# Implications of the latest experimental results for physics at a future Linear Collider

Georg Weiglein

DESY

Granada, 09 / 2011

- Introduction
- Physics of electroweak symmetry breaking
- New physics addressing the hierarchy problem
- Implications for Linear Collider physics

#### **Introduction**

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LHC results so far, executive summary:

Impressive rediscovery of the known ingredients of the Standard Model

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No evidence for new physics yet

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Latest experimental results:

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Latest experimental results:

There will be no discussion of LC implications of v > c neutrinos in this talk ...

#### LHC physics: exploring the Terascale

1 TeV  $\approx 1000 \times m_{\text{proton}} \Leftrightarrow 2 \times 10^{-19} \,\mathrm{m}$ 



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# What can we learn from exploring the new territory of TeV-scale physics?

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- How do elementary particles obtain the property of mass: what is the mechanism of electroweak symmetry breaking? Is there a Higgs boson (or more than one)?
- Do all the forces of nature arise from a single fundamental interaction?
- Are there more than three dimensions of space?
- Are space and time embedded into a "superspace"?
- What is dark matter? Can it be produced in the laboratory?
- Are there new sources of CP-violation? Can they explain the asymmetry between matter and anti-matter in the Universe?



Higgs: last missing ingredient of the Standard Model But: the Standard Model cannot be the ultimate theory

- The Standard Model does not include gravity  $\Rightarrow$  breaks down at the latest at  $M_{\text{Planck}} \approx 10^{19} \text{ GeV}$
- "Hierarchy problem": M<sub>Planck</sub>/M<sub>weak</sub> ≈ 10<sup>17</sup>
  How can two so different scales coexist in nature?
  Via quantum effects: physics at M<sub>weak</sub> is affected by physics at M<sub>Planck</sub>
  - $\Rightarrow$  Instability of  $M_{\text{weak}}$
  - ⇒ Would expect that all physics is driven up to the Planck scale
- Nature has found a way to prevent this The Standard Model provides no explanation Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 – p.5

## Hierarchy problem: how can the Planck scale be so much larger than the weak scale?

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#### Extra dimensions of space:

Fundamental Planck scale is  $\sim {\rm TeV}$  (large extra dimensions), hierarchy of scales is related to a "warp factor" ("Randall–Sundrum" scenarios)

### Supersymmetry (SUSY)

Supersymmetry: fermion ←→ boson symmetry, leads to compensation of large quantum corrections





# The Minimal Supersymmetric Standard Model (MSSM)

Superpartners for Standard Model particles:

 $\begin{bmatrix} u, d, c, s, t, b \end{bmatrix}_{L,R} \begin{bmatrix} e, \mu, \tau \end{bmatrix}_{L,R} \begin{bmatrix} \nu_{e,\mu,\tau} \end{bmatrix}_{L} \quad \text{Spin } \frac{1}{2}$   $\begin{bmatrix} \tilde{u}, \tilde{d}, \tilde{c}, \tilde{s}, \tilde{t}, \tilde{b} \end{bmatrix}_{L,R} \begin{bmatrix} \tilde{e}, \tilde{\mu}, \tilde{\tau} \end{bmatrix}_{L,R} \begin{bmatrix} \tilde{\nu}_{e,\mu,\tau} \end{bmatrix}_{L} \quad \text{Spin } 0$   $g \quad \underbrace{W^{\pm}, H^{\pm}}_{\tilde{\chi}_{1,2}} \quad \underbrace{\gamma, Z, H_{1}^{0}, H_{2}^{0}}_{\tilde{\chi}_{1,2,3,4}} \quad \text{Spin } 1 \text{ / Spin } 0$   $\underbrace{\tilde{g}} \quad \check{\chi}_{1,2}^{\pm} \quad \underbrace{\tilde{\chi}_{1,2,3,4}^{0}}_{\tilde{\chi}_{1,2,3,4}} \quad \text{Spin } \frac{1}{2}$ 

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 $\tilde{g} \qquad \tilde{\chi}_{1,2}^{\pm} \qquad \tilde{\chi}_{1,2,3,4}^{0} \qquad \qquad \mathsf{Spin}\ \frac{1}{2}$ 

Two Higgs doublets, physical states:  $h^0, H^0, A^0, H^{\pm}$ 

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Two Higgs doublets, physical states:  $h^0, H^0, A^0, H^{\pm}$ 

General parametrisation of possible SUSY-breaking terms  $\Rightarrow$  free parameters, no prediction for SUSY mass scale

