## Physics studies with polarisation

Mikael Berggren<sup>1</sup>

<sup>1</sup>DESY, Hamburg

BAW-II, SLAC, Jan 2011

Mikael Berggren (DESY)

Physics studies with polarisation

Outline

#### Outline

- Introduction.
- Examples
  - $\tilde{\tau}$  in SPS1a'.
  - Model independent WIMPs
  - TGCs and polarisation
  - Near degenerate  $\tilde{e}$
- Conclusions.

Polarisation is

#### Needed to analyse the chiral structure of interactions

• Key observable ALR.

• Relative error goes as that of  $P_{eff} = (P_{e^-} - P_{e^+})/(1 - P_{e^-}P_{e^+})$ 

- Useful to improve S/B:
  - Key-number  $\mathcal{L}_{eff} = (1 P_{e^-} P_{e^+}) \mathcal{L}/2$
  - Useful even if S and B depends on P in the same way: error ∝ S/√B!

• See:

- Overview in hep-ph/0507011, Phys.Rept., 460 (2008).
- Sabine's talk.

I'll talk about a few full simulation analyses demonstrating that the conclusions hold also when other beam-aspects and detector effects are included.

Polarisation is

- Needed to analyse the chiral structure of interactions
  - Key observable ALR.

• Relative error goes as that of  $P_{eff} = (P_{e^-} - P_{e^+})/(1 - P_{e^-}P_{e^+})$ 

- Useful to improve S/B:
  - Key-number  $\mathcal{L}_{eff} = (1 P_{e^-}P_{e^+})\mathcal{L}/2$
  - Useful even if S and B depends on P in the same way: error  $\propto S/\sqrt{B}!$

See:

- Overview in hep-ph/0507011, Phys.Rept., 460 (2008).
- Sabine's talk.

I'll talk about a few full simulation analyses demonstrating that the conclusions hold also when other beam-aspects and detector effects are included.

Polarisation is

- Needed to analyse the chiral structure of interactions
  - Key observable ALR.
  - Relative error goes as that of  $P_{eff} = (P_{e^-} P_{e^+})/(1 P_{e^-}P_{e^+})$
- Useful to improve S/B:
  - Key-number  $\mathcal{L}_{eff} = (1 P_{e^-}P_{e^+})\mathcal{L}/2$
  - Useful even if S and B depends on *P* in the same way: error  $\propto S/\sqrt{B}!$
- See:
  - Overview in hep-ph/0507011, Phys.Rept., 460 (2008).
  - Sabine's talk.

I'll talk about a few full simulation analyses demonstrating that the conclusions hold also when other beam-aspects and detector effects are included.

Polarisation is

- Needed to analyse the chiral structure of interactions
  - Key observable ALR.

• Relative error goes as that of  $P_{eff} = (P_{e^-} - P_{e^+})/(1 - P_{e^-}P_{e^+})$ 

- Useful to improve S/B:
  - Key-number  $\mathcal{L}_{eff} = (1 P_{e^-}P_{e^+})\mathcal{L}/2$
  - Useful even if S and B depends on *P* in the same way: error  $\propto S/\sqrt{B}!$
- See:
  - Overview in hep-ph/0507011, Phys.Rept., 460 (2008).
  - Sabine's talk.

I'll talk about a few full simulation analyses demonstrating that the conclusions hold also when other beam-aspects and detector effects are included.

## $\tilde{\tau}$ in SPS1a' and polarisation

(Work by J. List, P. Bechtle, P. Schade, M.B., PRD 82,no5 (2010), arXiv:0908.0876)

SPS1a' is a pure mSUGRA model, just outside what is excluded by LEP and low-energy observations. Compatible with WMAP, with  $\tilde{\chi}_1^0$  Dark Matter.

- At  $E_{CMS} = 500$  GeV:
  - All sleptons available.
  - No squarks.
  - Lighter bosinos, up to  $\tilde{\chi}^0_3$  (in  $e^+e^- \rightarrow \tilde{\chi}^0_1 \tilde{\chi}^0_3$ )

### Features of $\tilde{\tau}$ :s in SPS1a'

- In SPS1a', the  $\tilde{\tau}_1$  is the NLSP.
- For  $\tilde{\tau}_1$ :  $E_{\tau,min} = 2.6 \text{ GeV}$ ,  $E_{\tau,max} = 42.5 \text{ GeV}$ :  $\gamma\gamma - background \Leftrightarrow pairs - background$ .
- For  $\tilde{\tau}_2$ :  $E_{\tau,min} = 35.0 \text{ GeV}, E_{\tau,max} = 152.2 \text{ GeV}$ :  $WW \rightarrow l\nu l\nu$  background  $\Leftrightarrow$  *Polarisation*.
- $\tilde{\tau}$  NLSP  $\rightarrow \tau$ :s in most SUSY decays  $\rightarrow$  SUSY is background to SUSY.
- For pol=(-1,1):  $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$  and  $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-)$  = several hundred fb and BR(X $\rightarrow \tilde{\tau}$ ) > 50 %. For pol=(1,-1):  $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$  and  $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-) \approx 0$ .

