Electron Gun

Electron Gun

	Cathode	Extraction Field	Comments
Pierce type (thermionic DC)	Thermal	Static	Still conventional
Photo Cathode DC Gun	Photo-electron	Static	For special cathode
Photo-cathode RF Gun	Photo-electron	RF	Advanced
Thermionic RF Gun	Thermal	RF	Advanced

Thermionic DC gun

- Thermionic cathode with DC bias.
- It is a conventional gun widely used.
- Continous beam or
- Bunched beam by grid switching, but the bunch length is down to ~1ns.
- Energy at the gun exit is

$$K = (\gamma - 1)mc^2 = eV \qquad (3 - 1)$$



Thermionic Gun: A typical configuration (1)

- The beam emission is controlled by grid bias. The primary bunching and bunch repetition is determined by the grid pulse duration and repetition.
- The bunch length is shortened by bunchers for further acceleration.



Thermionic Gun: A typical configuration (2)

- Electron beam is extracted from thermionic gun continuously.
- RF cavity (Pre-Buncher) modulates the velocity of the electron beam. After some drift, the beam is bunched.
- The bunch repetition is determined by the Pre-Buncher frequency.
- Further bunching is made by SHB and Buncher.



Thermionic Cathode

- According to Richardson-Dushman equation, material with low work-function operated at high temperature is favor to generate high density electron beam.
- Practical operation temperature is limited by the operable temperature T_e which 10 atomic layers are lost per second at.
- Figure of merit of thermionic cathode is

$$\eta = \frac{\Phi}{T_e} \qquad (3-2)$$

Thermionic Cathode

- $\phi/T_e < 2.0$ is practically used as thermionic cathode.
- Impregnated type BaO cathode is widely used for conventional accelerator.
- CeB₆ and LaB₆ have advantage for high-brightness beam generation.

Material	φ (ev)	Te (K)	φ/T _e (x1E+3)
W	4.5	2860	1.57
Та	4.1	2680	1.53
Мо	4.2	2230	1.88
Cs	1.9	320	5.94
Th-W	2.6	1800	1.44
BaO	1.0	1400	0.71
CeB ₆	2.5	1400	1.79
LaB ₆	2.5	1400	1.79

Photo-Cathode DC Gun (1)

- Electron beam is generated by Photo-emission with laser.
- The bunch structure (repetition and duration) is determined by the laser.
- Beam extraction by a static electric field.
- Beam energy at the gun exit is determined by the bias voltage ,

 $K = (\gamma - 1)mc^2 = eV \qquad (3 - 3)$



Photo-Cathode DC Gun (2)

- Because the velocity at the gun exit is slow, the first cavity is "low β cavity", which synchronizes with the low speed beam.
- Time duration, in which the bunch travels cell length, L, has to be synchronize to the phase advance per cell.

$$L_{cell} = \frac{\Delta \phi}{2\pi f} \beta c \qquad (3-2)$$



Photo-Cathode DC Gun (3)

- In some case, SHB is placed for bunching.
- Low-beta cavity followed by accelerator.



HV Operation (1)

- Since the bunch length at the gun exit is determined by the space charge limit, higher voltage operation makes a higher peak current and bunch length can be shorter.
- Short bunch length from gun has merits
 - Simpler bunching section,
 - Energy spread after acceleration is smaller,
- For higher voltage operation, dark current by field emission from electrode surface should be suppressed.

FIV Operation (2) M. Yamamoto on behalf of F. Furuta



Only 7mm² area exposed maximum field gradient.

ILCWS2005, Snowmass

HV Operation (3)

- Future light source based on ERL employs PC-DC gun.
- For extremely low emittance, HV operation is essential.

HV operation (4)

- Segmented ceramic.
- Guard ring for FE electron.
- HV conditioning up to 550kV.
- 8 hour stable operation at 500kV without any damage.
- It is a world record of PC DC gun.

R. Nagai, RSI(81)033304(2010)

<u>RF Gun (1)</u>

- RF field for beam extraction.
- Electron beam is generated inside of the cavity.
- Laser photo-cathode type is popular.
- Beam from thermionic cathode type has a wide energy spred.

