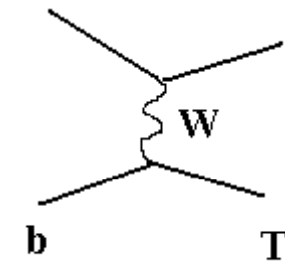
**Couplings:**  $\lambda_1(iQht_r + fT_L t_r hh^\dagger) + \lambda_2 f(T_L T_R)$

→ 3 parameters:  $m_t$ ,  $m_T$ , and  $\lambda_1/\lambda_2$

**Widths:**  $\Gamma(T \rightarrow th) = \Gamma(T \rightarrow tZ) = \frac{1}{2} \Gamma(T \rightarrow bW) = \frac{\kappa^2}{32\pi} M_T$

$$\kappa = \frac{\lambda_1^2}{\sqrt{\lambda_1^2 + \lambda_2^2}}$$

**Single production:**



[SN-ATLAS-2004-038](#)

[detail](#)

# E<sub>6</sub> isosinglet quarks

➤ quark sector in E<sub>6</sub> :  $\left[ \begin{pmatrix} u \\ d \end{pmatrix}_L, u_R, d_R, D_L, D_R \right] \times 3 \text{ families}$

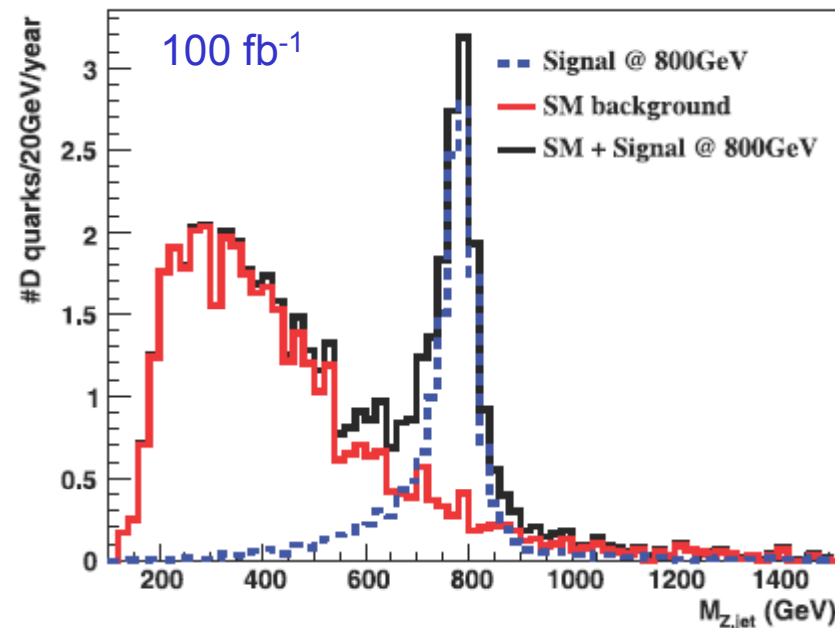
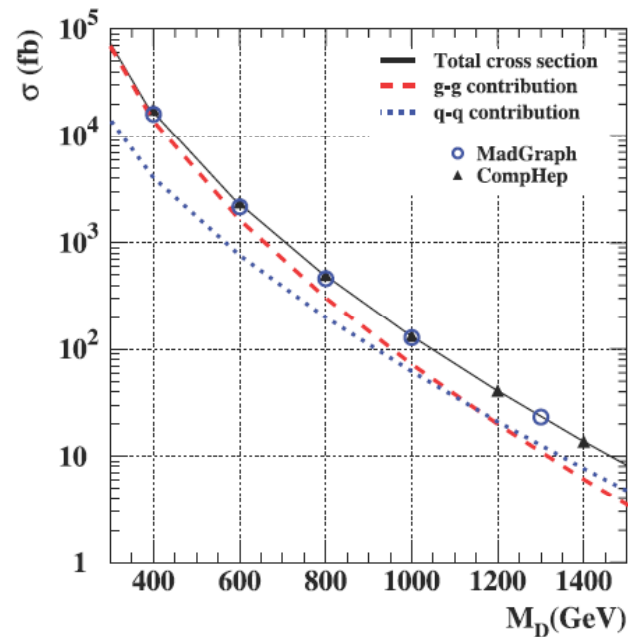
- present limit:  $m_D > 199 \text{ GeV}$  (PDG)
- 3 x 4 CKM matrix → constraint D-u mixing :  $\sin \varphi < 0.07$
- decays:  $D \rightarrow W u$  (50%);  $D \rightarrow Z d$  (25%);  $D \rightarrow d H$  (25%)

➤ **ATLAS study:**  $p p \rightarrow D \bar{D} \rightarrow Z d Z \bar{d} \rightarrow \ell \bar{\ell} d \ell \bar{\ell} \bar{d}$

- process implemented in CompHep and MADGRAPH

○ reach: 920 GeV with 300 fb<sup>-1</sup>

R. Mehdivev et al.. Eur Phys J C49 (2007) 613





- **Possibly different couplings to the 3<sup>rd</sup> family:**
  - LEP constraints weaker
  - topcolor
- **distinguish between scalar (A/H) from vector (Z')**
- **... but difficult to measure**
  - $Z' \rightarrow \tau \tau$  : poor resolution in mass reconstruction
    - method assumes collinear approximation:
      - neutrino in same direction as  $\tau$ , with the  $\tau$  massless
      - 2 x 3-vectors for charged particles, 1 x 2-vector for missing pT
      - 8 input + 2 constraints for  $4 \times 3 - 2 = 10$  deg. of freedom
      - works when  $\tau$ 's are not back to back (not too heavy  $Z'$ )
    - **good reconstruction at ILC, where  $\sqrt{s}$  and  $P_{\text{miss}}$  is known !**
  - $Z' \rightarrow bb$  : huge QCD background
  - $Z' \rightarrow tt$  : too much background from QCD production of top
    - except if resonance is from a strong interaction process
    - **good for ILC, but limited by mass**

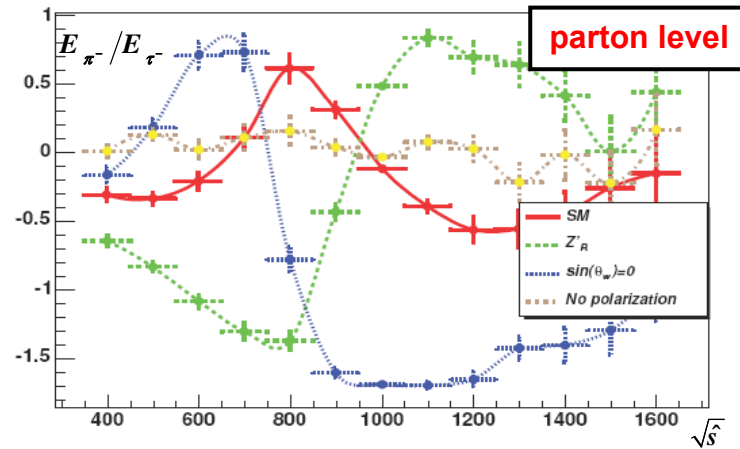
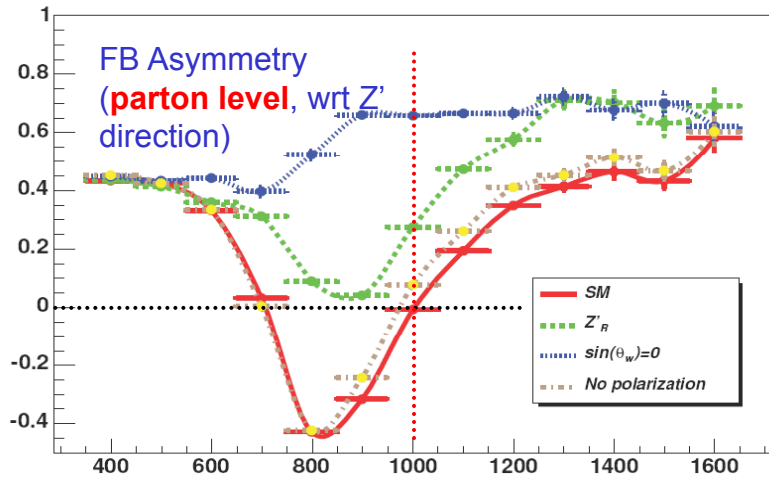
# further observables for 3<sup>rd</sup> generation

- ➔ can measure FB asymmetry wrt to Z' direction
- ➔ possibility to measure polarization or spin correlations through decay of  $\tau$  or  $t$

