

# Warszawa ECFA Alternatives, EW, Top (and QCD) Summary




On behalf of all speakers

ECFA 9-12 June Warszawa – Helenka Przysieczniak Frey @ U.of Montreal + CNRS

# Initial studies of top pair production (Andreas Moll, Alexei Raspereza)

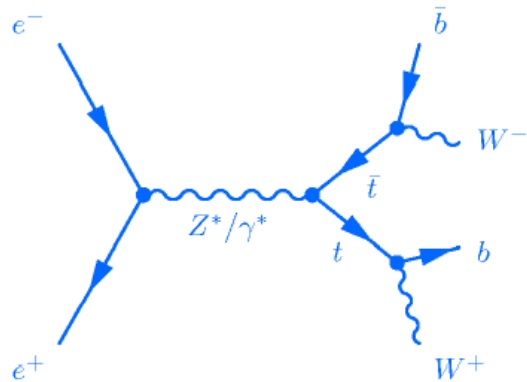
**Goal:** estimate statistical error on top mass+width for ILC from direct reconstruction of top decays. This method provides a consistency check and is complementary to the threshold scan technique.

 Tools used in this analysis:

**Mokka** – flexible geant4 based detector simulation framework

**Marlin** – modular and detector independent reconstruction software

- btag information from LCFIVertex package used
- kinematic fitting applied

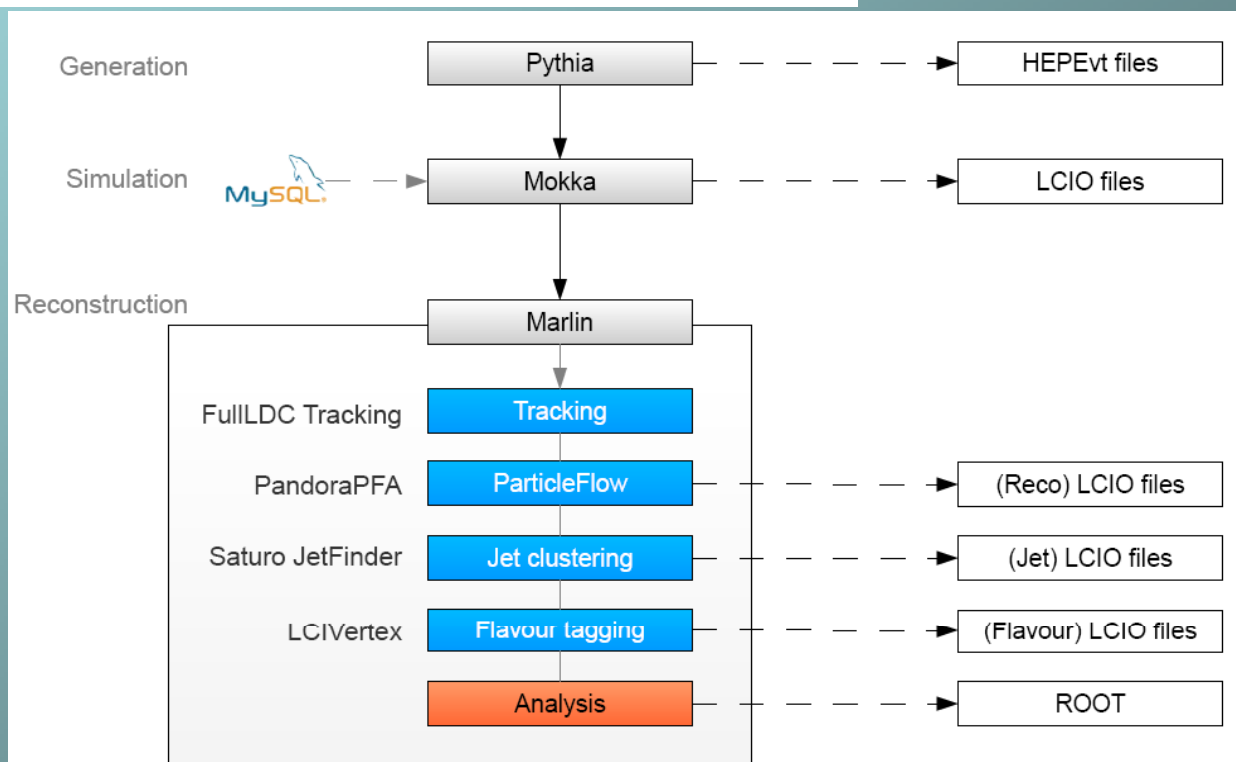


fully hadronic  $t\bar{t}$  decay mode used:

