

# Top Higgs Yukawa Coupling

## Analysis from

$$e^+ e^- \rightarrow \bar{t} tH \rightarrow \bar{b} W^- b W^+ \bar{b} b$$

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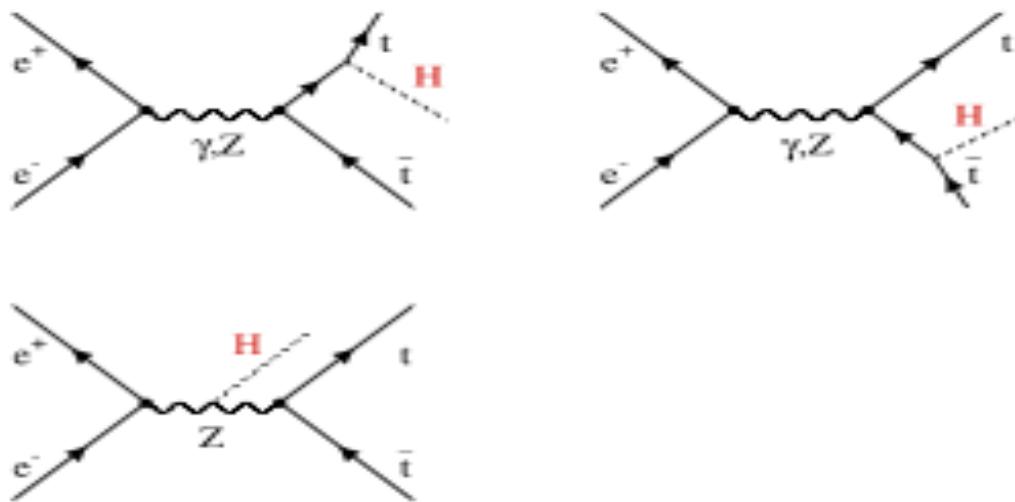


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# Introduction

- Once Higgs is found, its coupling with fermions is interesting to study.
- $g_{ffH} = \frac{m_f}{v}$ ,  $v$  is vacuum expectation value of Higgs field.
- Top is heaviest fermion, so top-Higgs Yukawa coupling is largest.
- The coupling of top to the Higgs is modified in the SUSY models
- This analysis was not done for LOI
- ILD software framework is used for this analysis



**Fig. 1.** Lowest order Feynman diagrams of the process  $e^+e^- \rightarrow ttH$

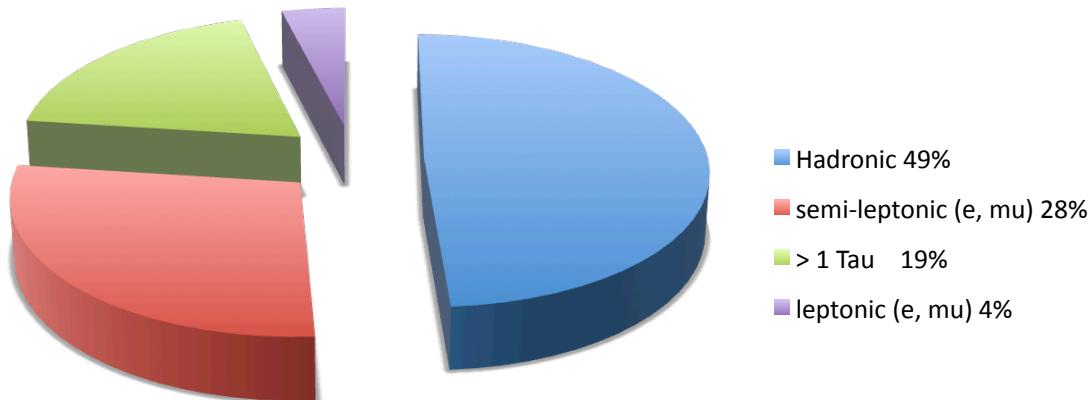
# Samples

- ILD\_00 centrally reconstructed sample with center of mass energy  $\sqrt{s} = 500 \text{ GeV}$ .
- $t\bar{t}$ -Higgs events with  $M_h = 120 \text{ GeV}/c^2$ ,  $M_t = 175 \text{ GeV}/c^2$ .
- Software version used is ILCSoft v01-06.

Process	$\sigma (\text{fb})$	Sample	$L (\text{ab}^{-1})$
$e^+ e^- \rightarrow t\bar{t}H$	0.577 [arXiv:hep-ph/0604166v2]	20,000	34
$e^+ e^- \rightarrow t\bar{t}$	521	1800000	3.4
$e^+ e^- \rightarrow t\bar{t}Z$	0.58	24,000	41
$e^+ e^- \rightarrow ZZ$	577.2		
$e^+ e^- \rightarrow W^-W^+$	7890		
$e^+ e^- \rightarrow q\bar{q}$	3951.8		

# Semi-Leptonic Channel

- $e^+ e^- \rightarrow \bar{t} tH \rightarrow \bar{b} W^- b W^+ \bar{b} b$
- Focus on semi-leptonic final state with one W decaying into lepton and neutrino and other W decaying into light jets
- Final state is 1 lepton, missing energy, 6 Jets with 4 b-jets
- Remove the leptons and force remaining particles into 6-jets (JetFinder Algorithm)
- High momentum Lepton and large missing momentum signature



# Filtering of Semi-Leptonic Channel for Monte Carlo Sample

- Initially 20,000 MC events with full final state where Higgs and W decay into anything
- Filter events with one lepton ( $\mu$ , e), and H decaying to  $b\bar{b}$ , 4466 events are left.

# Lepton Identification

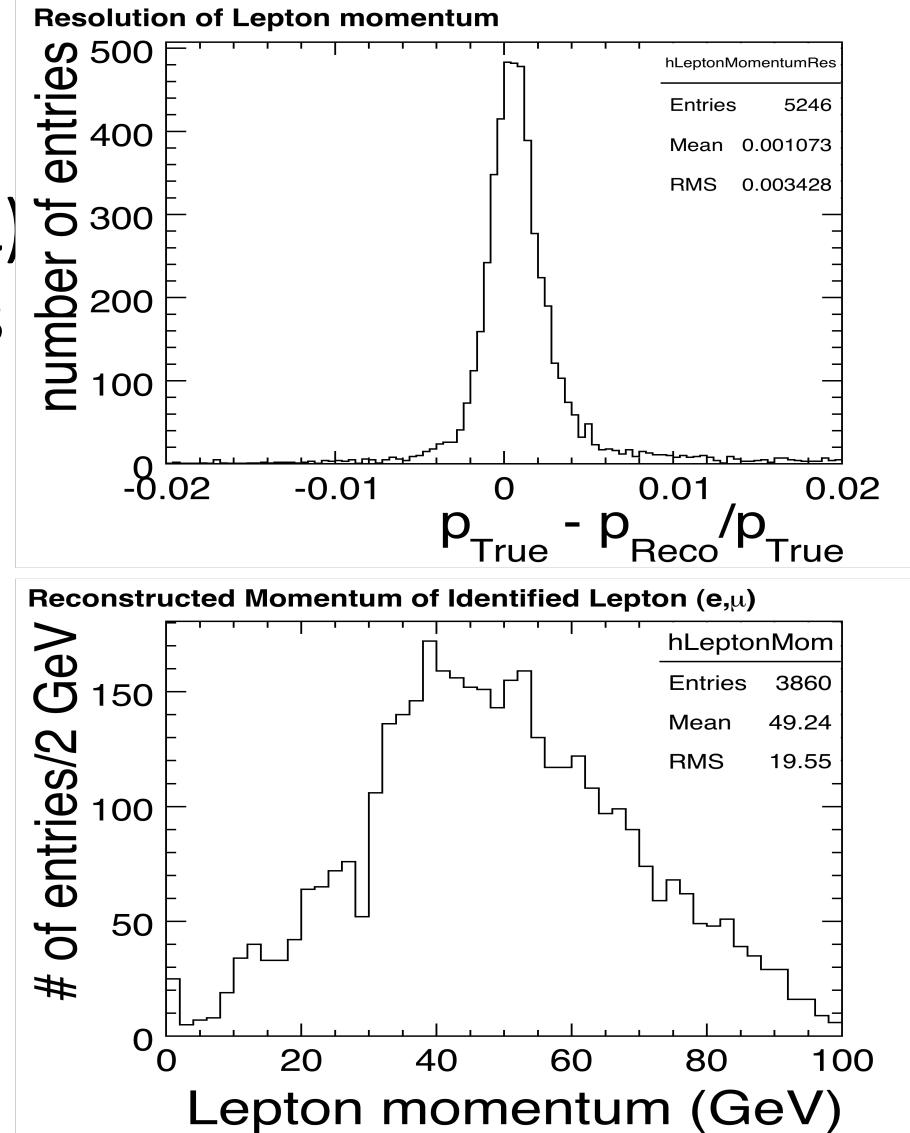
## (From study of Single MC Lepton with $P > 15 \text{ GeV}$ )

- Muon Identification:
  - Cut based selection is being used. Efficiency from single Muon sample is 98%.
    - (1)  $E_{\text{Ecal}} < 2.5 \text{ GeV}$
    - (2)  $E_{\text{Hcal}} < 15 \text{ GeV}$
    - (3)  $E_{\text{Ecal}} / E_{\text{Tot}} < 0.5$
    - (4)  $E_{\text{tot}} / p < 0.3$
- Electron Identification:
  - Cut-based selection on single Electron sample has showed that 98.57% electron are identified by using:
    - (1)  $E_{\text{Ecal}} / E_{\text{Tot}} > 0.6$
    - (2)  $E_{\text{Tot}} / p > 0.7$

efficiencies in %	electron cuts	muon cuts
$e$	$98.57 \pm 0.06$	$\sim 0$
$\mu$	$0.03 \pm 0.01$	$97.5 \pm 0.05$
$\pi$	$3.88 \pm 0.06$	$0.46 \pm 0.003$

# $W \rightarrow v\ell$ (Lepton Identification)

- we identify our reconstructed leptons ( $e, \mu$ ) using same cut variables as for single lepton case.
- Most of the reconstructed leptons are correctly identified as leptons.
- These lepton tracks are then removed from the PandorPFOs collection.