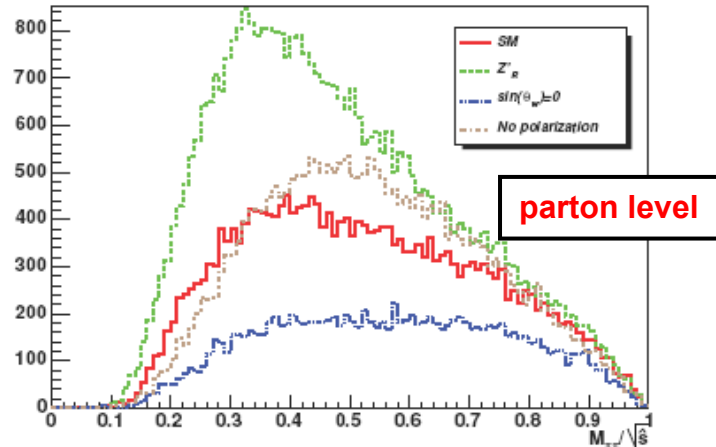
$$Z' \rightarrow \tau \tau \quad m(Z') = 1 \text{ TeV}$$

$$\quad \quad \quad \downarrow$$

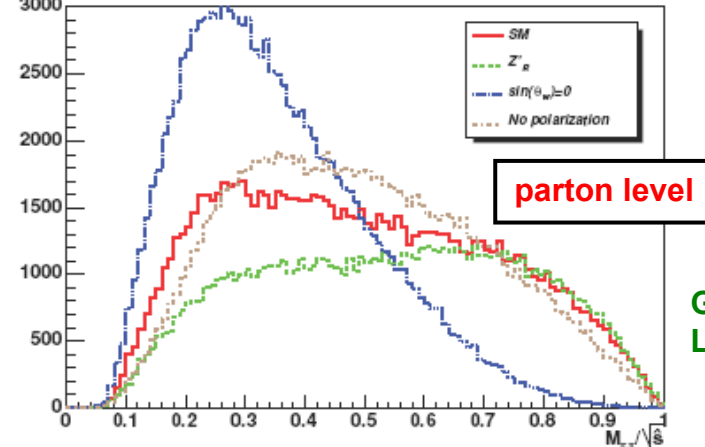
$$\quad \quad \quad \pi \nu$$



$M_{\pi\pi}/\sqrt{\hat{s}}, \sqrt{\hat{s}} < 800 \text{ GeV}$



$M_{\pi\pi}/\sqrt{\hat{s}}, \sqrt{\hat{s}} > 800 \text{ GeV}$

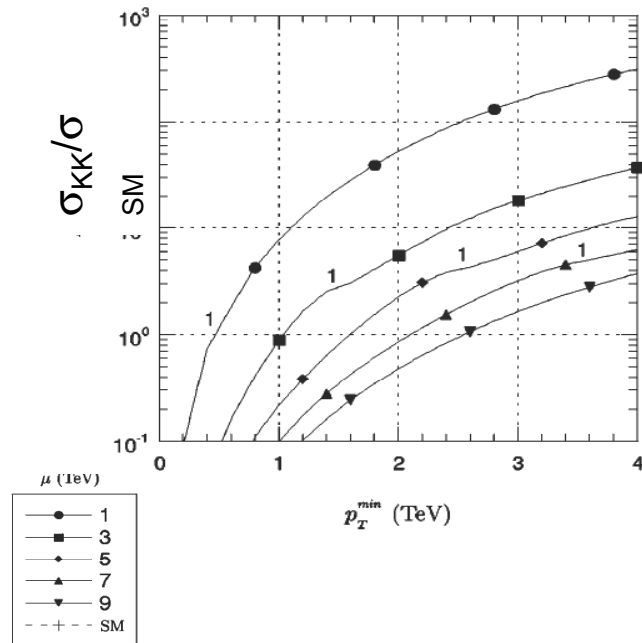


G.A. et al.,  
Les Houches 2005

# KK excitation of the gluon

## UED scenario

D.A. Dicus, C.D. McMullen and S. Nandi,  
PR D65 (2002) 076007



## top resonance from bulk RS KK gluon

- large overlap of KK gluon and top quark wave functions because both are localized towards TeV brane

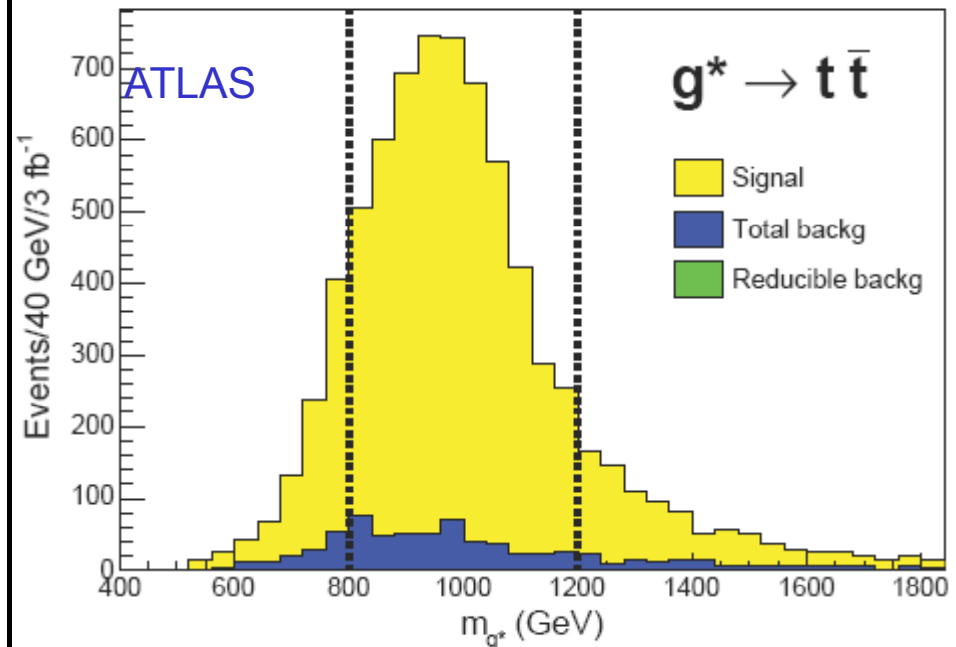
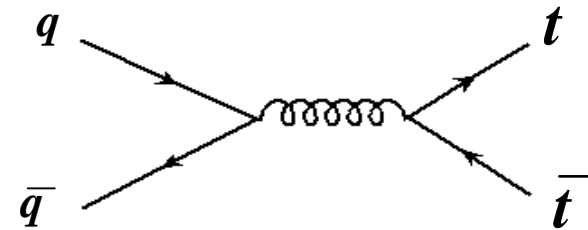
B. Lillie, L. Randall and L-T Wang, hep-ph/0701166

- can also measure spin correlations

- graviton resonance to top pairs (or WW)

Agashe et al., hep-ph/0701186

For heavy quark production, one diagram dominates:



March, L; Ros, E; Salvachúa, B; ATL-PHYS-PUB-2006-002

# Jet splitting

- ➔ **Highly boosted di-jet looks like one jet**
- **use algorithm to see if jet splits into two when narrower cone is used**

jets with  $p_T > 250$  GeV

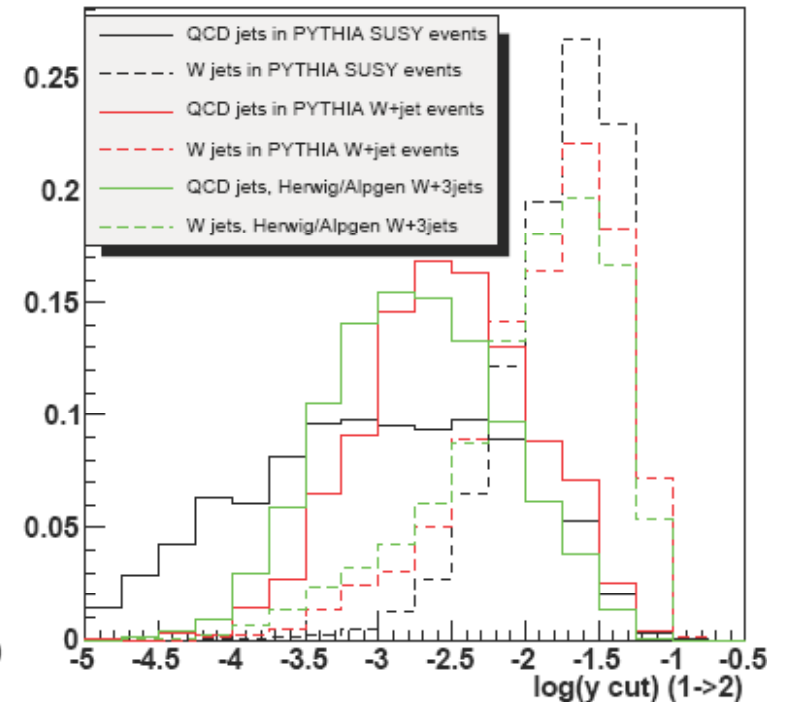
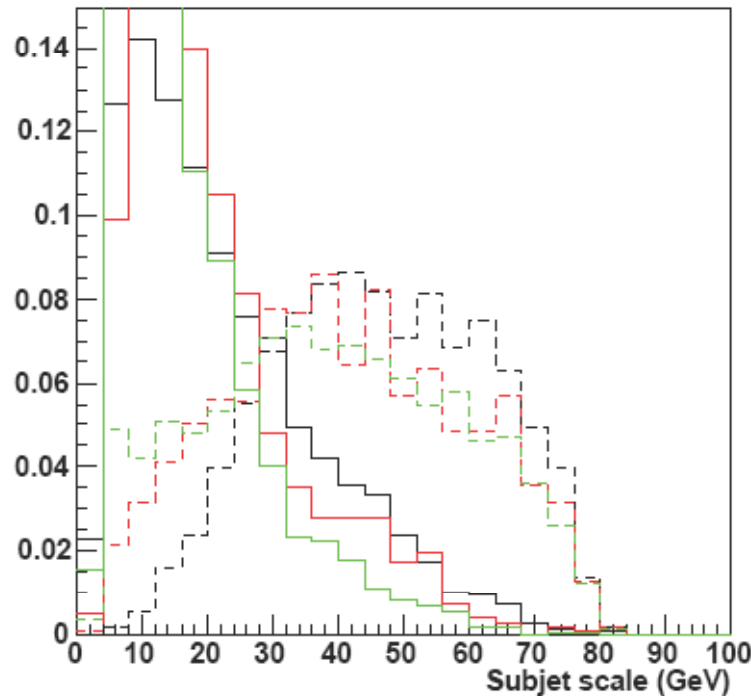
Once a jet is found, apply the inclusive kT algorithm to clusters composing it