RF Gun(2)

- Typical cavity configuration is 1.5 cells.
- TM01, pi-mode.
- Energy at the gun exit is given by

$$K = (\gamma - 1)mc^{2} = \int e E(z, t) c \beta(t) dt$$

= $\int e \sqrt{RP} \cos(\omega t - \phi) c \beta(t) dt$ (3-4)

P: RF input power, R:shunt impedance

Photo-cathode

- Quantum efficiency, η and temporal response are important property of Photo-cathode.
 - Quantum efficiency determines required laser pulse energy.
 - Temporal response should be even fast to form a short electron bunch, several 10s ps.
- Metal cathode (Cu, Mg) has low η and fast response.
 - η is typically 10⁻⁴~10⁻⁵, response is fast in fs.
- Alkali cathode (CsTe, CsKSb) high η and medium response.
 - η is typically 10⁻¹~10⁻², response is in sub ps.
- NEA GaAs cathode has high η and slow response.
 - η is typically 10⁻¹~10⁻², response is 10s ps

Electron source for Linear Colliders

Design Criteria

•Polarized Electron is essential for linear colliders; The electron must be polarized.

•Higher polarization is better. The world record is 90%, but the specification is determined from technical feasibility, i.e. reproduciability, stability, QE, etc.

•Depending on the polarization, the luminosity is decided to produce enough event rate for physics.

•The linear collider needs 2x10³⁴cm⁻²s⁻¹ luminosity with 80% polarization.

 $L = \frac{f_{rep} n_b N^2}{4\pi\sigma_x \sigma_y}$

ILC Requirements

Parameters		
Pulse length	0.9ms	
Pulse repetition <i>frep</i>	5Hz	
# of bunches in a pulse n _b	2625 (1310)	
Bunch separation	369(670)ns	
# of electrons in a bunch N	2x10 ¹⁰	
Micro bunch length at source	1ns	
Peak current	3.2A	
Electron Polarization	80%	

Basic Concept

- NEA GaAS cathode with circularly polarized laser is the only solution for polarized electrons.
- Beam extraction by a static electric field (DC photocathode gun) because RF gun is not compatible with GaAs cathode.

Electron Gun

E-source for LC

Laser

Bunch extraction Required Laser Pulse Energy

•From QE vs Polarization curve, required laser pulse energy is decided.

 $E_{L}[\mu J] = \frac{124 \times Q[nC]}{\eta[\%] \times \lambda[nm]}$

T. Nishitani et al., J of Appl. Phy. 97,094907(2005)

Electron Gun

E-source for LC

Lase

Bunch extraction Bunch shape

tь

S

•GaAs cathode is only operable in DC bias gun structure. The space charge limit gives possible charge density, J.

•Assuming a resonable spot size, the bunch length in time *t_b* is decided to extract 3.2nC bunch charge.

 $J[A/m^{2}] = 2.33 \times 10^{-6} \frac{V^{3/2}}{d^{2}}$

I[A] = JS

Injector Design

- If the bunch length at gun *t*^b is adequate for RF acceleration, any bunching section is not needed.
- Otherwise, we need a bunching section.
- RF period for the bunching should be long enough comparing to *t_b* for linear modulation.
- RF frequency for the bunching should be harmonics of bunch repetition, i.e. RF of the main linac.

 $T_{bunching} \gg t_b$

 $T_{bunching} = n T_{mainRF}$ $n \in N$

Electron Gun

E-source for LC

laser

Surface Charge Limit (1)

- For Linear colliders, multibunch electron beam are generated.
- Anomalous charge limit phenomena is observed (Surface Charge Limit) for high intensity and multi-bunch beam generation.

GaAs with a Be-dope 5E+18cm³

K. Togawa, NIM A 414 (1998) 431-445

Surface Charge Limit (2)

- The surface charge limit is caused by Photo-voltage effects;
- Some electrons, J_{surface} is captured at BBR(Band Bending Region).
- By the captured electrons, the effective vacuum level is increased.
- Photo-voltage effects decrease size of EA and limit the current.

K. Togawa, NIM A 414 (1998) 431-445

Surface Charge Limit (3)

- SCL was compensated by enhancement of the recombination of the captured electron.
- The recombination was boosted by increasing the positive carrier density in VB by high p-doping.
- Finally, 5.0A/cm² is achieved. It is more than the requirement of ILC gun.

(a) sample 1b(Na=0.5),
(b) sample 2a(Na=1.0), and
(c) sample 3(Na=2.0). The laser
intensity is 1 to 150 W/cm2.

