

A gauge mediation model for explaining the muon $g-2$ with gauge coupling unification

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in collaboration with

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Phys.Lett. B725 (2013) 339-343

arXiv:1311.1906

Introduction

2013 Physics Prize



Photo: G-M Greuel via Wikimedia Commons

Peter Higgs



Photo: Pnicolet via Wikimedia Commons

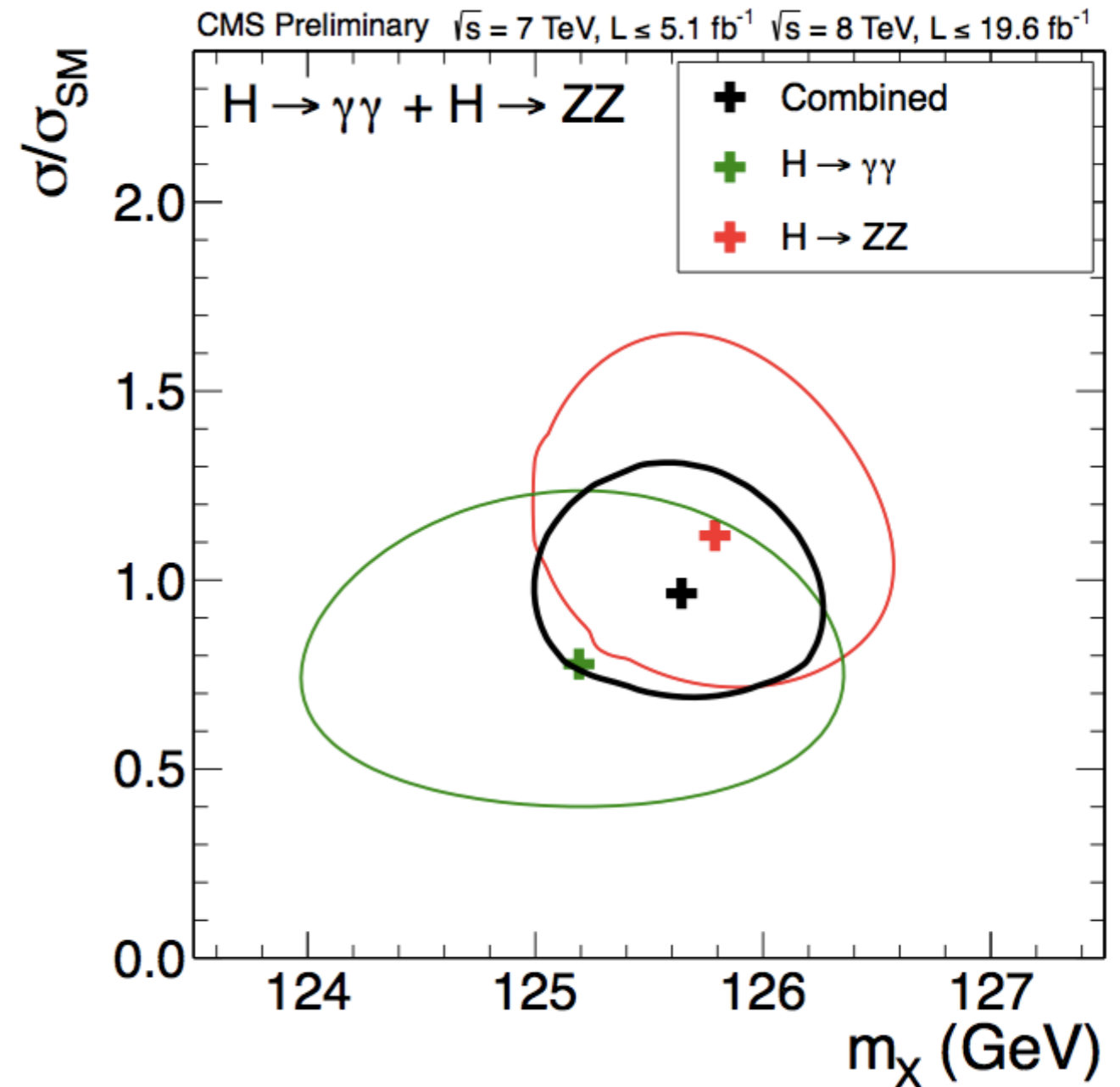
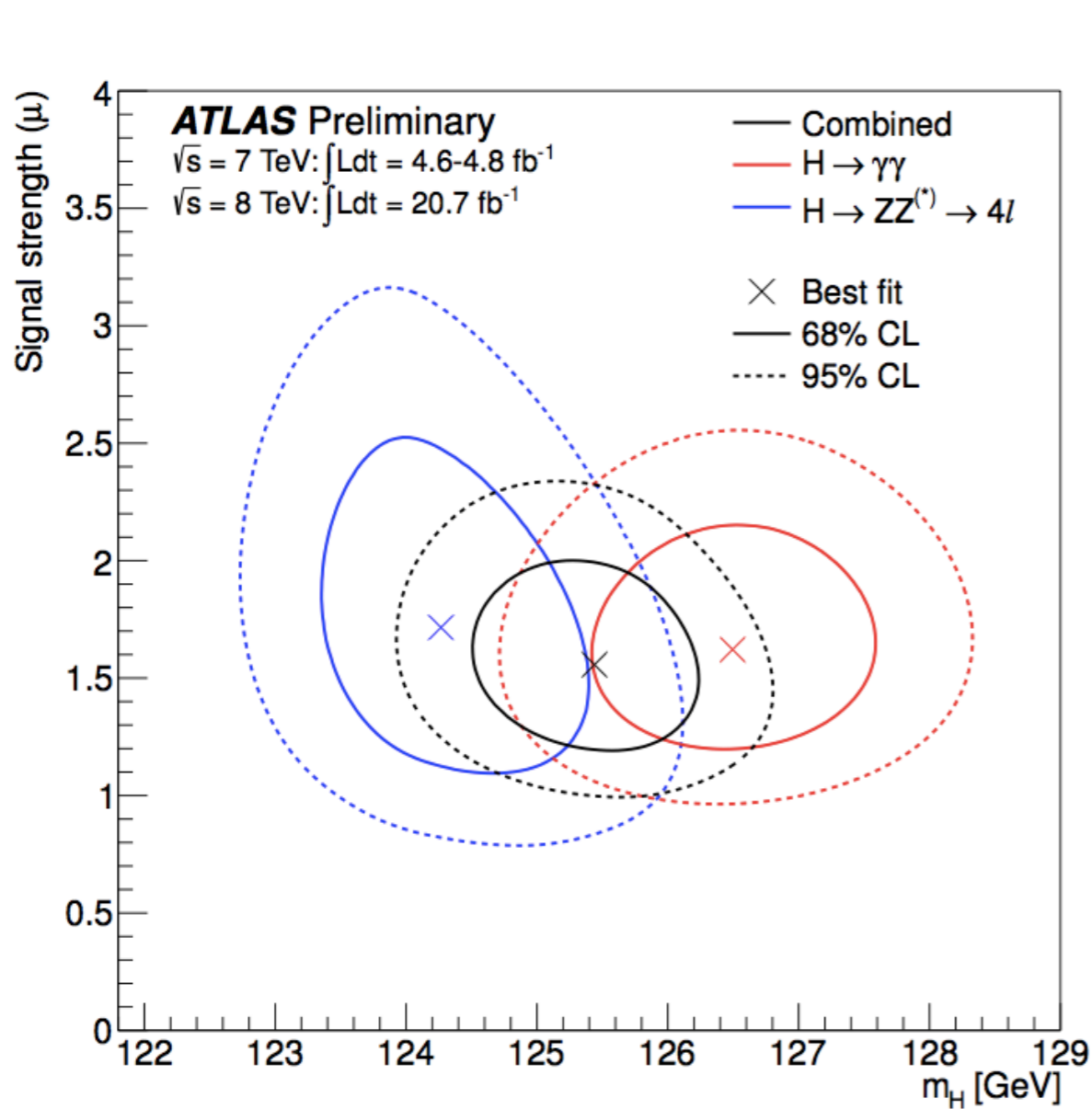
François Englert

Higgs boson has
been discovered!

[gallery](#)

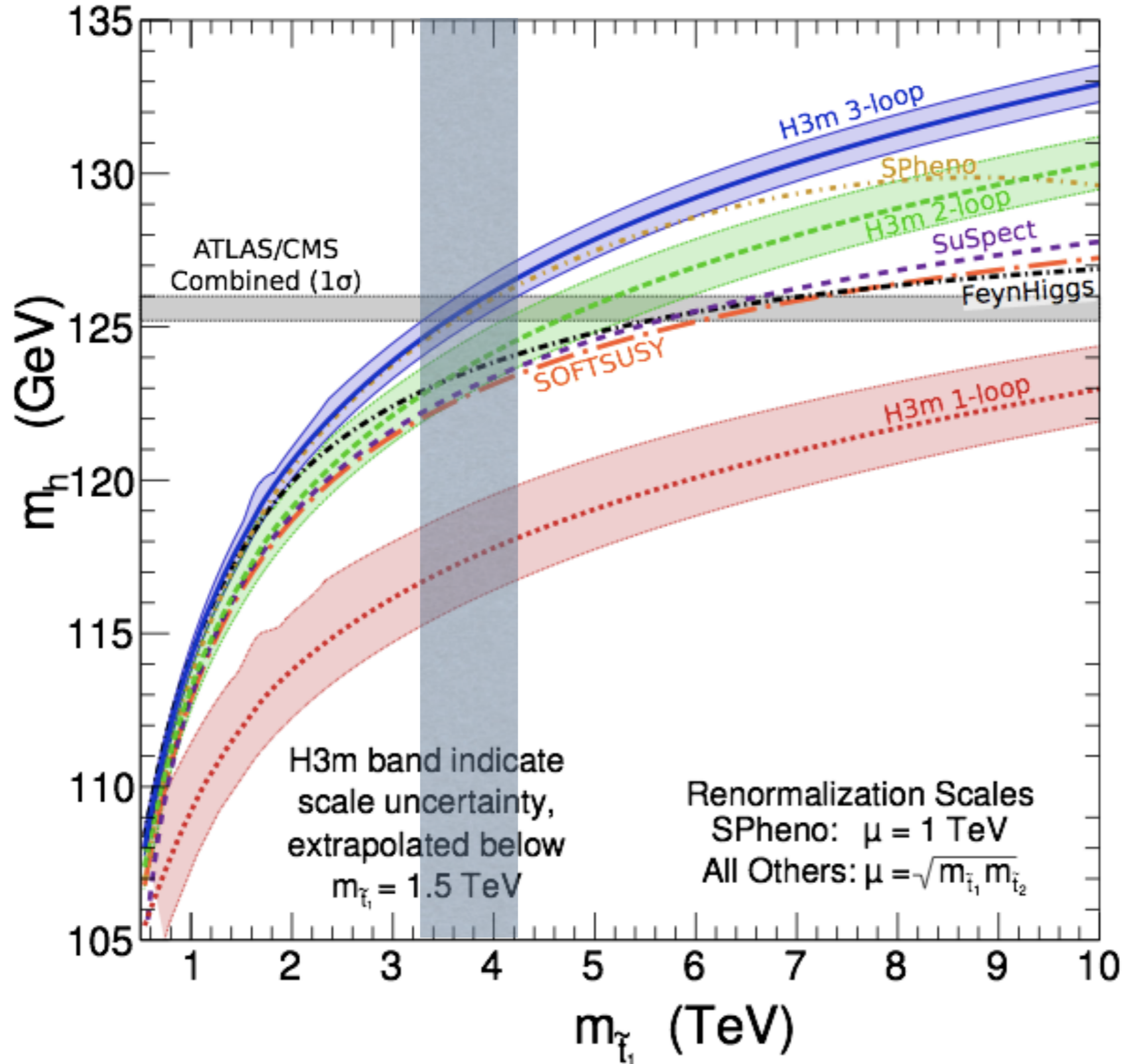
- ▶ [Listen to François Englert \(7:42\)](#)
- ▶ [Interview transcript](#)