Hierarchy problem  $\Rightarrow$  expect observable effects at TeV scale

#### How does SUSY breaking work?

Exact SUSY  $\Leftrightarrow m_{\rm e} = m_{\rm \tilde{e}}, \ldots$ 

⇒ SUSY can only be realised as a broken symmetry

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MSSM: no particular SUSY breaking mechanism assumed, parameterisation of possible soft SUSY-breaking terms

- ⇒ relations between dimensionless couplings unchanged
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- Most general case: 105 new parameters

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Strong phenomenological constraints on flavour off-diagonal and  $\mathcal{CP}$ -violating SUSY-breaking terms

 $\Rightarrow$  Good phenomenological description for universal SUSY-breaking terms ( $\approx$  diagonal in flavour space)

### Simplest ansatz: the Constrained MSSM (CMSSM)

Assume universality at high energy scale ( $M_{GUT}$ ,  $M_{Pl}$ , ...) renormalisation group running down to weak scale require correct value of  $M_Z$ 

 $\Rightarrow$  CMSSM characterised by

$$m_0^2, m_{1/2}, A_0, \tan\beta, \, \mathrm{sign}\,\mu$$

CMSSM has been the "favourite toy" for both theorists and experimentalists so far

CMSSM is in agreement with the experimental constraints from electroweak precision observables (EWPO) + flavour physics + cold dark matter density + ...

### SUSY-breaking scenarios

"Hidden sector": → Visible sector: SUSY breaking MSSM "Gravity-mediated": SUGRA "Gauge-mediated": GMSB "Anomaly-mediated": AMSB "Gaugino-mediated"

SUGRA: mediating interactions are gravitational

GMSB: mediating interactions are ordinary electroweak and QCD gauge interactions

AMSB, Gaugino-mediation: SUSY breaking happens on a different brane in a higher-dimensional theory

#### Models with extra dimensions of space





Brane-world Picture

Hierarchy between  $M_{\text{Planck}}$  and  $M_{\text{weak}}$  is related to the volume or the geometrical structure of additional dimensions of space

 $\Rightarrow$  observable effects at the TeV scale

### Gravity in a warped spacetime geometry

[L. Randall, LHC2TSP Workshop '11]

## Natural for gravity to be weak!



- Small probability for graviton to be near the Weakbrane
- If we live anywhere but the Gravitybrane, gravity will seem weak
- Natural consequence of warped geometry

 $ds^2 = g_{MN} dx^M dx^N = e^{-2\sigma} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$ ,

# Phenomenological consequences of extra dimensions

The wave function of a free particle must be  $2\pi R$  periodic



$$e^{ip.x_5} = e^{ip.(x_5 + 2\pi R)}$$

 $p = \frac{n}{R}$ 

 $\Rightarrow$  momentum is quantised

#### ⇒ Looks in 4-dim like a series of new, more massive partners associated with each known particle: "Kaluza–Klein tower"

### Workshop "Implications of LHC results for TeV-scale physics"

Kick-off meeting: 29/08/2011–02/09/2011, CERN,  $\gtrsim 200$  participants

⇒ Discuss impact of experimental results on future strategy for particle physics

Results will be summarised in a document to be submitted as input for the 2012 update of the European Strategy for Particle Physics (in time for "Orsay-type" meeting of strategy update, 09/2012)

Main organisers:

O. Buchmueller, P. De Jong, A. De Roeck, J. Ellis, C. Grojean,

S. Heinemeyer, J. Hewett, K. Jakobs, M. Mangano, F. Teubert, G. W. Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 – p.15

### Workshop "Implications of LHC results for TeV-scale physics"

Three working groups:

- WG1: Signals of electroweak symmetry breaking
  Conv.: S. Heinemeyer, M. Kado, C. Mariotti, G. W., A. Weiler
- WG2: Signatures with missing energy
  Conv.: R. Cavanaugh, J. Hewett, S. Kraml, G. Polesello
- WG3: Other signatures of possible BSM physics
  Conv.: C. Grojean, D. Martinez, J. Santiago Perez, P. Savard, S. Worm

 $\Rightarrow$  It is now the right time to join in to this activity!

### Physics of electroweak symmetry breaking

What is the mechanism of electroweak symmetry breaking?

