

No Higgs Models

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Fermilab, April 12, 2007



Life without a Higgs:

- The Standard Model has been tested very accurately, however...
- we still don't know the origin of the Electroweak symmetry breaking (and masses...)
- The Higgs mechanism is a nice description of such phenomenon, but it suffers from theoretical prejudices (hierarchy problem), and the Higgs boson has not been discovered yet.
- What if we don't see the Higgs at the (**early**) LHC?




The role of the Higgs:

- Weakly coupled $H = (h, W_L^\pm, Z_L)$
- The same dof's can arise as composite states of a strongly interacting sector (Technicolor) \rightarrow no fundamental scalar
- Can we decouple h from the theory? What would that imply?

$W_L W_L$ scattering violates Unitarity at ~ 1 TeV scale.

Strong coupling appears below 1 TeV \rightarrow broad resonances (techni- ρ , ...), theory out of control (precision EW tests, ...)



We would like to keep the theory under control up to a safe scale of 10 TeV:

- Let's consider the contribution of the ρ -mesons to the W_L scattering:

- Expanding the scattering amplitude for large Energy:

$$A(W_L-W_L) \sim A^{(4)} E^4 + A^{(2)} E^2 + \text{finite terms.}$$

- $A^{(4)}$ vanishes due to gauge invariance;


- $A^{(2)}$ receives contributions both from the vector resonances and the Higgs:

$$A^{(2)} \sim g_{WWWW}^2 - \frac{3}{4} \sum_k \frac{M_{Z^k}^2}{M_W^2} g_{WWZ^k}^2 - \frac{1}{4} \sum_k g_{WWH^k}^2$$

The violation of perturbative unitarity can be delayed up to 10 TeV!

We can smoothly go to a Higgsless model (gaugephobic Higgs model)

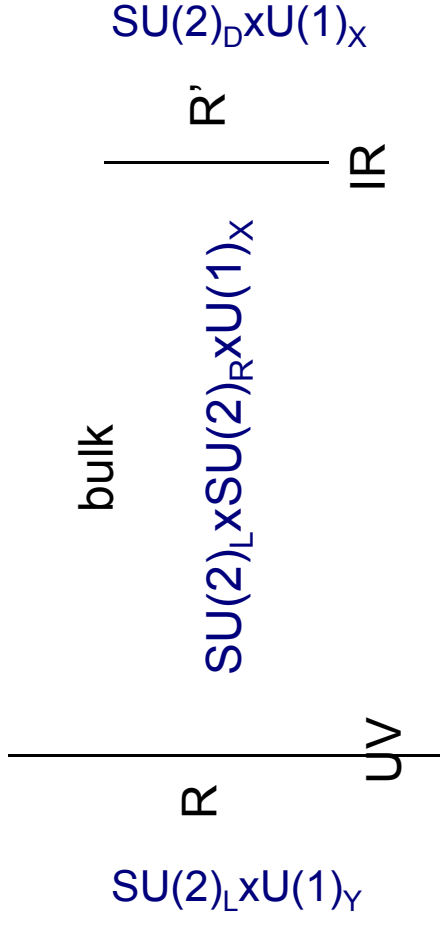
hep-ph/0611385



Extra dimensions: a holographic view into strongly coupled 4D theories

- conformal sector (strong)
- Mass gap
- CFT global symmetries
- Elementary sector
- Bound states
- Warped 5D
- IR brane
- Bulk gauge symmetries
- UV brane
- KK states

The model



- AdS₅ space: $ds^2 = \left(\frac{R}{z}\right)^2 (\eta_{\mu\nu} dx^\mu dx^\nu - dz^2)$
- UV brane at $z=R$, IR brane at $z=R'$;
- $SU(2)_L \times SU(2)_R \times U(1)_X$ in the bulk;
- $SU(2)_R \times U(1)_X \rightarrow U(1)_Y$ on the UV brane;
- $SU(2)_L \times SU(2)_R \rightarrow SU(2)_D$ on the IR brane.



Light fermions

- ▶ Light fermions are doublets of the $SU(2)$'s:
 - $\psi_L = (u_l, d_l) = (2, 1, 1/6)$ bulk mass c_L
 - $\psi_R = (u_r, d_r) = (1, 2, 1/6)$ bulk mass c_R
- ▶ The bulk masses control the localization of the zero modes.
- ▶ Delocalizing the zero modes in the bulk ($c_L \sim 0.46$) it is possible to minimize the EWPTs ($S \sim 0$).

