Something out of nothing:

Dark matter observation and mass determination in photon + missing energy events at the ILC

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Dark matter - WIMP

- Cosmological observations —> 20% of the energy density of Universe is due to non-relativistic, non-baryonic dark matter.
- WIMP most attractive candidate.
 - Mass $\longrightarrow 10 \text{ GeV} 1 \text{ TeV}$
 - Interactions strength \longrightarrow weak interactions of Standard Model.
- Have a relic abundance of the correct order of magnitude to account for the observed dark matter.
- Many BSM extensions contain suitable candidates for WIMP.
 - SUSY LSP ?
 - UED LKP ?
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Dark matter at collider

- Can we produce DM (χ) in collider?
 - We are in correct energy range
 - Expect "missing energy" signature
- Can we quantify these process : $X_i + \bar{X}_i \rightarrow \chi + \chi$ $(X_i = q, \ell, g, W, Z, \gamma, h)$
 - In a model independent way?
 - Using the known cosmological abundance of WIMP \rightarrow governed by scattering $\chi + \chi \rightarrow X_i + \bar{X}_i$
- We need inverse relation for any prediction @ collider

Relic density

- At early times, the DM particles are in thermal equilibrium with the SM particles : $\chi\chi \iff X_i \bar{X}_i$
- Freeze-out of thermal relics is described by the Boltzmann equation

$$\frac{dn_{\chi}}{dt} = -3Hn_{\chi} - \langle \sigma_A v \rangle \left(n_{\chi}^2 - n_{eq}^2\right)$$

- Total DM annihilation cross-section (σ_A) can be expanded: $\sigma_A v = \sum_i \sigma(\chi + \chi \to X_i + \overline{X}_i) v = a + bv^2 + O(v^4)$
- Approximate Relic density : $\Omega h^2 = 0.08 \frac{1pb}{a + (3b 0.75a)x_F}$
 - $x_F = T_F / M_\chi \sim 0.04$
 - T_F the freeze-out temperature.

Relic density

• WMAP: The present amount of dark matter $\sigma_{DM}h^2 = 0.111^{+0.011}_{-0.015}$

WMAP Coll. '06

- Two classes of DM candidates:
 - a = b: "s-annihilators"
 - $a < bx_F << b$: "p-annihilators".

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→ WIMP annihilation cross-section

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WIMP production cross-section

• Forward and backward WIMP production are related by detailed balancing:

$$\frac{\sigma(\chi + \chi \to X_i + \bar{X}_i)}{\sigma(X_i + \bar{X}_i \to \chi + \chi)} = 2.\frac{v_X^2 (2S_X + 1)^2}{v_\chi^2 (2S_\chi + 1)^2}$$

- Annihilation fraction: $\kappa_i = \frac{\sigma_i^{(J_0)}}{\sigma_{tot}}, \quad \sum_i \kappa_i = 1$
- Model-independent prediction of the WIMP production rate

$$\sigma(X_i + \bar{X}_i \to \chi + \chi) = 2^{2(J_0 - 1)} \kappa_i \sigma_{an} \frac{v_X^2 (2S_X + 1)^2}{v_\chi^2 (2S_\chi + 1)^2} \left(1 - \frac{4M_\chi^2}{s}\right)^{1/2 + J_0}$$

• Unknown parameters : κ_i , M_{χ} , S_{χ} , J_0

- But this is useless We need at least one detectable particle to trigger.
 - Tag the known initial state with a soft photon .. $e^+e^- \rightarrow 2\chi + \gamma$

WIMP production cross-section

• Soft/collinear photon factorization:

$$\frac{d\sigma(e+e-\rightarrow 2\chi+\gamma)}{dxd\cos\theta} = \frac{\alpha}{\pi} \frac{1+(1-x)^2}{x} \frac{1}{\sin\theta^2} \hat{\sigma}(e+e-\rightarrow 2\chi)$$