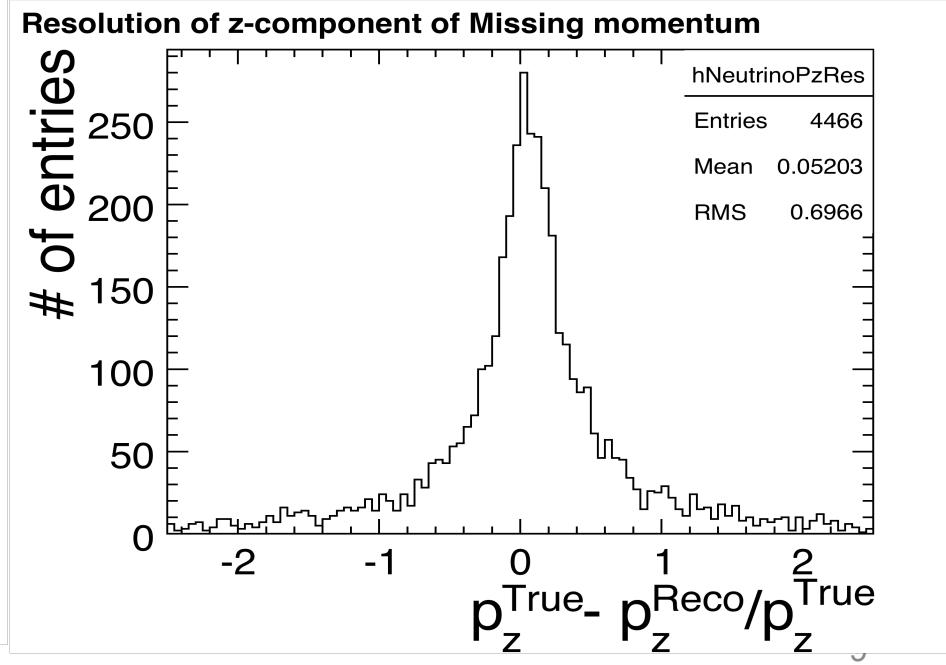
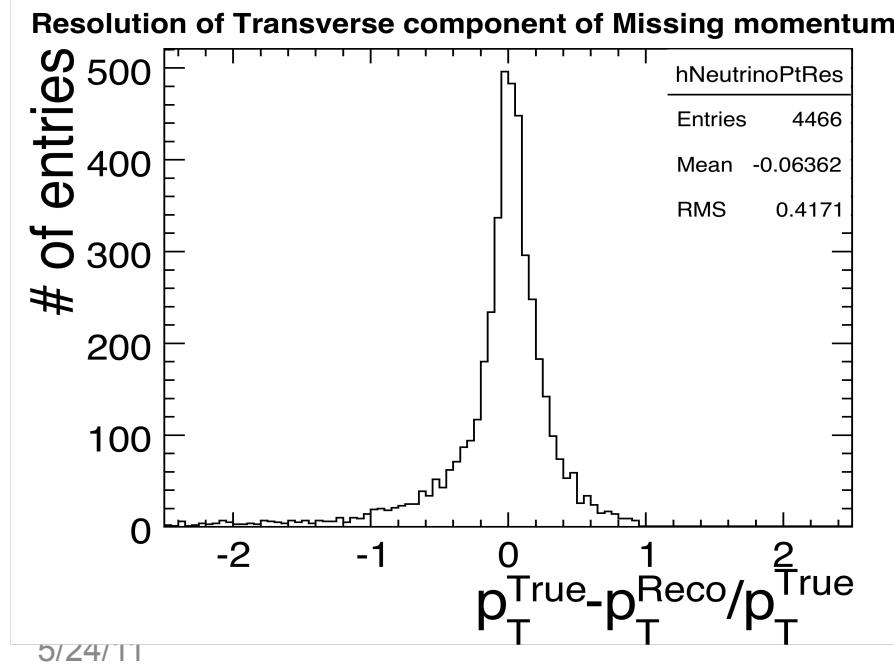


# $W \rightarrow v\ell$ (Missing Momentum)

- Using the information for all reconstructed particles, missing momentum is reconstructed.

$$p_x^{miss} = -\sum_i p_{xi}, \quad p_y^{miss} = -\sum_i p_{yi}, \quad p_z^{miss} = -\sum_i p_{zi}, \quad p_T^{miss} = \sqrt{(p_x^{miss})^2 + (p_y^{miss})^2}$$

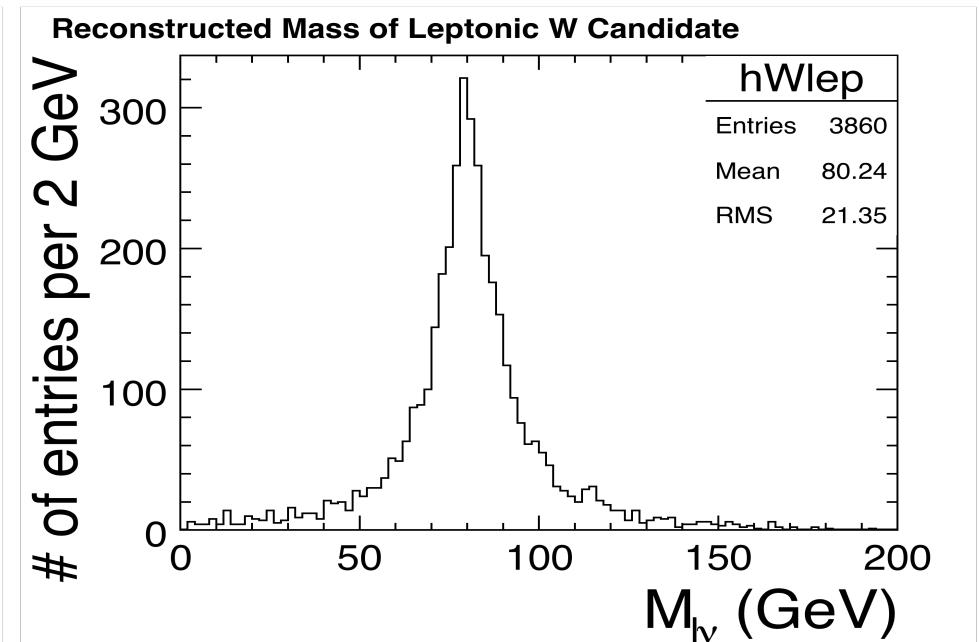
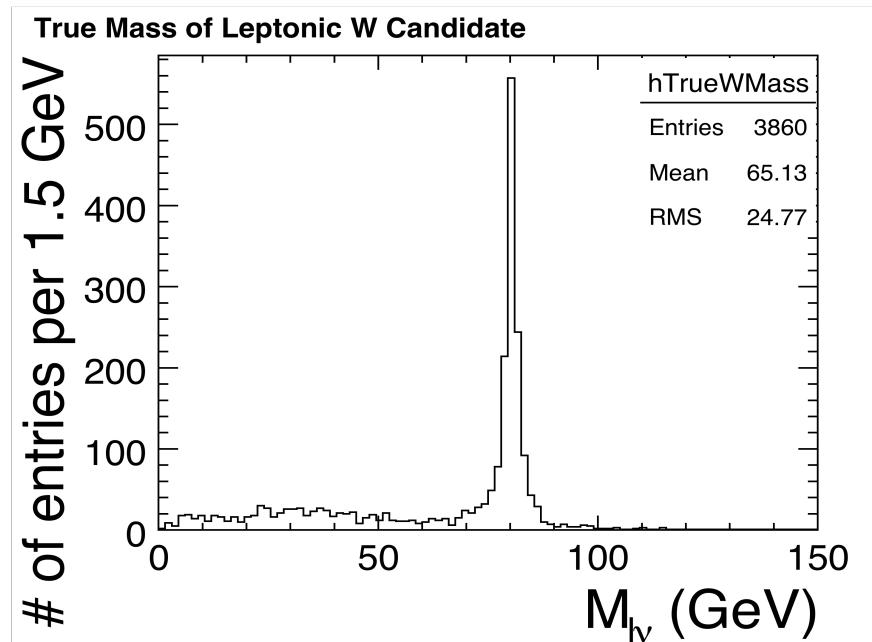
- Z-component of missing energy isn't as accurate as x and y components
- However, z-component can be used to reconstruct Semi-leptonic W mass.