- Standard Model (SM), SUSY, ...:
  Higgs mechanism, elementary scalar particle(s)
- Strong electroweak symmetry breaking: a new kind of strong interaction
- Higgsless models in extra dimensions: boundary conditions for SM gauge bosons and fermions on Planck and TeV branes in higher-dimensional space

#### $\Rightarrow$ New phenomena required at the TeV scale

### Higgs phenomenology: SM and beyond

Standard Model: a single parameter determines the whole Higgs phenomenology:  $M_{\rm H}$ 

Branching ratios of the SM Higgs:



⇒ dominant BRs:  $M_{\rm H} \lesssim 140$  GeV:  $H \rightarrow b\overline{b}$   $M_{\rm H} \gtrsim 140$  GeV:  $H \rightarrow W^+W^-, ZZ$ 

#### Production of a SM-like Higgs at the LHC

- SM Higgs production at the LHC:
- Dominant production processes:
- gluon fusion:  $gg \to H$ , weak boson fusion (WBF):  $q\bar{q} \to q'\bar{q}'H$



# Constraints on the SM Higgs from electroweak precision data

Indirect constraint on  $M_{H_{SM}}$ , no direct search limits included in the fit  $6^{July 2011}$   $m_{Limit} = 161 \text{ GeV}$  [LEPEWWG '11]



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### Higgs physics beyond the SM

In the SM the same Higgs doublet is used "twice" to give masses both to up-type and down-type fermions

- ⇒ extensions of the Higgs sector having (at least) two doublets are quite "natural"
- $\Rightarrow$  Would result in several Higgs states

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Many extended Higgs theories have over large part of their parameter space a lightest Higgs scalar with properties very similar to those of the SM Higgs boson

Example: SUSY in the "decoupling limit"

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Many extended Higgs theories have over large part of their parameter space a lightest Higgs scalar with properties very similar to those of the SM Higgs boson

Example: SUSY in the "decoupling limit"

But there is also the possibility that none of the Higgs bosons is SM-like

### Higgs physics in Supersymmetry

"Simplest" extension of the minimal Higgs sector:

Minimal Supersymmetric Standard Model (MSSM)

- Two doublets to give masses to up-type and down-type fermions (extra symmetry forbids to use same doublet)
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- SUSY imposes relations between the parameters
- $\Rightarrow$  Two parameters instead of one:  $\tan \beta \equiv \frac{v_u}{v_d}$ ,  $M_A$  (or  $M_{H^{\pm}}$ )

⇒ Upper bound on lightest Higgs mass,  $M_h$  (FeynHiggs): [S. Heinemeyer, W. Hollik, G. W. '99], [G. Degrassi, S. Heinemeyer, W. Hollik, P. Slavich, G. W. '02]  $M_h \lesssim 130 \,\mathrm{GeV}$ 

Very rich phenomenology

## **BSM** $\oplus$ Higgs phenomenology

- Large enhancement / suppression of standard search channels possible
  Example: large enhancement of Hb̄b coupling
  ⇒ large suppression of BR(h → γγ), BR(h → WW\*), ...
- New channels, different phenomenology:
  - Experimental evidence for dark matter
    - $\Rightarrow$  if dark matter particle is lighter than  $M_{\rm H}/2$
    - $\Rightarrow$  large branching fraction into invisible particles
    - $\Rightarrow$  large suppression of all other BRs
  - Higgs production in decays of BSM particles
  - $h_i \rightarrow h_j h_j$  decays
  - Higgs-radion mixing, ...
  - Higgses with nearly degenerate masses: large interference effects, resonance-type behaviour possible Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 p.:

# How to infer the underlying physics from the experimental signatures?

- A Higgs or not a Higgs?
- Fundamental or composite?
- SM, MSSM or beyond?
- Is there other new physics; what is it?
- How does the observed new physics fit into the global picture (ew precision observables, flavour physics, ...)?

⇒ Intense effort will be needed to identify the nature of electroweak symmetry breaking ATLAS SM Higgs search: combined upper limit normalised to the SM expectation (left) and observed result vs. expectation for a SM Higgs signal (right) [ATLAS Collaboration '11]



## SM Higgs search: combined CMS results

Combined confidence limit vs. expectation for a SM Higgs signal [CMS Collaboration '11]



• LHC excludes (at least at 90% C.L.) the range of  $145 \text{ GeV} \lesssim M_{\text{H}_{\text{SM}}} \lesssim 460 \text{ GeV}$ 

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A high mass Higgs?