Higgs boson mass



Higgs boson mass $\sim 125 \text{ GeV}$

heavy stop(3-4TeV) can be consistent with the observation



@3-loop
 $O(\alpha_t \alpha_s^2)$

J.L. Feng, P. Kant, S. Profumo and D. Sanford,
 1306.2318

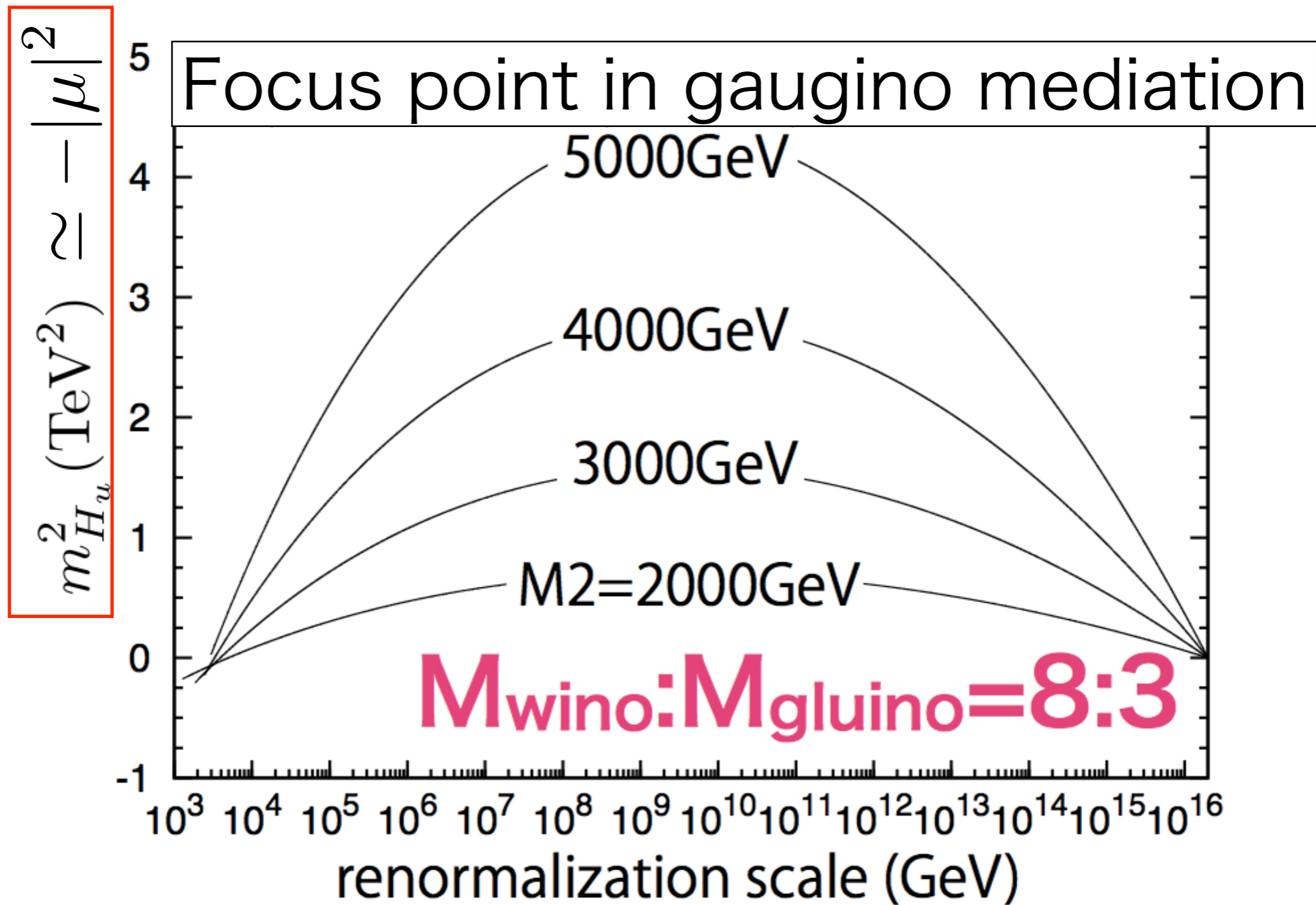
The H3m error corresponds to change of the renormalization scale from $M_s/2$ to $2M_s$

Light SUSY particles@ILC?

Light SUSY particles@ILC?