# Reconstructed Mass of leptonic $W$

- To reconstruct semi-leptonic  $W$ , we select events which have lepton momentum not equal to zero and one lepton with highest momentum
- The momentum of lepton and Missing Momentum are used to reconstruct the mass of  $W$ .

$$M_W = \sqrt{(E_\nu + E_l)^2 - (p_{\nu x} + p_{lx})^2 - (p_{\nu y} + p_{ly})^2 - (p_{\nu z} + p_{lz})^2}$$

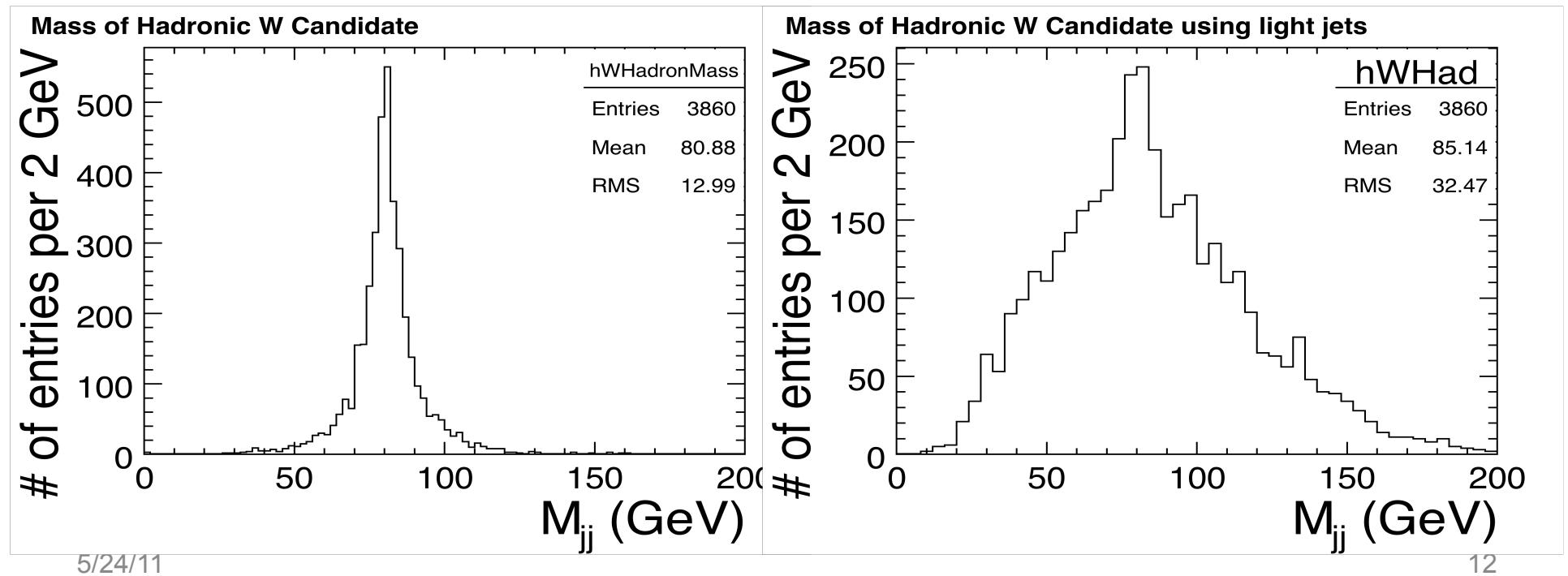


# Jets reconstruction

- Identified leptons are removed from the sample
- remaining particles are forced into 6 Jets using JetFinder algorithm
- Jets pass LCFIVertex reconstruction [[arXiv:0908.3019v1](#)]
- LCFI flavour tagging is used to separate light and b-jets
- Jets are sorted in descending order of b-tag value
- top four jets with highest b-tag value are selected as b-jets
- Light jets are used to reconstruct hadronic W

# Reconstructed Mass of Hadronic $W$

- Used the light jets (lowest b-tag value jets) momentum to reconstruct  $W$  mass (Right figure)
- Wrong particles are picked as light jets
- Make combinations of jets which have b-tag value less than 0.09 and pick the combination which is closest to  $W$  mass (81.4 GeV) (Left figure)



# Final State Reconstruction

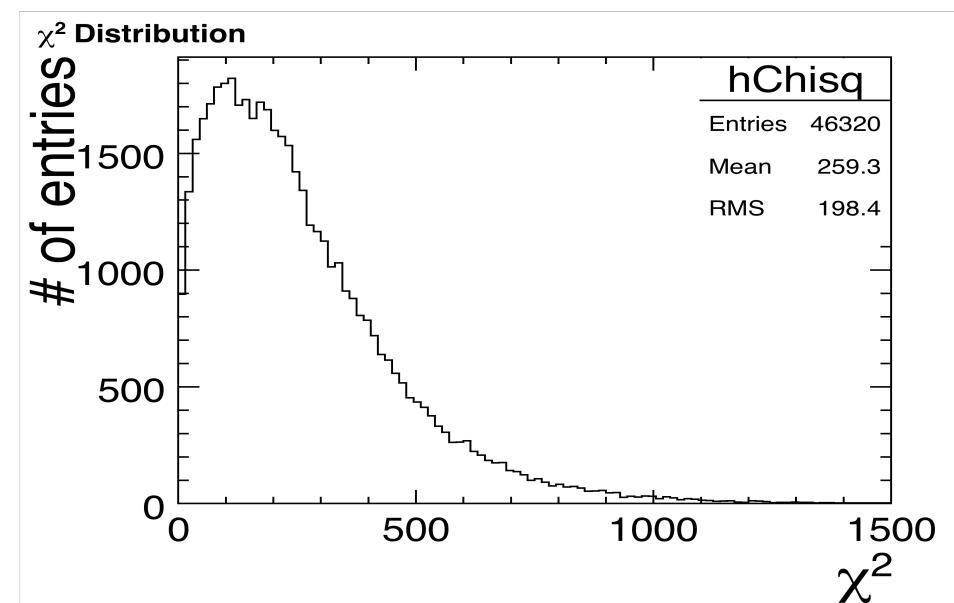
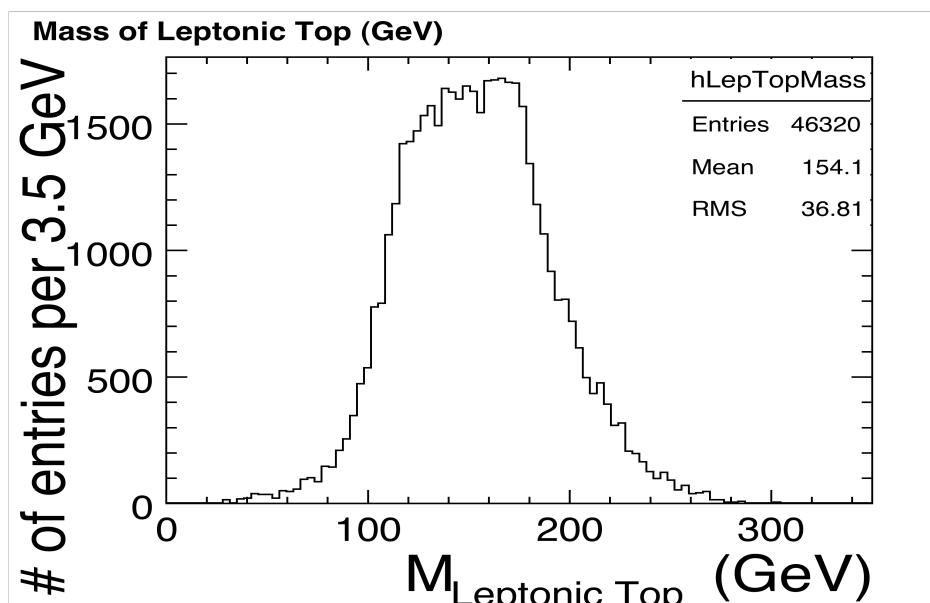
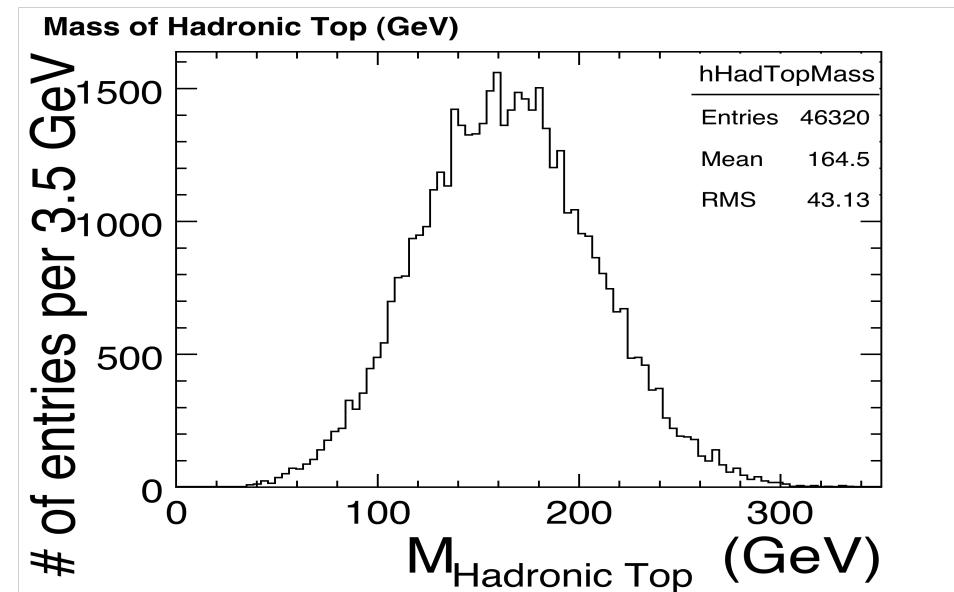
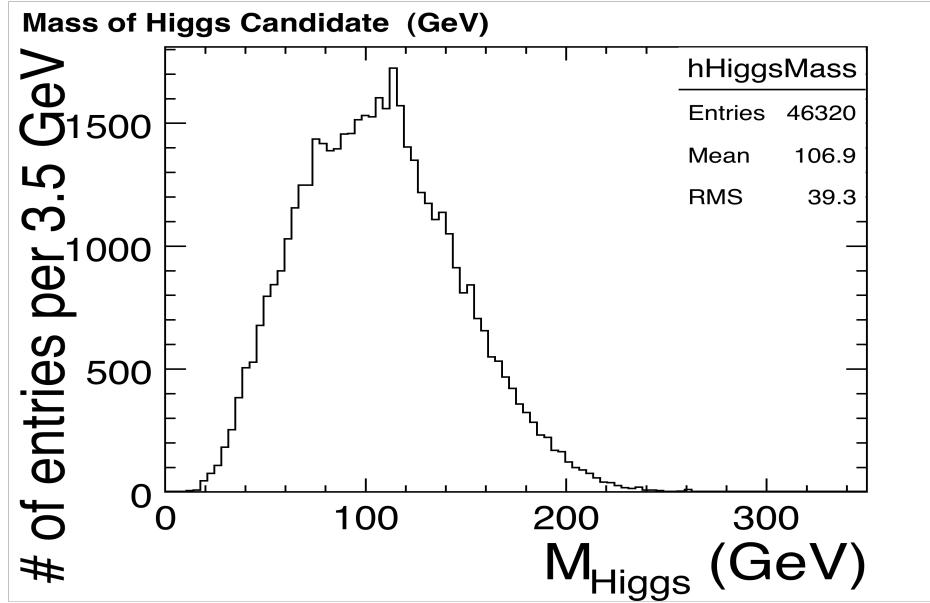
- As we have a good reconstruction of z-component of the missing momentum, it is used to reconstruct the Mass of semi-leptonic W

$$M_W = \sqrt{(E_\nu + E_l)^2 - (p_{\nu x} + p_{lx})^2 - (p_{\nu y} + p_{ly})^2 - (p_{\nu z} + p_{lz})^2}$$