$$t\bar{t} \rightarrow (W^+ b)(W^- \bar{b})$$

$W \rightarrow q\bar{q}$  branchingratio 44.4 %

$t\bar{t} \rightarrow 6 \text{ jets}$



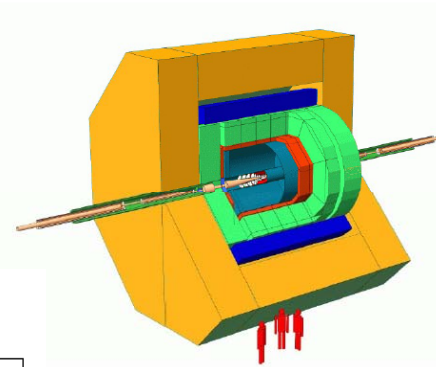
## Initial studies of top pair production (Andreas Moll, Alexei Raspereza)

- Pandora Particle flow algorithm used
  - Tracking in VTX/TPC/SiT/FTD
  - Clustering of hits in ECAL/HCAL
    - Selections cuts
      - B-tag
      - Kinematic fit

**Mokka:** geant4 based framework for full detector simulation

Detector used for simulation with Mokka: **LDCPrime\_02Sc\_p01**

⇒ Interpolation between the two detector concepts GLD and LDC



**Magnetic field:** 3.5 T

**Tracking:**

VTX (inner radius=1.5 cm)

TPC (R=1.7 m, L=4.4 m)

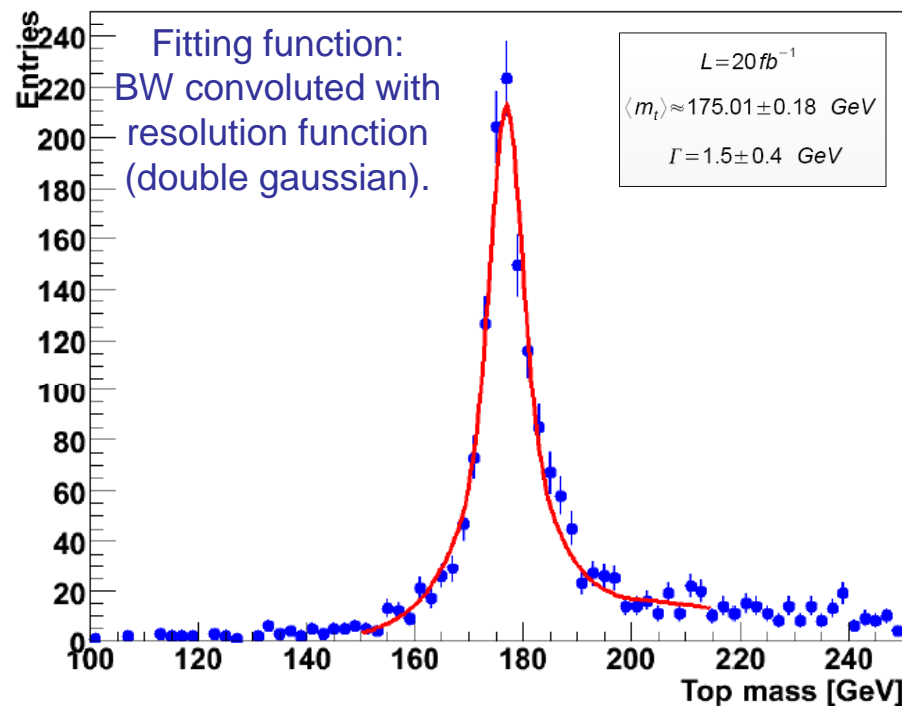
FTD (acceptance down to 7 degrees in polar angle)

**High granularity calorimeters:**

ECAL W-Si, 23x<sub>o</sub>, 1x1 cm<sup>2</sup>

HCAL Iron-Scintillator, ~4-6λ, 3x3 cm<sup>2</sup>

### # top quark invariant mass



## Conclusion

Top quark invariant mass reconstructed!!!