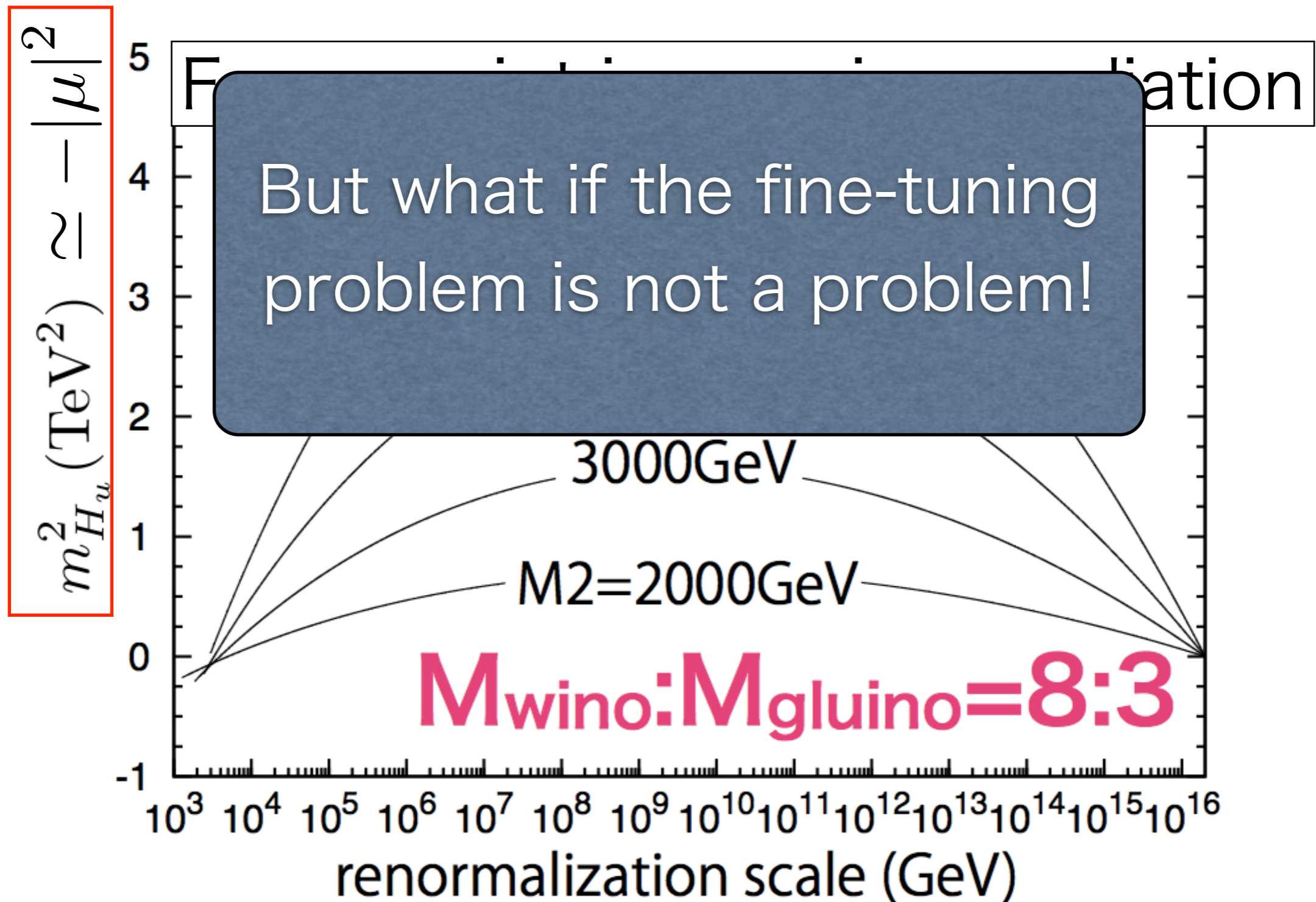
Light Higgsino ← Fine-tuning problem

If the model has a focus point behavior



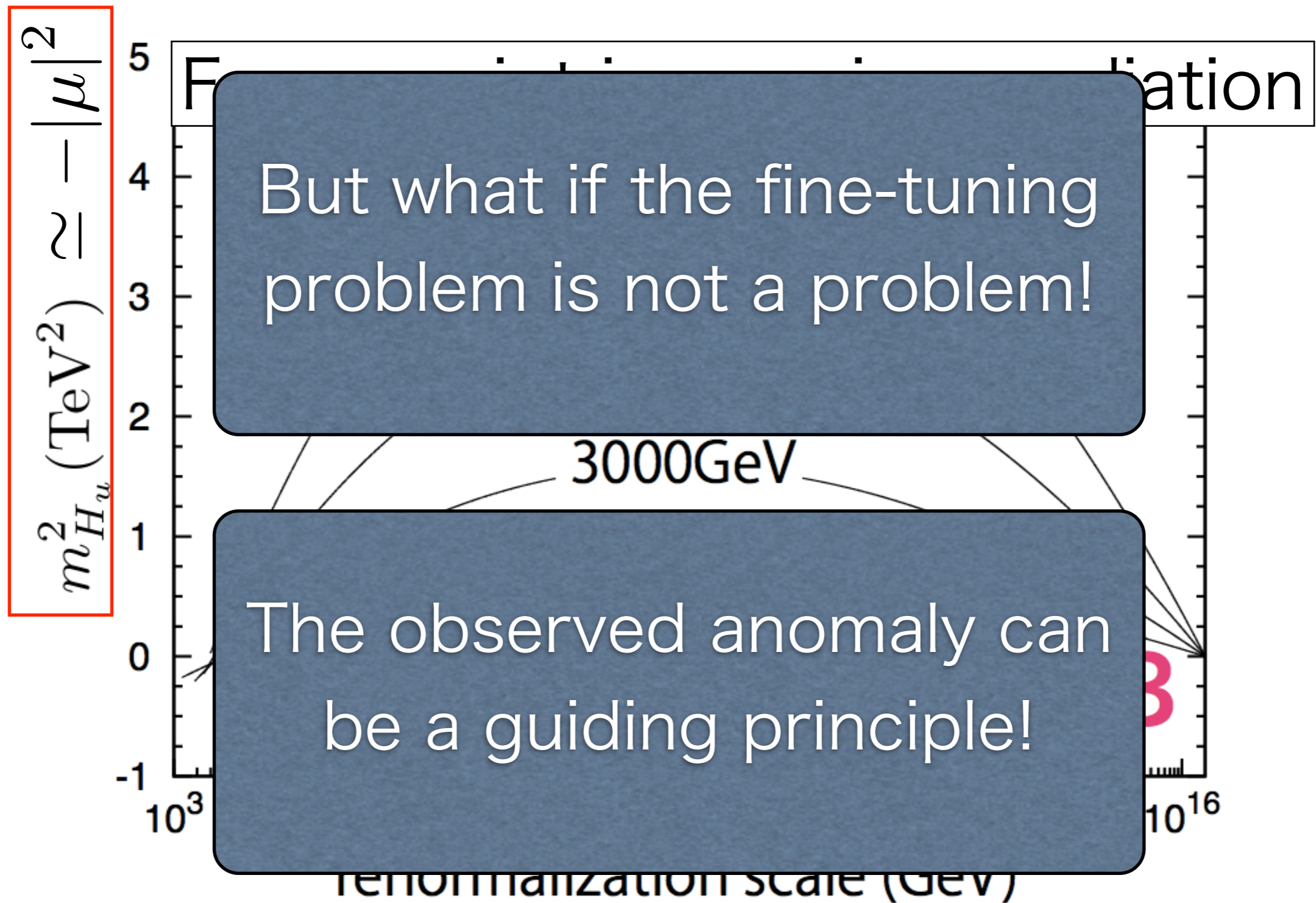
The Higgsino of $O(100)\text{GeV}$ can be observed

If the model has a focus point behavior



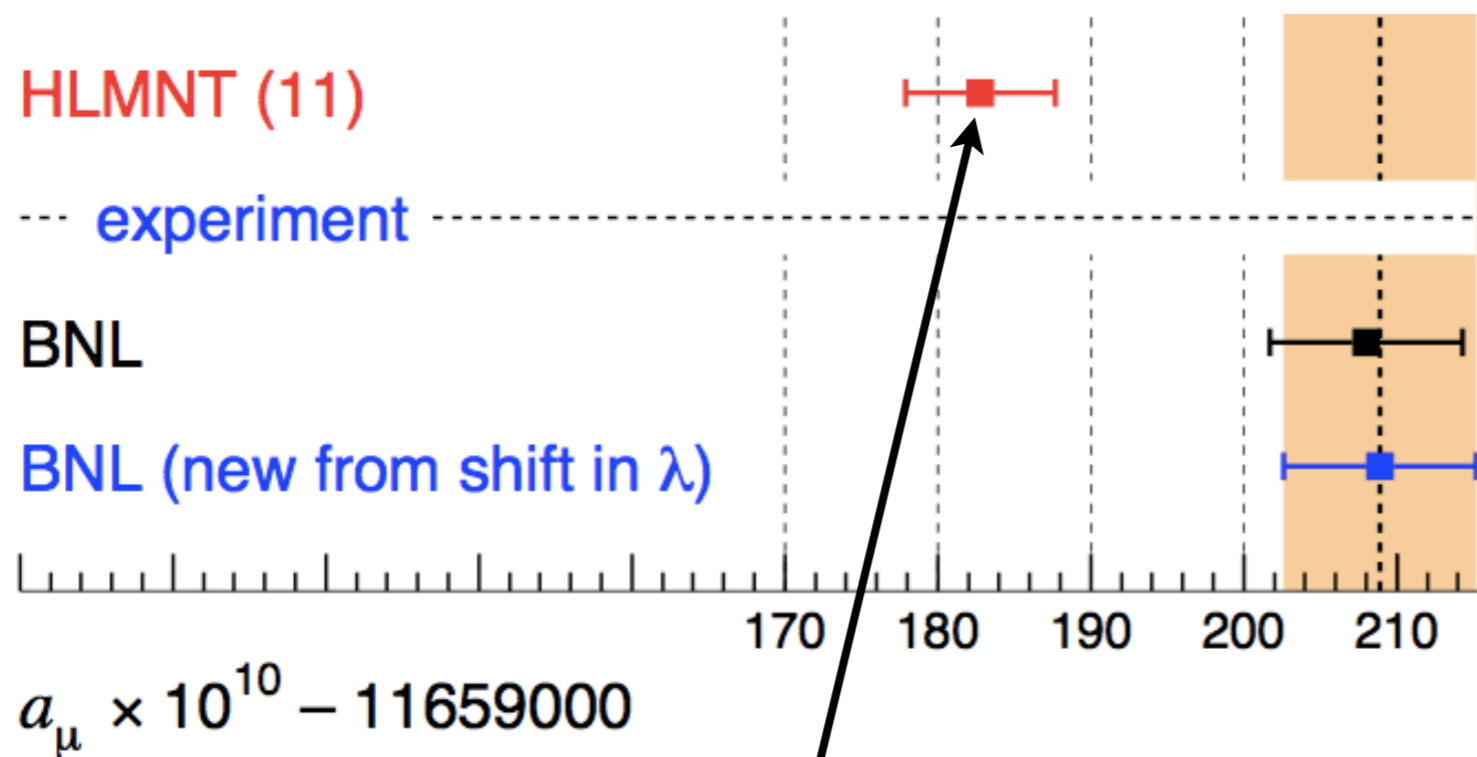
The Higgsino of $O(100)\text{GeV}$ can be observed

If the model has a focus point behavior



The Higgsino of $O(100)\text{GeV}$ can be observed

Muon $g-2$ anomaly



**>3 σ deviation
from SM
prediction!**

[Hagiwara, Liao, Martin, Nomura, Teubner, J.Phys. G38 (2011) 085003]

SM value

$$a_\mu = \frac{g - 2}{2}$$

$$V(\vec{x}) = -\vec{\mu} \cdot \vec{B}(\vec{x}) \quad \vec{\mu} = g \left(\frac{e}{2m_\mu} \right) \vec{S}$$

**What is a expected NP mass scale to
explain the muon $g-2$?**

What is a expected NP mass scale to explain the muon g-2 ?

$$\Delta(a_\mu)_{\text{NP}} \sim \frac{g^2}{16\pi^2} \frac{m_\mu^2}{m_{\text{NP}}^2}$$
$$= 20.7 \times 10^{-10} \left(\frac{g}{0.65}\right)^2 \left(\frac{120\text{GeV}}{m_{\text{NP}}}\right)^2$$

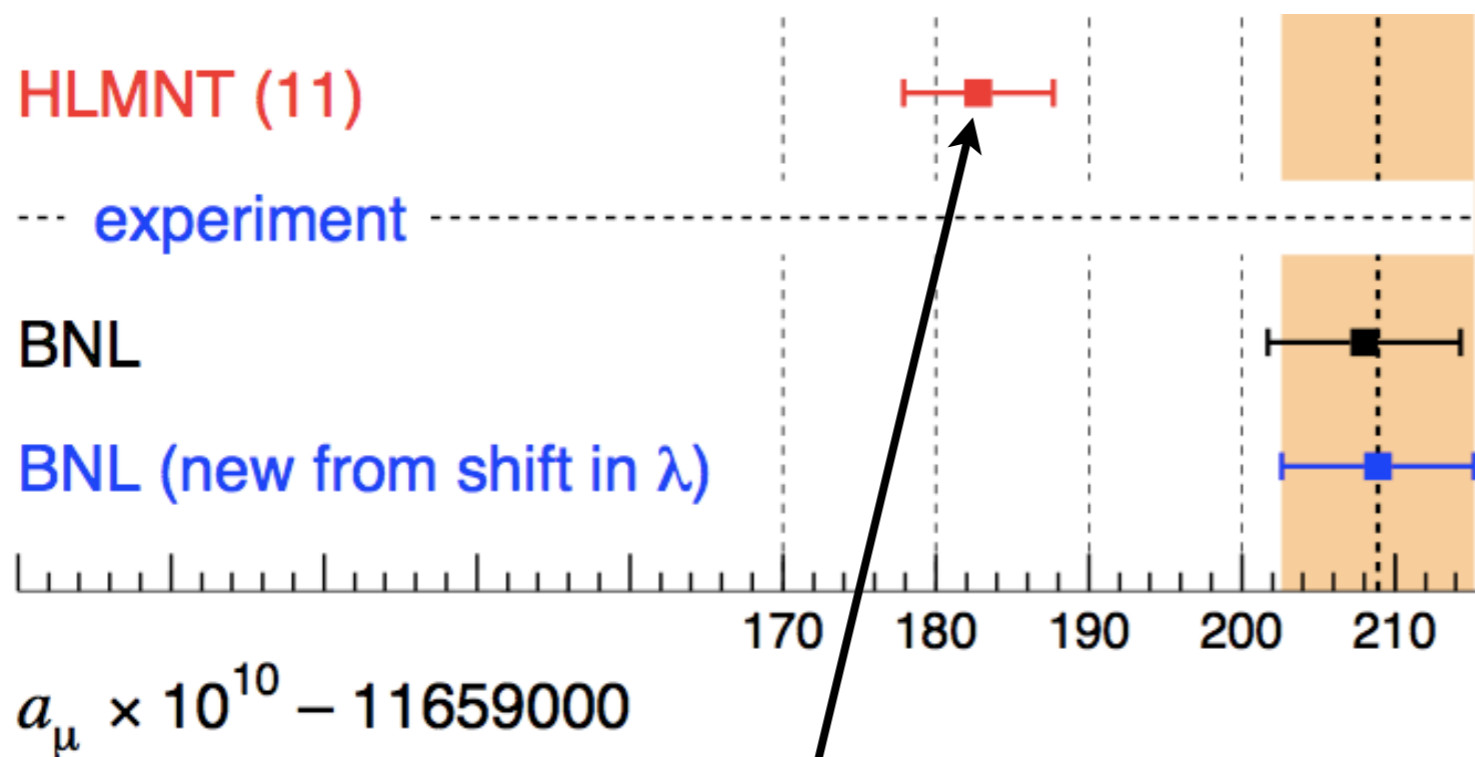
New coupling

Mass scale of new physics

We need $O(100)\text{GeV}$ new particles.