- Four b-jets are used to reconstruct two tops and Higgs particle
- To reduce combinatorial backgrounds, minimisation of  $\chi^2$  technique is used

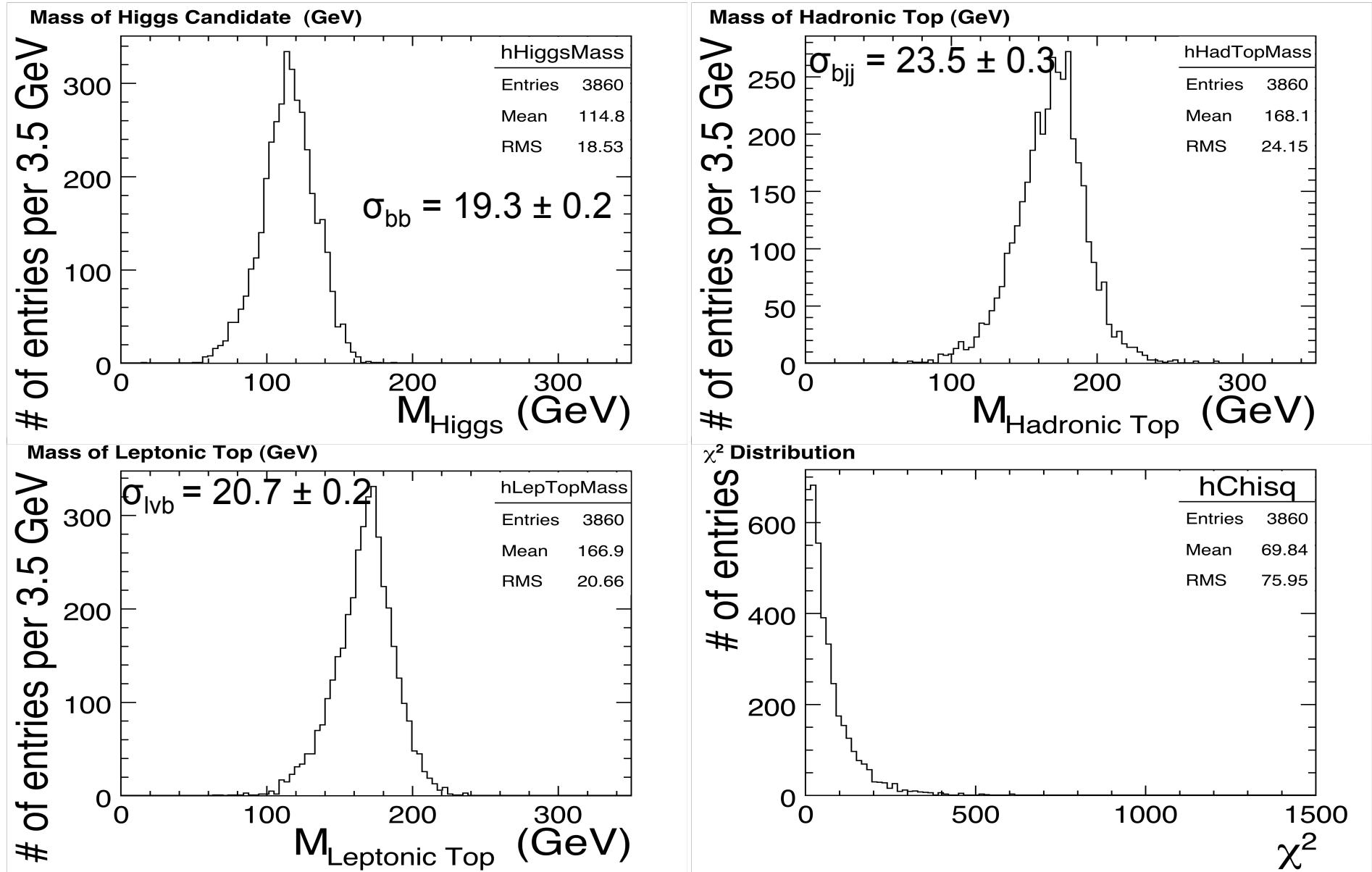
$$\chi^2 = \frac{(M_{l\nu b} - M_t)^2}{\sigma_{l\nu b}^2} + \frac{(M_{jjb} - M_t)^2}{\sigma_{jjb}^2} + \frac{(M_{bb} - M_{Higgs})^2}{\sigma_{bb}^2}$$

# Reconstructed Final State

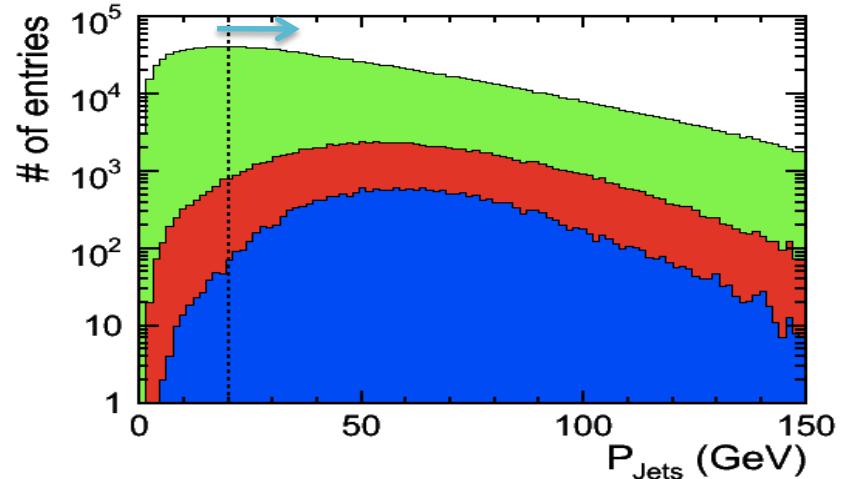
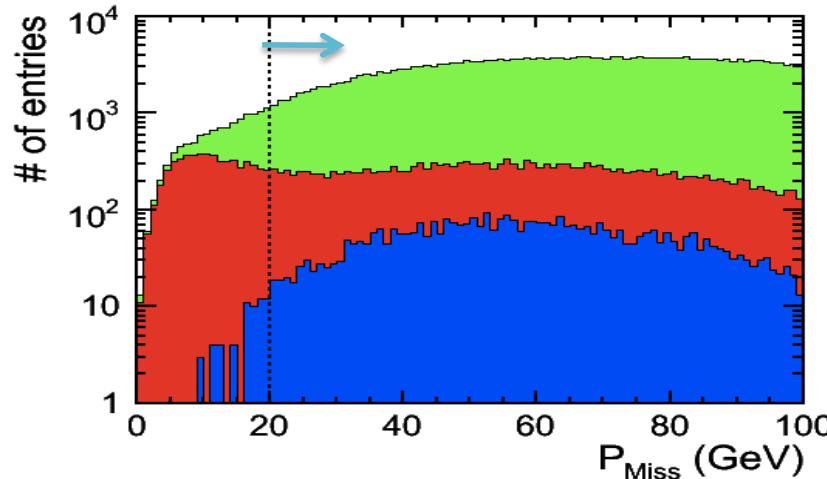
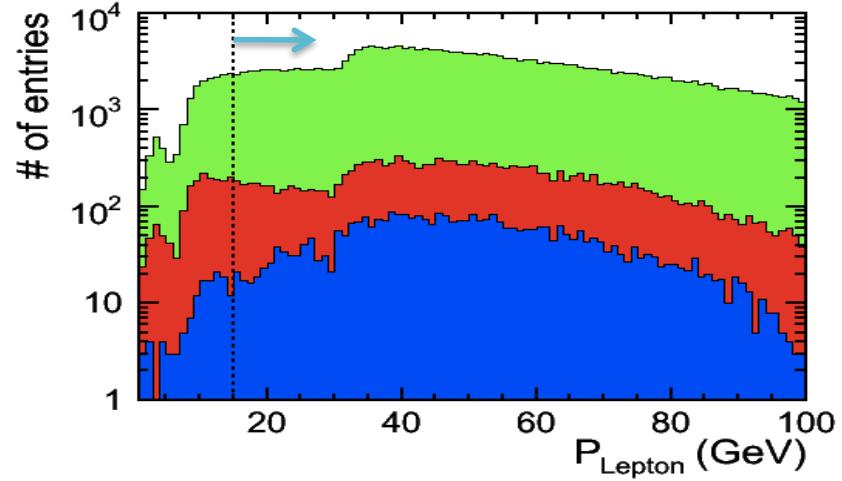
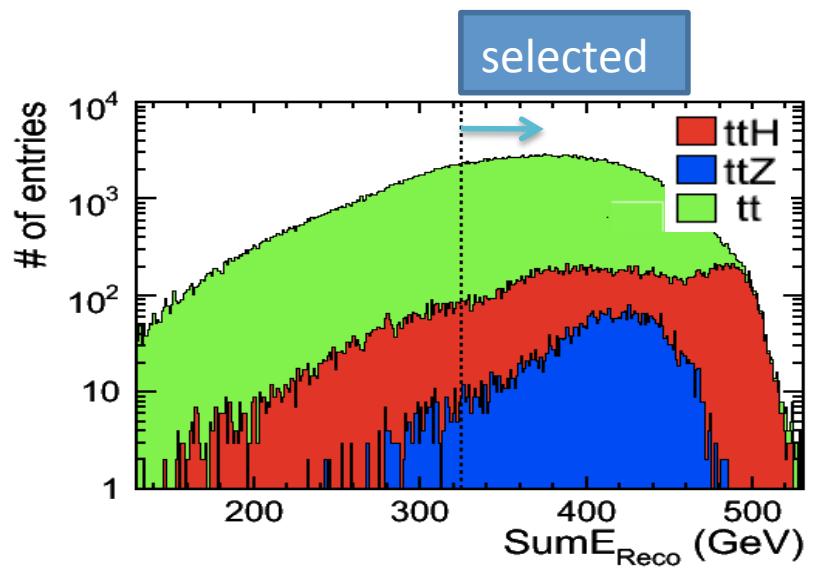