G.A. Mulhollan, Phy. Lett. A 282 (2001)

Electron Emission

nission Rela

Related Physics

Lase

Surface Charge Limit (4)

- Super-lattice Cathode has an advantage against SCL.
- J_{escape} is proportional to the size of NEA.
- The effective size of NEA in Super-lattice cathode is larger than that of bulk-GaAs.
- The escape probability, J_{esacape}/J_{total} is larger for Superlattice cathode. SCL current should be higher for Superlattice cathode.

Super-Lattice Cathode

- GaAs/GaAsP super lattice cathode for high polarization (90%) and high QE (0.5%).
- Heavy P (Zn) -doped GaAs surface layer to suppress SCL.
- Cathode is operated in Space charge limit regime.

Bunching(1)

- According Child-Langmuir law, peak current of ILC Electron gun (120kV, d~5cm, and 1cm diameter) is ~3A.
- To generate ILC bunch (3.2nC), 1.1ns is necessary.
- It is significantly longer than RF acceleration and should be shorten down to 10ps.
- A special section for this purpose is placed at downstream of Electron gun: Bunching section
- SHB(Sub Harmonic Buncher)
 - 216.7 MHz (1/6 of 1.3 GHz)
 - 433 Mhz (1/3 of 1.3 GHz)
 - Buncher : 1.3 G Hz NC tube.

Bunching (2)

- Bunch length is 1ns at the exit of Electron gun.
- Velocity bunching to shorten the bunch length for RF acceleration.
- Acceleration by high gradient RF cavity for the whole bunch, compensates the velocity modulation and the beam becomes rigid.

Energy Compression (1)

- According to a simulation, the energy spread is 2%, which is larger than DR acceptance, 1%.
- Energy compressor by de/acceleration at the dispersive area is added before the DR.
- After the energy compression, the energy spread is 0.5%, which is in tolerance.

Before RF Compression

After RF Compression

ILC and CLIC comparison

Accelerator Beam parameter	CLIC (ACC.)	ILC
Pulse length	156ns	0.86
Pulse reputation	50Hz	5Hz
# of micro bunches in a pulse	312	2625
Bunch separation	500ps	369ns
Bunch charge	0.9nC	3.2nC
Polarization	80%	80%
Bunch length at gun	100ps	1ns
Peak current	9A	3A

A similar system to ILC based on Polarized electron source with GaAs cathode is assumed.

Less bunch charge, but high repetition rate and high average current in a pulse are challenging.

Laser for Photo-cathode

Laser for Photo-Cathode

- Laser is one of the most important element of the photocathode gun.
- Beam properties depend on the laser.
 - Temporal structure : 3MHz repetition, 0.9 ms macro pulse.
 - Beam emittance : 10 μrad.
 - Polarization :circular polarization and tunable around 700nm.
- A laser system, which meets fully LC requirements, is not available commercially.
- Several candidates for ILC.
 - Ti:Al₂O₃ : baseline
 - **Yb fiber laser : possible alternative**

$Ti:Al_2O_3$

- Spontaneous mode-locking by Carr effect, bunch length > 17fs 。
- Wide band width for lasing (700-1100nm), wave length tuneability by filtering.
- Require 488nm light for pumping; SH of Nd:YAG/YLF is employed limiting the efficiency from the pumping power to the laser light.
- Luminescence time is 3.2 ms, which is not suitable to form a long macro pulse.

ILC Baseline Design

- Ti:Al₂O₃ mode lock + 3MHz pulse picker by Pockels cell makes a pulse train.
- Macro-pulse amplification by Ti:Al₂O₃ regenerative amplifier pumped by SH of Nd:YAG.
- Wave length is tunable. The stability is a challenging issue.

E-source fo

Laser

Pulse stretcher

UN.

DES

Pulse from Model-lock laser (200fs) is stretched to 1ns for ILC Bunch.

 $L = f_{M1}$

2200 g/mm, BW 5 nm bW, OPD 516ps OPTICAL SYSTEM LAYOUT

202

M1

Retro reflector

Yb: YAG fiber laser + OPA

- Yb:YAG mode lock and Pockels cell Pulse Picker generate 3MHz pulse train.
- Yb: fiber laser amplifies the pulse train.
- NOPA (Non-collinear Optical Parametric Amplification) realize the wavelength tune-ability around 700nm.
- It could be LD pumped-full solid stable laser.