Polarisation = (+,-) assumed.

▲□▶ ▲圖▶ ▲目▶ ▲目▶ 三回日 のへ⊙

## Features of $\tilde{\tau}$ :s in SPS1a'

- In SPS1a', the  $\tilde{\tau}_1$  is the NLSP.
- For  $\tilde{\tau}_1$ :  $E_{\tau,min} = 2.6 \text{ GeV}$ ,  $E_{\tau,max} = 42.5 \text{ GeV}$ :  $\gamma\gamma - background \Leftrightarrow pairs - background$ .
- For  $\tilde{\tau}_2$ :  $E_{\tau,min} = 35.0 \text{ GeV}, E_{\tau,max} = 152.2 \text{ GeV}$ :  $WW \rightarrow l\nu l\nu$  background  $\Leftrightarrow$  *Polarisation*.
- For pol=(-1,1):  $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$  and  $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-)$  = several hundred fb and BR(X $\rightarrow \tilde{\tau}$ ) > 50 %. For pol=(1,-1):  $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$  and  $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-) \approx 0$ .

Polarisation = (+,-) assumed.

▲□▶ ▲圖▶ ▲目▶ ▲目▶ 三回日 のへ⊙

# Extracting the $\tilde{\tau}$ properties

From decay kinematics:

- $M_{\tilde{\tau}}$  from  $M_{\tilde{\chi}_1^0}$  and end-point of spectrum =  $E_{\tau,max}$ .
- Need to measure end-point of spectrum.
- Must get  $M_{\tilde{\chi}^0_{+}}$  from other sources.

From cross-section:

• 
$$\sigma_{\widetilde{\tau}} = A(\theta_{\widetilde{\tau}}, \mathcal{P}_{beam}) \times \beta^3 / s$$
, so

• 
$$M_{\tilde{\tau}} = E_{beam} \sqrt{1 - (\sigma s/A)^{2/3}}$$
: no  $M_{\tilde{\chi}_1^0}$ !

- Only the upper end-point is relevant.
- Background subtraction:
  - *τ˜*<sub>1</sub>: Important SUSY background,but region above 45 GeV is signal free. Fit exponential and extrapolate.
  - *τ˜*<sub>2</sub>: ~ no SUSY background above 45 GeV. Take background from SM-only simulation and fit exponential.
- Fit line to (data-background fit).

- Only the upper end-point is relevant.
- Background subtraction:
  - *τ˜*<sub>1</sub>: Important SUSY
     background,but region
     above 45 GeV is signal free.
     Fit exponential and
     extrapolate.
  - *τ˜*<sub>2</sub>: ~ no SUSY background above 45 GeV. Take background from SM-only simulation and fit exponential.
- Fit line to (data-background fit).



・ 何 ト ・ ヨ ト ・ ヨ ト ・ ヨ

- Only the upper end-point is relevant.
- Background subtraction:
  - *τ˜*<sub>1</sub>: Important SUSY
     background,but region
     above 45 GeV is signal free.
     Fit exponential and
     extrapolate.
- Fit line to (data-background fit).



(B)

- Only the upper end-point is relevant.
- Background subtraction:
  - *τ˜*<sub>1</sub>: Important SUSY
     background,but region
     above 45 GeV is signal free.
     Fit exponential and
     extrapolate.
  - <sup>˜</sup><sub>2</sub>: ~ no SUSY background above 45 GeV. Take background from SM-only simulation and fit exponential.
- Fit line to (data-background fit).



4 3 > 4 3

## Fitting the $\tilde{\tau}$ mass: Cross-section

- Poorly known SUSY background is most important contribution to uncertainty.
- Select region where is is as low as possible.

## Fitting the $\tilde{\tau}$ mass: Cross-section

- Poorly known SUSY background is most important contribution to uncertainty.
- Select region where is is as low as possible.



## Fitting the $\tilde{\tau}$ mass: Cross-section

- Poorly known SUSY background is most important contribution to uncertainty.
- Select region where is is as low as possible.