There are 12 entries for each event due to different combinations

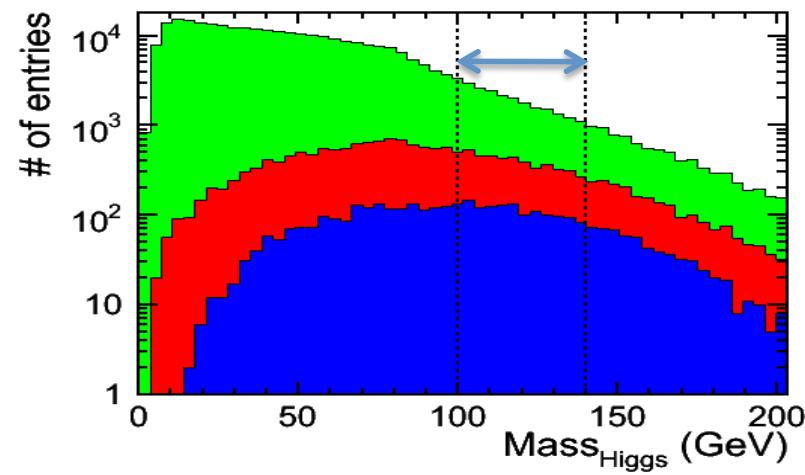
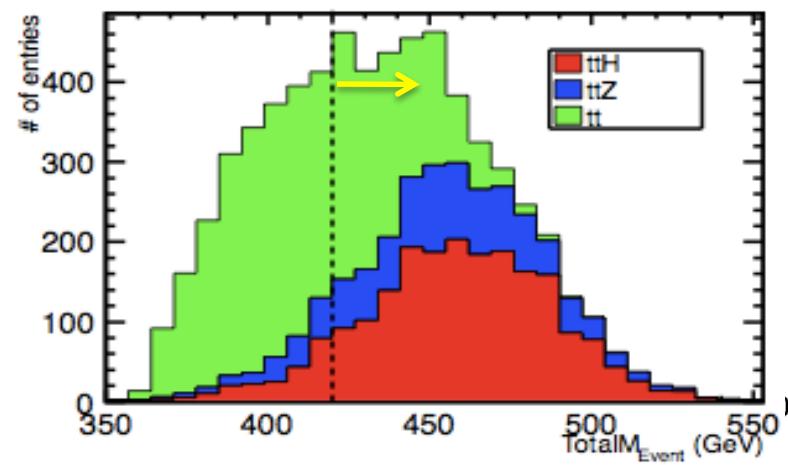
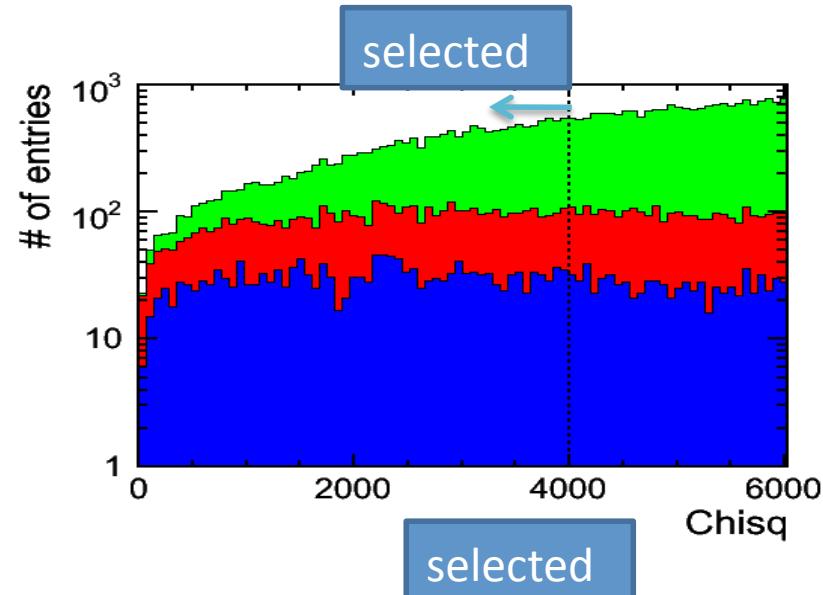
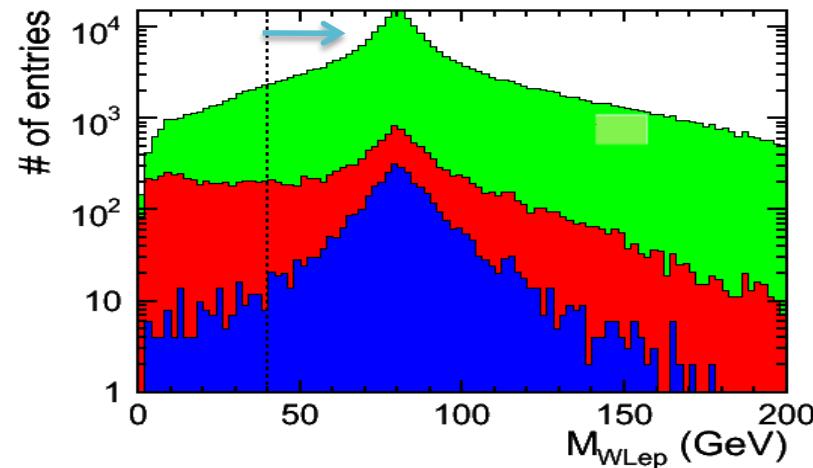
# Reconstructed Final State after Minimizing $\chi^2$



# Selection variables (I)



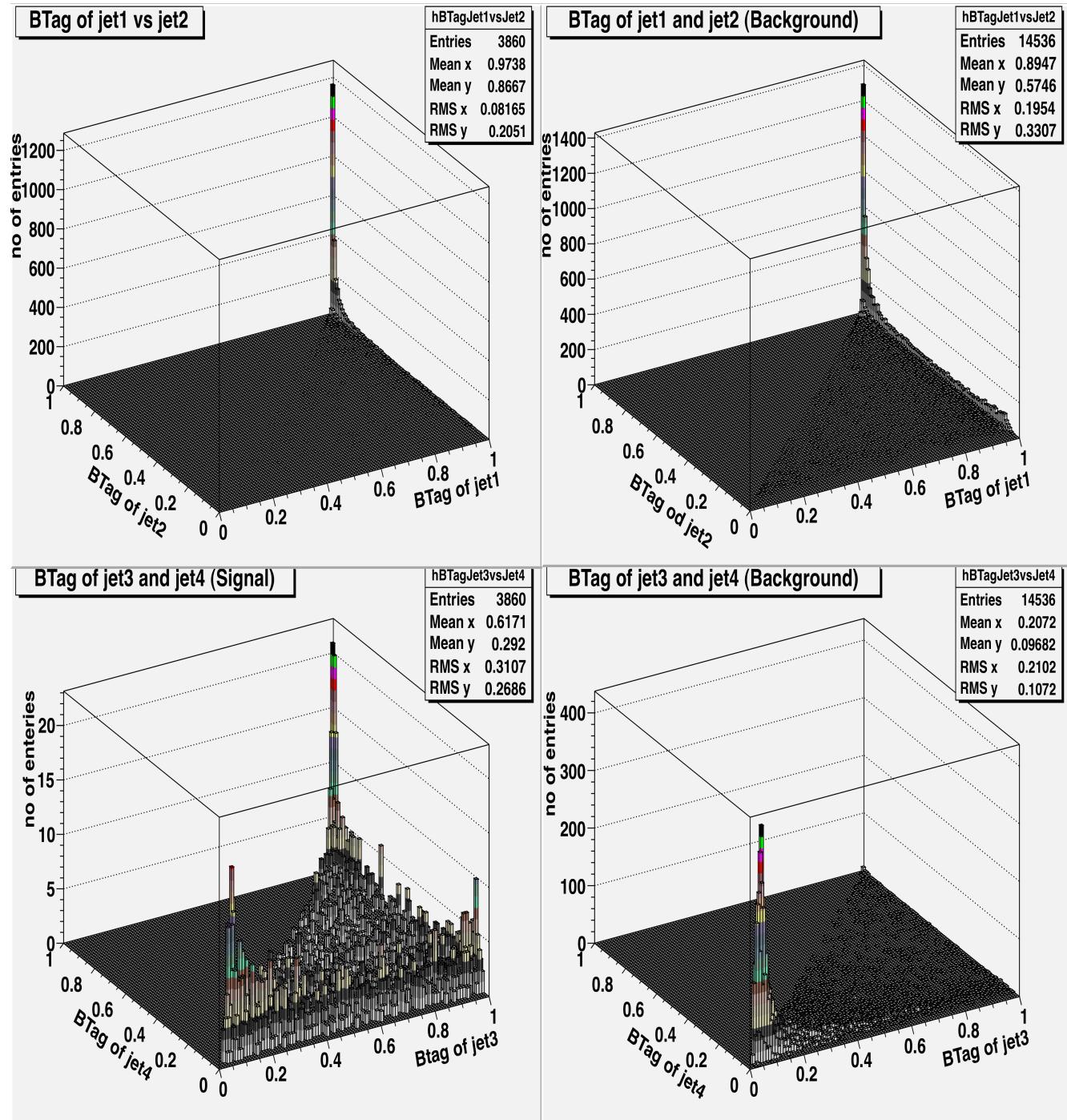
# Selection variables (II)