Both representations and localization are preferred by precision measurements!

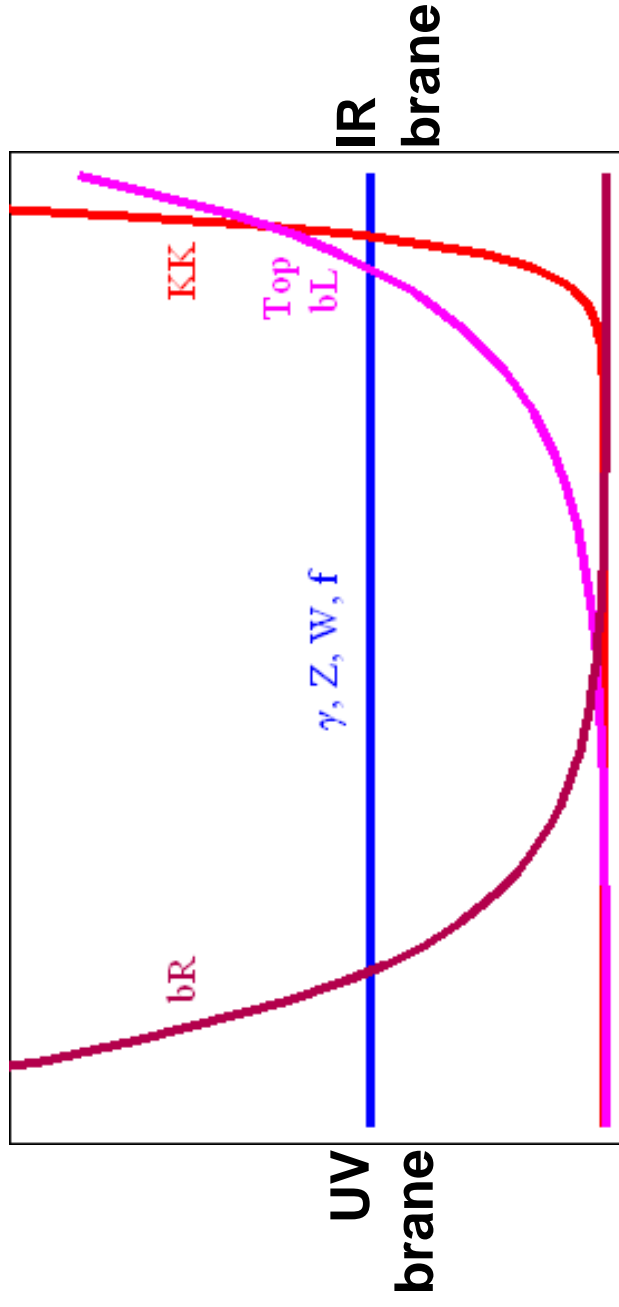


Third generation: m_{top} vs $Zb\bar{b}$


Csaki et al, hep-ph/0607146

- In Higgsless models it is a challenge to have a heavy top and small corrections to the $Zb\bar{b}$ coupling (below 1%).
- The top must be composite due to its heaviness: the b_L is also composite.
- The $SU(2)_R$ doublet ψ_R contains a left-handed b' that mixes with the b_1 via the top Yukawa.
- If we extend the custodial symmetry to $O(4) \sim SU(2)_L \times SU(2)_R \times P_{LR}$, we can use P_{LR} to protect $Zb\bar{b}$: **need to use new representations that respect $O(4)$** . (Agashe, Contino, daRold, Pomarol)
- **Or, we can exile top and bottom to a different AdS throat** (like, we have two independent CFT sectors – Csaki et al, hep-ph/0505001 and 0604218)

Group portrait without the Lady:



- Gauge bosons and light fermions are flat
- Top, bL and KK resonances are localized near the IR brane -> highly composite
- bR is localized near the UV brane -> elementary



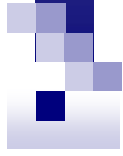
Summary of the models:

- The only free parameter is the mass of the KK states M_{KK} . (minimal model)
- All the other parameters are fixed by SM values (M_Z , G_F , α , ...) or precision tests.
- Reasonable mass range (unitarity):
 $600 \text{ GeV} < M_{\text{KK}} < 1 \text{ TeV}$.
- Perturbative and under control up to $\sim 10 \text{ TeV}$ (safe for EWPTs).