## RDR , SB2009 and $\tilde{\tau}$ :s

Potential effects on the  $\tilde{\tau}$ -channels:

- Decrease of P(e<sup>+</sup>): Less signal, more background for τ
  <sub>1</sub>, and more signal, but still more background for τ
  <sub>2</sub>
- Incoming energy-spread grows: end-point blurred.
- Luminosity within 1 % of nominal reduced: lower signal.
- Twice as much beam-strahlung:
  - more overlayed tracks (real or fake): Destroys  $\tau$  topology.
  - Twice as much energy in BeamCal: More  $\gamma\gamma$ .
- Higher probability for a γγ event in the same BX as the physics event (this effect has not yet been studied).
- Also: Total luminosity decrease for SB2009 w/o TF.

Impossible to re-do fully simulated samples - months of Grid-time - , so a transformation from RDR to SB2009 is needed !

- BeamCal: Double the energy density wrt. numbers from RDR simuation.
- Tracking: Fully simulated and reconstructed BX:es available both for RDR and SB2009. Choose which set to overlay.
- Beam-spectrum Reweight events with Lumi distributions from GP (A. Hartin, T Barklow) (avaiable both for RDR and SB2009) to modify the fully simulated samples to an SB2009 one.
- Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

Impossible to re-do fully simulated samples - months of Grid-time - , so a transformation from RDR to SB2009 is needed !

- BeamCal: Double the energy density wrt. numbers from RDR simuation.
- Tracking: Fully simulated and reconstructed BX:es available both for RDR and SB2009. Choose which set to overlay.
- Beam-spectrum Reweight events with Lumi distributions from GP (A. Hartin, T Barklow) (avaiable both for RDR and SB2009) to modify the fully simulated samples to an SB2009 one.
- Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

Impossible to re-do fully simulated samples - months of Grid-time - , so a transformation from RDR to SB2009 is needed !

- BeamCal: Double the energy density wrt. numbers from RDR simuation.
- Tracking: Fully simulated and reconstructed BX:es available both for RDR and SB2009. Choose which set to overlay.
- Beam-spectrum Reweight events with Lumi distributions from GP (A. Hartin, T Barklow) (avaiable both for RDR and SB2009) to modify the fully simulated samples to an SB2009 one.
- Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

Impossible to re-do fully simulated samples - months of Grid-time - , so a transformation from RDR to SB2009 is needed !

- BeamCal: Double the energy density wrt. numbers from RDR simuation.
- Tracking: Fully simulated and reconstructed BX:es available both for RDR and SB2009. Choose which set to overlay.
- Beam-spectrum Reweight events with Lumi distributions from GP (A. Hartin, T Barklow) (avaiable both for RDR and SB2009) to modify the fully simulated samples to an SB2009 one.
- Polarisation: Straight-forward relative weighting of generated samples with *P*=(-1,1) and *P*=(1,-1).

## RDR , SB2009 and $\tilde{\tau}$ :s: Effect on end-results

Errors on end-point (GeV)

case	#	$\tilde{\tau}_1$	$ ilde{ au}_2$
RDR	1	0.129	1.83
+SB bck	2	0.144	2.02
+SB ppol	3	0.153	2.06
+SB spect	4	0.152	2.10
+SB noTF	5	0.179	2.42

Errors on cross-section (%)

case	#	$ ilde{ au}_1$	$ ilde{ au}_2$
RDR	1	2.90	4.24
+SB bck	2	3.03	4.72
+SB ppol	3	3.31	4.77
+SB spect	4	3.52	5.09
+SB noTF	5	3.79	5.71



Red: cross-section, Blue: end-point, Solid :  $\tilde{\tau}_1$ , Dashed:  $\tilde{\tau}_2$ .

# Search for WIMPS and polarisation

(Work by C. Bartels, presented at IWLC 2010)

WIMP Dark Matter

- Masses of 0.1–1 TeV.
- In thermal equilibrium with SM soup after inflation.
- Weak interactions naturally give observed relic density.
- In SUSY with conserved R-Parity: LSP:  $\tilde{\chi}_1^0$  or  $\tilde{G}$ .
- Here: no model assumptions.

#### Pair production at ILC

- $e^+e^- \rightarrow \chi \chi$ , WIMPs not detected.
- Detection via ISR:  $e^+e^- \rightarrow \chi \chi \gamma$ .
- Missing ₽.
- Dominant background:  $e^+e^- \rightarrow \nu\nu(N)\gamma$ .

Search for WIMPS and polarisation

## Model Independent Production Cross Section



#### Parameters:

- $\kappa_e(P_e, P_p)$ : Helicity dependent annihilation fraction to  $e^+e^-$ .
- $S_{\chi}$ : Spin, scale factor.
- $M_{\chi}$ ,  $J_0 \rightarrow$  shape,  $J_0$  dominant partial wave.