# b-tagging

## ttH vs ttZ

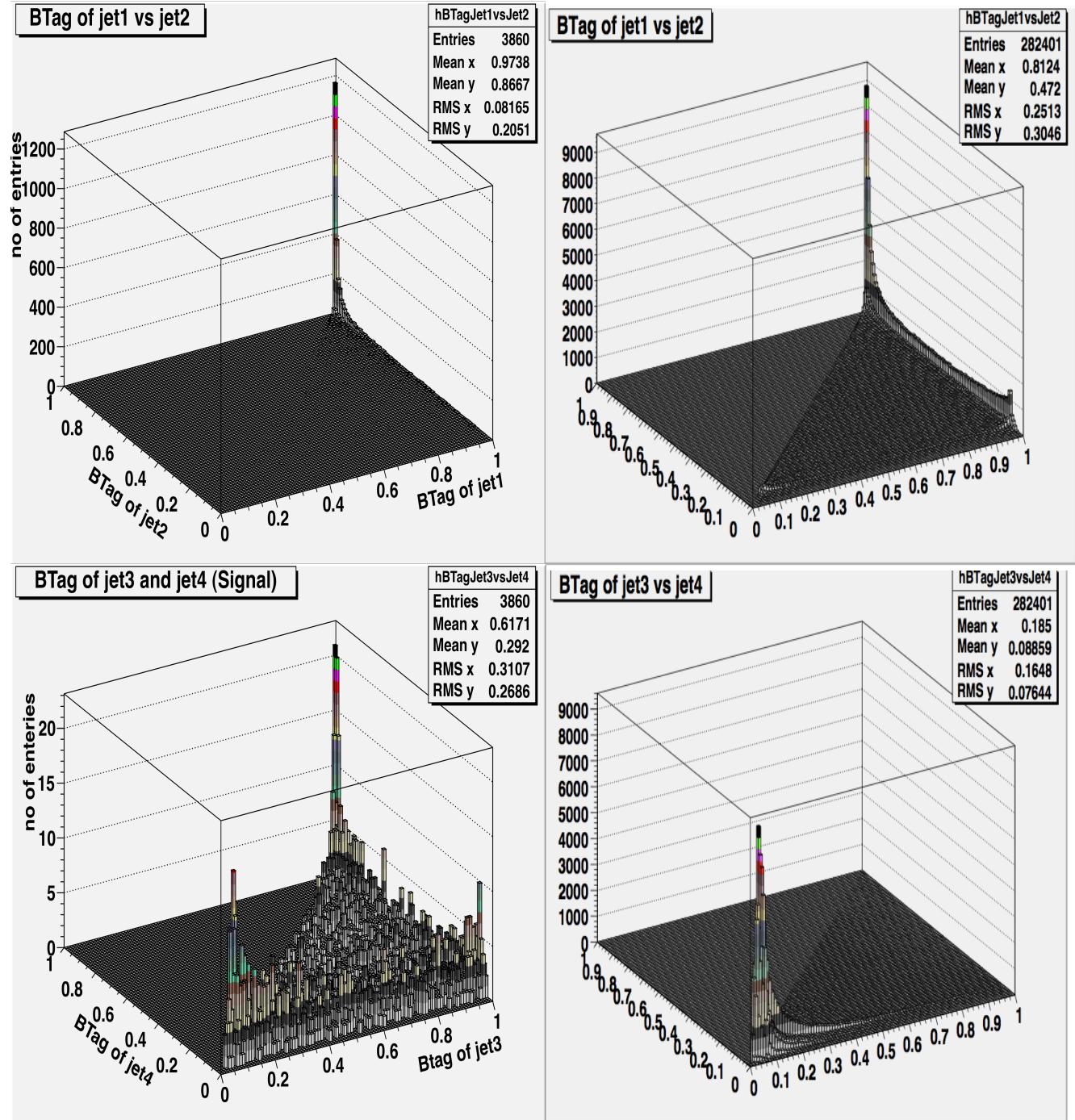
- B-tag of all 6 jets are looked at
- 4 with largest b-tag value are considered as b-jets
- Cut on b-tag values of jet3 and jet4 can reduce a significant number of background events



# b-tagging

## tH vs tt

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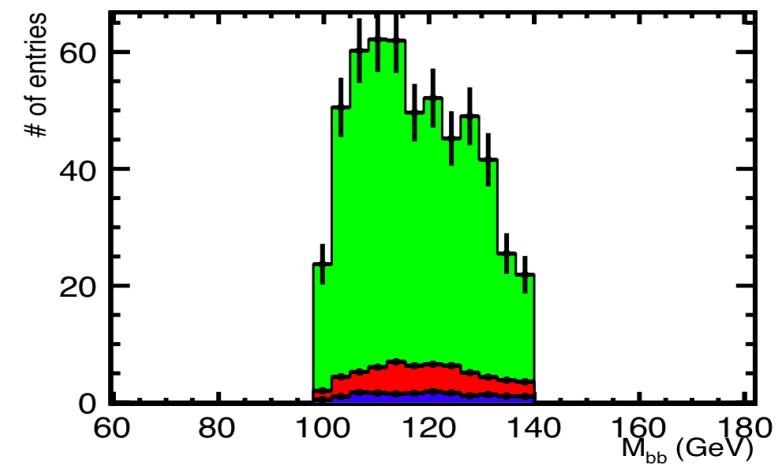
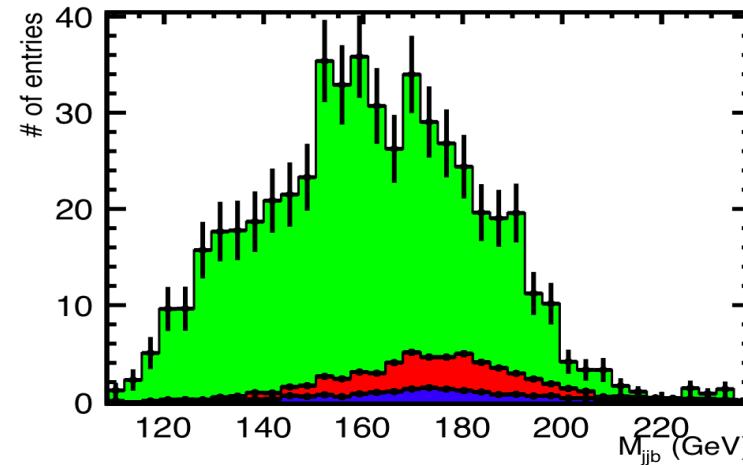
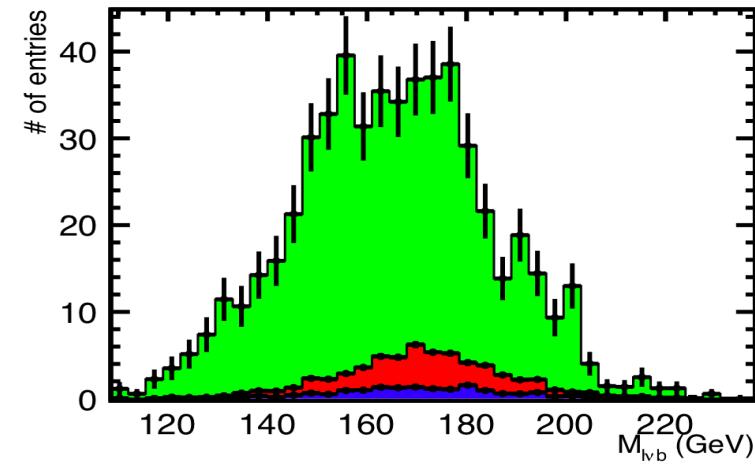
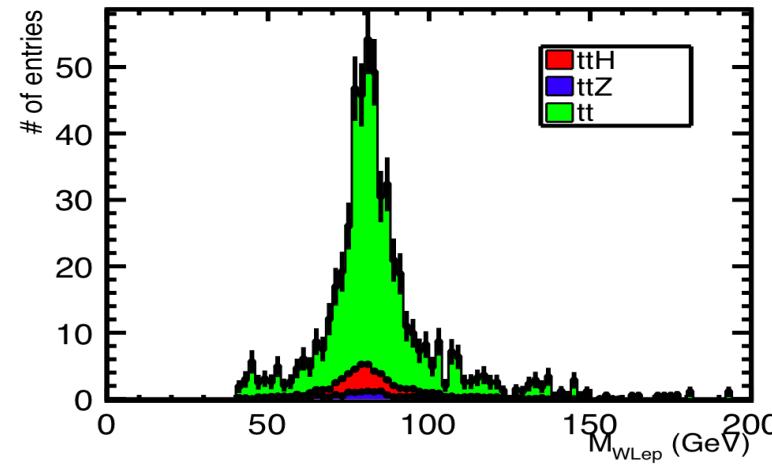
# b-tag efficiency (3<sup>rd</sup> and 4<sup>th</sup> Jets)

b-tag (both jets)	No of Signal events (S)	No of bkg events (B)	S/ $\sqrt{S+B}$
0	3097	7304	30.36
> 0.08	2410	2371	34.84
> 0.09	2237	1666	35.80
> 0.1	2154	1525	35.51
> 0.15	1708	771	34.30
> 0.2	1429	515	32.41
Btag (3 <sup>rd</sup> jet) > 0.09 Btag (4 <sup>th</sup> jets) > 0.08	2406	2282	35.14