Accessible at the ILC!

Muon $g-2$ anomaly

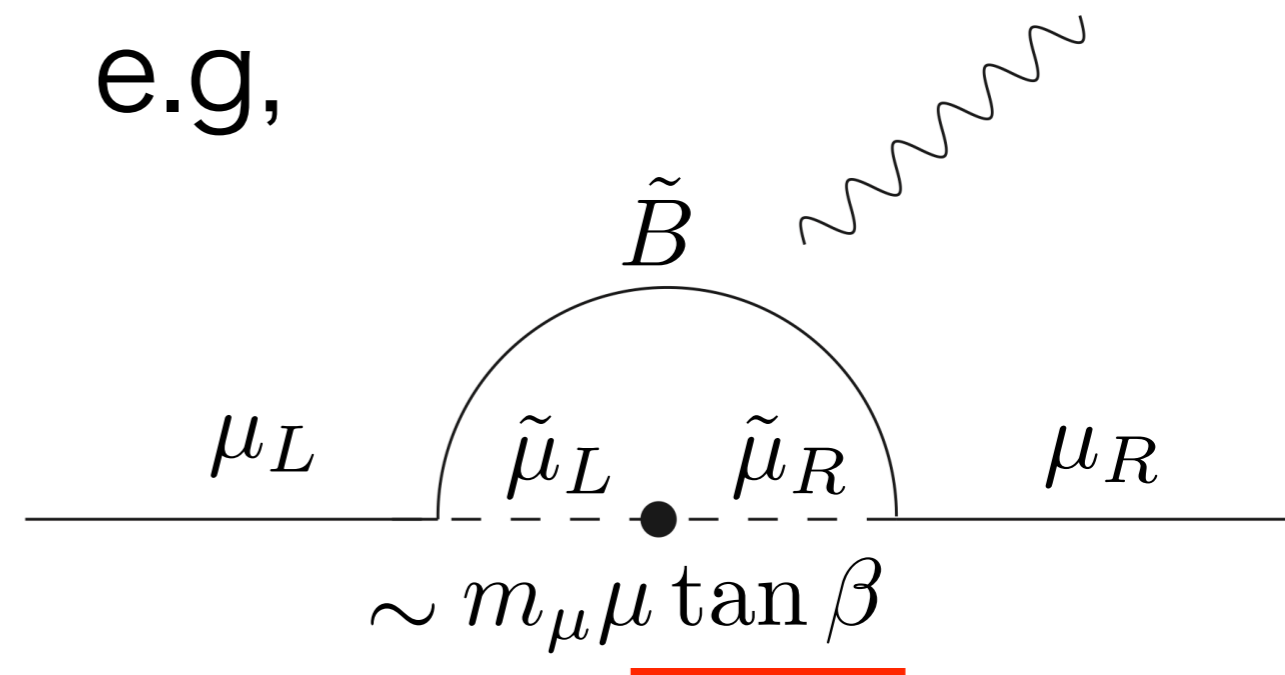


**>3 σ deviation
from SM
prediction!**

[Hagiwara, Liao, Martin, Nomura, Teubner, J.Phys. G38 (2011) 085003]

SM value

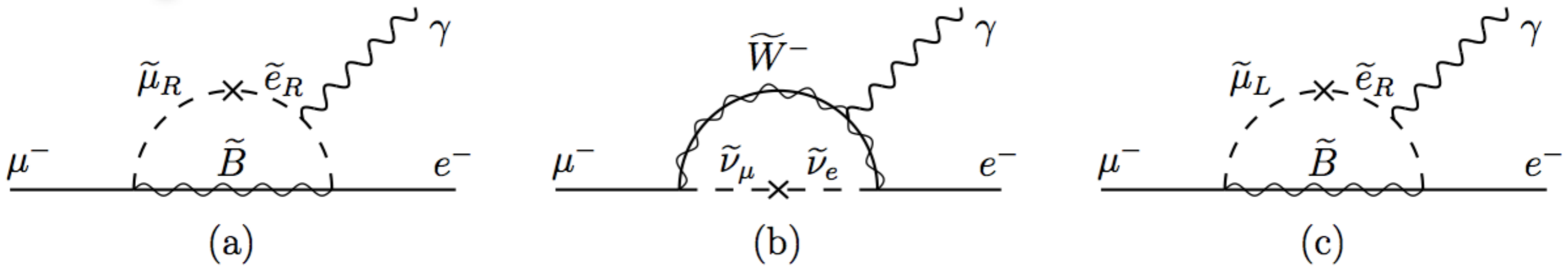
e.g,



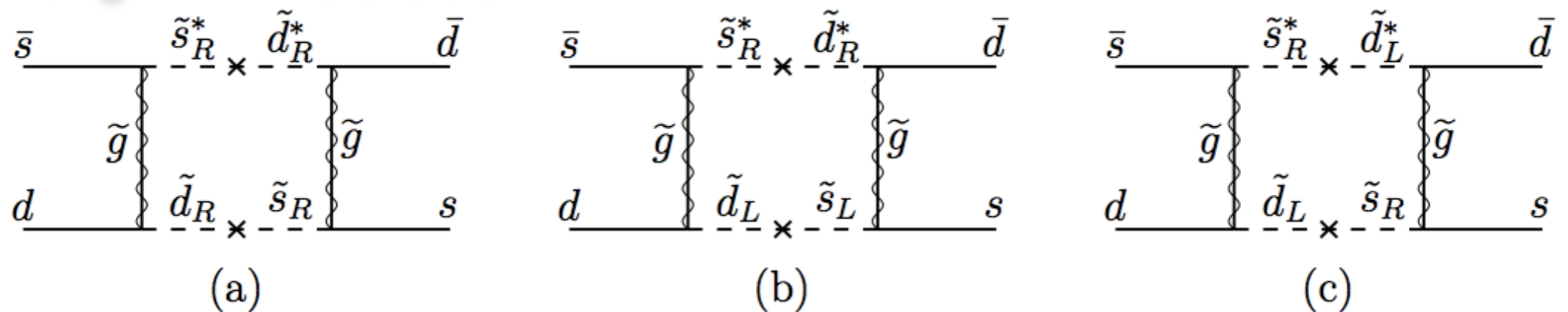
**light non-colored
SUSY particle can
explain this deviation**

Severe constraints from flavor violating processes

slepton sector



squark sector



We can not take them to be light without thinking

Sever constraints from flavor violating processes

slepton sector $\sim 10\text{TeV}$ for $O(1)$ mixing

(a) Slepton sector diagram: A muon line (μ^-) emits a $\tilde{\mu}_R$ and a \tilde{B} , which then interact to produce an electron (e^+) and a photon (γ).

(b) Squark sector diagram: A muon line (μ^-) emits a \tilde{d}_R and a \tilde{s} , which then interact via a gluon (\tilde{g}) to produce a muon (μ^-) and a photon (γ).

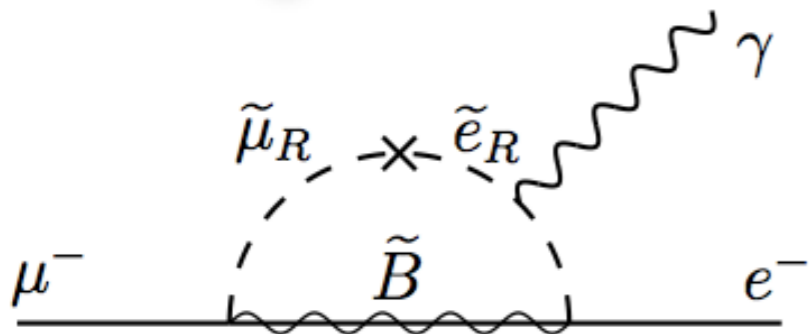
(c) Squark sector diagram: A muon line (μ^-) emits a \tilde{d}_R and a \tilde{s} , which then interact to produce a muon (μ^-) and a photon (γ).

In conclusion the MEG experiment has so established the most stringent upper limit to date on the branching ratio of the $\mu^+ \rightarrow e^+ \gamma$ decay, $\mathcal{B} < 5.7 \times 10^{-13}$ at 90% C.L. using data collected between 2009 and 2011, which improves the previous best upper limit by a factor of four. Further data have also been acquired in 2012 with an additional three-month run scheduled for 2013; the final number of stopped muons is expected to be almost twice that of the sample analyzed so far. Currently an upgrade

We can not take them to be light without thinking

Sever constraints from flavor violating processes

slepton sector



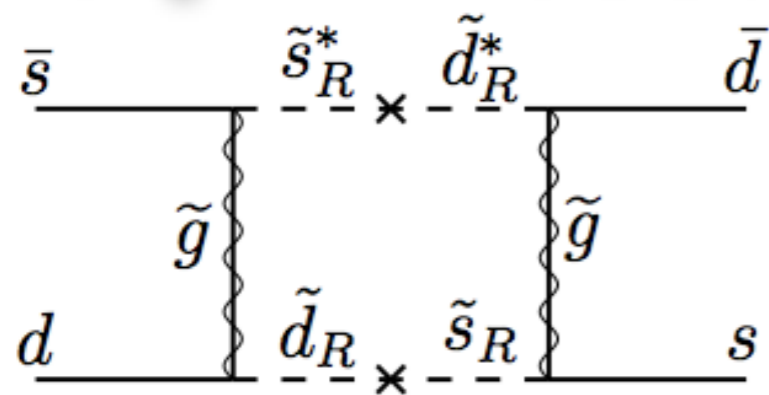
(a)