## Model Independent Production Cross Section

Signal shape at threshold provides information on partial wave (s- or p-wave).



#### Parameters:

- $\kappa_e(P_e, P_p)$ : Helicity dependent annihilation fraction to  $e^+e^-$ .
- $S_{\chi}$ : Spin, scale factor.
- $M_{\chi}$ ,  $J_0 \rightarrow$  shape,  $J_0$  dominant partial wave.

# Model Independent Production Cross Section

Crossover for s-wave and p-wave signal with same cross section.  $(\Rightarrow \text{ important later})$ 



#### Parameters:

- $\kappa_e(P_e, P_p)$ : Helicity dependent annihilation fraction to  $e^+e^-$ .
- $S_{\chi}$ : Spin, scale factor.
- $M_{\chi}$ ,  $J_0 \rightarrow$  shape,  $J_0$  dominant partial wave.

## **Polarised Cross Sections**



• 500 fb<sup>-1</sup> with  $(P_{e^-}/P_{e^+}) = (0.8/-0.6), (-0.8/0.6), (0.8/0.6), (-0.8/-0.6)$  (devided as 200, 200, 50, 50 fb<sup>-1</sup>).

150 GeV p-wave WIMP

고나님

< 17 >

## **Polarised Cross Sections**



- 500 fb<sup>-1</sup> with  $(P_{e^-}/P_{e^+}) = (0.8/-0.3), (-0.8/0.3), (0.8/0.3), (-0.8/-0.3)$  (devided as 200, 200, 50, 50 fb<sup>-1</sup>)
- 150 GeV p-wave WIMP

12

< A >

#### Mass Determination, $P_{e^-} = 80\% P_{e^+} = 0\%$



#### Mass Determination, $P_{e^-} = 80\% P_{e^+} = -30\%$



 $\sigma_{RR} = \sigma_{LL} = \sigma_{LR} = \sigma_{RL}$ 

Only small change in resolution with positron polarisation.

 $(P_{p}/P_{p}) = (0.8/-0.3)$ Luminosity = 500.0 fb  $J_0 = 1; \sigma_0 = 100.0 \text{ fb}$ Coup. struct. = 1 140 160 180 200 220 M<sub>v</sub> [GeV]

 $\sigma_{LR} = \sigma_{RL}; \sigma_{RR} = \sigma_{LL} = \mathbf{0}$ 

• Additional resolution increase by 3/4.

## TGC:s and Polarisation

(Work by I. Marchesini, presented at IWLC 2010)

Polarisation measurement from data with the Blondel scheme:

$$\sigma = \sigma_{u} \left[ 1 - P_{e^{+}} P_{e^{-}} + A_{LR} (P_{e^{+}} - P_{e^{-}}) \right], \tag{1}$$

hence

$$\mathbf{P}_{e^{\pm}} = \sqrt{\frac{(\sigma_{+-} + \sigma_{-+} - \sigma_{++} - \sigma_{--})(\mp \sigma_{-+} \pm \sigma_{+-} - \sigma_{++} + \sigma_{--})}{(\sigma_{-+} + \sigma_{+-} + \sigma_{++} + \sigma_{--})(\mp \sigma_{-+} \pm \sigma_{+-} + \sigma_{++} - \sigma_{--})}}$$

However: 100:s of  $fb^{-1}$  needed to get to 0.2 %.

## WW and Polarisation

Enter *WW* production : a high cross-section, polarisation dependent process



Ideally suited to make polarisation measurements, with less data than the Blondel scheme.

## TGC:s in WW

There is a catch, however:

Triple Gauge Couplings

TGC:s :

- 14 complex parameters, 8 CP conserving.
- In the SM: only 4 real parameters non-zero, all equal to unity
- Deviations from SM loop-corrections and beyond SM physics

Deviations from the SM still allowed (by LEP), affecting the polarisation measurement up to the % level.

TGC:s and Polarisation

## TGC:s+Polarisation in WW

TGC:s modifies angular diff. cross-sections  $\rightarrow$  % level corrections to polarisation measurement  $\rightarrow$  fit simultaneously.





If individually C and P conserving and real: 6 TGC:s, but one fixed by EM-gauge invariance. Gauge conditions: some relations  $\rightarrow$  3 TGC parameters + 2 polarisations to fit.

