

## 4 SOLENOID HOMEWORK

1. Consider a 200 MeV/c particle starting on the axis with a transverse momentum of 20 MeV/c in an axial solenoidal field of 3.33 T.
  - (a) What is its motion in the lab frame and out to what transverse distance from the axis does it get.
  - (b) What is the distance along the axis before it first returns to that axis?
  - (c) What is the wavelength  $\lambda$  in the Larmor frame?
  - (d) what is the lattice parameter  $\beta_{\perp}$  for that particle

## Answers

(a) It's motion is helical with radius in the lab frame from eq.25 of

$$\rho = \frac{[pc/e]_{\perp}}{B c} = \frac{20 \cdot 10^6}{3 \cdot 10^8 \cdot 3.33} = 2 \cdot 10^{-2} \text{ m, or 2 cm. The helix pitch angle is } 0.1 \text{ radians}$$

Its furthest out will be at  $2 \times \rho = 4 \text{ cm}$ .

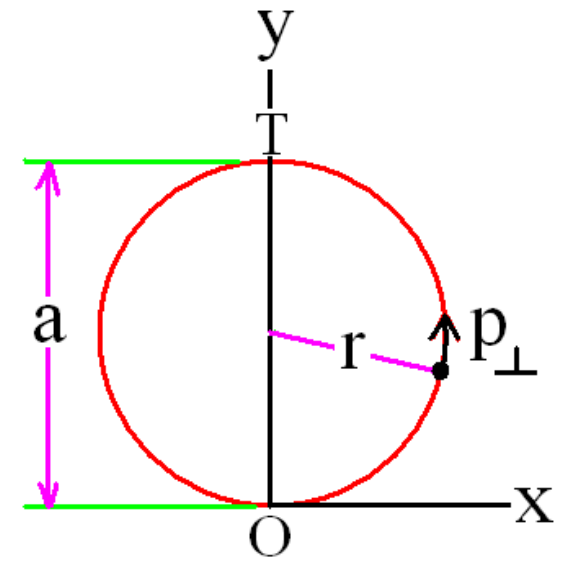
(b) This distance is the helix period in the lab frame from eq.30

$$\lambda_{\text{helix}} = 2\pi\rho \frac{p_z}{p_{\perp}} = 2\pi \cdot 0.02 \cdot \frac{200}{20} = 1.256 \text{ m}$$

(c) From eq.31  $\lambda_{\text{Larmor}} = 2 \lambda_{\text{helix}} = 2.512 \text{ m}$

(d) From eq.32:  $\beta_{\perp} = \frac{\lambda_{\text{Larmor}}}{2\pi} = \frac{2 [pc/e]_z}{B c} = 0.4 \text{ m}$

2. Consider again a  $200 \text{ MeV}/c$  particle starting on the  $z$  axis with a transverse momentum of  $20 \text{ MeV}/c$  in an axial solenoidal field of  $3.33 \text{ T}$ . After a distance
- A) corresponding to  $1/2$  a helix rotation, or
  - B) corresponding to a full helix rotation determine, the field abruptly doubles to  $6.66 \text{ T}$ . In the two cases determine:
- (a) The shape of the motion projected onto the  $x, y$  plane
  - (b) The following length in  $z$  for one helix rotation ( $\lambda_{\text{helix}}$ )



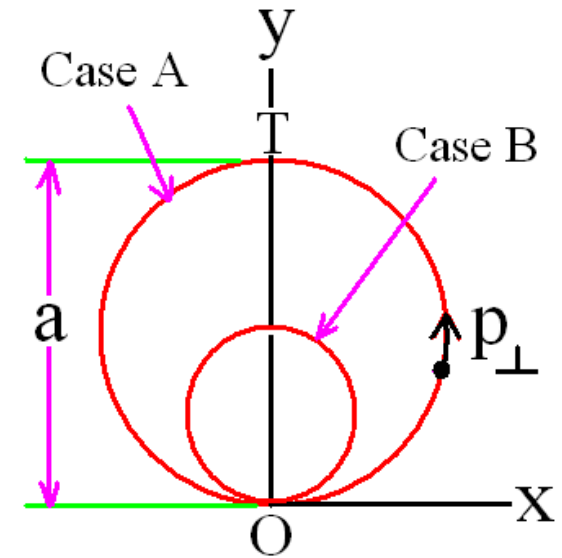
## Answers

- (a) In case A) the field doubles as the muon is at the top (T) of its rotation. At that point the particle receives a kick (eq.19) of

$$\Delta p_{\perp} = ac \frac{\Delta B_z}{[pc/e]_z} = 20 \text{ MeV}/c$$

so now  $p_{\perp} = 2 \times 20 = 40 \text{ MeV}/c$

At the same time the axial field doubles, so the transverse bending radius remains the same = 2 cm