CPV in $K \bar{K}$ mixing

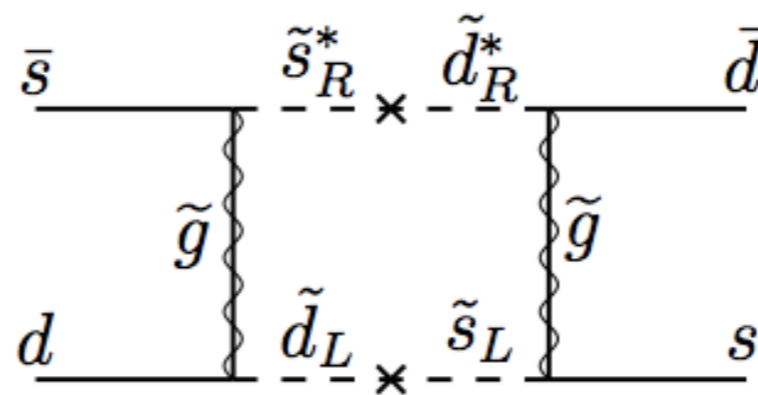
Marco Ciuchini¹, Livio Conti², Andrea Donini³, Enrico Franco⁴, Vicent Gimenez⁵, Leonardo Giusti⁶, Vittorio Lubicz¹, Guido Martinelli⁴, Antonio Masiero⁷, Ignazio Scimemi⁴, Luca Silvestrini⁸, Mauro Talevi⁹, Anastassios Vladikas², 1998

squark sector

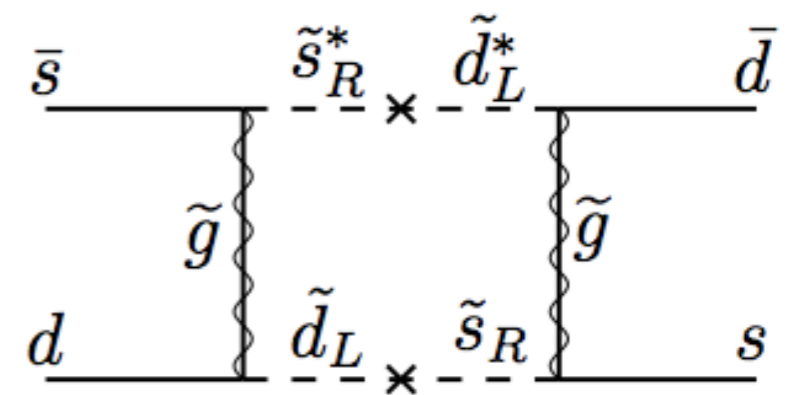
~1000TeV! for O(1) mixing



(a)



(b)

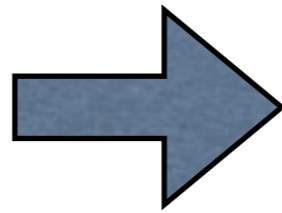


(c)

We can not take them to be light without thinking

What we should do?

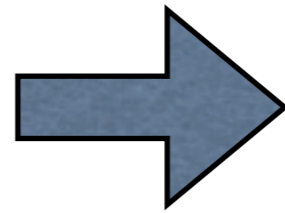
Higgs boson
mass~125GeV



heavy colored
SUSY particles

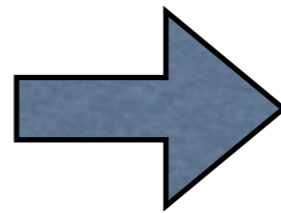
What we should do?

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heavy colored
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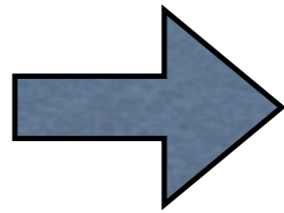
Anomaly of
the muon $g-2$



light non-colored
SUSY particles

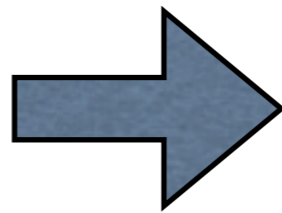
What we should do?

Higgs boson
mass $\sim 125\text{GeV}$



heavy colored
SUSY particles

Anomaly of
the muon $g-2$

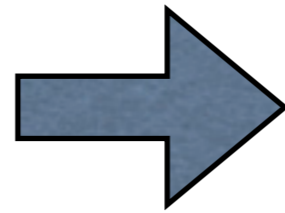


light non-colored
SUSY particles

FCNC should be avoided

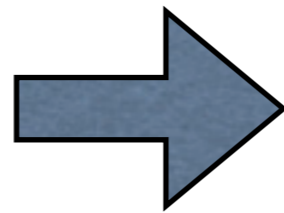
What we should do?

Higgs boson
mass $\sim 125\text{GeV}$



heavy colored
SUSY particles

Anomaly of
the muon $g-2$



light non-colored
SUSY particles

FCNC should be avoided

**answer: colored and non-colored mass
splitting in gauge mediation**

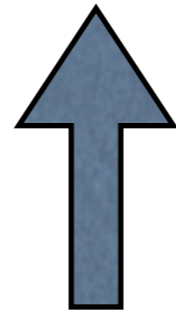
Our model

Adjoint messenger gauge mediation model

Adjoint Messenger Model

$$W = (M_8 + \lambda_8 F \theta^2) \text{Tr} \Sigma_8^2 + (M_3 + \lambda_3 F \theta^2) \text{Tr} \Sigma_3^2 + (M_5 + \lambda_5 F \theta^2) \Psi_5 \Psi_{\bar{5}}$$

No hyper charge



To give masses to right handed sleptons/bino

Octet/triplet messenger does not give masses to the bino and right handed sleptons.

Gauge Coupling Unification

With MSSM matter contents, gauge coupling unification is satisfied. (Of course!)

Even if we include color octet and SU(2) triplet, the gauge coupling unified non-trivially

Σ_8 :SU(3) octet

Σ_3 :SU(2) triplet

Renormalization group equations change

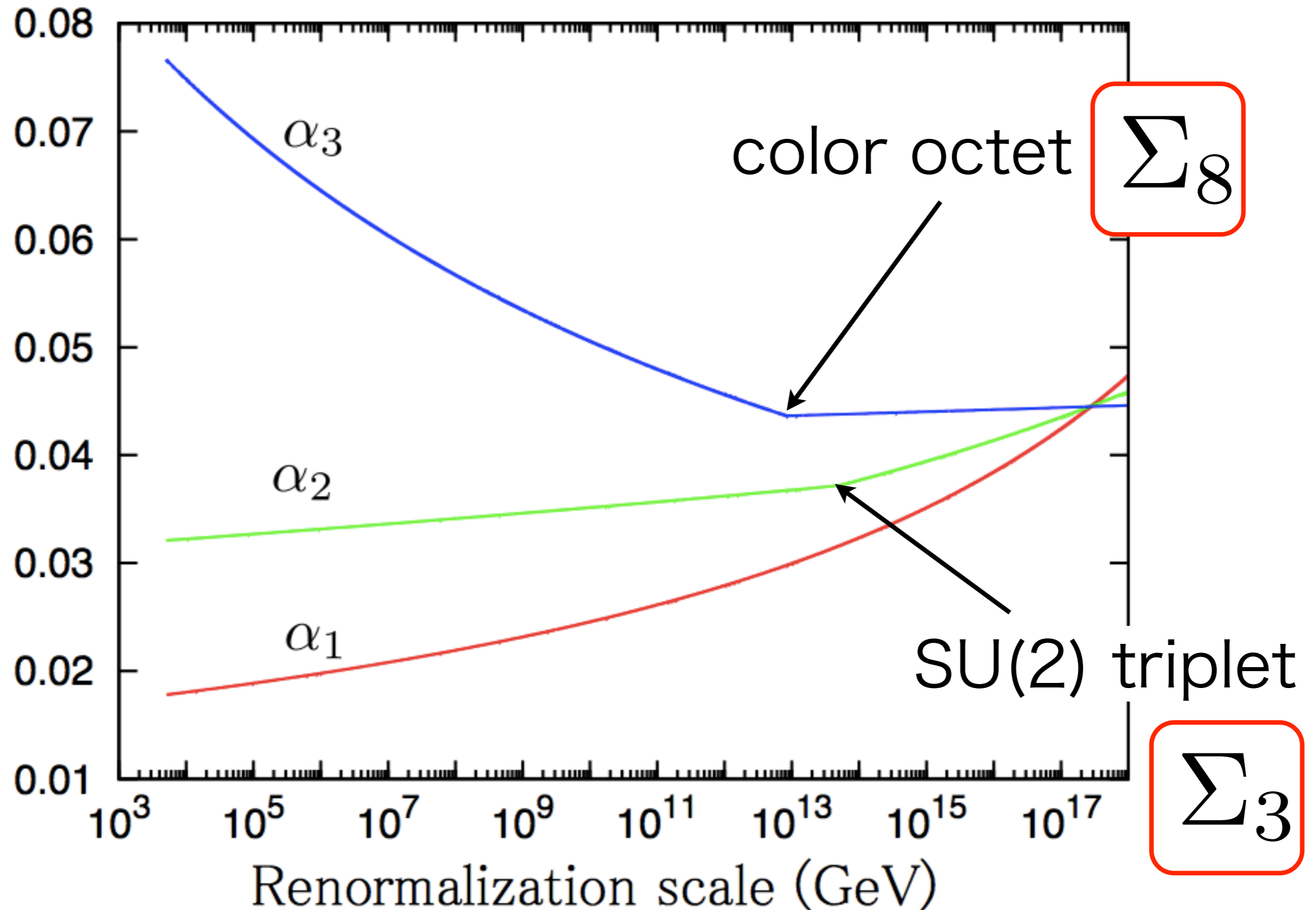
$$\begin{aligned}\alpha_1^{-1}(M_{\text{str}}) &= \alpha_1^{-1}(m_{\text{SUSY}}) - \frac{b_1}{2\pi} \ln \frac{M_{\text{str}}}{m_{\text{SUSY}}}, \\ \alpha_2^{-1}(M_{\text{str}}) &= \alpha_2^{-1}(m_{\text{SUSY}}) - \frac{b_2}{2\pi} \ln \frac{M_{\text{str}}}{m_{\text{SUSY}}} - \frac{2}{2\pi} \ln \frac{M_{\text{str}}}{M_3}, \\ \alpha_3^{-1}(M_{\text{str}}) &= \alpha_3^{-1}(m_{\text{SUSY}}) - \frac{b_3}{2\pi} \ln \frac{M_{\text{str}}}{m_{\text{SUSY}}} - \frac{3}{2\pi} \ln \frac{M_{\text{str}}}{M_8}.\end{aligned}$$