$$d_{kl} = \min(p_{T_k}^2, p_{T_l}^2) R_{kl}^2 / R^2$$

$$y \equiv \frac{d_{kl}}{(p_T^{jet})^2}$$

$y_{cut}$  at which jet splits into 2 subjets

$$\text{scale: } p_T^{jet} \sqrt{y} = \sqrt{d_{kl}}$$

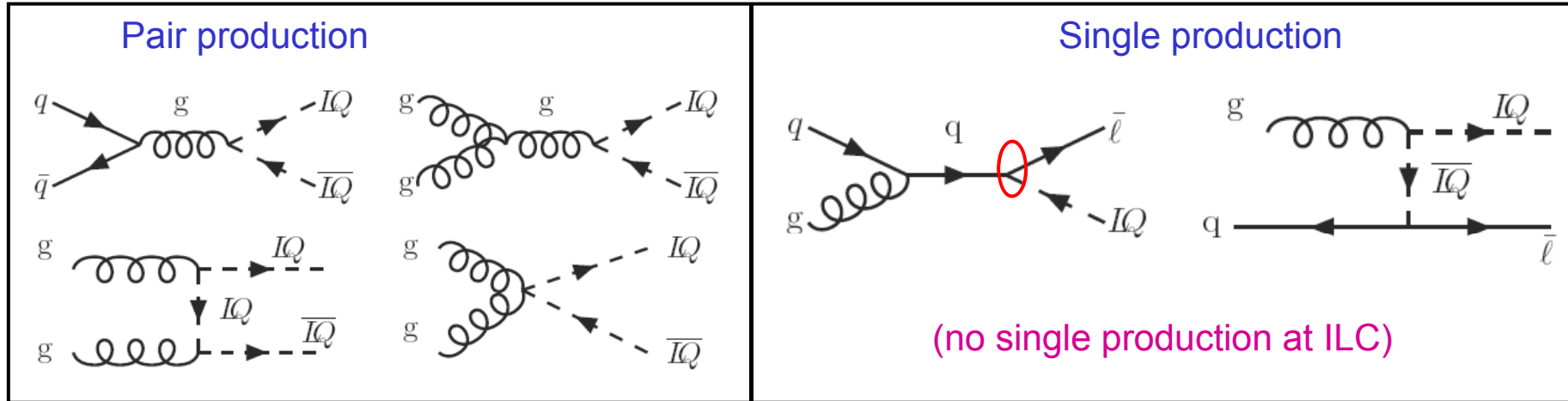


J.M. Butterworth, J.R. Ellis, A.R. Raklev, hep-ph/0702150

# Leptoquarks

implemented in CompHep/CalcHep

A. Belyaev et al., *J. High Energy Phys.* 09 (2005) 005



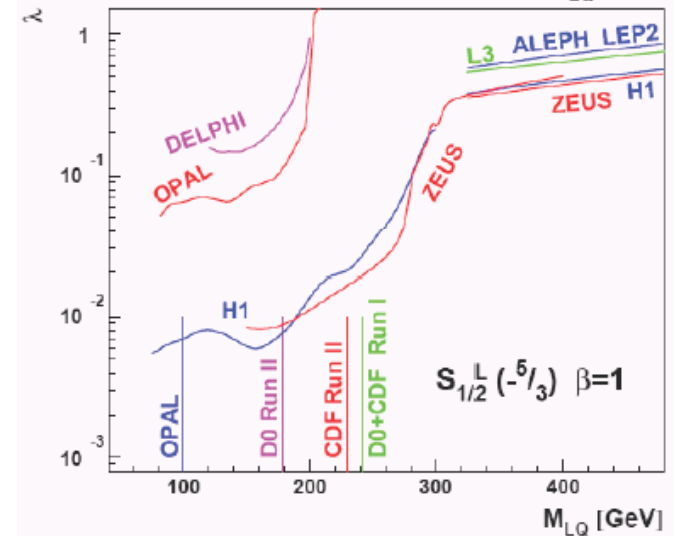
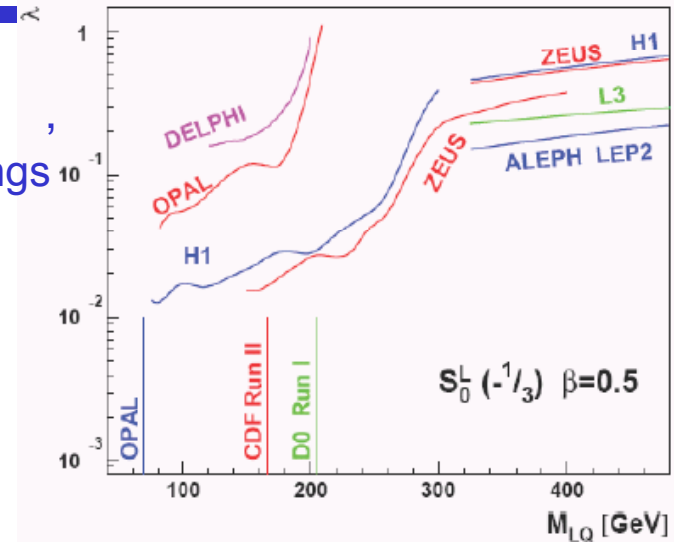
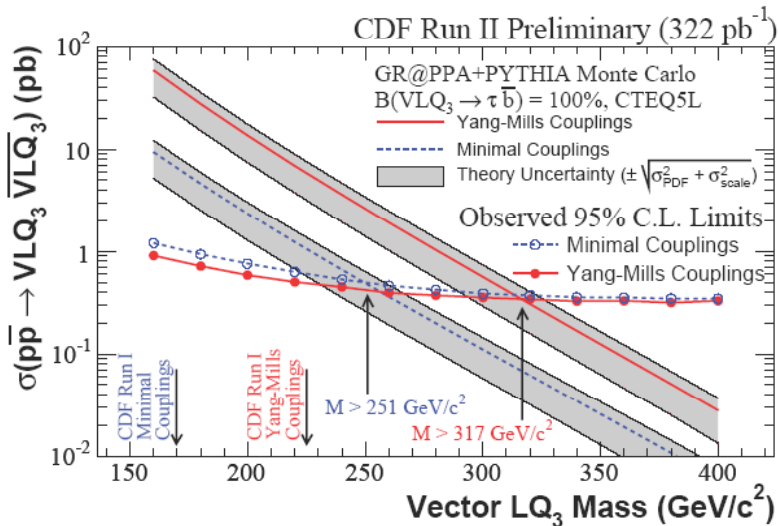
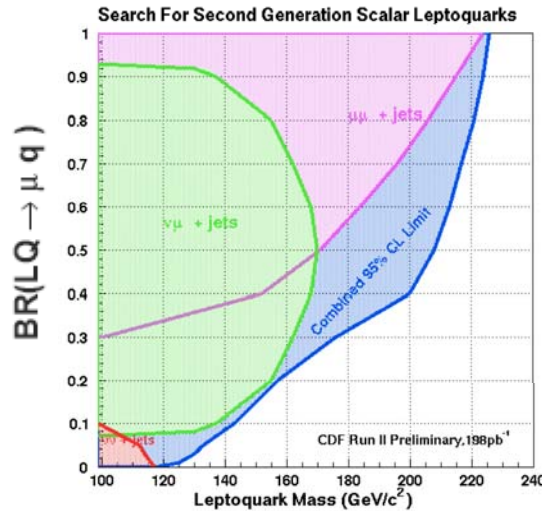
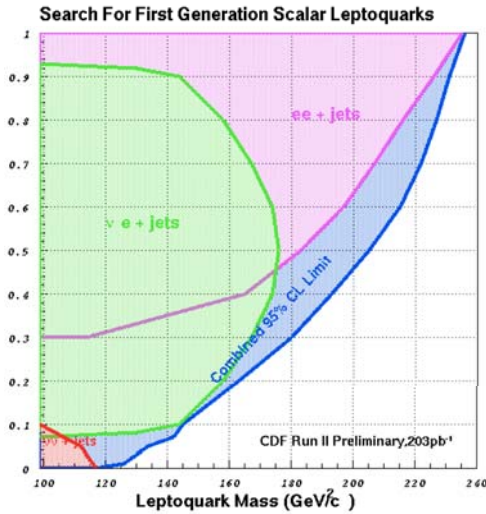
Characterized by: Fermion number: 0 ( $\ell^+q$ ) or 2 ( $\ell^-q$ )  
 spin: scalar or vector ( $g\text{-}\Phi^\mu\text{-}\Phi_\mu$  coupling depends on  $\kappa_g$  and  $\lambda_G$ )  
 isospin: 0,  $\frac{1}{2}$ , 1  
 charge:  $\pm 1/3, \pm 2/3, -4/3, -5/3$