## Simultaneous fit : TGC:s

Result of simultaneous fit: TGCs



-

4 3 > 4 3

< 17 ▶
# Simultaneous fit : TGC:s

Result of simultaneous fit: TGCs



-

4 3 > 4 3

< 17 ▶

# Simultaneous fit : TGC:s

Result of simultaneous fit: TGCs



TGC:s and Polarisation

# Simultaneous fit: Polarisation

Result of simultaneous fit: polarisation



315

3 > 4 3

< 🗇 🕨

TGC:s and Polarisation

# Simultaneous fit: Polarisation

Result of simultaneous fit: polarisation



315

B > 4 B >

< 6 b

# Simultaneous fit: Polarisation

Result of simultaneous fit: polarisation



- Outperforms Blondel scheme
- Much gain with positron polarisation



3

(B) (A) (B) (A)

< 🗇 🕨

# Near Degenerate *ẽ* and polarisation

(Preliminary work by M.B., G. Moortgat-Pick)

SUSY associates scalars to chiral (anti)fermions



What if  $M_{\tilde{e}_L} \approx M_{\tilde{e}_R}$ , so that thresholds can't separate  $e^+e^- \rightarrow \tilde{e}_L \tilde{e}_L$ ,  $\tilde{e}_R \tilde{e}_R$  and  $\tilde{e}_R \tilde{e}_L$ ?

Model: SPS1a' like, but:

 $M_{\tilde{e}_{L}}$  = 200 GeV and  $M_{\tilde{e}_{R}}$  = 195 GeV. Both decay 100 % to  $\tilde{\chi}_{1}^{0} e$ .

Even with  $P_{e^-} \ge +90\%$ : No separation of  $\tilde{e}^+_{\rm L} \tilde{e}^-_{\rm R}$  and  $\tilde{e}^+_{\rm R} \tilde{e}^-_{\rm R}$ : Ratio of the cross sections  $\approx$  constant.

Model: SPS1a' like, but:

 $M_{\tilde{e}_{L}}$  = 200 GeV and  $M_{\tilde{e}_{R}}$  = 195 GeV. Both decay 100 % to  $\tilde{\chi}_{1}^{0} e$ .

Even with  $P_{e^-} \ge +90\%$ : No separation of  $\tilde{e}_L^+ \tilde{e}_R^-$  and  $\tilde{e}_R^+ \tilde{e}_R^-$ : Ratio of the cross sections  $\approx$  constant.



Model: SPS1a' like, but:

 $M_{\tilde{e}_{\rm L}}$  = 200 GeV and  $M_{\tilde{e}_{\rm R}}$  = 195 GeV. Both decay 100 % to  $\tilde{\chi}_1^0 e$ .

Even with  $P_{e^-} \ge +90\%$ : No separation of  $\tilde{e}_L^+ \tilde{e}_R^-$  and  $\tilde{e}_R^+ \tilde{e}_R^-$ : Ratio of the cross sections  $\approx$  constant.



### Polarised positrons a must !

Background and efficiency from Full-sim SPS1a' sample, kinematics from Whizard simulation of the model.

Background and efficiency from Full-sim SPS1a' sample, kinematics from Whizard simulation of the model.

- The ẽ signal was extracted from the same sample as was used for the SPS1a' τ̃study, using the same cuts except
  - Demand exactly two well identified electrons.
  - Reverse the τ̃anti-SUSY background cut
  - Some cuts could be loosened
- Almost background-free !

김 글 제 김 글 제 글 날

Background and efficiency from Full-sim SPS1a' sample, kinematics from Whizard simulation of the model.

- The ẽ signal was extracted from the same sample as was used for the SPS1a' τ̃study, using the same cuts except
  - Demand exactly two well identified electrons.
  - Reverse the *τ* anti-SUSY background cut
  - Some cuts could be loosened
- Almost background-free !



Background and efficiency from Full-sim SPS1a' sample, kinematics from Whizard simulation of the model.

For the signal:

- Generate (with Whizard 1.95) the modified model.
- Apply the kinematic cuts used for the full simulation analysis.
- Scale down the over-all event-weight so that the efficiency agrees with the full simulation.

The handle: Opposite polarisation beams produces  $\tilde{e}$ :s in both s- and t-channel. Same polarisation produces  $\tilde{e}$ :s in t-channel only  $\Rightarrow$ 

### Modification of $\Theta$ distribution with changed positron polarisation

However, the effect is small since t-channel always dominates !  $\tilde{e}$ :s are heavy (and are scalars)  $\Rightarrow$  t- and s- channel kinematic distributions of the electrons are not very different.

