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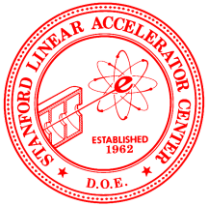
# Linear Normal Form and "Ascript"

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**November 6-17, 2011**

*6<sup>th</sup> International Accelerator School for  
Linear Collider,  
Pacific Grove, California, USA*



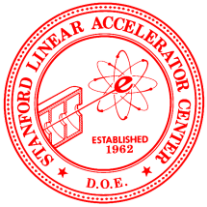
# "Ascript"

- Definition

- "Ascript" is a symplectic transformation from the normal to physical coordinates

- Why ascript?

- Only have to deal with real matrix and TPSA
- Relate to a rotation
- Closer to conventional treatment such as Courant-Snyder parameters
- Natural extension from one-dimensional case
- Include coupling and effects of errors



# Courant-Snyder Parameters

One-turn matrix:

$$M = \begin{pmatrix} \cos \mu + \alpha \sin \mu & \beta \sin \mu \\ -\gamma \sin \mu & \cos \mu - \alpha \sin \mu \end{pmatrix}$$

Rotation matrix:

$$R = \begin{pmatrix} \cos \mu & \sin \mu \\ -\sin \mu & \cos \mu \end{pmatrix}$$

We have:

$$M = ARA^{-1}$$

where  $A^{-1}$  is a transformations from physical to normalized coordinates:

$$A^{-1} = \begin{pmatrix} \frac{1}{\sqrt{\beta}} & 0 \\ \frac{\alpha}{\sqrt{\beta}} & \sqrt{\beta} \end{pmatrix}, A = \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix}$$

$A$  is an "ascript" and is not unique. Since two-dimensional rotational group is commutative  $AR(q)$  is also an ascript. Courant and Snyder choose to have  $A_{12}=0$ .



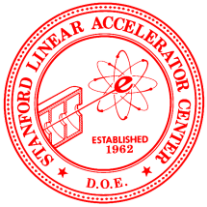
# Symplectic Matrix

$M$  is a symplectic matrix if it has the property that

$$MJM^T = J,$$

where  $J$  is

$$J = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & -1 & 0 \end{pmatrix}$$



# How to Construct "Ascript"

We use eigen vectors to construct a complex symplectic matrix

$$U = [E_I, iE_{-I}, E_{II}, iE_{-II}, E_{III}, iE_{-III}],$$

which is symplectic and has the property that

$$U^{-1}MU = \Lambda = \text{diag}(e^{i2\pi\nu_I}, e^{-i2\pi\nu_I}, e^{i2\pi\nu_{II}}, e^{-i2\pi\nu_{II}}, e^{i2\pi\nu_{III}}, e^{-i2\pi\nu_{III}})$$

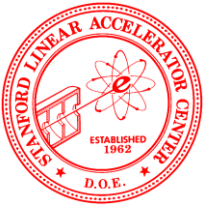
"Ascript" is defined as  $A=UK$  has the property that

$$A^{-1}MA = R = K^{-1}AK$$

Further more  $A$  is symplectic and real.

Clearly, it is an extension of one dimension case.

$$K = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i & 0 & 0 & 0 & 0 \\ -i & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -i & 0 & 0 \\ 0 & 0 & -i & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & -i \\ 0 & 0 & 0 & 0 & -i & 1 \end{pmatrix}$$



# A Solution of Ascript

Explicitly, ascript can be written

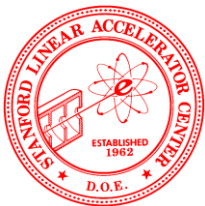
$$A = \sqrt{2}[\text{Re } E_I, \text{Im } E_I, \text{Re } E_{II}, \text{Im } E_{II}, \text{Re } E_{III}, \text{Im } E_{III}],$$

The eigen vectors are normalized as

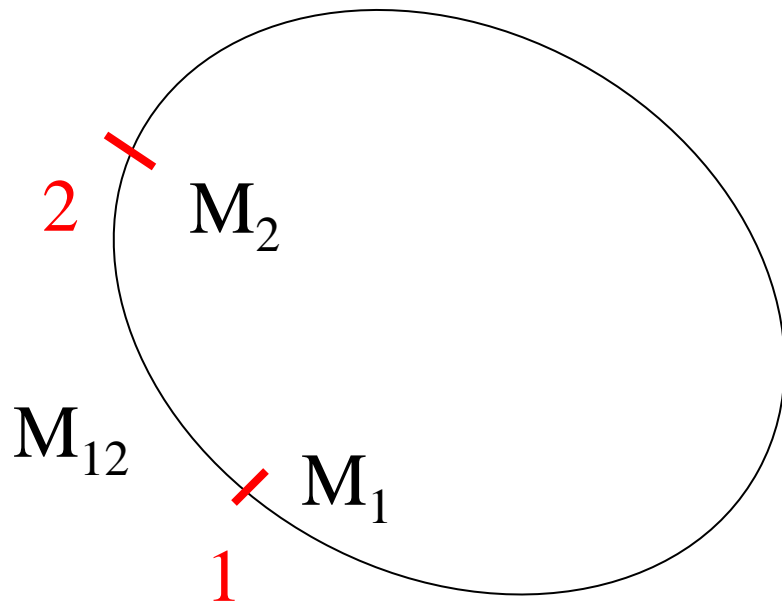
$$E_{I,II,III}^{T*} J E_{I,II,III} = i,$$