Benchmark points: two scenarios

$$M_{\text{KK}} \sim 700 \text{ GeV}$$

- Scenario 1: topless
- The third generation decouples from the EWSB (i.e. “a brane on their own”, [hep-ph/0505001](#))
- Top and bottom decouple from the (light) KK states
- Scenario 2: top
- “New custodian” to protect $Zb_L b_L$: [hep-ph/0607146](#)
- KK states do couple to t, b
- KK states of the third generation (effects of the new custodian)

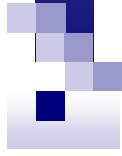


Phenomenology: Gauge Bosons

[GeV]	0	1	2	3
W	80.4	699	1105	1583
Z	91.2	694	1110	1578
γ (A)	0	718	-	1603
G	0	718	-	1603

Couplings, an example:

Z1 W W	0.034	$\sim 10\%$ SM
Z1 $f_L f_L$	$(-0.048) T_{3L} + (0.096) Y$	$\sim 10\%$ SM
Z1 tt	$(0.46) P_L + (0.44) P_R$	\sim SM

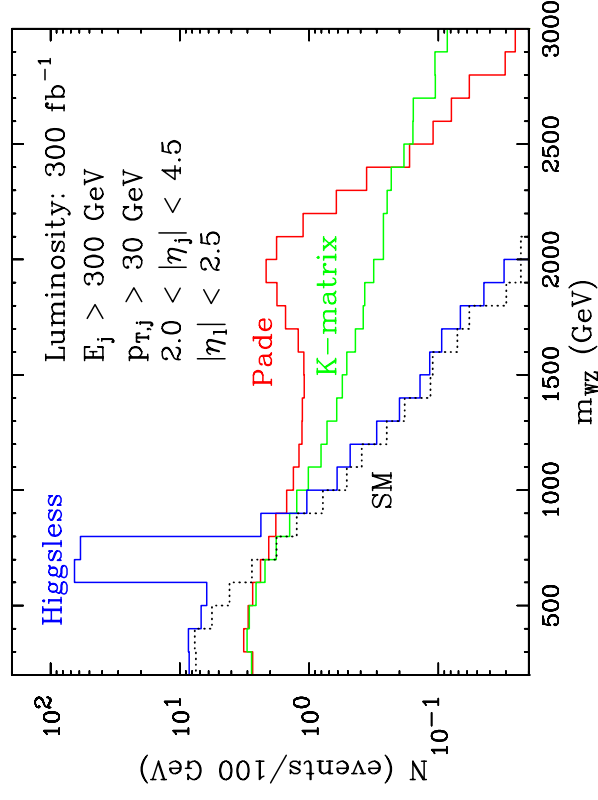


	Scenario 1	Scenario 2
W1	$\Gamma = 16 \text{ GeV}$	$\Gamma = 54 \text{ GeV}$
	99% in ZW	30% in ZW
		34% in bt; 35% in heavy tops
	1% jets; 0.1% leptons	0.2% jets; 0.04% leptons
Z1	$\Gamma = 9 \text{ GeV}$	$\Gamma = 60 \text{ GeV}$
	97% in WW	15% in WW
		67% in bb; 18% in tt
	1% jets; 0.15% l^+l^- ; 1.5% inv	0.02% jets; 0.025% l^+l^- ; 0.3% inv
A1	$\Gamma = 7 \text{ GeV}$	$\Gamma = 21 \text{ GeV}$
	90% in WW	31% in WW
		59% in tt; 6% in bb
	5% jets; 2% l^+l^-	2% jets; 0.6% l^+l^-

Vector Boson Fusion

- Probes the couplings unitarizing longitudinal GB scattering, i.e. the core of the new EWSB mechanism!
- Studied in [Perelstein-Matchev-Birkedal hep-ph/0412278](#) (assuming 100% BR in GBs)

- $\sigma(pp \rightarrow W1 jj) = 200 \text{ fb}$
- decay channel $W1 \rightarrow WZ$
 - $jj + 3l + \text{missing } E_T$
- With 10 fb^{-1} , probe $< 550 \text{ GeV}$
- To probe up to 1 TeV mass, 60 fb^{-1} are needed!
- BR in tb makes it worse!
- $Z1, A1 \rightarrow WW$ is even more challenging.