$$\mathcal{L}_{|F|=0}^f = (h_{2L}\bar{u}_R\ell_L + h_{2R}\bar{q}_L i\tau_2 e_R)R_2 + \tilde{h}_{2L}\bar{d}_R\ell_L\tilde{R}_2 + (h_{1L}\bar{q}_L\ell_L + h_{1R}\bar{d}_R\gamma^\mu e_R)U_{1\mu} + \tilde{h}_{1R}\bar{u}_R\gamma^\mu e_R\tilde{U}_{1\mu} + h_{3L}\bar{q}_L\vec{\tau}\gamma^\mu\ell_L\vec{U}_{3\mu} + \text{h.c.},$$

$$\mathcal{L}_{|F|=2}^f = (g_{1L}\bar{q}_L^c i\tau_2\ell_L + g_{1R}\bar{u}_R^c e_R)S_1 + \tilde{g}_{1R}\bar{d}_R^c e_R\tilde{S}_1 + g_{3L}\bar{q}_L^c i\tau_2\vec{\tau}\ell_L\vec{S}_3 + (g_{2L}\bar{d}_R^c\gamma^\mu\ell_L + g_{2R}\bar{q}_L^c\gamma^\mu e_R)V_{2\mu} + \tilde{g}_{2R}\bar{u}_R^c\gamma^\mu\ell_L\tilde{V}_{2\mu} + \text{h.c.},$$

# LQ limits

In effect, only 3 parameters:  $M_{LQ}$ ,  $\lambda_L$  and  $\lambda_R$   
 total width depends on  $\lambda_{eff}^2 = \lambda_L^2(lq) + \lambda_R^2(lq) + \lambda_L^2(vq)$ ,  
 where  $\lambda$  are linear combinations of  $g, h, \tilde{g}, \tilde{h}$  couplings



from CDF:

[http://ncdf70.fnal.gov:8001/lq/LQ\\_comb/Combinations.html](http://ncdf70.fnal.gov:8001/lq/LQ_comb/Combinations.html)

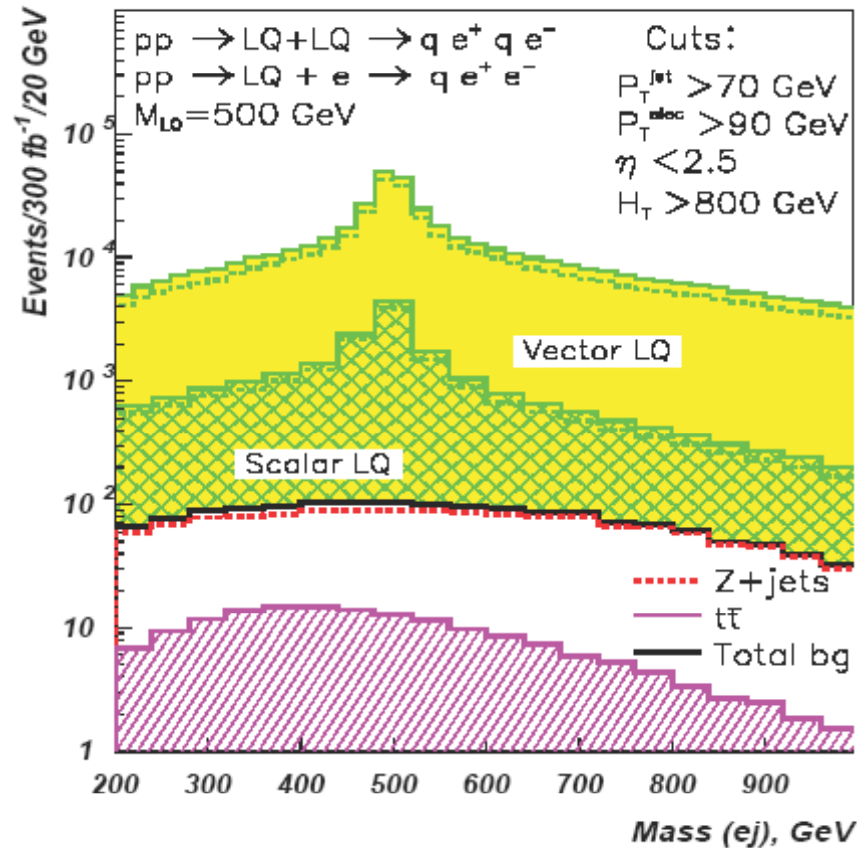
[http://www-cdf.fnal.gov/physics/joint\\_physics/public/ichep/CDF\\_Physics\\_2006.html](http://www-cdf.fnal.gov/physics/joint_physics/public/ichep/CDF_Physics_2006.html)

similar results from D0:

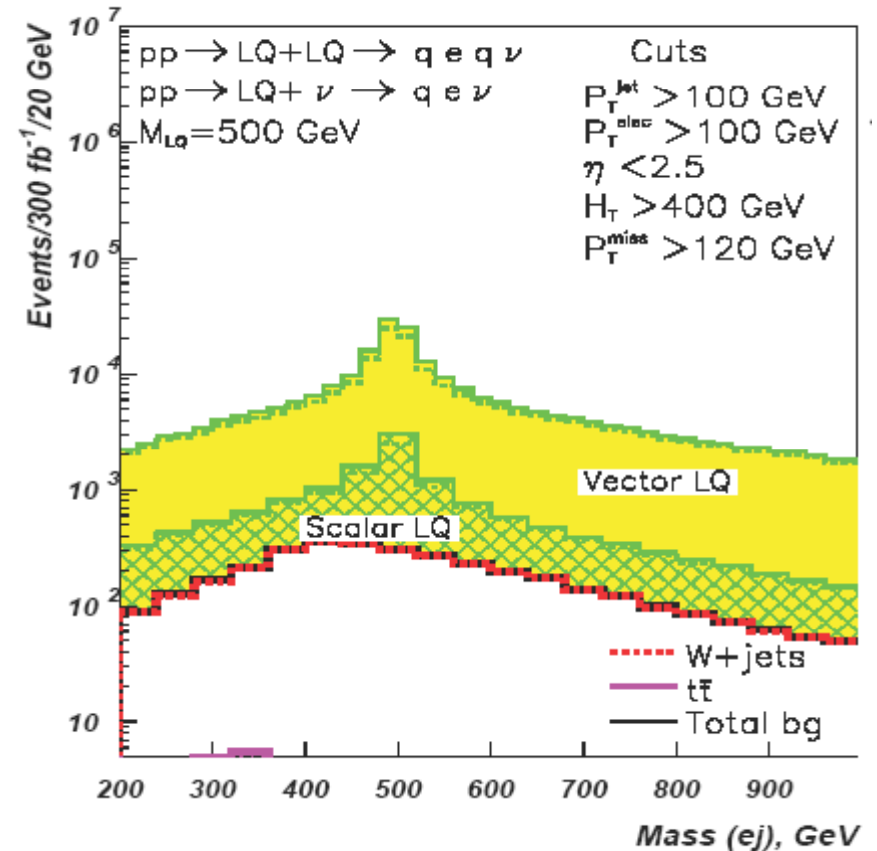
<http://www-d0.fnal.gov/Run2Physics/WWW/results/np.htm>

# ATLAS study

“type 1”: 2 leptons + jets



“type 2”: 1 lepton + jets +  $E_T^{\text{miss}}$



Combine single and pair production

A. Belyaev et al., [J. High Energy Phys. 09 \(2005\) 005](#)

Reach with 10 fb<sup>-1</sup>: ~ 800 and 900 GeV for S (types 1 and 2 respect.)