- full ILD detector simulation
- full hadronic bgd
- complete reco

Result extrapolated for higher lumi:

$$\Delta m_{\text{top}} = 50 \text{ MeV (stat) for } 300 \text{ fb}^{-1}$$

Analysis will be used for  
detector optimization  
and performance studies  
and should be included in LOI

# A 4<sup>th</sup> generation scenario (François Richard)

Expect the unexpected (at the LHC) ! Like a heavy Higgs, or heavy fermions without SUSY.

Why not through a 4<sup>th</sup> generation? It is possible (!!):

The 4<sup>th</sup> chiral generation is excluded by S/T constraints only for the mass degenerate case !!

$$\Delta T \sim \frac{\Delta m^2}{(150 \text{ GeV})^2}$$

$$\Delta S = \frac{N_c}{6\pi} \left( 1 - 2Y \ln \frac{m_u^2}{m_d^2} \right)$$

hence when  $\Delta T=0$  and  $\Delta S=3/4\pi$ .

So one can play with the positive correlation between these variables and easily pass the constraints when fermions are partially degenerate in mass

## The main motivation for more than 3 generations is baryogenesis

Baryogenesis needs C+CP violation & strong EW 1st order transition

SM alone ?

Not enough CPV and insufficient EW transition

MSSM alone ?

New phases but severely constrained by EDM and EW transition not strong enough  
→ extra particles needed, strongly coupled to the Higgs field, scalars or fermions  
(cf. Carena et al. hep-ph 0410352)

With a 4<sup>th</sup> generation :

- CPV fine
  - Large Yukawa couplings to Higgs field (**strong** at scale  $\sim \text{TeV}$ )
    - However not enough to get the right EW transition
- Hence include SUSY → **4MSSM** (R. Fok G. Kribs [arXiv:0803.4207](https://arxiv.org/abs/0803.4207))

## A 4<sup>th</sup> generation scenario (François Richard)

### Predictions and present day results

- $300 < m_{t',b'} < 450$  GeV + lighter leptons
- Squarks ~ mass degenerate with quarks
- **Higgs could be heavy** through radiative corrections
- Spectacular & early signals at LHC
- Accessible at a TeV LC (heavy leptons)
- Interesting Tevatron results : searches exclude  $m_{t'} < 260$  GeV &  $140 < m_H < 180$  GeV and b-sector results could point to 4<sup>th</sup> gen.

$$m_h^2 = \sum_{f=t,t',b'} \frac{3}{2\pi^2} \frac{m_f^4}{v^2} \ln \frac{m_f^2}{m_{f'}^2}.$$

### Outlook and Conclusion

- At LHC an early discovery of the new fermions + SUSY squarks
- In addition to a heavy Higgs, which can easily be observed
- ILC very powerful in particular for leptons and for a light Higgs
  
- Requires SUSY but cannot be extrapolated to GUT because of the large Yukawa constants
- Allows for a **heavy Higgs** within SUSY
- Rich physics for a TeV LC
- → Could serve as an illustration of LHC/LC complementarity



# STU in mHDM: Precision constraints on multi-Higgs doublet models (Per Osland based on work with W.Grimus, L.Lavoura, OM OGREID)

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W}$$

where  $\rho - 1 = \Delta\rho \approx \alpha T$

## Definitions of S, T, U

- M. Peskin & T. Takeuchi, PRL 65 (1990) 964, PRD 46 (1992) 381
- G. Altarelli, R. Barbieri, PL B 253 (1991) 161
- G. Altarelli, R. Barbieri, S. Jadach, NPB 369 (1992) 3
- I. Maksymyk, C.P. Burgess, D. London, PRD 50 (1994) 529
- R. Barbieri, A. Pomarol, R. Rattazzi, A. Strumia, NPB 703 (2004) 127

Recent papers: more observables

Early papers defined S, T, U in terms of first derivatives of self energy function at the origin

Maksymyk et al define S, T, U, V, W, X in terms of differences  
BPRS introduce second derivatives