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However: a heavy SM-like Higgs appears to be theoretically questionable

# Prospects for SM Higgs searches in the high mass region

[W. Murray, LHC2TSP Workshop '11]



#### $\Rightarrow$ Large increase in coverage expected

# SM Higgs search: ATLAS and CMS results in the low mass region

Combined upper limit normalised to the SM expectation, low mass region

#### [ATLAS Collaboration '11]

[CMS Collaboration '11]



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# CMS results for SM Higgs searches: local p value and observed best-fit signal strength



 $\Rightarrow$  With LEE: probability to see an excess at least as large as the one observed in the data is  $\approx 0.4$ 

Best compatibility with a SM Higgs for  $M_{\rm H_{SM}} \lesssim 125~{\rm GeV}$ 

# ATLAS results for SM Higgs searches: local p value vs. expectation for a SM Higgs signal

[ATLAS Collaboration '11]



⇒ Best compatibility with a SM Higgs for  $M_{\rm H_{SM}} \lesssim 130 {\rm ~GeV}$ Slight deficit w.r.t. SM expectation

## SM Higgs search: Tevatron results, CDF + D0

CDF + D0 combined upper limit normalised to the SM expectation [CDF and D0 Collaborations '11]



Tevatron Run II Preliminary,  $L \le 8.6 \text{ fb}^{-1}$ 

#### $\Rightarrow$ Broad excess in low mass region

• Search for SM Higgs has narrowed down to  $114 \text{ GeV} \lesssim M_{\text{H}_{\text{SM}}} \lesssim 135 \text{ GeV}$  (+ high mass region)

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- Broad excess in low mass region observed by ATLAS, CMS and the Tevatron
   Note: H → WW\* channel has very limited mass resolution
- What about a Higgs with M<sub>H</sub> ≈ 145 GeV with somewhat reduced σ × BR(H → WW\*) compared to SM case?
  Difficult to get sufficiently large BR(H → WW\*) in the MSSM, can better be accomodated in the NMSSM

#### CMS excess in $H \rightarrow \gamma \gamma$ search

#### [W. Murray, LHC2TSP Workshop '11]



#### $\Rightarrow 1.6\sigma$ excess at $M_{\rm H} \approx 140 \ {\rm GeV}$ after taking into account LEE

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# Search for the heavy SUSY Higgs bosons H, A: *limits in the* $M_A$ -tan $\beta$ *plane*

[ATLAS Collaboration '11]

[CMS Collaboration '11]



CMS Preliminary 2011 1.6 fb<sup>-1</sup>

# Search for the heavy SUSY Higgs bosons *H*, *A*: cross section limit from CMS

[CMS Collaboration '11]



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## Fundamental or composite Higgs?

Renewed interest in composite Higgs models, mostly from extra dimensions

[N. Arkani-Hamed, A. Cohen, H. Georgi '01]

[K. Agashe, R. Contino, A. Pomarol '05], ...

Composite Higgs: light remnant of a strong force

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Relation extra dimensions  $\Leftrightarrow$  new strong forces?

- Correspondence (AdS/CFT):
- Warped gravity model  $\Leftrightarrow$  Technicolour-like theory in 4D

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Signatures at LHC: new resonances, W', Z', t', KK excitations Under pressure from electroweak precision tests

# Effective field-theory description of a composite Higgs

- Agreement with electroweak precision data can be improved if there is a strongly interacting light Higgs, e.g.
- Little Higgs [N. Arkani-Hamed, A. Cohen, E. Katz, A. Nelson '02] Holographic Higgs [R. Contino, Y. Nomura, A. Pomarol '03], [K. Agashe, R. Contino, A. Pomarol '05], ...
- Effective Lagrangian formalism for model-independent analysis of effects of a Strongly-Interacting Light Higgs (SILH) [*G. Giudice, C. Grojean, A. Pomarol, R. Ratazzi '07*]
- ⇒ Specific pattern of modified Higgs couplings Strong WW scattering at high energies despite light Higgs
- ⇒ Need precision measurement of Higgs couplings
  + test of longitudinal gauge-boson scattering

# Strongly-Interacting Light Higgs: deviation of $\sigma \times BR$ from the case of a SM Higgs

[G. Giudice, C. Grojean, A. Pomarol, R. Ratazzi '07]



Sensitivity at LHC: 20–40%, ILC: O(1%) $\Rightarrow$  ILC can test scales up to  $\sim 30$  TeV

#### Further prospects

Prospects for searches for a 114 GeV SM-like Higgs:

[W. Murray, LHC2TSP Workshop '11]