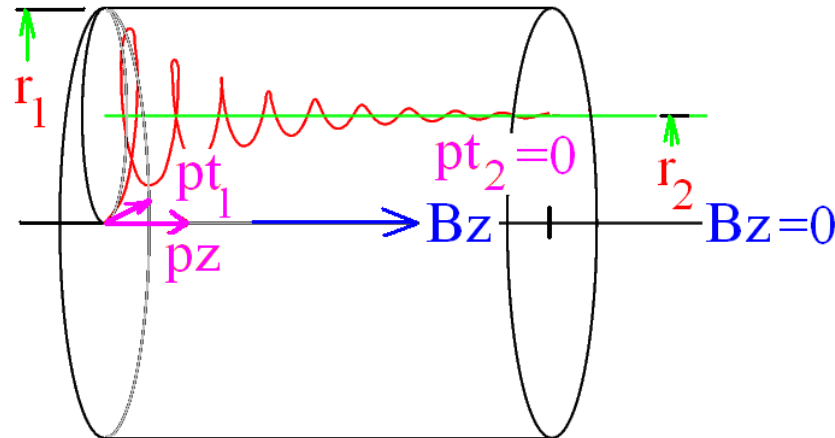


In case B The field doubles when the particle is on the magnet axis and there is no transverse momentum kick, and the transverse momentum remains at 20 MeV/c. Now with double the axial field, the transverse radius is halved to 1 cm.

- (b) For both cases A and B, the longitudinal momentum is approximately unchanged, and the wavelengths (eq.30)  $\propto 1/B_z$  the helix wavelength is halved to  $1.256/2 = 0.628 \text{ m}$ .

## 7 IONIZATION COOLING HOMEWORK

1. Again consider a solenoid with  $B_z = 3.33$  T and a muon with starting on axis with  $p_t = 20$  MeV/c and  $p_z = 200$  MeV/c. Imagine an ideal transverse cooling system with continuous energy loss and re-acceleration so that all transverse momenta are reduced to near zero, then the above particle will settle at half its maximum distance from the axis  $r_2 = r_1/2$  and pass straight down the field lines at  $p_z$ .



- (a) What now is its motion in the Larmor frame?
- (b) If now the field  $B_z$  suddenly stops, what is the further motion of the muon?
- (c) Taking the initial phase space to be  $\Phi_1 = p_{t1} r_1$ , what is the final phase space if it is defined as  $\Phi_2 = p_{t3} r_3$ ?
- (d) This (from Bush's Law) is not good. What could one do to avoid it ?

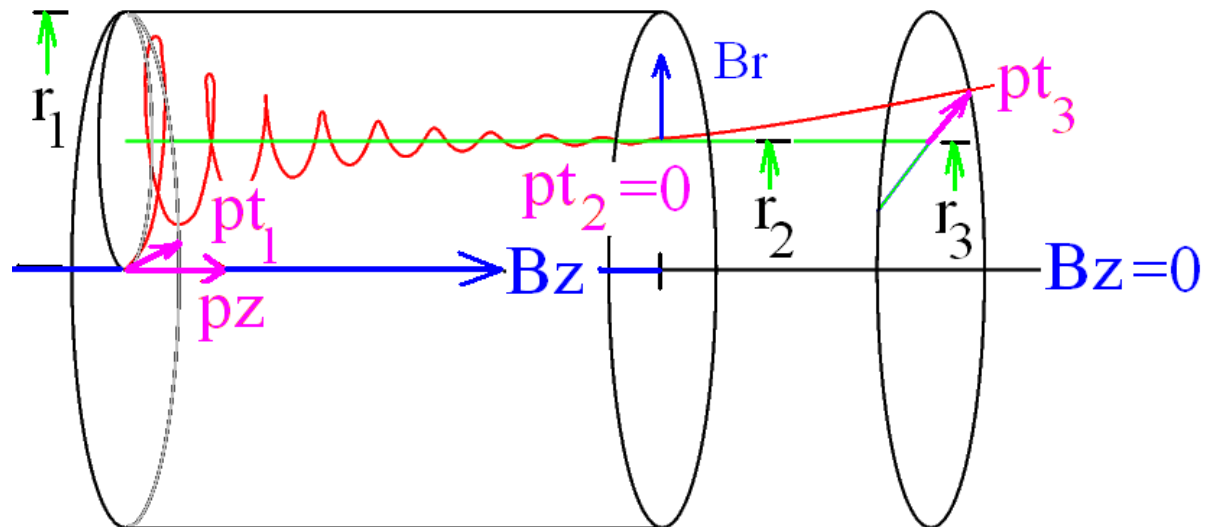
## Answers

- (a) The motion in the Larmor plane will be a helix about the axis with radial amplitude  $r_2 = r_1/2$ , wavelength 2.5 m, and transverse momentum of  $pt$
- (b) When it leaves the magnet the muon will see radial field components that will give it a lateral kick

$$[ptc/e]_3 = \Delta[ptc/e] = + \frac{r_2 c B_z}{2} = \frac{[ptc/e]_1}{2} = 10 \text{ MeV}/c$$

Thus  $\epsilon_3 = \frac{\epsilon_1}{2}$  not a very good gain.