This part

@1-loop level

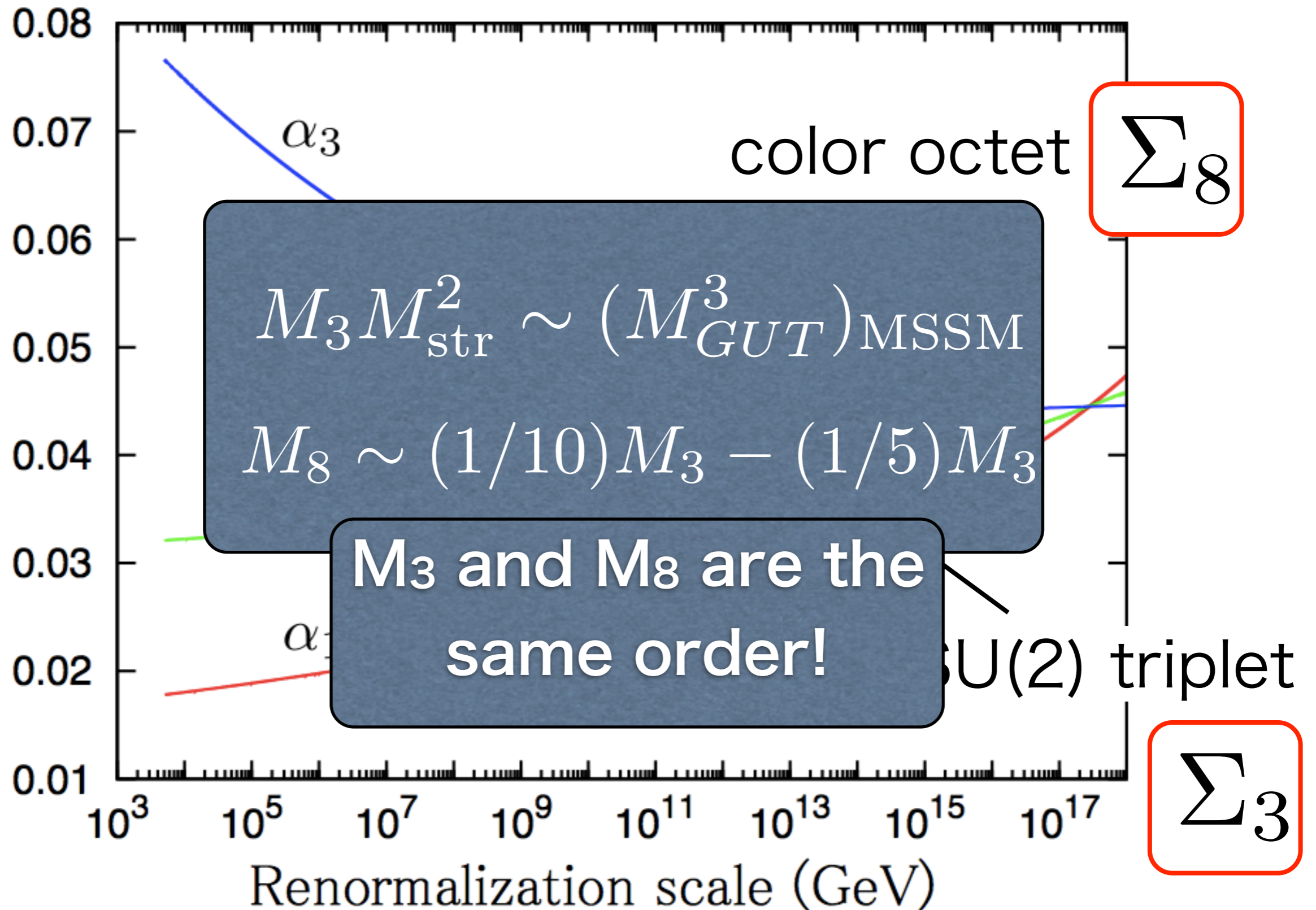
Two-loop RGE analysis

$$M_3 = 5 \times 10^{13} \text{ GeV}$$



Two-loop RGE analysis

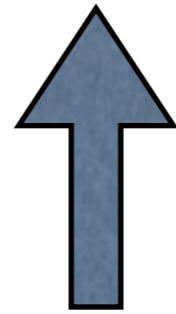
$$M_3 = 5 \times 10^{13} \text{ GeV}$$



Adjoint Messenger Model

$$W = (M_8 + \lambda_8 F \theta^2) \text{Tr} \Sigma_8^2 \\ + (M_3 + \lambda_3 F \theta^2) \text{Tr} \Sigma_3^2 \\ + (M_5 + \lambda_5 F \theta^2) \Psi_5 \Psi_{\bar{5}}$$

No hyper charge



To give masses to right handed sleptons/bino

Octet/triplet messenger does not give masses
to the bino and right handed sleptons.

Gaugino mass

$$M_1 \simeq \frac{\alpha_1}{4\pi} \Lambda_5, \quad M_2 \simeq \frac{\alpha_2}{4\pi} (2\Lambda_3 + \Lambda_5), \quad M_3 \simeq \frac{\alpha_3}{4\pi} (3\Lambda_8 + \Lambda_5),$$

Sfermion mass

$$m_{\tilde{Q}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 + \Lambda_5^2) + \frac{3}{4} \alpha_2^2 (2\Lambda_3^2 + \Lambda_5^2) + \frac{1}{60} \alpha_1^2 \Lambda_5^2 \right]$$

$$m_{\tilde{U}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 + \Lambda_5^2) + \frac{4}{15} \alpha_1^2 \Lambda_5^2 \right],$$

$$m_{\tilde{D}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 + \Lambda_5^2) + \frac{1}{15} \alpha_1^2 \Lambda_5^2 \right],$$

$$m_{\tilde{L}}^2 = m_{H_u}^2 = m_{H_d}^2 \simeq \frac{1}{8\pi^2} \left[\frac{3}{4} \alpha_2^2 (2\Lambda_3^2 + \Lambda_5^2) + \frac{3}{20} \alpha_1^2 \Lambda_5^2 \right],$$

$$m_{\tilde{E}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{3}{5} \alpha_1^2 \Lambda_5^2 \right],$$

Definition

$$\Lambda_8 \equiv \frac{\lambda_8 F}{M_8}, \quad \Lambda_3 \equiv \frac{\lambda_3 F}{M_3}, \quad \Lambda_5 \equiv \frac{\lambda_5 F}{M_5}.$$

Gaugino mass

$$M_1 \simeq \quad M_2 \simeq \frac{\alpha_2}{4\pi} (2\Lambda_3 \quad), \quad M_3 \simeq \frac{\alpha_3}{4\pi} (3\Lambda_8 \quad),$$

Sfermion mass

$$m_{\tilde{Q}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 \quad + \frac{3}{4} \alpha_2^2 (2\Lambda_3^2 \quad) \right]$$

$$m_{\tilde{U}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 \quad) \right],$$

$$m_{\tilde{D}}^2 \simeq \frac{1}{8\pi^2} \left[\frac{4}{3} \alpha_3^2 (3\Lambda_8^2 \quad) \right],$$

$$m_{\tilde{L}}^2 = m_{H_u}^2 = m_{H_d}^2 \simeq \frac{1}{8\pi^2} \left[\frac{3}{4} \alpha_2^2 (2\Lambda_3^2 \quad) \right],$$

$$m_{\tilde{E}}^2 \simeq$$

Definition

$$\Lambda_8 \equiv \frac{\lambda_8 F}{M_8}, \quad \Lambda_3 \equiv \frac{\lambda_3 F}{M_3}, \quad \Lambda_5 \equiv \frac{\lambda_5 F}{M_5}.$$

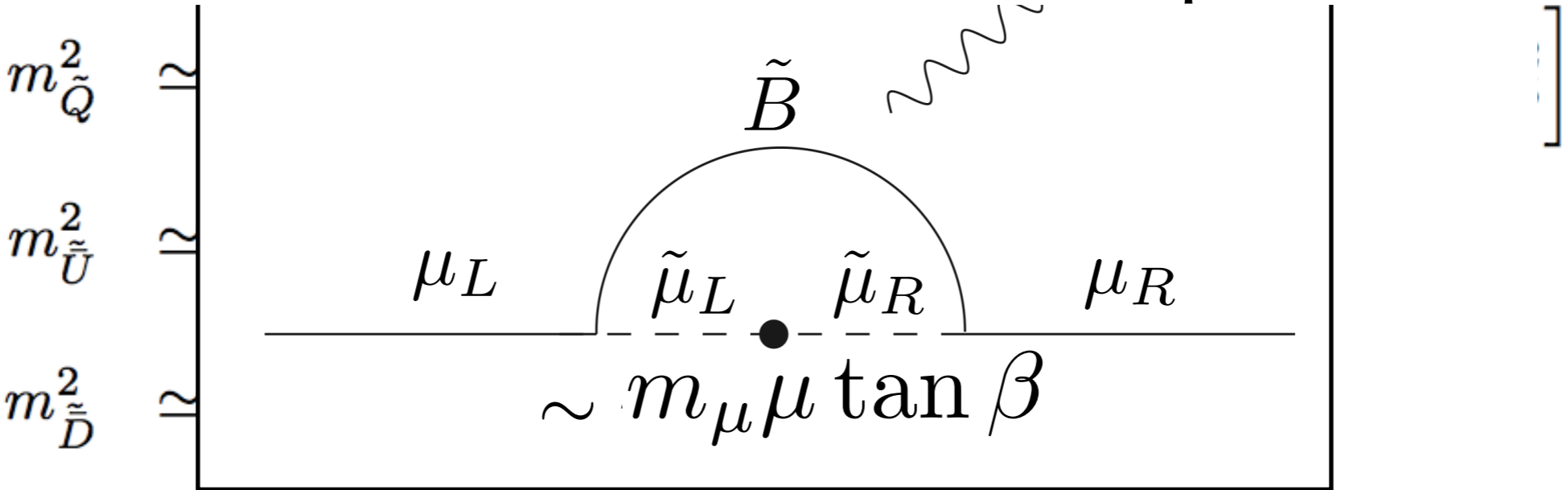
Gaugino mass

$$M_1 \simeq$$

$$M_2 \simeq \frac{\alpha_2}{4\pi} (2\Lambda_3), \quad M_3 \simeq \frac{\alpha_3}{4\pi} (3\Lambda_8),$$

Fermion mass

Bino contributions are important



$$m_{\tilde{L}}^2 = m_{H_u}^2 = m_{H_d}^2 \simeq \frac{1}{8\pi^2} \left[\frac{3}{4} \alpha_2^2 (2\Lambda_3^2) \right],$$