The handle: Opposite polarisation beams produces  $\tilde{e}$ :s in both s- and t-channel. Same polarisation produces  $\tilde{e}$ :s in t-channel only  $\Rightarrow$ 

Modification of  $\Theta$  distribution with changed positron polarisation

However, the effect is small since t-channel always dominates !  $\tilde{e}$ :s are heavy (and are scalars)  $\Rightarrow$  t- and s- channel kinematic distributions of the electrons are not very different.

Analyse assuming  $100 \text{ fb}^{-1}$  for each of the polarisations configurations.

- P(e<sup>-</sup>)= +80 % and ..
- P(e<sup>+</sup>) = ± 22 % ...
- P(e<sup>+</sup>) = ± 30 % ...
- P(e<sup>+</sup>) = ± 60 % ...
- ... and for  $P(e^{-})=\pm 80 \%$  $P(e^{+})=0$

Analyse assuming 100  $\rm fb^{-1}$  for each of the polarisations configurations.

P(e<sup>-</sup>)= +80 % and ..
P(e<sup>+</sup>) = ± 22 % ...
P(e<sup>+</sup>) = ± 30 % ...
P(e<sup>+</sup>) = ± 60 % ...
... and for P(e<sup>-</sup>)= ± 80 % P(e<sup>+</sup>) = 0



Analyse assuming 100  $\rm fb^{-1}$  for each of the polarisations configurations.

P(e<sup>-</sup>)= +80 % and ..
P(e<sup>+</sup>) = ± 22 % ...
P(e<sup>+</sup>) = ± 30 % ...
P(e<sup>+</sup>) = ± 60 % ...
... and for P(e<sup>-</sup>)= ± 80 %



Analyse assuming 100  $\rm fb^{-1}$  for each of the polarisations configurations.

P(e<sup>-</sup>)= +80 % and ..
P(e<sup>+</sup>) = ± 22 % ...
P(e<sup>+</sup>) = ± 30 % ...
P(e<sup>+</sup>) = ± 60 % ...
... and for P(e<sup>-</sup>)= ± 80 % P(e<sup>+</sup>) = 0



Analyse assuming 100  $\rm fb^{-1}$  for each of the polarisations configurations.

P(e<sup>-</sup>)= +80 % and ..
P(e<sup>+</sup>) = ± 22 % ...
P(e<sup>+</sup>) = ± 30 % ...
P(e<sup>+</sup>) = ± 60 % ...
... and for P(e<sup>-</sup>)= ± 80 % P(e<sup>+</sup>) = 0



Analyse assuming 100  $\rm fb^{-1}$  for each of the polarisations configurations.



### Conclusions

### Also when full simulation of both detector and beams:

### • The $\mathcal{L}_{eff}$ effect is seen in $\tilde{\tau}$ and WIMP:s

- Strongly in in 
   <sup>7</sup><sub>1</sub>: signal and background have opposite P dependence.
- But also in  $\tilde{\tau}_2$ , even though they have the same.

### • The polarisation measurement using WW was shown to

- More powerful than the Blondel scheme.
- Profit strongly from positron polarisation.
- The preliminary determination of the chiral structure of near-degenerate ess
  - Cannot be done without positron polarisation.
  - Profits largely even from a modest increase (22 % to 30 % ↔ doubling the luminosity)

### Higher positron polarisation enhances the physics potential of the ILC

### Conclusions

Also when full simulation of both detector and beams:

### • The $\mathcal{L}_{\textit{eff}}$ effect is seen in $\tilde{\tau}$ and WIMP:s

- Strongly in in (\tilde{\tau}\_1: signal and background have opposite P) dependence.
- But also in  $\tilde{\tau}_2$ , even though they have the same.
- The polarisation measurement using WW was shown to
  - More powerful than the Blondel scheme.
  - Profit strongly from positron polarisation.
- The preliminary determination of the chiral structure of near-degenerate ess
  - Cannot be done without positron polarisation
  - Profits largely even from a modest increase (22 % to 30 % ↔ doubling the luminosity)

### Higher positron polarisation enhances the physics potential of the ILC

### Conclusions

Also when full simulation of both detector and beams:

### • The $\mathcal{L}_{\textit{eff}}$ effect is seen in $\tilde{\tau}$ and WIMP:s

- Strongly in in (\tilde{\tau}\_1: signal and background have opposite P) dependence.
- But also in  $\tilde{\tau}_2$ , even though they have the same.
- The polarisation measurement using WW was shown to
  - More powerful than the Blondel scheme.
  - Profit strongly from positron polarisation.
- The preliminary determination of the chiral structure of near-degenerate ess
  - Cannot be done without positron polarisation
  - Profits largely even from a modest increase (22 % to 30 % ↔ doubling the luminosity)