$$E_{-I,-II,-III}^{T*} J E_{-I,-II,-III} = -i$$

How to get ascript directly from the one-turn matrix? Given ascript, we have  $U=AK^{-1}$ , which we should use in our map analysis. How about propagation of  $U$ ?  $A_2=T_{12}^*A_1$  leads to  $U_2=T_{12}^*U_1$ . But that implies we need to write force in complex, That is rather "dangerous". Therefore, we should use the complex coordinates only in the analysis.



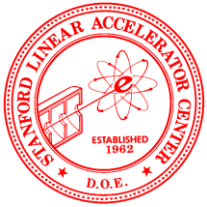
# Propagation of "Ascript"



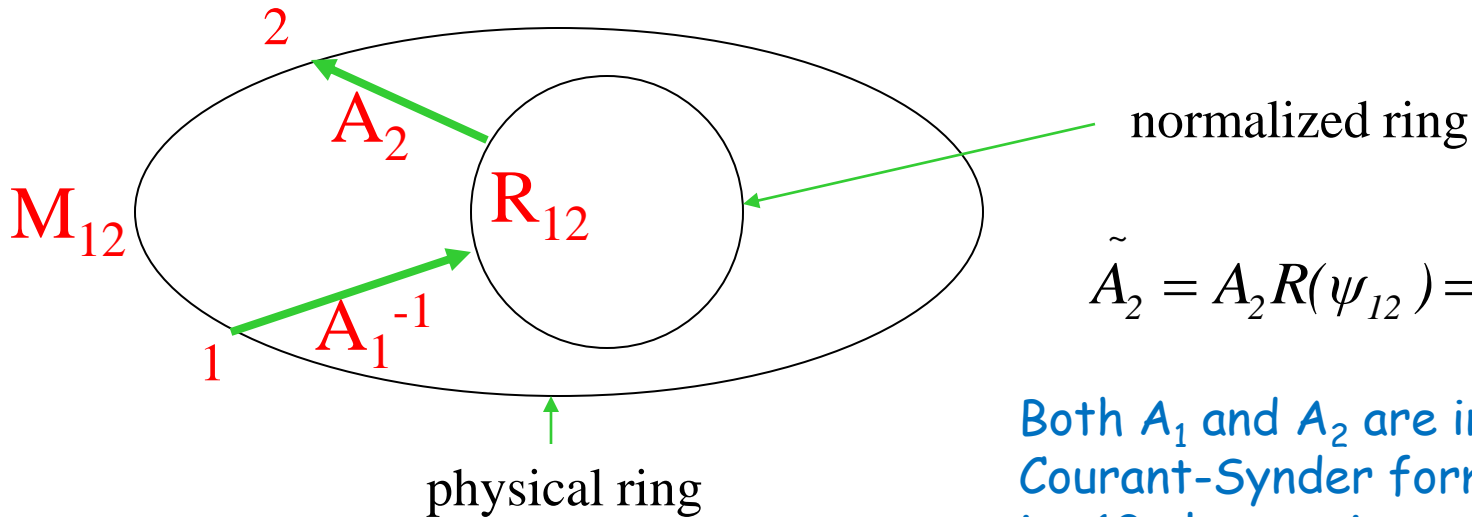
$M_1$  and  $M_2$  are one-turn matrices at position 1 and respectively.  $M_{12}$  is the transport matrix from 1 to 2. It is easy to show

$$M_2 = M_{12} M_1 M_{12}^{-1}$$

As a result of this identity,  $\tilde{A}_2 = M_{12} A_1$  is an "ascript" At position 2 if  $A_1$  is an "ascript" at position 1. We do not need to solve eigen vectors at every position in the ring.



# Propagation of "Lattice Functions"



$$\tilde{A}_2 = A_2 R(\psi_{12}) = M_{12} A_1$$

Both  $A_1$  and  $A_2$  are in Courant-Snyder form, namely its 12 element is zero.