## Maksymyk et al (1994)

$$\frac{\alpha}{4s_W^2 c_W^2} S = \frac{A_{ZZ}(m_Z^2) - A_{ZZ}(0)}{m_Z^2} - \frac{\partial A_{\gamma\gamma}(q^2)}{\partial q^2} \Big|_{q^2=0} + \frac{c_W^2 - s_W^2}{c_W s_W} \frac{\partial A_{\gamma Z}(q^2)}{\partial q^2} \Big|_{q^2=0}$$

$$\alpha T = \frac{A_{WW}(0)}{m_W^2} - \frac{A_{ZZ}(0)}{m_Z^2},$$

$$\frac{\alpha}{4s_W^2} U = \frac{A_{WW}(m_W^2) - A_{WW}(0)}{m_W^2} - c_W^2 \frac{A_{ZZ}(m_Z^2) - A_{ZZ}(0)}{m_Z^2} - s_W^2 \frac{\partial A_{\gamma\gamma}(q^2)}{\partial q^2} \Big|_{q^2=0} + 2c_W s_W \frac{\partial A_{\gamma Z}(q^2)}{\partial q^2} \Big|_{q^2=0},$$

$$\alpha V = \frac{\partial A_{ZZ}(q^2)}{\partial q^2} \Big|_{q^2=m_Z^2} - \frac{A_{ZZ}(m_Z^2) - A_{ZZ}(0)}{m_Z^2},$$

$$\alpha W = \frac{\partial A_{WW}(q^2)}{\partial q^2} \Big|_{q^2=m_W^2} - \frac{A_{WW}(m_W^2) - A_{WW}(0)}{m_W^2},$$

$$\frac{\alpha}{s_W c_W} X = \frac{\partial A_{\gamma Z}(q^2)}{\partial q^2} \Big|_{q^2=0} - \frac{A_{\gamma Z}(m_Z^2)}{m_Z^2}.$$

There are 3 kinds of scalars

- scalar SU(2) doublets (1 charged and 1 neutral component)
  - charged SU(2) singlets
  - neutral SU(2) singlets

where the neutral fields have vevs

## STU in mHDM: Precision constraints on multi-Higgs doublet models (Per Osland based on work with W.Grimus, L.Lavoura, OM OGREID)

Scalars couple to  $W$  and  $Z$  via covariant derivatives, except for the neutral singlets  $S, T, U$  expressions were shown in the talk while  $V, W, X$  given in arXiv0802.4353

### Applications with 2HDM :

- In the limit of  $m_Z = m_W$ , with partial degeneracy  $m_2 = (\text{charged}) = \mu_3 = \mu_4$  (neutrals),  $T \sim 0$
- In the limit of CP conservation while applying the custodial symmetry  $m_2 = \mu_4$ ,  $T \sim 0$ 
  - In twisted custodial symmetry,  $T$  does not vanish but can be made close to 0 by a suitable choice of masses.
- For inert or dark scalar, with  $\mu_4$  being the SM Higgs  $T$  will vanish in the limits of  $m_2 = \mu_3$  or  $m_2 = \mu_4$

### Summary

- $S, T, U, V, W, X$  calculated for general scalar sector, consisting of doublets and singlets
- Result expressed in terms of **two** mixing matrices and simple functions of masses:
  - $T$ : **one** function
  - $S, U$ : **2** functions
- In general ( $S, T, U, V, W, X$ ): **5** functions of masses

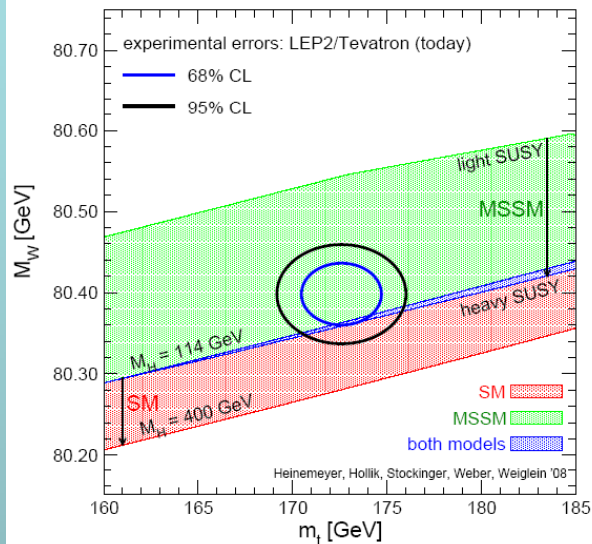
# Constraining SUSY with Electroweak Precision Observables (Sven Heinemeyer)

based on collaboration with W.Hollik, AM Weber and G.Weiglein)