~ 1000 and 1300 GeV for V (types 1 and 2 respect.)

assuming  $\Lambda_{\text{eff}} \sim \alpha$  and minimal couplings for vector LQs

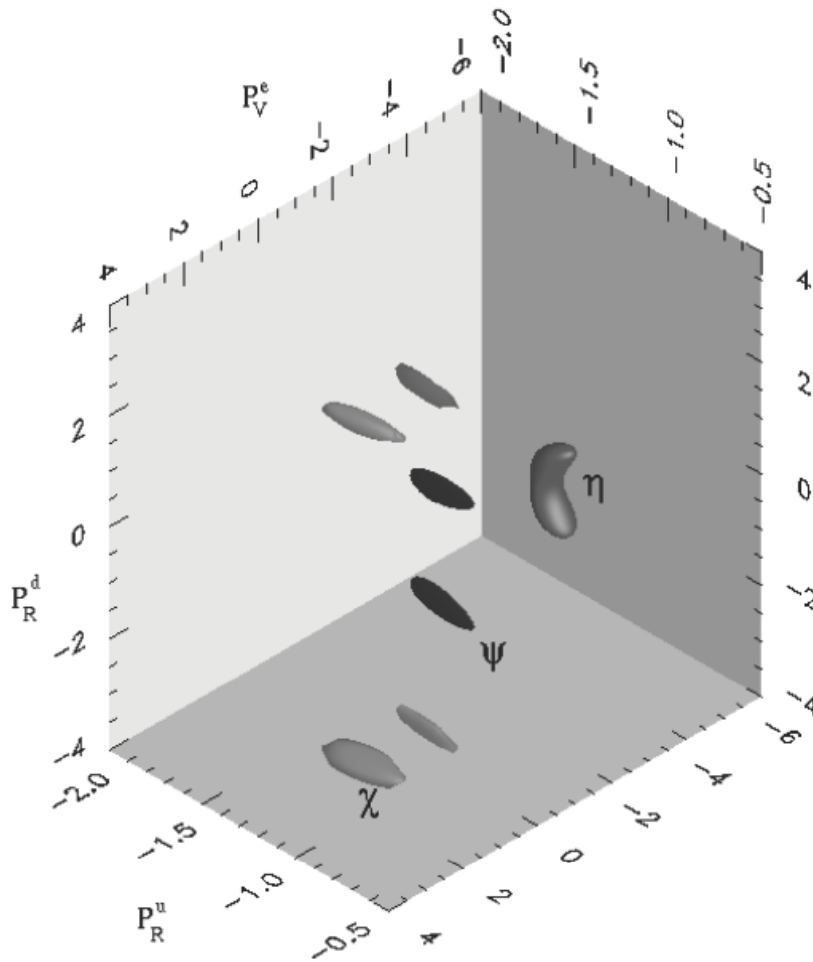
## Conclusions

- **All possible types of 2-body resonances are predicted by new models**
- **In general, exotic physics probes high masses with high luminosity**
  - push for highest energy machine, or at least possibility of upgrade
  
- **but**
  - possibilities of weak couplings buried in background at LHC
  - higher precision attainable at ILC



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## Backups

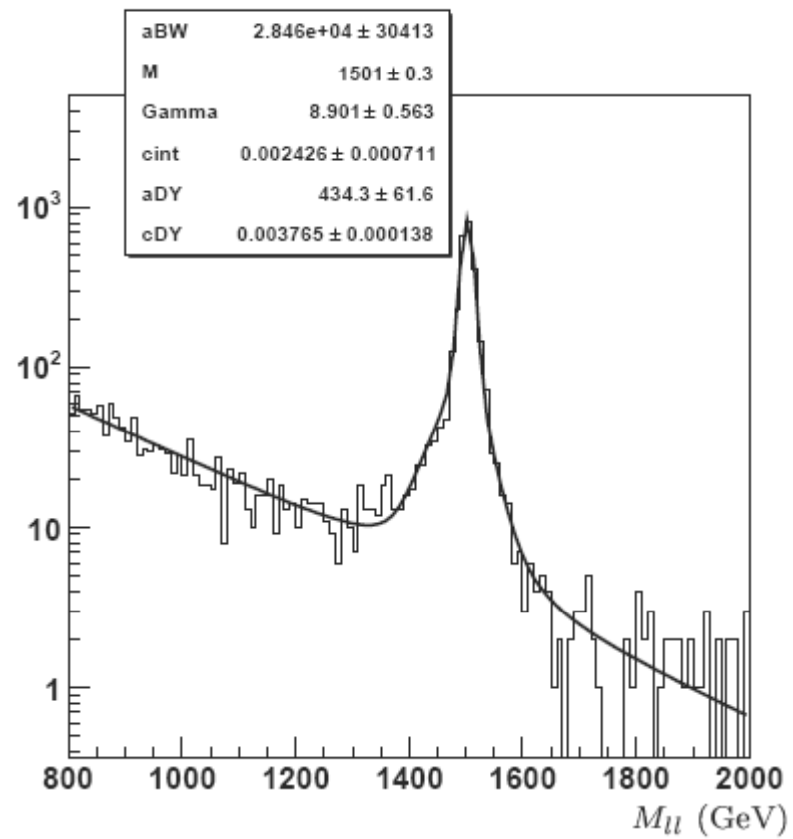


from Cvetič and Godfrey,  
[hep-ph/9404216](https://arxiv.org/abs/hep-ph/9404216)

Figure 6: 90% confidence level ( $\Delta\chi^2 = 6.3$ ) regions for the  $\chi$ ,  $\psi$ , and  $\eta$  models with  $M_{Z'} = 1$  TeV are plotted for  $P_R^u$  versus  $P_R^d$  versus  $P_V^e$  at the LHC. The figure reflects a few-fold ambiguity in the determination of these couplings at the LHC.

## excess fermions

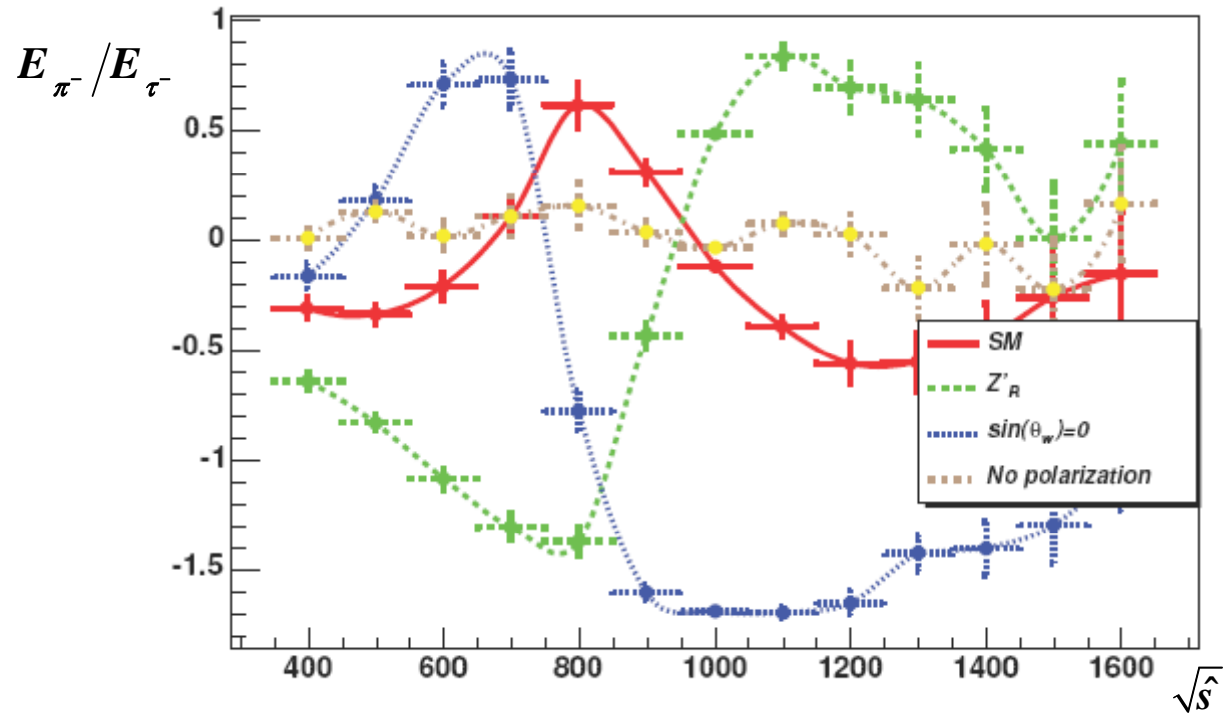
- ➔ **contact interaction**
- ➔ **virtual gravitons**
- ➔ **torsion**