- $x = 2E_{\gamma}/\sqrt{s}$
- $\hat{\sigma}$ calculated at $\hat{s} = (1 x)s$
- Applicable when $E_{\gamma} << \sqrt{s} M_{\chi}$

VIMP production cross-section

• Soft/collinear photon factorization:

$$\frac{d\sigma(e+e-\rightarrow 2\chi+\gamma)}{dxd\cos\theta} = \frac{\alpha}{\pi} \frac{1+(1-x)^2}{x} \frac{1}{\sin\theta^2} \hat{\sigma}(e+e-\rightarrow 2\chi)$$



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- We look at the signal for 500 GeV ILC and $\mathcal{L} = 500 f b^{-1}$
- Main irreducible background from $e^+e^- \rightarrow \nu \bar{\nu} \gamma$
- Background from $e^+e^- \rightarrow e^+e^-\gamma$ can be eliminated using lower cut on the $p_{T,\gamma}$
- Photon energy distribution for Signal and Background:



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- $\sin \theta_{\gamma} > 0.1$
- $p_{T_{\gamma}} > 7.5 \ GeV$ (suppress Bhabha : mask calorimeter acceptance of 1°)
- χ non-relativistic and E_{γ} below threshold-



- Discovery reach for $\mathcal{L} = 500 f b^{-1} \text{ ILC}$
- No polarization.
- For p-annihilator WIMP
- (red, blue) band include a systematic uncertainty of 0.3%

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- $\sin \theta_{\gamma} > 0.1$
- $p_{T_{\gamma}} > 3.0 \ GeV$ (suppress Bhabha : BeamCAL acceptance of 0.38°)
- χ non-relativistic and E_{γ} below threshold-



- Discovery reach for $\mathcal{L} = 500 f b^{-1} \text{ ILC}$
- No polarization.
- For p-annihilator WIMP
- (red, blue) band include a systematic uncertainty of 0.3%
- Revised from better beam-CAL acceptance.

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- $\sin \theta_{\gamma} > 0.1$
- $p_{T_{\gamma}} > 3.0 \ GeV$ and $7.5 \ GeV$ (suppress Bhabha : BeamCAL acceptance of 0.38^o)
- χ non-relativistic and E_{γ} below threshold-

 $\frac{\sqrt{s}}{2} \left(1 - \frac{8M_{\chi}^2}{s} \right) \le E_{\gamma} \le \frac{\sqrt{s}}{2} \left(1 - \frac{4M_{\chi}^2}{s} \right)$



- Discovery reach for $\mathcal{L} = 500 f b^{-1}$ ILC
- No polarization.
- For p-annihilator WIMP
- (red, blue) band include a systematic uncertainty of 0.3%
- Improvement in 5σ level.

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- Once WIMP observed in ILC, mass measurement next step.
- Photon energy distribution can be used to reconstruct WIMP mass
- Both unpolarized case and polarized case can be considered

- Once WIMP observed in ILC, mass measurement next step.
- Photon energy distribution can be used to reconstruct WIMP mass
- Both unpolarized case and polarized case can be considered
- Reconstruction done scanning over the (M_{χ}, κ_e) parameter space and calculating χ^2 .
- Our sample signal point is M_{χ} = 150 GEV and κ_e =0.6
- Spin $\frac{1}{2}$ and p-annihilator WIMP considered
- Smeared with energy resolution $14.4\%/\sqrt{E} + 0.5\%$



• For $\mathcal{L} = 500 f b^{-1}$ ILC

• No polarization.

- For p-annihilator, Spin $\frac{1}{2}$ WIMP
- Smeared with energy resolution $14.4\%/\sqrt{E} + 0.5\%$



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Summary & Conclusions

- Model independent approach based on less assumptions over model parameters.
- ILC can efficiently resolve the mass determination of WIMP.
- Polarized beam significantly improves the result.
- Some similar analysis using recoil mass distribution indicates to measure the mass quite precisely.
 Bartels, List '07