Search for indirect effects of SUSY via Precision Observables (POs)  
by comparing EW POs ( $m_W, \sin^2\theta_{\text{eff}}, \Gamma_Z, \dots$ ) with theory (SM, MSSM, ...)  
Test of the theory at quantum level by its sensitivity to loop corrections.  
Very high accuracy of measurements and theoretical predictions needed.

Example: Prediction for  $M_W$  in the SM and the MSSM :

[S.H., W. Hollik, D. Stockinger, A.M. Weber, G. Weiglein '07]



MSSM band:  
scan over  
SUSY masses

overlap:  
SM is MSSM-like  
MSSM is SM-like

SM band:  
variation of  $M_H^{\text{SM}}$

Results for EWPO fit in the SM: [LEPEWWG '08]

Pull distributions:

Variable	Measurement	Fit	$ \text{O}^{\text{meas}} - \text{O}^{\text{fit}}  / \sigma^{\text{meas}}$
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	$0.02758 \pm 0.00035$	0.02768	0.1
$m_Z$ [GeV]	$91.1875 \pm 0.0021$	91.1875	0.0
$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	2.4957	0.2
$\sigma_{\text{had}}^0$ [nb]	$41.540 \pm 0.037$	41.477	1.7
$R_l$	$20.767 \pm 0.025$	20.744	0.9
$A_{\text{fb}}^{0,l}$	$0.01714 \pm 0.00095$	0.01645	0.7
$A_1(P_e)$	$0.1465 \pm 0.0032$	0.1481	0.5
$R_b$	$0.21629 \pm 0.00066$	0.21586	0.6
$R_c$	$0.1721 \pm 0.0030$	0.1722	0.0
$A_{\text{fb}}^{0,b}$	$0.0992 \pm 0.0016$	0.1038	2.8
$A_{\text{fb}}^{0,c}$	$0.0707 \pm 0.0035$	0.0742	1.0
$A_b$	$0.923 \pm 0.020$	0.935	0.6
$A_c$	$0.670 \pm 0.027$	0.668	0.1
$A_1(\text{SLD})$	$0.1513 \pm 0.0021$	0.1481	1.5
$\sin^2\theta_{\text{eff}}^{\text{lept}}(Q_{\text{fb}})$	$0.2324 \pm 0.0012$	0.2314	0.8
$m_W$ [GeV]	$80.398 \pm 0.025$	80.374	0.9
$m_t$ [GeV]	$170.9 \pm 1.8$	171.3	0.2
$\Gamma_W$ [GeV]	$2.140 \pm 0.060$	2.091	0.8