# Signal and Background separation

Cuts	$t\bar{t}H$	$t\bar{t}Z$	$t\bar{t}$
initial	20000	24000	376276
# Lep > 0	3860 ** After semi-leptonic selection	14536	282404
E_Reco > 325 GeV	3600	8021	68439
P_Lep > 15 GeV	3167	7128	55206
P_miss > 20 GeV	3119	5610	54488
P_Jet > 20 GeV	2978	4837	33909
3 <sup>rd</sup> & 4 <sup>th</sup> jet b-tag > 0.09	2215	1544	11017
Chisq < 400	2161	1487	1822
M_Lep > 40 GeV	2135	1330	1778
TotalM_Event > 420 GeV	1871	873	711
100 GeV < M_bb < 140 GeV	1513	662	464
Efficiency	7.57	2.76	0.11

# Scaled Signal and Background distributions after all selection cuts



# Measuring top-Higgs Yukawa coupling (Eur.Phys.J.C 49, 489-497(2007))

- The Yukawa coupling is scaled to the fermion mass:  $g_{ffH} = \frac{m_f}{v}$ ,  $v$  is the vacuum expectation value of the Higgs field = 246 GeV
- For selection efficiency of the signal ( $\varepsilon$ ) and purity of the selected sample ( $\rho$ ), systematic and statistical uncertainties are given by:

$$\left(\frac{\Delta g_{t\bar{t}H}}{g_{t\bar{t}H}}\right)_{stat} \approx \frac{1}{S_{stat}(g_{t\bar{t}H}^2) \sqrt{\epsilon_{signal}^{sel} \rho_{sample}^{sel} L}} \quad \left(\frac{\Delta g_{t\bar{t}H}}{g_{t\bar{t}H}}\right)_{syst} \approx \frac{1}{S_{syst}(g_{t\bar{t}H}^2)} \frac{1 - \rho_{sample}^{sel}}{\rho_{sample}^{sel}} \frac{\Delta \sigma_{eff}^{BG}}{\sigma_{eff}^{BG}}$$

- $\Delta\sigma/\sigma$  is the uncertainty in the residual background normalisation mainly from  $t\bar{t}$  pairs. In our case it is 5%.
- The sensitivity factors  $S_{stat}(g_{t\bar{t}H}^2) = \frac{1}{\sqrt{\sigma_{t\bar{t}H}}} \left| \frac{d\sigma_{t\bar{t}H}}{d(g_{t\bar{t}H}^2)} \right|$  and  $S_{syst}(g_{t\bar{t}H}^2) = \frac{1}{\sigma_{t\bar{t}H}} \left| \frac{d\sigma_{t\bar{t}H}}{d(g_{t\bar{t}H}^2)} \right|$  express dependence of cross section on the coupling square which is inversely proportional to the square of  $g_{t\bar{t}H}^2$  due to small cross section of the Higgs radiating off the Z.

# Results

- Expected uncertainty on the coupling measurement. We used  $S_{stat} = 1.5 \text{ fb}^{1/2}$  and  $S_{syst} = 1.98$  with Luminosity  $L = 1000 \text{ fb}^{-1}$

Final State	$\epsilon_{sel}$ (%)	$\sigma_{eff}$ ( $\text{fb}^{-1}$ )
$t\bar{t}H$	$7.57 \pm 0.19$	0.04
$t\bar{t}$	$0.11 \pm 0.00$	0.29
$t\bar{t}Z$	$2.76 \pm 0.12$	0.02

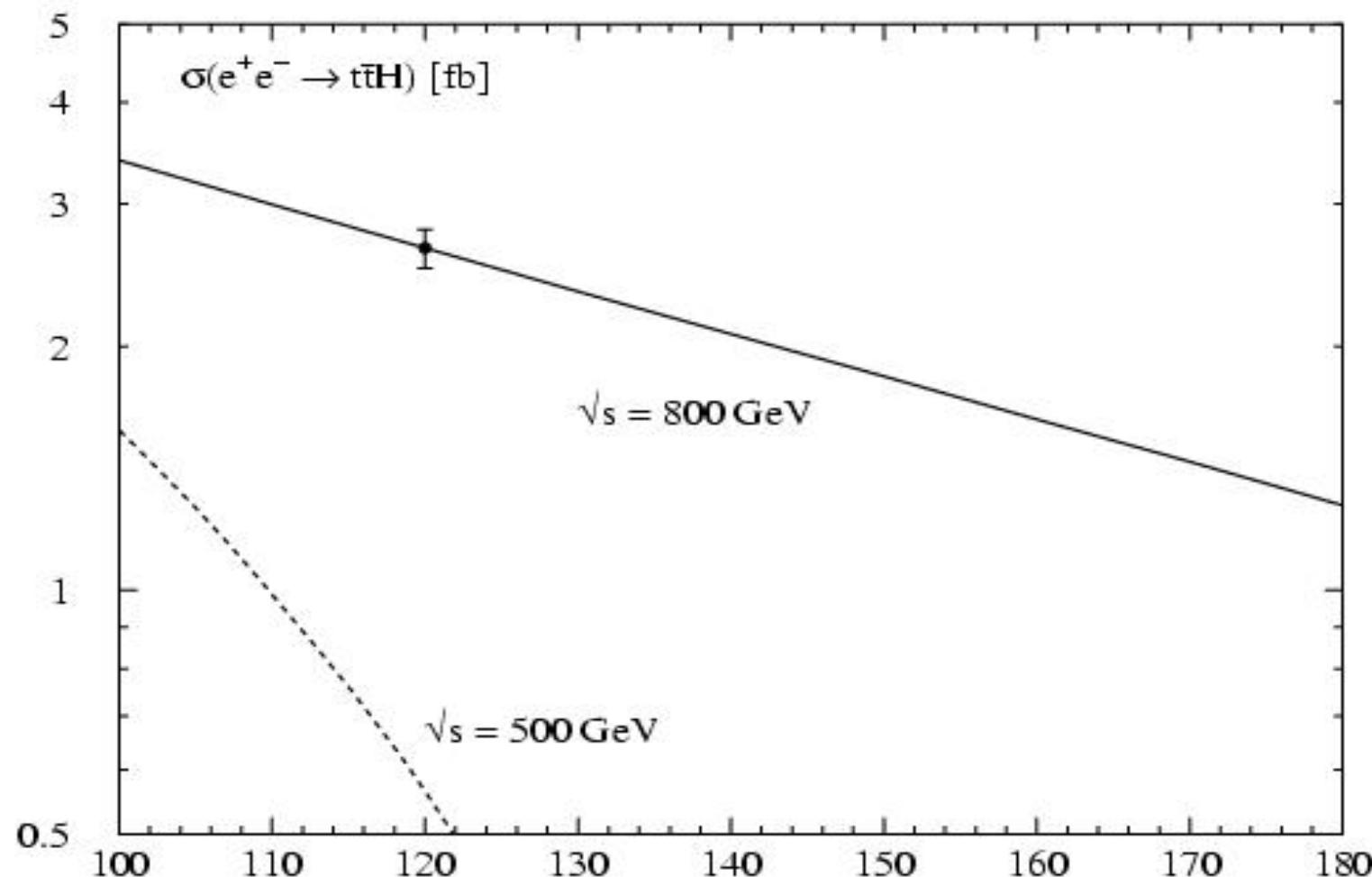
Parameter	value (%)
$\frac{\Delta\sigma_{eff}^{BG}}{\sigma_{eff}^{BG}}$	5
$\epsilon_{sel}$	$7.6 \pm 0.2$
$\rho_{sample}^{sel}$	$12.5 \pm 0.3$
$\left( \frac{\Delta g_{t\bar{t}H}}{g_{t\bar{t}H}} \right)_{stat}$	21.6
$\left( \frac{\Delta g_{t\bar{t}H}}{g_{t\bar{t}H}} \right)_{syst}$	17.6
$\frac{\Delta g_{t\bar{t}H}}{g_{t\bar{t}H}}$	27.9

# Conclusion/Discussion

- For  $L = 1000 \text{ fb}^{-1}$ ,  $\sqrt{s}=500 \text{ GeV}$ , the measured uncertainty in the top-Higgs-Yukawa coupling is 27.9%
- Improvements in hadronic W reconstruction, can refine the results
- Multivariate can be used to improve the background suppression
- The largest uncertainty is due to the uncertainty  $\Delta\sigma/\sigma$  on the background normalisation  $\sigma_{\text{tt}}$ , which is taken to be 5%. Updated values of this uncertainty can change the results.
- It will be necessary to extend this study to a higher energy but b-tagging will be needed to be optimised.