### Higher positron polarisation enhances the physics potential of the ILC

### Conclusions

Also when full simulation of both detector and beams:

### • The $\mathcal{L}_{\textit{eff}}$ effect is seen in $\tilde{\tau}$ and WIMP:s

- Strongly in in (\tilde{\tau}\_1: signal and background have opposite P) dependence.
- But also in  $\tilde{\tau}_2$ , even though they have the same.
- The polarisation measurement using WW was shown to
  - More powerful than the Blondel scheme.
  - Profit strongly from positron polarisation.
- The preliminary determination of the chiral structure of near-degenerate ess
  - Cannot be done without positron polarisation
  - Profits largely even from a modest increase (22 % to 30 %  $\leftrightarrow$  doubling the luminosity)

### Higher positron polarisation enhances the physics potential of the ILC

### Conclusions

Also when full simulation of both detector and beams:

### • The $\mathcal{L}_{\textit{eff}}$ effect is seen in $\tilde{\tau}$ and WIMP:s

- Strongly in in (\tilde{\tau}\_1: signal and background have opposite P) dependence.
- But also in  $\tilde{\tau}_2$ , even though they have the same.
- The polarisation measurement using WW was shown to
  - More powerful than the Blondel scheme.
  - Profit strongly from positron polarisation.
- The preliminary determination of the chiral structure of near-degenerate ess
  - Cannot be done without positron polarisation
  - Profits largely even from a modest increase (22 % to 30 %  $\leftrightarrow$  doubling the luminosity)

### Higher positron polarisation enhances the physics potential of the ILC

### Backup

# Main items for physics • Half RF power:

- - To keep L: decrease beam-size.
  - But: .  $\rightarrow$  increases  $\delta_{BS}$
  - Doubled luminosity/BX → doubled probability for a  $\gamma\gamma$  event in the same BX.

### Undolator move :

# Main items for physics • Half RF power:

- - To keep L: decrease beam-size.
  - But: .  $\rightarrow$  increases  $\delta_{BS}$
  - Doubled luminosity/BX → doubled probability for a  $\gamma\gamma$  event in the same BX.

### Undolator move :

- Higher energy-spread at 500 GeV.
- Lower positron polarisation at 500

イロト 不得 トイヨト イヨト 正言 ろくの

# Main items for physics • Half RF power:

- - To keep L: decrease beam-size.
  - But: .  $\rightarrow$  increases  $\delta_{BS}$
  - Doubled luminosity/BX → doubled probability for a  $\gamma\gamma$  event in the same BX.

### Undolator move :

- Higher energy-spread at 500 GeV.



< ロ > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 >

# Main items for physics • Half RF power:

- - To keep L: decrease beam-size.
  - But: .  $\rightarrow$  increases  $\delta_{BS}$
  - Doubled luminosity/BX → doubled probability for a  $\gamma\gamma$  event in the same BX.

### Undolator move :

- Higher energy-spread at 500 GeV.
- Lower positron polarisation at 500 GeV.



### $\text{RDR} \rightarrow \text{SB2009 procedure}$

- BeamCal: From our studies: SB2009(TF) ≈ SB2009(noTF) ≈ 2 × RDR. → Multiply BeamCal Energy-density map(RDR) by 2. Use same function p([E<sub>e</sub> at x, y], [pairs energy density at x, y]) as the probability to detect an electron of energy E entering the BeamCal at (x, y) would be seen.
- Tracking Fully simulated and reconstructed BX:es available both for RDR and SB2009 (no TF). Use method outlined above. NB: optimistic when applied to SB2009 with TF !
- Beam-spectrum Lumi distributions from GP (A. Hartin, T Barklow) for RDR and SB2009 and *E*<sub>beam1,2</sub> used to calculate even-by-event weights, to modify the existing fully simulated sample to an RDR one.

 Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

Mikael Berggren (DESY)

January 19, 2011 30 / 28

### $\text{RDR} \rightarrow \text{SB2009 procedure}$

- BeamCal: From our studies: SB2009(TF) ≈ SB2009(noTF) ≈ 2 × RDR. → Multiply BeamCal Energy-density map(RDR) by 2. Use same function p([E<sub>e</sub> at x, y], [pairs energy density at x, y]) as the probability to detect an electron of energy E entering the BeamCal at (x, y) would be seen.
- Tracking Fully simulated and reconstructed BX:es available both for RDR and SB2009 (no TF). Use method outlined above. NB: optimistic when applied to SB2009 with TF !
- Beam-spectrum Lumi distributions from GP (A. Hartin, T Barklow) for RDR and SB2009 and *E*<sub>beam1,2</sub> used to calculate even-by-event weights, to modify the existing fully simulated sample to an RDR one.
- Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