Ledroit atl-phys-pub-2005-010

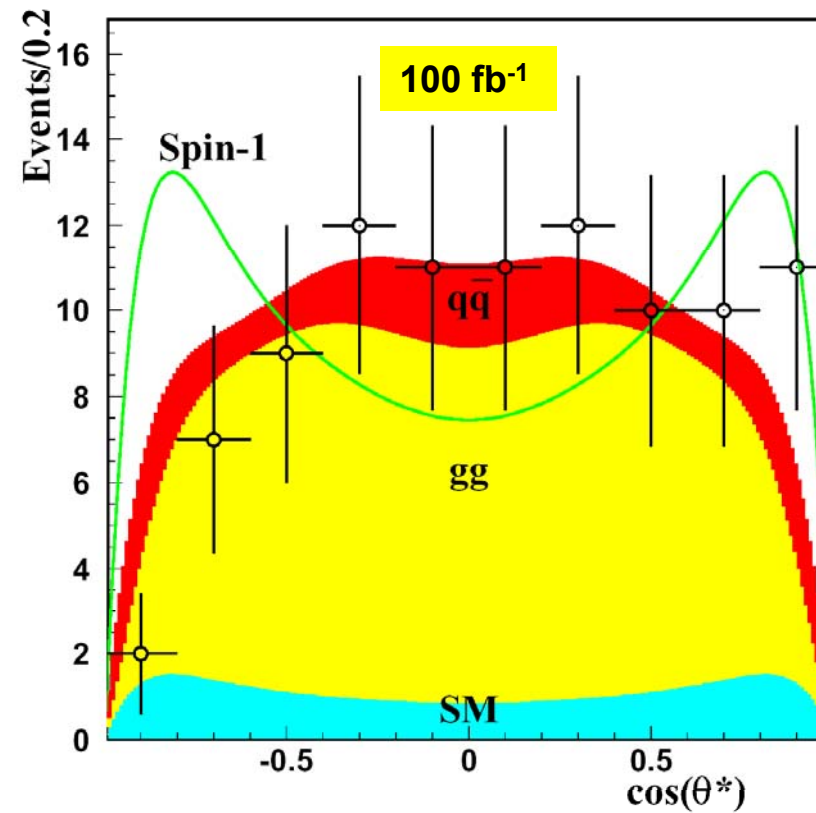
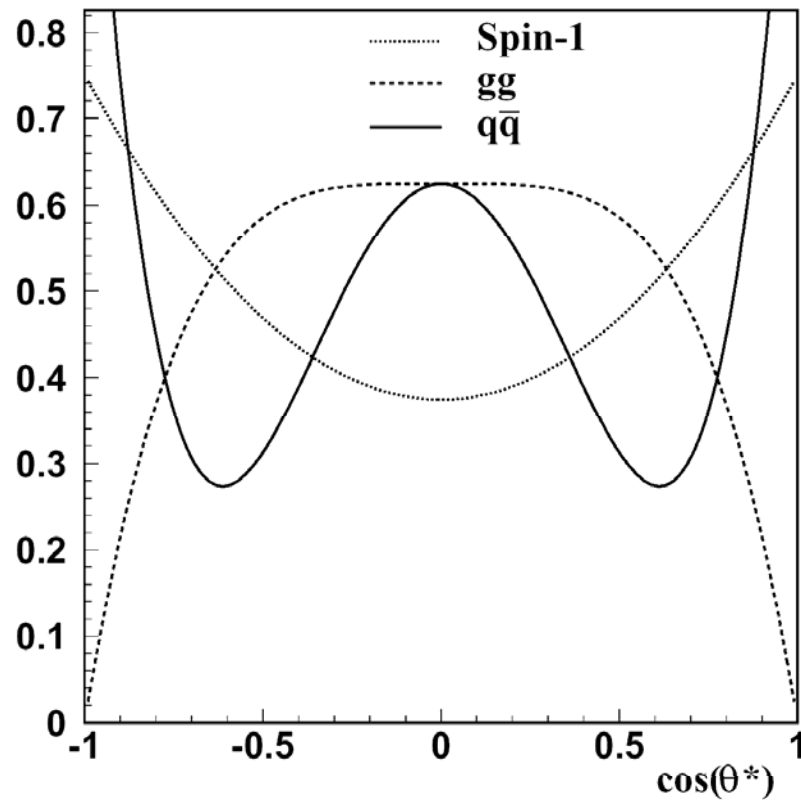


$Z' \rightarrow \tau\tau$   
 $\quad \quad \quad \searrow \pi\nu$   
 $m(Z') = 1\text{TeV}$



# Narrow Graviton Resonance

## Spin determination



spin-2 could be determined (spin-1 ruled out) with 90% CL up to graviton mass of **1720 GeV**

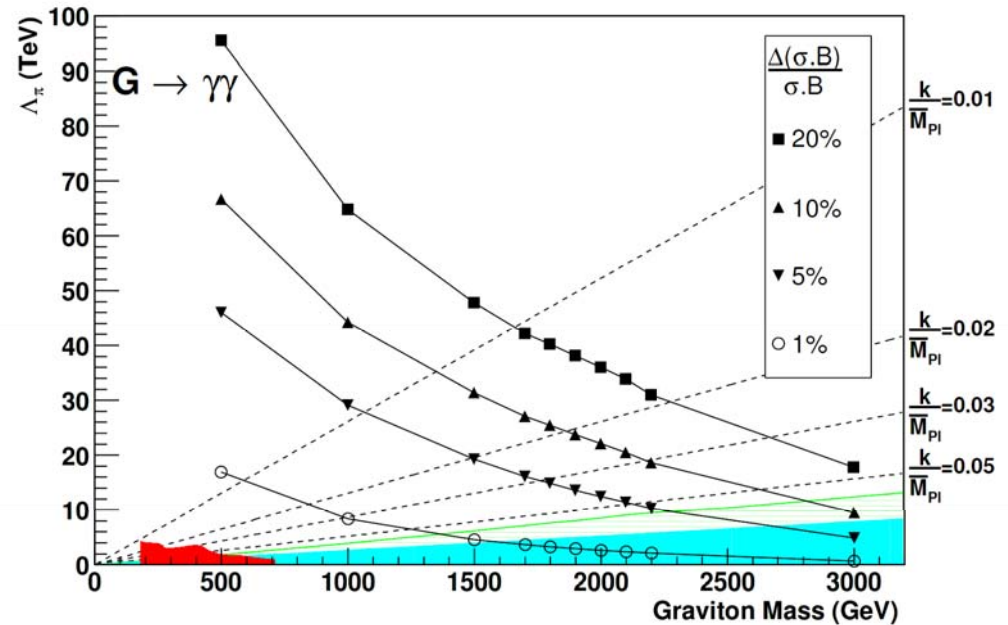
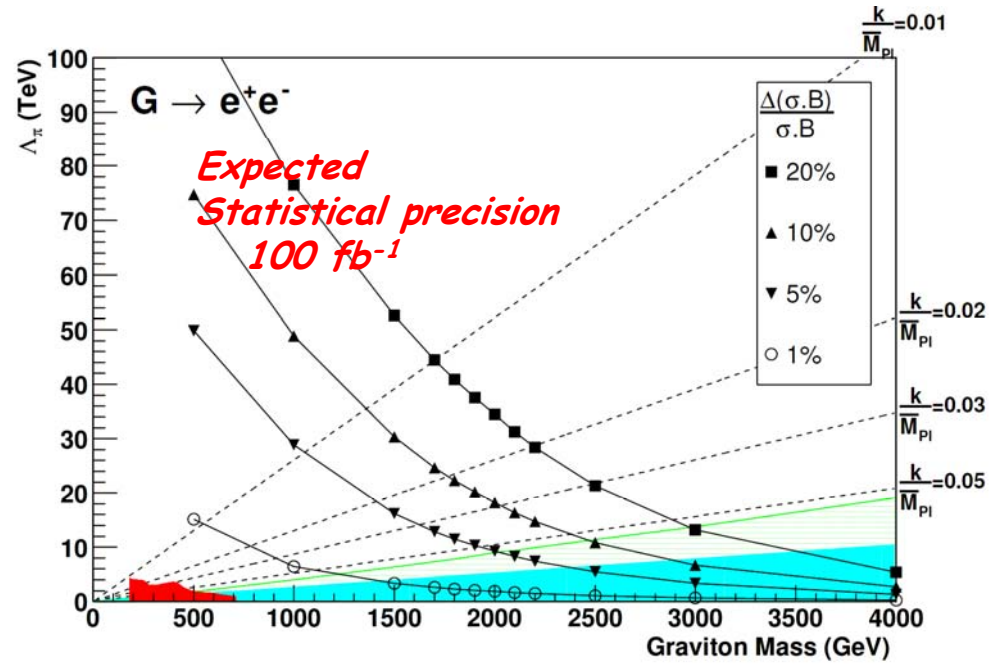
Allanach et al., ATL-PHYS-2002-031

- also  $G \rightarrow WW, ZZ, jj, mm, tt, hh$

$$\Lambda = \frac{m_G}{x_1 \cdot k / M_{Pl}},$$

$x_1 = 3.83 = 1^{st}$  root of Bessel function

e.g.: for a resonance observed at  
 $m_G = 1.5$  TeV in ee channel  
 $\Delta m_G < 10.5$  GeV (energy scale error)  
 $\Delta \sigma.B \sim 18\%$   
 if  $k/r_c = 0.01$  (pessimistic)  
 $\Rightarrow r_c = (82 \pm 7) \times 10^{-33}$  m !!