# Extra slides

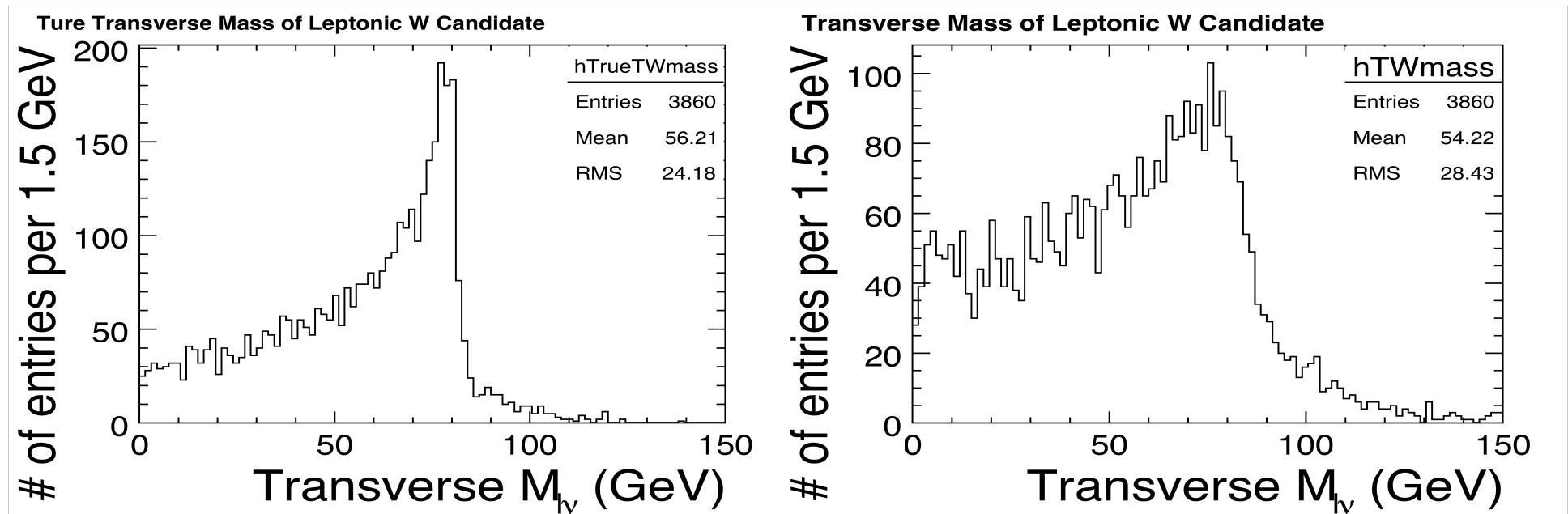
# Cross section of $\bar{t} t H$



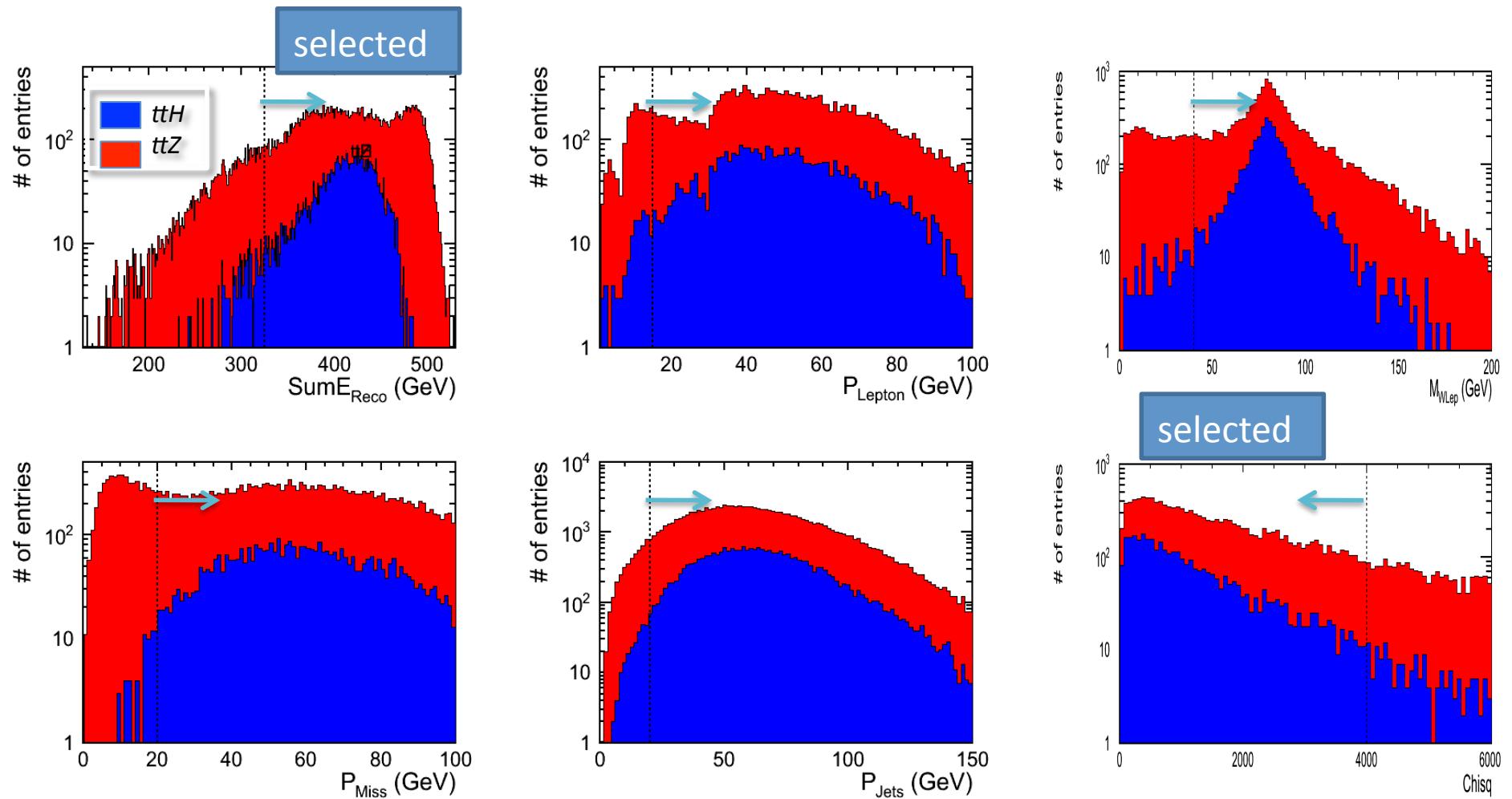
# $W \rightarrow \nu\ell$ (Transverse Mass)

- It is worth trying to reconstruct the transverse mass of W.
- The transverse momentum of lepton and Missing Momentum are used to reconstruct the transverse mass of W.

$$M_{WT} = \sqrt{2p_T^l p_T^\nu (1 - \cos(\phi^l - \phi^\nu))}$$

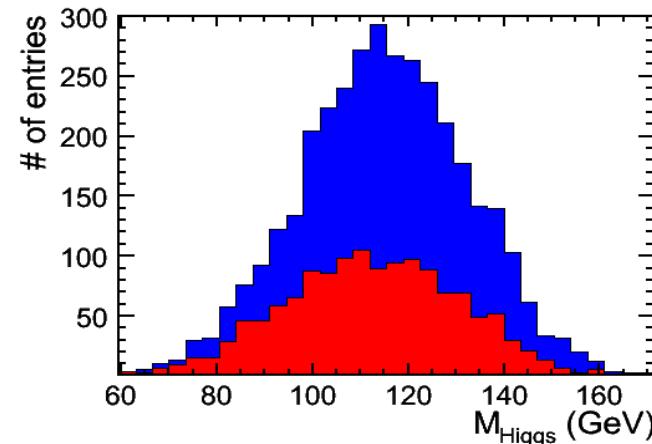
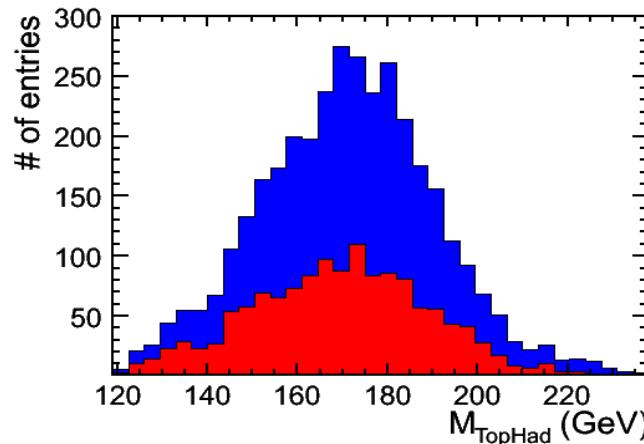
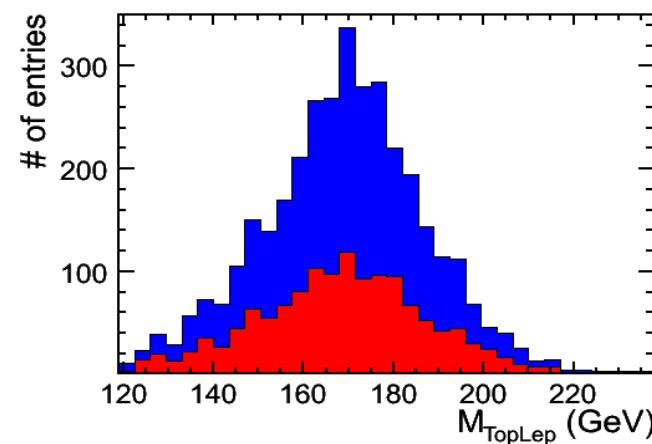
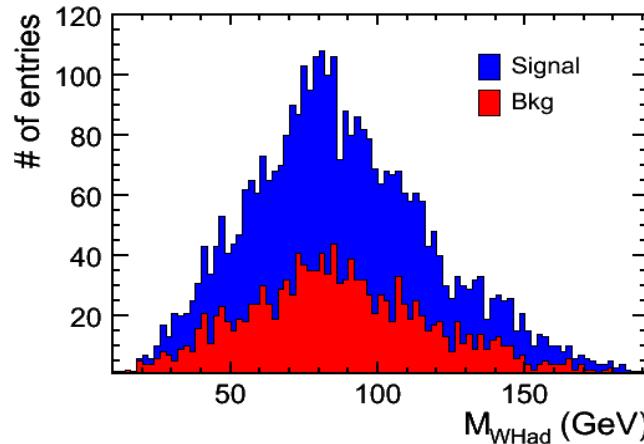


# Selection variables



Signal and background are arbitrarily normalised  
Cuts are selected by optimising  $S/\sqrt{S+B}$

# Signal and Background Final State after applying selection cuts



5/24/11  
Signal and background are arbitrarily normalised