 $\Rightarrow$  2011 data, when combined between ATLAS + CMS, should provide  $2\sigma$  sensitivity down to  $M_{\rm H} = 114~{\rm GeV}$ 

#### Prospects for the signal significance

#### [W. Murray, LHC2TSP Workshop '11]



 $\Rightarrow \text{With 2012 data, ATLAS} + \text{CMS combined:} \\ \text{expect sensitivity of at least } 3.5\sigma \\ \text{Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09/2011 - p.41} \\ \end{array}$ 

### New physics addressing the hierarchy problem

SUSY production cross sections at the LHC with 7  $\mathrm{TeV}$ :



⇒ Highest cross section for gluino and squarks of the first two generations

Squark and gluino couplings  $\sim \alpha_s$ ; cross sections mainly determined by  $m_{\tilde{q},\tilde{g}}$ , small residual model dependence

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### SUSY searches at the LHC

Dominated by production of coloured particles: gluino, squarks (mainly first two generations)

Very large mass reach in the searches for jets + missing energy

 $\Rightarrow$  gluino, squarks accessible up to 2–3 TeV at LHC (14 TeV)

Coloured particles are usually heavier than the colour-neutral ones

⇒ long decay chains possible; complicated final states

**e.g.:** 
$$\tilde{g} \to \bar{q}\tilde{q} \to \bar{q}q\tilde{\chi}_2^0 \to \bar{q}q\tilde{\tau}\tau \to \bar{q}q\tau\tau\tilde{\chi}_1^0$$

Many states produced at once, difficult to disentangle

#### **Pre-LHC:** Fit results for the CMSSM

#### from precision data

Comparison: preferred region in the  $m_0-m_{1/2}$  plane vs. prospective CMS 95% C.L. reach for 0.1, 1 fb<sup>-1</sup> at 7 TeV

[O. Buchmueller, R. Cavanaugh, A. De Roeck, J. Ellis, H. Flächer, S. Heinemeyer, G. Isidori, K. Olive, P. Paradisi, F. Ronga, G. W. '10]



#### $\Rightarrow$ Best fit point was within the 95% C.L. reach with 1 fb<sup>-1</sup>

### SUSY search results for the CMSSM



400

300

200

#### [CMS Collaboration '11]



 $\Rightarrow$  High sensitivity from search for jets + missing energy Previous best-fit point is excluded CMSSM starts to get under pressure Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 - p.45
# Interpretation of SUSY search result in "simplified model"

"Simplified model": squarks of first two generations, gluino + massless neutralino (LSP), all other SUSY particles heavy [ATLAS Collaboration '11]



# Limits for gluinos and squarks in simplified models, LSP mass varied from 0 to $m_{\tilde{g}} - 200 \text{ GeV}$



[CMS Collaboration '11]

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# **Excitement at SUSY '11**

#### Stop production in gluino decays

[ATLAS Collaboration '11]



 $\Rightarrow$  Observed limit decreased with 30  $\times$  more luminosity 1.2  $\sigma$  excess in both electron and muon channels

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# Global fit in the CMSSM including 2011 LHC data ( $1 \text{ fb}^{-1}$ ) and XENON100 results

68% and 95% CL contours, pre- and post-LHC

[O. Buchmueller, R. Cavanaugh, A. De Roeck, M. Dolan, J. Ellis, H. Flächer,

S. Heinemeyer, G. Isidori, D. Martínez Santos, K. Olive, S. Rogerson, F. Ronga,

G. W. '11]



 ⇒ Preferred region "opens up", overall χ<sup>2</sup> worsened Shift towards higher mass scales, higher values of tan β
Comparison: GMSB yields much larger splitting between coloured and colour-neutral part of the spectrum

- Search for jets (+ leptons) + missing energy
  - $\Rightarrow$  Bounds on gluino and squarks of first two generations of  $\mathcal{O}(\text{ TeV})$

- Search for jets (+ leptons) + missing energy
  - $\Rightarrow$  Bounds on gluino and squarks of first two generations of  $\mathcal{O}({\rm ~TeV})$
  - $\Rightarrow$  The constrained scenario CMSSM starts to get under some tension: direct search limits vs.  $(g-2)_{\mu}$

- Search for jets (+ leptons) + missing energy
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Reduced sensitivity to compressed spectra

 Limited sensitivity to 3rd generation squarks Hardly any direct constraints from the LHC on colour neutral SUSY particles up to now