Drell-Yan production

- Sensitive to couplings to light quarks ($S=0$)
- Channel ignored in the previous literature

$$\sigma(p p \rightarrow W1^+) = 0.95 \text{ pb}$$

$$\sigma(p p \rightarrow W1^-) = 0.48 \text{ pb}$$

$$\sigma(p p \rightarrow Z1) = 0.7 \text{ pb}$$

$$\sigma(p p \rightarrow A1) = 2.5 \text{ pb}$$

Compare with Higgs:

$$\sigma(g g \rightarrow h) = 0.8 \text{ pb}$$

$$@ m_H = 700 \text{ GeV}$$

- $\sigma(Z1, A1 \rightarrow W^+ W^-) = 2.9 \text{ pb}$ (0.9 pb): better discovery potential than a 700 GeV Higgs \rightarrow discovery at 10 fb^{-1} ! Impossible to distinguish Z1-A1.
- $\sigma(W1 \rightarrow W Z) = 1.4 \text{ pb}$ (0.4 pb): discovery at 10 fb^{-1} !
- top-bottom channels: swamped by gluon (see in a few slides)
- $\sigma(Z1, A1 \rightarrow l^+ l^-) = 43 \text{ fb}$ (15 fb): hundreds of dilepton events, easy to reconstruct the resonance!

Top-Bottom associated production

- Sensitive to couplings to the third generation (Scenario 2 only!)

$$\sigma(p p \rightarrow t \bar{b} W1^+) = 0.64 \text{ pb}$$

$$\sigma(p p \rightarrow b \bar{t} W1^-) = 0.64 \text{ pb}$$

$$\sigma(p p \rightarrow t \bar{t} Z1) = 26 \text{ fb}$$

$$\sigma(p p \rightarrow t \bar{t} A1) = 26 \text{ fb}$$


$$\sigma(p p \rightarrow b \bar{b} Z1) = 5.7 \text{ pb}$$

$$\sigma(p p \rightarrow b \bar{b} A1) = 0.2 \text{ pb}$$

- Large cross sections: good chance to be observed in W,Z!
- t, b channels swamped by the gluon.
- Other channels: pair production, associated production with gauge bosons, jets... (work in progress)

Gluon

- Large couplings to top-bottom, large production σ (strong interactions).
- Broad resonance: $\Gamma = 110$ GeV
- tt – 66%; bb – 30%; jj – 4%
 - $\sigma(pp \rightarrow G1) = 145$ pb $\sigma(pp \rightarrow G1 G1) = 46$ pb
 - $\sigma(pp \rightarrow t t G1) = 75$ pb $\sigma(pp \rightarrow b b G1) = 3$ nb
- Discovery at $> 5\sigma$ in tt channel (i.e. Barger, Han, Walker, [hep-ph/0612016](https://arxiv.org/abs/hep-ph/0612016))
- Associated production: see Han, Valencia, Wang, [hep-ph/0405055](https://arxiv.org/abs/hep-ph/0405055) ($pp \rightarrow tt G1 \rightarrow tt bb$ and $pp \rightarrow bb G1 \rightarrow bb tt$)
- Pollution for the Z1 and A1 decay channels in tt and bb !
- However, the gluon is not directly related to the EWSB...

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- **tt channel:** Barger, Han, Walker, [hep-ph/0612016](#)
Semileptonic decay ($t \bar{t} \rightarrow l \nu b + j j b$), small background;
at 10 fb^{-1} , for $m_V = 800 - 1000 \text{ GeV}$, 5σ discovery if

$$\sigma \geq 1/5 \sigma(\text{SM-Z'})$$

Not enough for our Z1, but yes for G1!
Interesting to study tb final state (no G1 pollution!)

- **Associated production:** Han, Valencia, Wang, [hep-ph/0405055](#)
For a Z' with SM couplings with t and b:
 5σ discovery for $m_V < 1200 \text{ GeV}$ in $t\bar{t} V \rightarrow t\bar{t} b\bar{b}$;
Discovery guaranteed for the G1, possible for Z1 and A1.
Interesting to look at a resonance in tb channel (W1)!