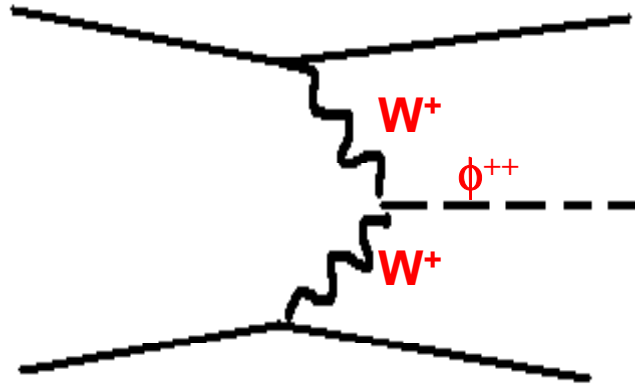
LQ	$F$	$I_3$	$Q_{em}$	decay	coupling	$\beta$
$S_0$	2	0	-1/3	$e_L^- u_L$	$\lambda_{LS_0}$	$\frac{\lambda_{LS_0}^2 + \lambda_{RS_0}^2}{2\lambda_{LS_0}^2 + \lambda_{RS_0}^2}$
				$e_R^- u_R$	$\lambda_{RS_0}$	
				$\nu_e d_L$	$-\lambda_{LS_0}$	
$\tilde{S}_0$	2	0	-4/3	$e_R^- d_R$	$\lambda_{R\tilde{S}_0}$	1
				1	2/3	$\nu_e u_L$
$S_1$	2	0	-1/3	$\nu_e d_L$	$-\lambda_{LS_1}$	1/2
				$e_L^- u_L$	$-\lambda_{LS_1}$	
				-1	-4/3	$e_L^- d_L$
$S_{1/2}$	0	1/2	-2/3	$\nu_e \bar{u}_L$	$\lambda_{LS_{1/2}}$	$\frac{\lambda_{RS_{1/2}}^2}{\lambda_{LS_{1/2}}^2 + \lambda_{RS_{1/2}}^2}$
				$e_R^- \bar{d}_R$	$-\lambda_{RS_{1/2}}$	
				-1/2	-5/3	$e_L^- \bar{u}_L$
				$e_R^- \bar{u}_R$	$\lambda_{RS_{1/2}}$	
$\tilde{S}_{1/2}$	0	1/2	1/3	$\nu_e \bar{d}_L$	$\lambda_{L\tilde{S}_{1/2}}$	0
				-1/2	-2/3	$e_L^- \bar{d}_L$

LQ	$F$	$I_3$	$Q_{em}$	decay	coupling	$\beta$
$V_0$	0	0	-2/3	$e_L^- \bar{d}_R$	$\lambda_{LV_0}$	$\frac{\lambda_{LV_0}^2 + \lambda_{RV_0}^2}{2\lambda_{LV_0}^2 + \lambda_{RV_0}^2}$
				$e_R^- \bar{d}_L$	$\lambda_{RV_0}$	
				$\nu_e \bar{u}_R$	$\lambda_{LV_0}$	
$\tilde{V}_0$	0	0	-5/3	$e_R^- \bar{u}_L$	$\lambda_{R\tilde{V}_0}$	1
				1	1/3	$\nu_e \bar{d}_R$
$V_1$	0	0	-2/3	$\nu_e \bar{u}_R$	$\lambda_{LV_1}$	1/2
				$e_L^- \bar{d}_R$	$-\lambda_{LV_1}$	
				-1	-5/3	$e_L^- \bar{u}_R$
$V_{1/2}$	2	1/2	-1/3	$\nu_e d_R$	$\lambda_{LV_{1/2}}$	$\frac{\lambda_{RV_{1/2}}^2}{\lambda_{LV_{1/2}}^2 + \lambda_{RV_{1/2}}^2}$
				$e_R^- u_L$	$\lambda_{RV_{1/2}}$	
				-1/2	-4/3	$e_L^- d_R$
				$e_R^- d_L$	$\lambda_{RV_{1/2}}$	
$\tilde{V}_{1/2}$	2	1/2	2/3	$\nu_e u_R$	$\lambda_{L\tilde{V}_{1/2}}$	0
				-1/2	-1/3	$e_L^- u_R$

from LQ - OPAL



# Triplet Higgs

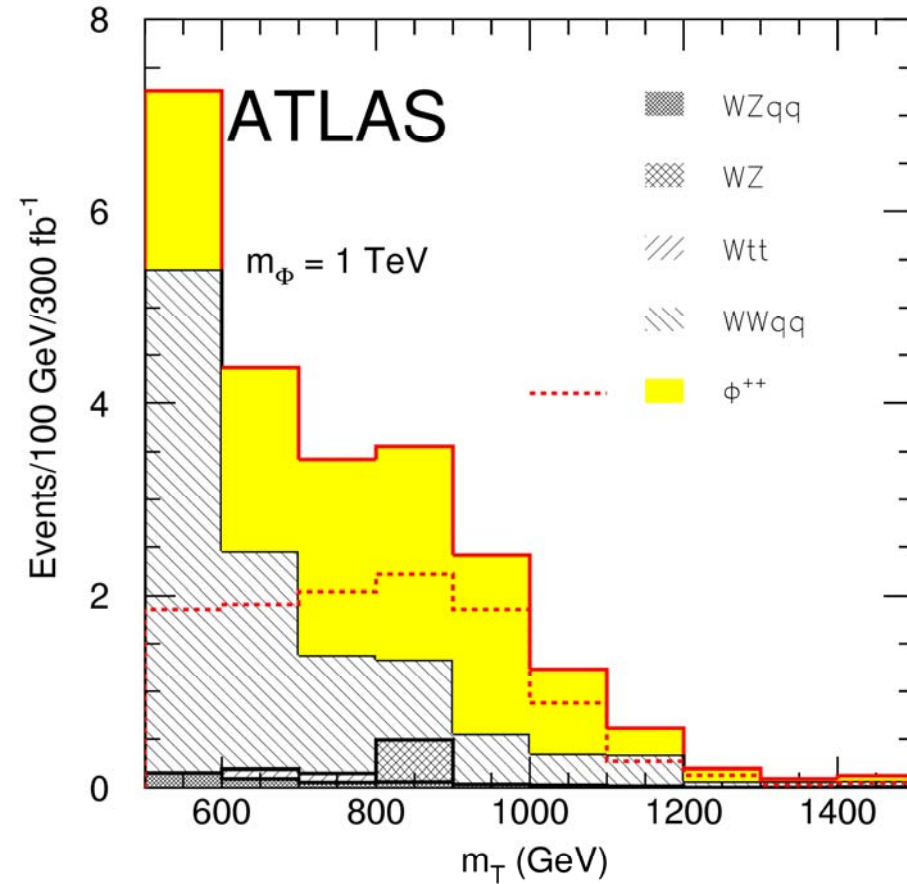


Single production

Main background from  $W_T W_T$  scattering

$$qq \rightarrow qq \phi^{++}$$

$$\searrow W^+ W^+ \rightarrow \ell^+ \nu \ell^+ \nu$$



# Right-handed interactions?

## Z' : first sign of extended gauge group ?

### Left-Right Symmetric Model: triplet Higgs

$$\begin{array}{ccccc}
 SU(2)_L \times SU(2)_R \times U(1)_{B-L} & \xrightarrow{\Delta} & SU(2)_L \times U(1)_Y & & \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbf{W}_L^i & \mathbf{W}_R^i & \mathbf{C} & \mathbf{W}_L^i & \mathbf{B} \\
 \mathbf{g}_L & \mathbf{g}_R & \mathbf{g}' & \mathbf{g}_L & \mathbf{g}_Y
 \end{array}$$