Probability: 15%



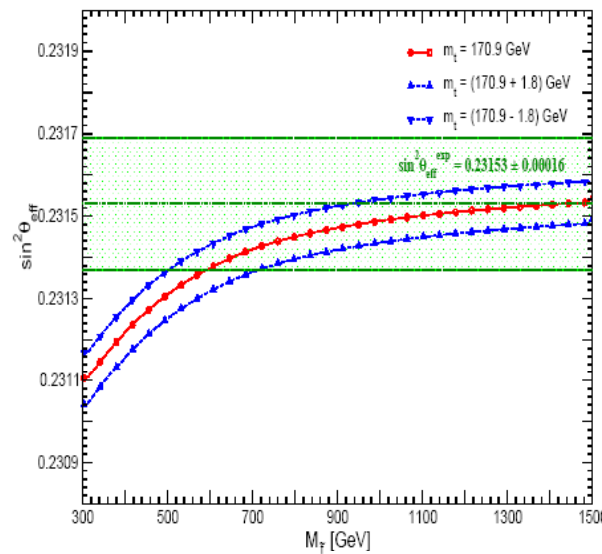
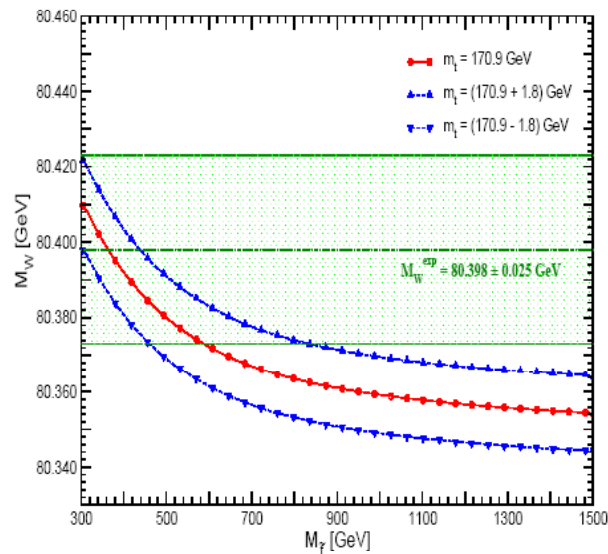
# Constraining SUSY with Electroweak Precision Observables (Sven Heinemeyer based on collaboration with W.Hollik, AM Weber and G.Weiglein)

Perform fit  
with MSSM predictions  
of EW POs

observable	central exp. value	$\sigma \equiv \sigma^{\text{today}}$	$\sigma^{\text{LHC}}$	$\sigma^{\text{ILC/GigaZ}}$
$M_W$ [GeV]	80.398	0.025	0.015	0.007
$\sin^2 \theta_{\text{eff}}$	0.23153	0.00016	0.00020–0.00014	0.000013
$\Gamma_Z$ [GeV]	2.4952	0.0023	—	0.001
$R_t$	20.767	0.025	—	0.01
$R_b$	0.21629	0.00066	—	0.00014
$\sigma_{\text{had}}^0$	41.540	0.037	—	0.025

⇒ ILC/GigaZ precision yields a very strong improvement

## A) $M_{\text{SUSY}}$ and $m_t$ dependence (I)



⇒ strong  $M_{\text{SUSY}}$  dependence  
⇒ important  $m_t$  dependence

$m_W$ ,  $\sin^2 \theta_{\text{eff}}$  strong dependence  
while  $\Gamma_Z$  relevant dependence  
on  $m_t$  and  $m_{\text{SUSY}}$

$m_W$ ,  $\sin^2 \theta_{\text{eff}}$ ,  $\Gamma_Z$   
strong dependence  
on  $\mu$  and on  $M_2$  for small  $M_2$

Also studied complex phases  
in the squark sector  
dependence  
of EW POs predictions

Split SUSY model predictions  
not at all dependant.

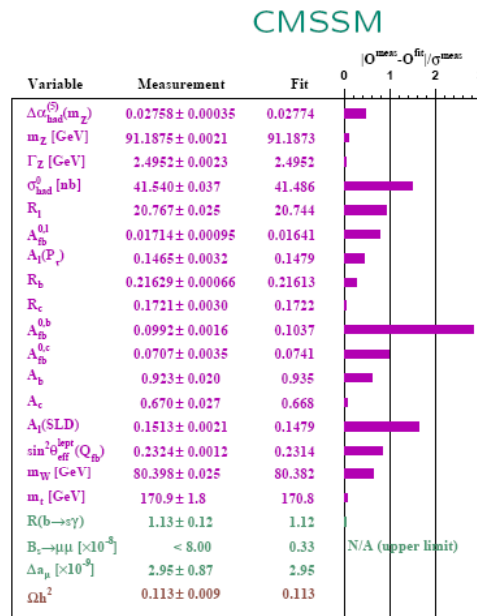
# Constraining SUSY with Electroweak Precision Observables (Sven Heinemeyer based on collaboration with W.Hollik, AM Weber and G.Weiglein)

In a scenario with no SUSY particles at the LHC,  
looking at  $\sin^2\theta_{\text{eff}}$   
the ILC (1000)/GigaZ could detect SUSY directly or indirectly.