- (c) The problem arises because the cooling reduces the lab angular momentum inside the solenoid. This causes the canonical angular momentum to grow (eq.23). Periodic reversals of  $B_z$  during the cooling avoids a build up of canonical angular momentum



2. Assume a transverse cooling channel with, at the absorber,  $\beta_{\perp} = 0.18$  m,  $\beta_v = v/c = 0.85$ , and the dimensionless constant defined by eq.35:  $C(mat, E) = 45 \cdot 10^{-4}$ , and emittance  $\epsilon = 2.2$  (pi mm rad). These are the values for a lattice given in Feasibility study 2. The value of  $C(mat, E)$  differs somewhat from the value for pure liquid hydrogen ( $C \approx 35 \cdot 10^{-4}$ ) because of the effect of the hydrogen and rf windows.

a) calculate the theoretical equilibrium emittance.

b) calculate the final differential rate of cooling  $d\epsilon_{\perp}/\epsilon_{\perp}$  compared with its ideal value  $dp/p$ .

c) This rate would be improved if the lattice parameter  $\beta_{\perp}$  were lowered. What is the likely problem that will be encountered if this approach were followed?

## answers

a) calculate the theoretical equilibrium emittance.

eq.37

$$\epsilon_{x,y}(\min) = \frac{\beta_{\perp}}{\beta_v} C(\text{mat}, E)$$

so

$$\epsilon_{\perp}(\min) = \frac{45 \cdot 10^{-4} \cdot 0.18}{0.85} = 950 \text{ } (\pi \text{ mm mrad})$$

b) calculate the final differential rate of cooling  $d\epsilon_{\perp}/\epsilon_{\perp}$  compared with its ideal value  $dp/p$ .

eq.38

$$\frac{d\epsilon}{\epsilon}(\text{end}) = \left(1 - \frac{\epsilon(\min)}{\epsilon}\right) \frac{dp}{p} \approx 0.57 \frac{dp}{p}$$

c) This rate would be improved if the lattice parameter  $\beta_{\perp}$  were lowered. What is the likely problem that will be encountered if this approach were followed?

This rate can be improved if  $\frac{\epsilon(\min)}{\epsilon}$  is reduced by decreasing  $\epsilon(\min)$ , and this can be accomplished by reducing the lattice parameter  $\beta_{\perp}$ . The main problem is that if one does this, is that the angular spread of the resulting beam, and thus required angular aperture rises because eq.39:

$$\sigma_{\theta} = \sqrt{\frac{\epsilon_{\perp}}{\beta_{\perp} \beta_v \gamma}}$$



3. In the longitudinal cooling section, we describe a cooling ring that, with emittance exchange in wedges cools all 6 dimensions. Assume  $\beta_{\perp} = 0.4$  m, dispersion at the hydrogen wedge  $D = 7$  cm, the length of the wedge on axis  $\ell = 28.6$  cm, and the height from the axis to the apex of the wedge  $h = \frac{\ell}{2 \tan(100^\circ/2)} = 12$  cm. Assume that the sum of partition functions  $\sum J_i \approx 2.0$ ,  $C(mat, E) = 38 \cdot 10^{-4}$ , and assume good mixing between  $x$  and  $y$ . As before,  $\beta_v = 0.85$ .

- a) What are the three partition functions in this case?
- b) What is the expected equilibrium transverse emittance?

answer

a) What are the three partition functions in this case?

The longitudinal partition function is thus

$$J_z = \frac{D}{h} = 0.58$$

and since there is good mixing between  $x$  and  $y$ ,  $J_x = J_y$

From  $\sum J_i \approx 2.0$ :

$$J_x + J_y + 0.58 = 2.0$$

$$J_x = J_y \approx \frac{2 - 0.58}{2} = 0.71$$

i.e. The wedge angle gave somewhat less cooling (0.58) in the longitudinal direction than in the other two (.71) .

b) What is the expected equilibrium transverse emittance?

The equilibrium emittances, with values of  $J_{x,y} \neq 1$  from eq.45

$$\epsilon_{\perp}(\text{min}) = \frac{C \beta_{\perp}}{J_{x,y} \beta_v} = \frac{38 \cdot 10^{-4} \cdot 0.4}{0.71 \cdot 0.85} = 2.5 \quad (\pi \text{ mm})$$