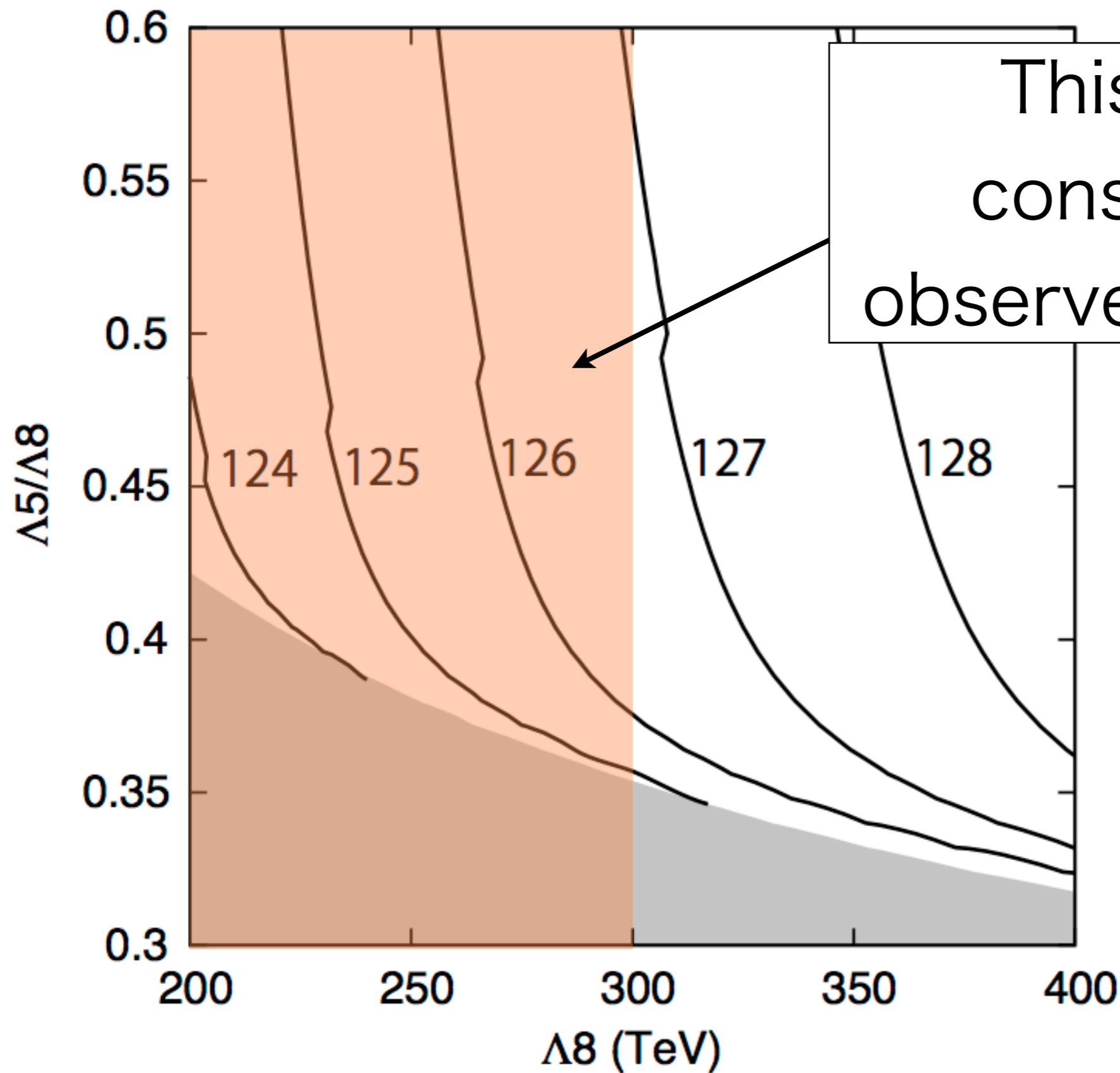
$$m_{\tilde{E}}^2 \simeq$$

Definition

$$\Lambda_8 \equiv \frac{\lambda_8 F}{M_8}, \quad \Lambda_3 \equiv \frac{\lambda_3 F}{M_3}, \quad \Lambda_5 \equiv \frac{\lambda_5 F}{M_5}.$$

Results

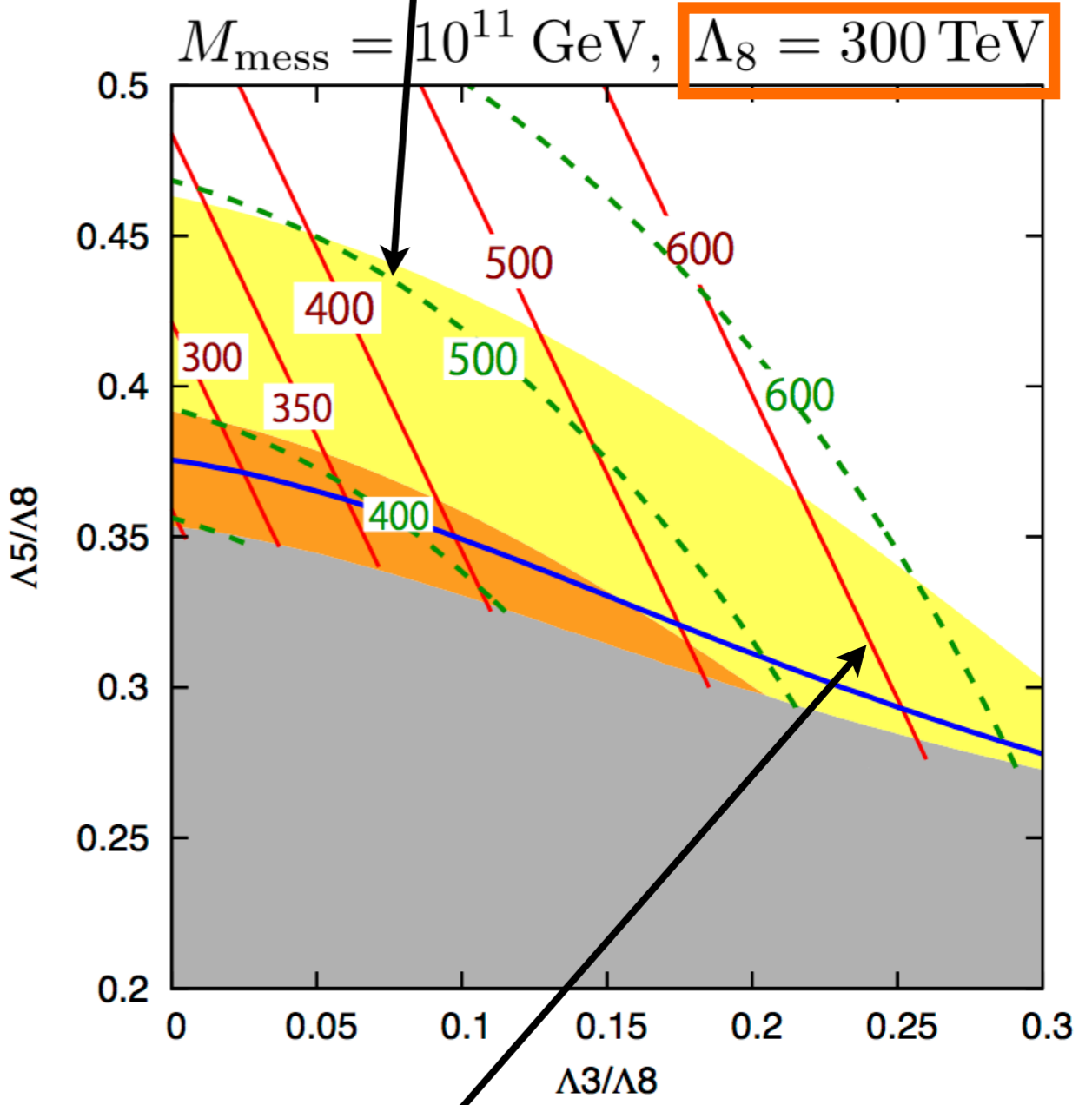
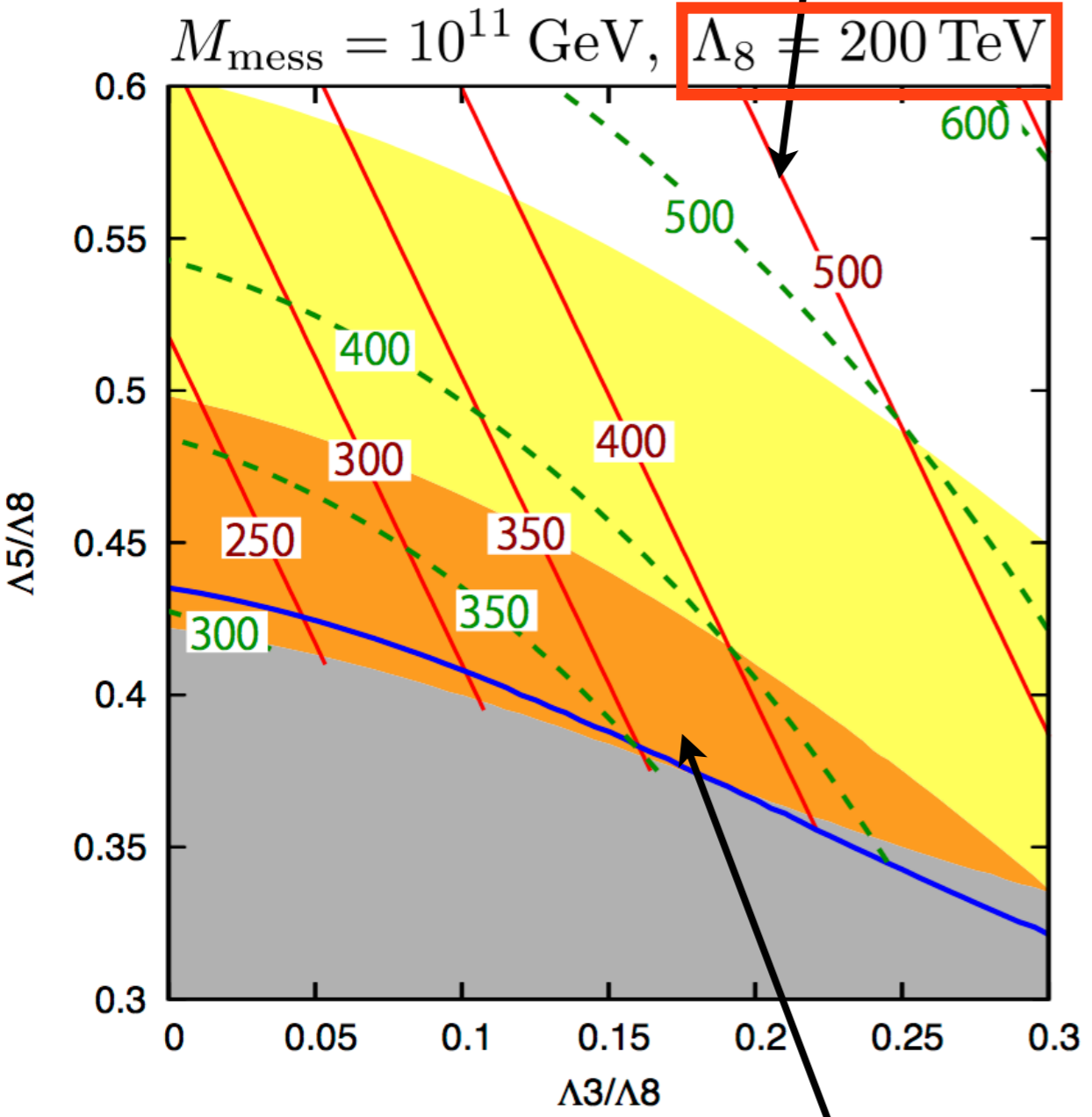
$$M_{\text{mess}} = 10^{11} \text{ GeV}, \tan \beta = 10, \Lambda_3 = 0$$