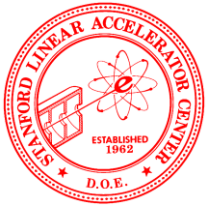
lattice functions at location 2:

$$\beta = \tilde{A}_{11}^2 + \tilde{A}_{12}^2, \alpha = -(\tilde{A}_{11} \tilde{A}_{21} + \tilde{A}_{12} \tilde{A}_{22}), \gamma = \tilde{A}_{21}^2 + \tilde{A}_{22}^2$$

phase advance:

$$\psi_{12} = \tan^{-1} \tilde{A}_{12} / \tilde{A}_{11}$$





# Edwards-Teng Coupling Parameters

Given an one-turn matrix  $M$ , we can decouple it with a symplectic transformation:

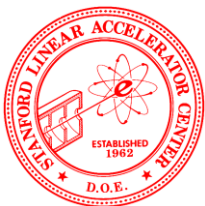
$$M = \begin{pmatrix} gI & \bar{w} \\ -w & gI \end{pmatrix} \begin{pmatrix} u_1 & 0 \\ 0 & u_2 \end{pmatrix} \begin{pmatrix} gI & -\bar{w} \\ w & gI \end{pmatrix},$$

$\swarrow$   $C_{ET}$

where  $u_1$  and  $u_2$  can be parameterized as if no coupling case and  $w$  is a symplectic matrix:

$$u_1 = \begin{pmatrix} \cos \mu_1 + \alpha_1 \sin \mu_1 & \beta_1 \sin \mu_1 \\ -\gamma_1 \sin \mu_1 & \cos \mu_1 - \alpha_1 \sin \mu_1 \end{pmatrix},$$
$$u_2 = \begin{pmatrix} \cos \mu_2 + \alpha_2 \sin \mu_2 & \beta_2 \sin \mu_2 \\ -\gamma_2 \sin \mu_2 & \cos \mu_2 - \alpha_2 \sin \mu_2 \end{pmatrix},$$
$$w = \begin{pmatrix} w_{11} & w_{12} \\ w_{21} & w_{22} \end{pmatrix}.$$

There are **ten** independent parameters. Bar notes symplectic conjugate.  
 $g^2 = 1 - \det(w)$ .

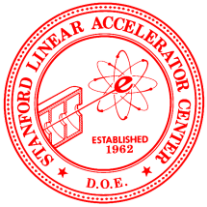


# "Ascript" for Coupled Lattices

$$A = C_{ET} A_{CS} = \begin{pmatrix} g\sqrt{\beta_1} & 0 & \frac{w_{12}\alpha_2 + w_{22}\beta_2}{\sqrt{\beta_2}} & -\frac{w_{12}}{\sqrt{\beta_2}} \\ -\frac{g\alpha_1}{\sqrt{\beta_1}} & \frac{g}{\sqrt{\beta_1}} & -\frac{w_{11}\alpha_2 + w_{21}\beta_2}{\sqrt{\beta_2}} & \frac{w_{11}}{\sqrt{\beta_2}} \\ \frac{w_{12}\alpha_1 - w_{11}\beta_1}{\sqrt{\beta_1}} & -\frac{w_{12}}{\sqrt{\beta_1}} & g\sqrt{\beta_2} & 0 \\ \frac{w_{22}\alpha_1 - w_{21}\beta_1}{\sqrt{\beta_1}} & -\frac{w_{22}}{\sqrt{\beta_1}} & -\frac{g\alpha_2}{\sqrt{\beta_2}} & \frac{g}{\sqrt{\beta_2}} \end{pmatrix}$$

$$g = \sqrt{I - (w_{11}w_{22} - w_{12}w_{21})}, A^{-1} = -JA^T J$$

$A$  is symplectic and its presentation is far from unique. In fact, there are two independent angles. There are eight independent parameters.



# "Symplectic Dispersion Matrix"

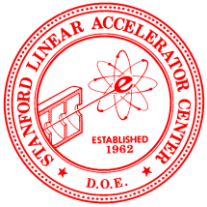
by Ohmi, Hirata, and Oide

$$H = \begin{pmatrix} \left(1 - \frac{|h_x|}{1+a}\right)I & -\frac{h_x \bar{h}_y}{1+a} & h_x \\ -\frac{h_y \bar{h}_x}{1+a} & \left(1 - \frac{|h_y|}{1+a}\right)I & h_y \\ -\bar{h}_x & -\bar{h}_y & aI \end{pmatrix}$$

$h_x$  and  $h_y$  are 2x2 matrices and parameter  $a$  is related to their determinates by

$$a^2 + |h_x| + |h_y| = 1$$

$H$  has 8 independent parameters. Four parameters describe dispersions and the other fours for "crab dispersions"



# A Symplectic Factorization of "Ascript"

$$A = H_{OHO} C_{ET} A_{CS} R(\psi_1, \psi_2, \psi_3)$$

- $H_{OHO}$  is a dispersion matrix by Ohmi, Hirata, and Oide (8 independent parameters)
- $C_{ET}$  is coupling matrix by Edwards and Teng (4 independent parameters)
- $A_{CS}$  is "three two-dimensional ascripts" in Courant-Snyder form (6 independent parameters)
- $R(\psi_1, \psi_2, \psi_3)$  are "three rotation matrix" for phase advances (3 independent parameters)
- $A$  has 21 independent parameters, which is the dimensionality of 6x6 symplectic matrix