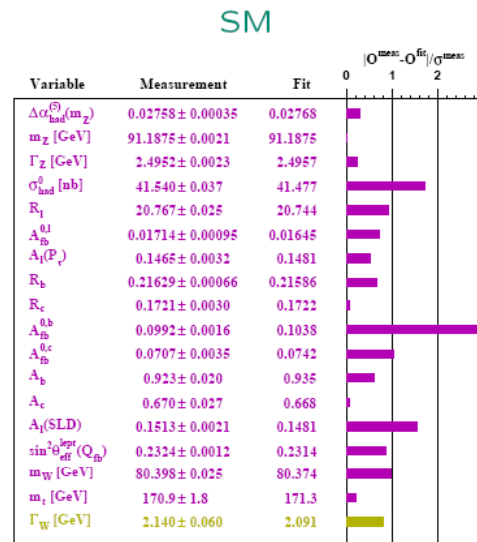
Made studies of the CMSSM/mSUGRA

## Pull distributions:

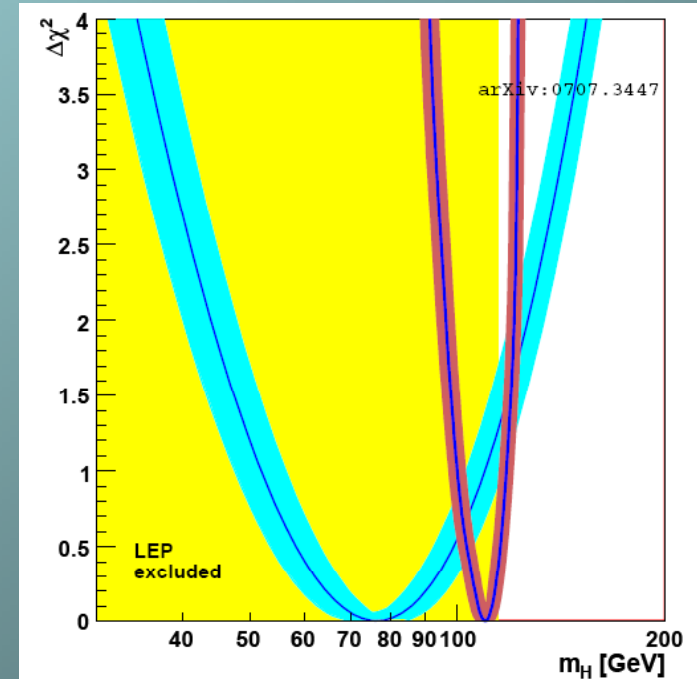
[Buchmüller, Cavanaugh, de Roeck, S.H., Isidori, Paradisi, Ronga, Weber, Weiglein '07]



Probabilities: 24% / 20%



15% / 15% (incl. / excl.  $M_h$ )



Red band (CMSSM)  
versus blue band (SM) !!!

# Constraining SUSY with Electroweak Precision Observables (Sven Heinemeyer based on collaboration with W.Hollik, AM Weber and G.Weiglein)

## 4. Conclusions

- Precision observables
  - can give valuable information about the “true” Lagrangian
  - can provide bounds on SUSY parameter space
- SM: Blue band plot:  $\Rightarrow M_H^{\text{SM}} = 87_{-27}^{+36}$  GeV (too light for LEP bounds?)
- electroweak precision observables (EWPO):  
 $\mathcal{O} = M_W, \sin^2 \theta_{\text{eff}} (A_{\text{FB}}^{b,c}, A_{\text{LR}}^{e,\mu}), R_l, R_b, \sigma_0^{\text{had}}, \dots$
- best MSSM prediction = full (available) SM result
  - + all existing MSSM corrections
  - e.g. full 1L incl. complex phases
  - double counting
- SUSY dependencies:
  - strong dependence only for  $M_W, \sin^2 \theta_{\text{eff}}, \Gamma_Z$
  - strong dependence on  $M_{\text{SUSY}}, \mu, M_2, m_t, \dots$
  - strong dependence on  $\phi_{A_t}$
  - strong dependence on  $\phi_{A_b}$  for large  $\tan \beta$
- CMSSM/mSUGRA: Red band plot:  $\Rightarrow M_h^{\text{CMSSM}} = 110 \pm 8 \pm 3$  GeV

## Personal conclusions:

very nice session with extremely clear talks  
and strong dynamical audience dependence