Light fermions KK states

- $SU(2)_L$ doublet is flat \rightarrow KK state degenerate with A1:
 $m_{fL} = 720 \text{ GeV}$.
- Flatness is required by EWPTs ($S=0$)
- General facts (common to leptons and light quarks):
 - $\Gamma = 1.2 \text{ GeV}$
 - $BR(W) \sim 2/3$; $BR(Z) \sim 1/3$
- Leptons: single production via a W or Z \rightarrow tens of ab!
Not enough at 10 fb^{-1} .
- Pair production: few fb. Not enough at 10 fb^{-1} ,
but discovery guaranteed at higher luminosity!
(multi-lepton events!)
Good chances to be studied at ILC.



KK quarks

$$\sigma(pp \rightarrow q \bar{q}) = 310 \text{ fb} \quad (W, Z)$$

$$\sigma(pp \rightarrow q \bar{q}) = 3.3 \text{ pb} \quad (\text{strong})$$

- Signatures in W_j, Z_j
- Similar studies in Little Higgs models: good chances for discovery at 10 fb^{-1} .
- See for example, [Azuelos et al, hep-ph/0402037](#)
- Need of a detailed study at those lower masses!

The new custodian for the bottom.

$$\Psi_L \sim (\mathbf{2}, \mathbf{2})_{2/3} \supset (t_\ell, b_\ell)$$

$$t_R \sim (\mathbf{1}, \mathbf{1})_{2/3} \supset t_r$$

$$\Psi_R \sim (\mathbf{1}, \mathbf{3})_{2/3} \supset b_r$$

They will contain extra (heavy) degrees of freedom:

$$\Psi_L = \begin{pmatrix} q_L & Q_L \end{pmatrix} = \begin{pmatrix} t_L & X_L \\ b_L & T_L \end{pmatrix}, \quad \Psi_R = \begin{pmatrix} X_R \\ T_R \\ b_R \end{pmatrix}, \quad t_R$$

The top Yukawa (M1) does not involve the bottom!
The dangerous mixing is absent.

$$\mathcal{L}_{IR} = M_3 \left[\frac{1}{\sqrt{2}} T_R (t_L + T_L) + b_R b_L + X_R X_L \right] + \frac{M_1}{\sqrt{2}} t_R (t_L - T_L) + h.c.$$

$$Q(b_L, b_R) = -1/3; \quad Q(t_L, t_R, T_L, T_R) = 2/3; \quad Q(X_L, X_R) = 5/3.$$

B(-1/3)	4.5	-	670	943
T(2/3)	172.5	450	622	835
X(5/3)	-	435	-	812

X1) $\Gamma = 6 \text{ GeV}$; $W t - 100\%$

T1) $\Gamma = 1.5 \text{ GeV}$; $Z t - 76\%$ $W b - 24\%$

$\sigma(pp \rightarrow X1 X1) = 11.3 \text{ pb}$ $\sigma(pp \rightarrow T1 T1) = 9.3 \text{ pb}$ (strong)

- Interesting decay chain $X1 \rightarrow W^+ t \rightarrow W^+ W^+ b$
- Final state: $4W+2b\text{-jet}$
- Studied in [Dennis, Unel, Servant, Tseng, hep-ph/0701158](#):
discovery $>5\sigma$ guaranteed, looked at 1 leptonic decay of W s.
- Interesting channel: 2 same sign W 's decaying leptonically.
It allows to reconstruct a peak in the $4j+b$ channel!



Conclusions and outlook

- EWSB via a strong sector is a viable alternative to the (weakly coupled) Higgs mechanism.
- Extra dimensions are a new handle on the strong sector: perturbative control up to 10 TeV.
- Realistic No-Higgs models are now available.
- Rich phenomenology: $>5\sigma$ discovery at 10 fb^{-1} for $W1$, $Z1$, $A1$, $q1$, $T1$, $G1$...and more!
- “Indirect” probes: $Z t t$, $W t b$ ($\sim 10\%$ deviation) [single top production], WWZ ($\sim 2\%$ deviation)
- ILC necessary to study the couplings, and discover other particles (like $L1$).
- More detailed studies are necessary! (work in progress)