- right-handed fermions in doublets → heavy Majorana  $\nu_R = N$ 
  - explains low mass of  $\nu_L$  (see-saw mechanism)
- $\mathbf{W}_R, \mathbf{Z}_R$  associated with right-handed sector  $\mathbf{W}_R \rightarrow eN, \mathbf{Z}_R \rightarrow ee$

### larger GUT groups (includes LRSM)

$$\begin{aligned}
 E_6 &\rightarrow SO(10) + U(1)_\psi \\
 &\quad \hookrightarrow SU(5) + U(1)_\chi
 \end{aligned}$$

$$Q_{E_6} = \cos \beta Q_\chi + \sin \beta Q_\psi \quad \Rightarrow \quad Q_\eta = \sqrt{\frac{3}{8}} Q_\chi - \sqrt{\frac{5}{8}} Q_\psi$$

	$\chi$ -Model		$\eta$ -Model		$\psi$ -Model	
	$g_V$	$g_A$	$g_V$	$g_A$	$g_V$	$g_A$
$Z' \bar{e}e$	$-\frac{g \tan \theta_W}{\sqrt{6}}$	$-\frac{g \tan \theta_W}{2\sqrt{6}}$	$-\frac{g \tan \theta_W}{4}$	$-\frac{g \tan \theta_W}{12}$	0	$\frac{-g \tan \theta_W \sqrt{10}}{12}$
$Z' \bar{\nu}_e \nu_e$	$-\frac{3g \tan \theta_W}{4\sqrt{6}}$	$-\frac{3g \tan \theta_W}{4\sqrt{6}}$	$-\frac{g \tan \theta_W}{12}$	$-\frac{g \tan \theta_W}{12}$	$-\frac{g \tan \theta_W \sqrt{10}}{24}$	$-\frac{g \tan \theta_W \sqrt{10}}{24}$
$Z' \bar{u}u$	0	$\frac{g \tan \theta_W}{2\sqrt{6}}$	0	$\frac{g \tan \theta_W}{3}$	0	$-\frac{g \tan \theta_W \sqrt{10}}{12}$
$Z' \bar{d}d$	$\frac{g \tan \theta_W}{\sqrt{6}}$	$-\frac{g \tan \theta_W}{2\sqrt{6}}$	$\frac{g \tan \theta_W}{4}$	$\frac{g \tan \theta_W}{12}$	0	$-\frac{g \tan \theta_W \sqrt{10}}{12}$

TABLE III: Couplings between  $Z'$  and fermions in  $\psi$ ,  $\eta$  and  $\chi$  models where  $\theta_W$  is the Weinberg angle.

from Almeida et al., arXiv:hep-ph/0702137v1

# The littlest Higgs Model - particles

In  $SU(2)_L \times U(1)_Y$  basis,  $(2\tau+1)_{Y/2} \rightarrow \left. \begin{matrix} X \\ 1_0 \end{matrix} \right\} \rightarrow$  eaten  $\Rightarrow$  4 heavy bosons:  $Z_H, \gamma_H, W_H^\pm$

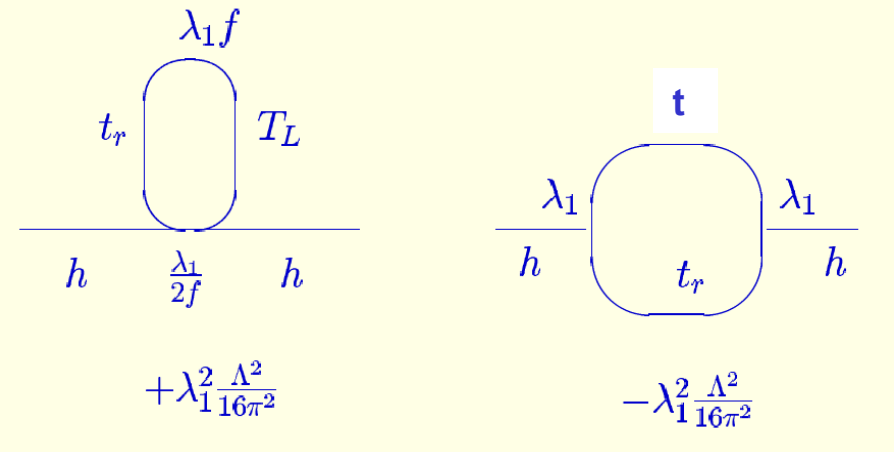
$Q = \tau_3 + \frac{Y}{2} \rightarrow \left. \begin{matrix} h \\ 2_{1/2} \end{matrix} \right\} \rightarrow$  complex higgs doublet, massless

$\left. \begin{matrix} \phi \\ 3_1 \end{matrix} \right\} \rightarrow$  complex triplet  
 acquires mass from one-loop gauge interactions

triggers EW symm. breaking  
 $\rightarrow$  mass to Z, W, h massless

1-loop gauge interactions:

To cancel the top loop, introduce  $SU(2)_L$  singlet quark  $T_L$ , and  $T_R$

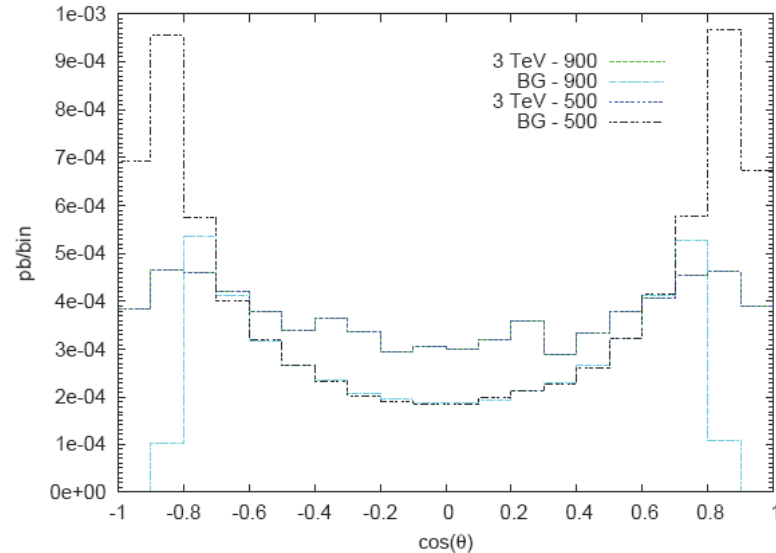
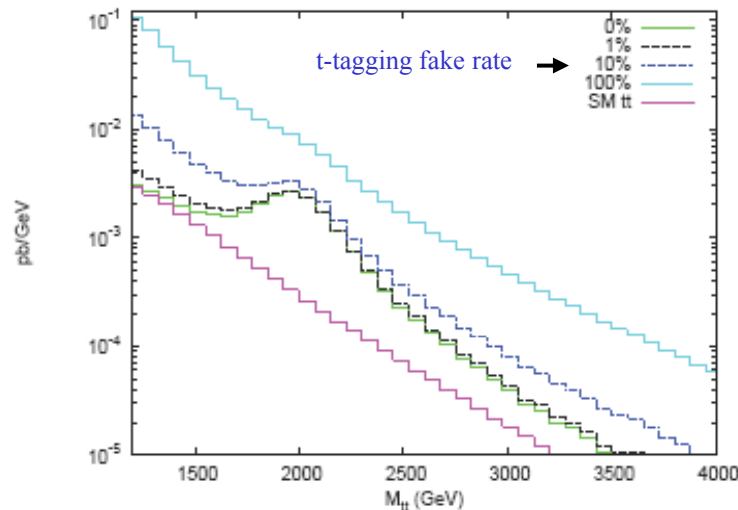


$$\lambda_1 (iQht_r + fT_L t_r h h^\dagger) + \lambda_2 f (T_L T_R)$$

# top resonance from bulk RS KK gluon

B. Lillie, L. Randall and L-T Wang, hep-ph/0701166

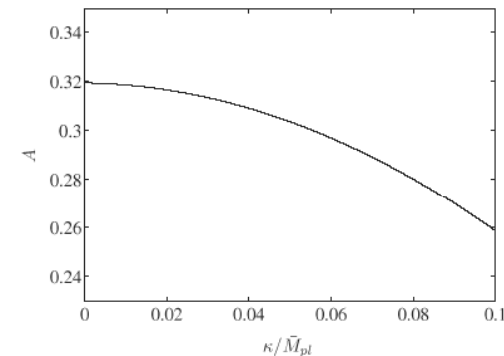
- large overlap of KK gluon and top quark wave functions because both are localized towards TeV brane



- can also measure spin correlations ( $t_R$  coupling?)
- experimental issues:
  - b-tagging
  - strong collimation of jets from top and from W's
    - jet mass can be used, as in:

W. Skiba, D. Tucker-Smith hep-ph/0701247

- can also have graviton resonance to top pairs (or WW), but higher mass (Agashe et al., hep-ph/0701186)



spin asymmetry vs  $k/M_{Pl}$  (in classical RS model),

M. Arai et al., hep-ph/0701157