## SUSY searches: what next?

#### [S. Padhi (CMS), LHC2TSP Workshop '11]

Assuming colored particles (1<sup>st</sup> and 2<sup>nd</sup> generation squarks and gluinos) are beyond the LHC range:

#### a) Need dedicated exclusive studies to constrain stops and sbottoms

- With and without the cross section help from the colored particles
- See also M. Papucci's EPS-2011 talk
- http://indico.in2p3.fr/contributionDisplay.py?contribId=904&sessionId=6&confId=5116

#### b) Need dedicated activity on EWK inos

- Current limits on Chargino/neutralinos are low
- Explore LHC reach for the electroweak sector (See also Shufang Su SUSY-11 talk)



# Search for dilepton resonances: ATLAS

#### [ATLAS Collaboration '11]



Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 - p.52

# Search for dilepton resonances: CMS

[CMS Collaboration '11]

# Limits with dimuons, dielectrons



$\Lambda_T$ [TeV] (GRW)	$M_s$ [TeV] (HLZ)					
	n = 2	n = 3	n = 4	n = 5	n = 6	n = 7
	ADD k-f	actor: 1.	0			
2.62	2.58	3.12	2.62	2.36	2.20	2.08
2.56	2.58	3.10	2.56	2.27	2.09	1.95
	ADD k-f	actor: 1.	3			
2.70	2.72	3.22	2.70	2.44	2.28	2.16
2.66	2.72	3.20	2.66	2.37	2.17	2.02

EXO-11-019

 $Z'_{SSM}$ : 1940 GeV  $Z'_{\psi}$ : 1620 GeV KK: 1450 GeV ( $\frac{k}{M} = 0.05$ ) KK: 1780 GeV ( $\frac{k}{M} = 0.1$ ) Exclusion limits for SSM, superstring-inspired, RS KK (1.5-2 TeV, as well as ADD models for several parameters (2-3 TeV)

EXO-11-039

# Search for dijet resonances: ATLAS



### Search for dijet resonances: CMS

[CMS Collaboration '11]

# Resonances: limits with dijets



Derived limits for several models, with excluded masses up to 4 TeV

Model	Excluded Mass (TeV)			
	Observed	Expected		
String Resonances	4.00	3.90		
$E_6$ Diquarks	3.52	3.28		
Excited Quarks	2.49	2.68		
Axigluons/Colorons	2.47	2.66		
W' Bosons	1.51	1.40		

arXiv.1107.4771 (submitted to PLB) EXO-11-015

Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 – p.55

### Status of search for heavy resonances

#### No evidence for a signal

• Limits of 
$$\mathcal{O}(1-4 \text{ TeV})$$

### Search for the rare decay $B_s \rightarrow \mu^+ \mu^-$



Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 – p.57

# $BR(B_s \rightarrow \mu^+ \mu^-)$ : combined result from LHCb and CMS

 $B_s \rightarrow \mu^+\mu^-$ : combination with CMS



This is ~ 3 times the SM BR

⇒ Very good agreement with SM expectation (so far) Implications of the latest experimental results for physics at a future Linear Collider, Georg Weiglein, LCWS11, Granada, 09 / 2011 – p.58

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⇒ Preference for low-mass region has further strengthened

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  - coloured states of new physics
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- - - -

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 $\Rightarrow$  Experimental situation for colour–neutral sector  $+~{\rm \tilde{t}},~{\rm \tilde{b}}$  is essentially unchanged compared to the pre-LHC case

# LC implications (personal view)

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In view of the current experimental situation the strongest physics case exists for a

"Higgs factory": LC with  $\sqrt{s} \lesssim 250 \text{ GeV}$ (+ running at WW threshold and Z resonance)

or a

- "Higgs + top factory": LC with  $\sqrt{s} \lesssim 350 \text{ GeV}$
- $\Rightarrow$  Decay-mode independent search (recoil method)
  - Absolute measurement of ZZH coupling
  - Precise measurement of Higgs mass, width, couplings,
  - branching ratios (also Higgs  $\rightarrow$  invisible)
  - Determination of Higgs spin and quantum numbers, ...
  - Precision top + electroweak physics

• Top Yukawa coupling: ILC500 yields 10% acc. for  $1ab^{-1}$ Best sensitivity for 120 GeV Higgs at  $\sqrt{s} \approx 700$  GeV

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Results from the LHC run 2011 / 2012 will provide important further information  $\Rightarrow$  Let's stay tuned!