### $\text{RDR} \rightarrow \text{SB2009 procedure}$

- BeamCal: From our studies: SB2009(TF) ≈ SB2009(noTF) ≈ 2 × RDR. → Multiply BeamCal Energy-density map(RDR) by 2. Use same function p([E<sub>e</sub> at x, y], [pairs energy density at x, y]) as the probability to detect an electron of energy E entering the BeamCal at (x, y) would be seen.
- Tracking Fully simulated and reconstructed BX:es available both for RDR and SB2009 (no TF). Use method outlined above. NB: optimistic when applied to SB2009 with TF !
- Beam-spectrum Lumi distributions from GP (A. Hartin, T Barklow) for RDR and SB2009 and *E*<sub>beam1,2</sub> used to calculate even-by-event weights, to modify the existing fully simulated sample to an RDR one.

 Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

### $\text{RDR} \rightarrow \text{SB2009 procedure}$

- BeamCal: From our studies: SB2009(TF) ≈ SB2009(noTF) ≈ 2 × RDR. → Multiply BeamCal Energy-density map(RDR) by 2. Use same function p([E<sub>e</sub> at x, y], [pairs energy density at x, y]) as the probability to detect an electron of energy E entering the BeamCal at (x, y) would be seen.
- Tracking Fully simulated and reconstructed BX:es available both for RDR and SB2009 (no TF). Use method outlined above. NB: optimistic when applied to SB2009 with TF !
- Beam-spectrum Lumi distributions from GP (A. Hartin, T Barklow) for RDR and SB2009 and *E*<sub>beam1,2</sub> used to calculate even-by-event weights, to modify the existing fully simulated sample to an RDR one.
- Polarisation: Straight-forward relative weighting of generated samples with P=(-1,1) and P=(1,-1).

# RDR , SB2009 and $\tilde{\tau}$ :s: Signal and background

### Events after cuts, end-point analysis

case	$ ilde{ au}_1$			$ ilde{ au}_2$		
	SM	SUSY	signal	SM	SUSY	signal
RDR	317	998	10466	1518	241	1983
SB09(TF)	814	956	8410	1346	223	1555
SB09(nTF)	611	717	6308	1009	167	1166

### Events after cuts, cross-section analysis

case	$ ilde{ au}_1$			$ ilde{ au}_2$		
	SM	SUSY	signal	SM	SUSY	signal
RDR	17.6	47.7	2377	1362	33.7	1775
SB09(TF)	17.6	45.7	1784	1194	32.4	1366
SB09(nTF)	13.2	34.3	1337	895	24.3	1025
#### Model-independent search Birkedal *et al.* [hep-ph/0403004]

#### Model independence

- Assume only one DM candidate, no co-annihilation.
- Constrain WIMP pair annihilation XSec from observation.
- Crossing Symmetry (annihilation  $\Rightarrow$  production).
- ISR.



## Mass Determination, $P_{e^-} = 0\% P_{e^+} = 0\%$



Physics studies with polarisation

## Mass Determination, $P_{e^-} = 0\% P_{e^+} = 0\%$





M = 220.

## Mass Determination, $P_{e^-} = 0\% P_{e^+} = 0\%$



# RDR , SB2009 and Near Degenerate $\widetilde{e}$

Need to reconstruct the  $\tilde{e}$  direction:

- 8 unknown  $\tilde{\chi}_1^0$  momentum components
- Assume  $M_{ ilde{e}}$  and  $M_{ ilde{\chi}^0_1}$  known ightarrow
- 8 constraints (E and p conservation, 4 mass-relations)
- WORKS !

# RDR , SB2009 and Near Degenerate ẽ

Need to reconstruct the  $\tilde{e}$  direction:

- 8 unknown  $\tilde{\chi}_1^0$  momentum components
- Assume  $M_{\widetilde{e}}$  and  $M_{\widetilde{v}_{+}^{0}}$  known  $\rightarrow$
- 8 constraints (E and p conservation, 4 mass-relations)
- WORKS !



< E

# RDR , SB2009 and Near Degenerate ẽ

Need to reconstruct the  $\tilde{e}$  direction:

- 8 unknown  $\tilde{\chi}_1^0$  momentum components
- Assume  $M_{\widetilde{e}}$  and  $M_{\widetilde{\chi}^0_{+}}$  known  $\rightarrow$
- 8 constraints (E and p conservation, 4 mass-relations)
- WORKS !



< 17 ▶

315

4 3 > 4 3