Yb fiber laser (1)

- Double clad-core optical fiber.
- Light from InGaAs LD (940nm) is introduce to 1st clad for pumping. Direct pumping by LD is very efficient and stable.
- Signal propagates in the inner core, where Yb ion is doped, and is amplified by stimulated emission.
- Due to the long structure, power density can be low and large limit on the high power operation.

J. Limpert

Yb fiber laser (2)

- The gain per length is low, but the propagation is quite efficient and the total gain can be quite high.
- High efficiency, low-loss, high-power, very stable.
- 2kW CW amplification is achieved.

E-source for LC

Fig. 4: Power evolution of cw double-clad fiber lasers with diffraction-limited beam quality over the last decade

Related Physics

Electron Gu

Nonlinear Optics (1)

Non-linear polarization is induced by intense laser field in material. In usual linear regime, the electric polarization is

 $\mathbf{P} = \varepsilon_0 \boldsymbol{\chi} \mathbf{E} \qquad (5-1)$

Non-linear polarization (up to second order) is $\mathbf{P} = \varepsilon_0 \chi^{(1)} \mathbf{E} + \varepsilon_0 \Big[\chi^{(2)} (2 \omega = \omega + \omega) + \chi^{(2)} (0 = \omega - \omega) \Big] \mathbf{E}^2 \quad (5-2)$ Sum frequency 0 frequency

Second harmonics (2ω) and 0 frequency mode are induced. That corresponds to that second order of the fundamental mode is expressed with 2ω mode and 0 mode.

$$\mathbf{P^{(2)}} \propto \cos^2(\omega t) = \frac{1}{2} \cos(2\omega t) + \frac{1}{2} \qquad (5-3)$$

SH 0

Nonlinear Optics (2)

The phase velocity of polarization and SH is

$$v_1 = \frac{2k_1}{2\omega} = \frac{n_1}{c}$$
 (5-4) $v_2 = \frac{k_2}{2\omega} = \frac{n_2}{c}$ (5-5)

 k_1 and k_2 are wave number, n_1 and n_2 are refractive index. The phase velocity should be same for efficient SH generation, because the growth is expressed as

$$E_{2}|^{2} = \frac{\omega_{2}^{2}}{4\epsilon_{0}^{2}n_{2}^{2}c^{2}}|P^{(2\omega)}|^{2}\frac{\sin^{2}\left(\frac{\Delta kz}{2}\right)}{\left(\frac{\Delta kz}{2}\right)^{2}}z^{2} \quad (5-6)$$

which is maximized by $\Delta k \equiv 2k_1 \cdot k_2 = 0$, when the phase velocity is same for both modes. Usually, material shows normal dispersion, that $n_1 > n_2$ for $\omega_1 > \omega_2$ and the condition is never satisfied. It is satisfied only with birefringence material.

SHG: Second Harmonic Generation

By focusing laser light in birefringence material, second harmonics is generated.

 $\omega_1 + \omega_1 = \omega_2 \qquad (5-7)$

The phase matching condition should be satisfied for efficient conversion.

 $n_1 \omega_1 + n_1 \omega_1 = n_2 \omega_2$ (5-8)

<u>OPA (1)</u>

As reverse process of harmonic generation, high energy photon can be split into two low energy photons,

 $\omega_1 = \omega_2 + \omega_3 \qquad (5-9)$

When intense ω_1 laser is given, ω_2 and ω_3 light are amplified. This is Optical Parametric Amplifier (OPA). The phase matching condition is

$$n_1 \omega_1 = n_2 \omega_2 + n_3 \omega_3$$
 (5-10)

 $ω_1$ is called as "driver" in OPA. When $ω_2$ is what we want (signal), $ω_3$ is called as "idler".

OPA(2)

► Yb fiber amplifier generate powerful pulse train in 1030nm.

► The fundamental mode is converted to SH, 515nm by SHG.

► The 515nm is driver in OPA. It can be converted to 800nm signal and 1500nm Idler.

► The phase matching condition can be modified by changing the angle between crystal axis and light direction. Wavelength is tune-abile.

Summary

- Fundamentals of electro-emission and electron gun are explained.
- Polarized electron is generated by photo-emission from NEA GaAs cathode with circularly polarized laser.
- ILC and CLIC electron sources are DC bias gun with NEA GaAs.
- The beam property depends on the emission process, external condition (surface field, etc.), beam transport and acceleration.

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