$$m_{\text{stop}} \sim 3.6 - 5.1 \text{ TeV}$$

chargino mass

Left-handed slepton mass



g-2 1σ

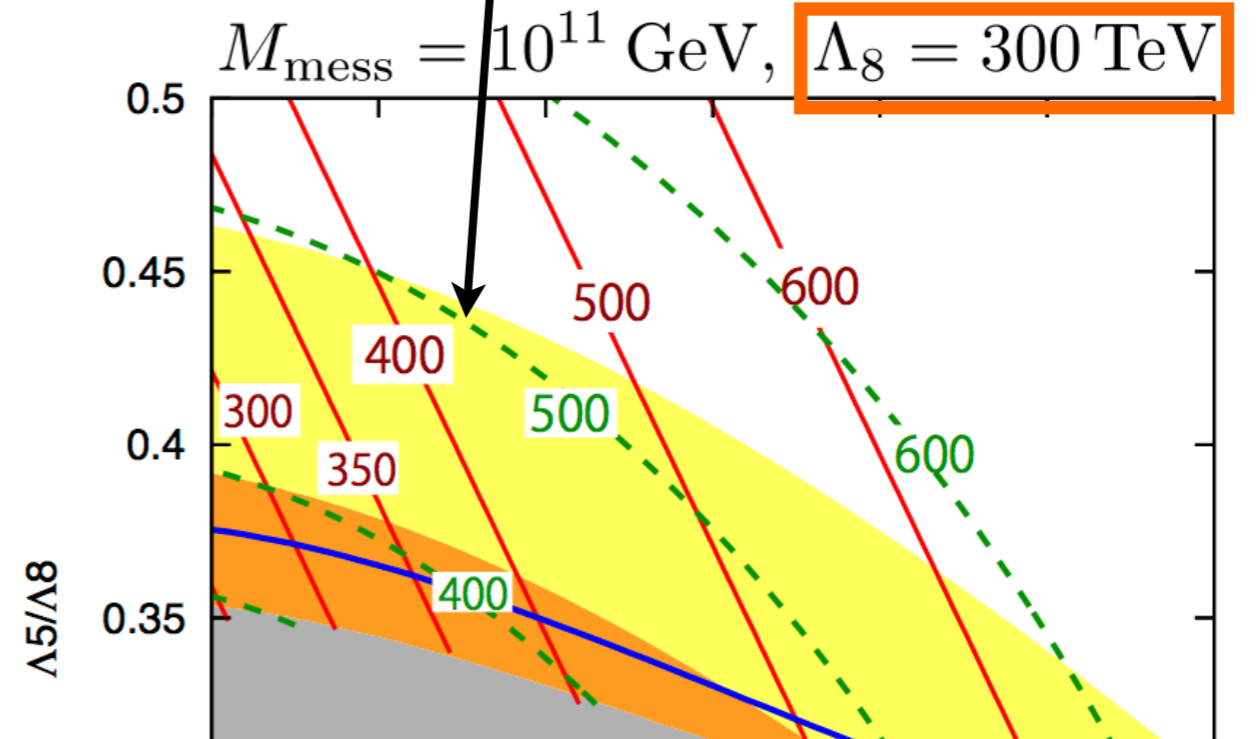
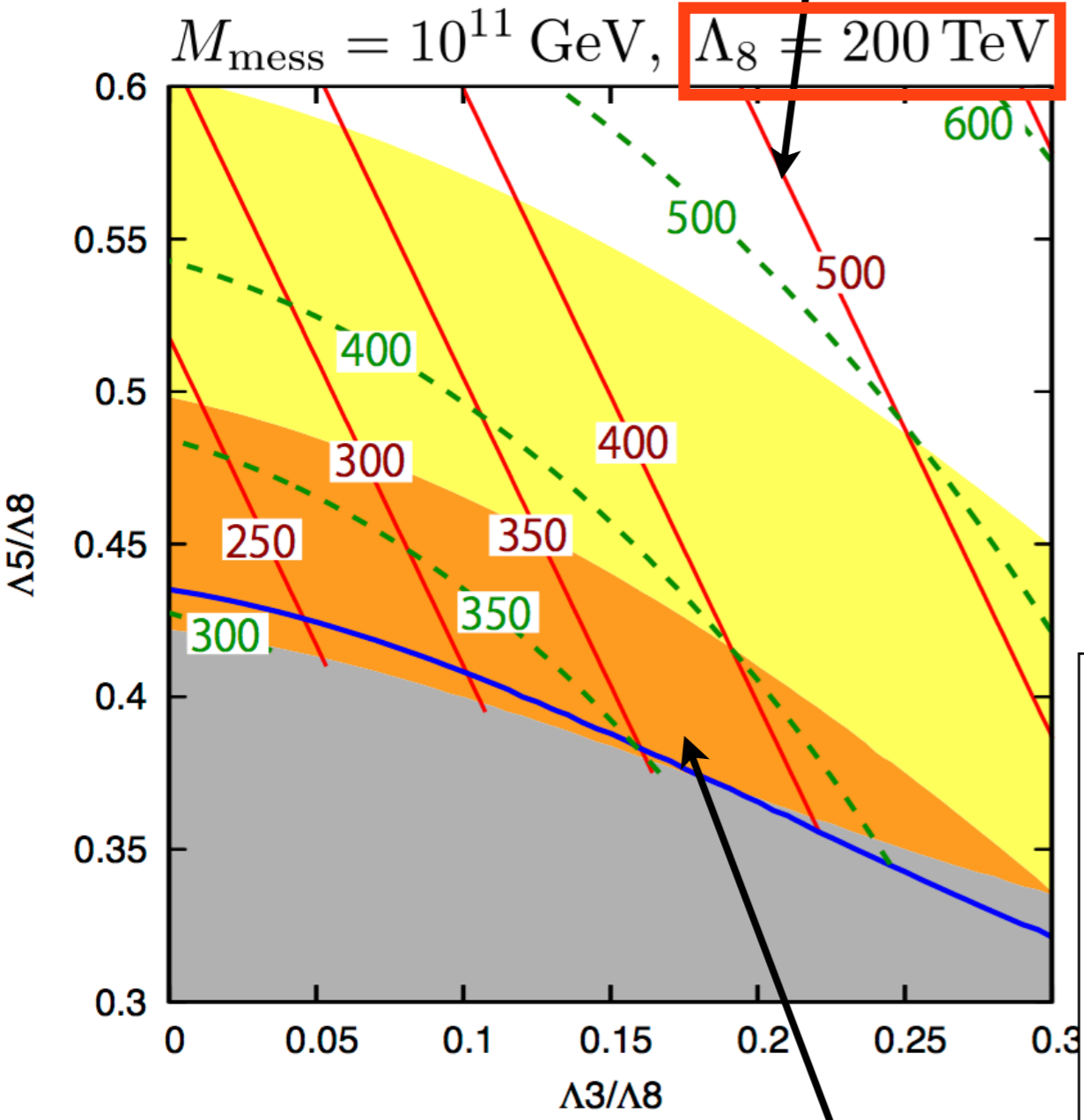
g-2 2σ

[g-2: FeynHiggs]

The stau NLSP needs to be heavier than 340GeV

chargino mass

Left-handed slepton mass



neutralino NLSP

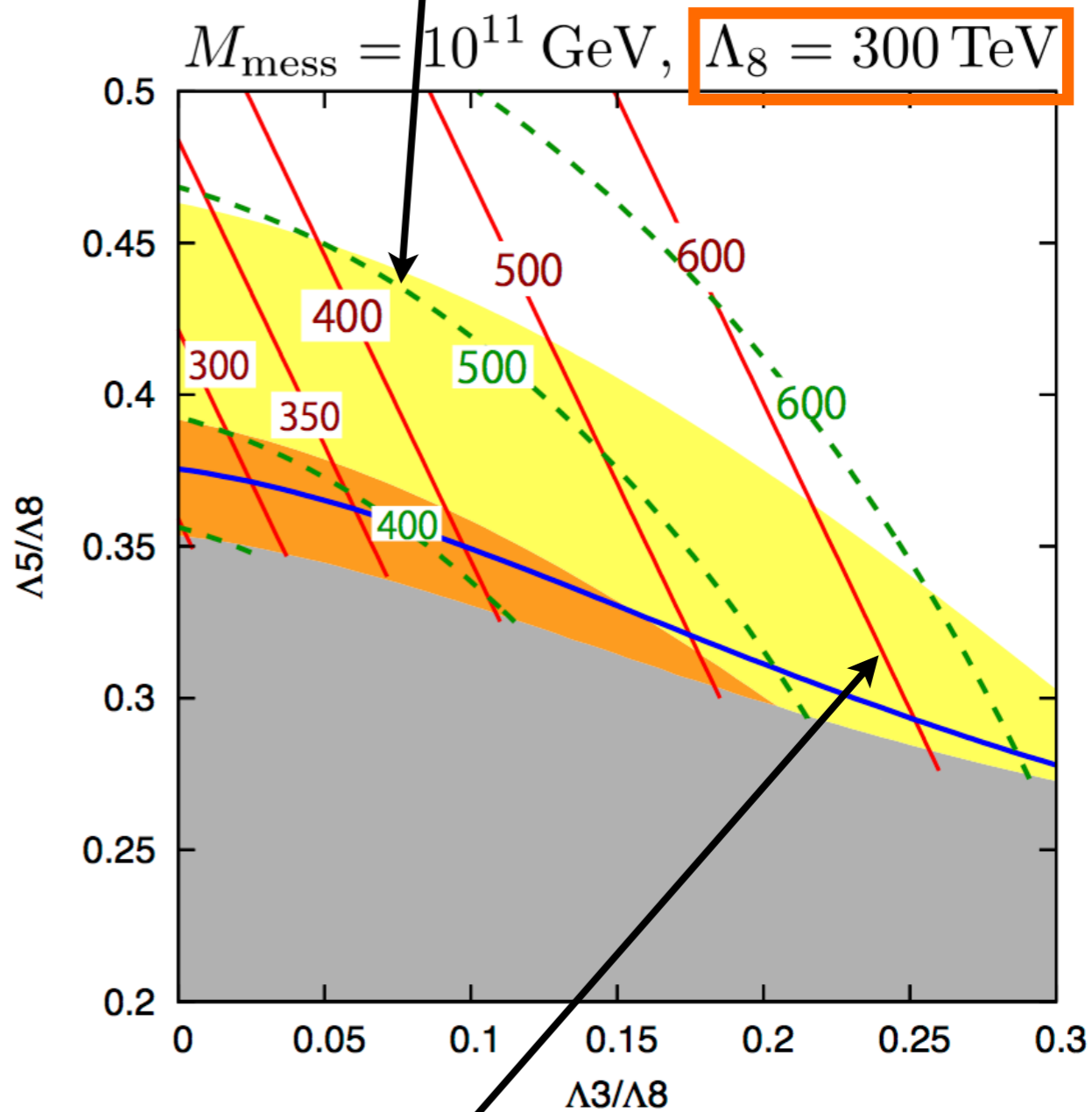
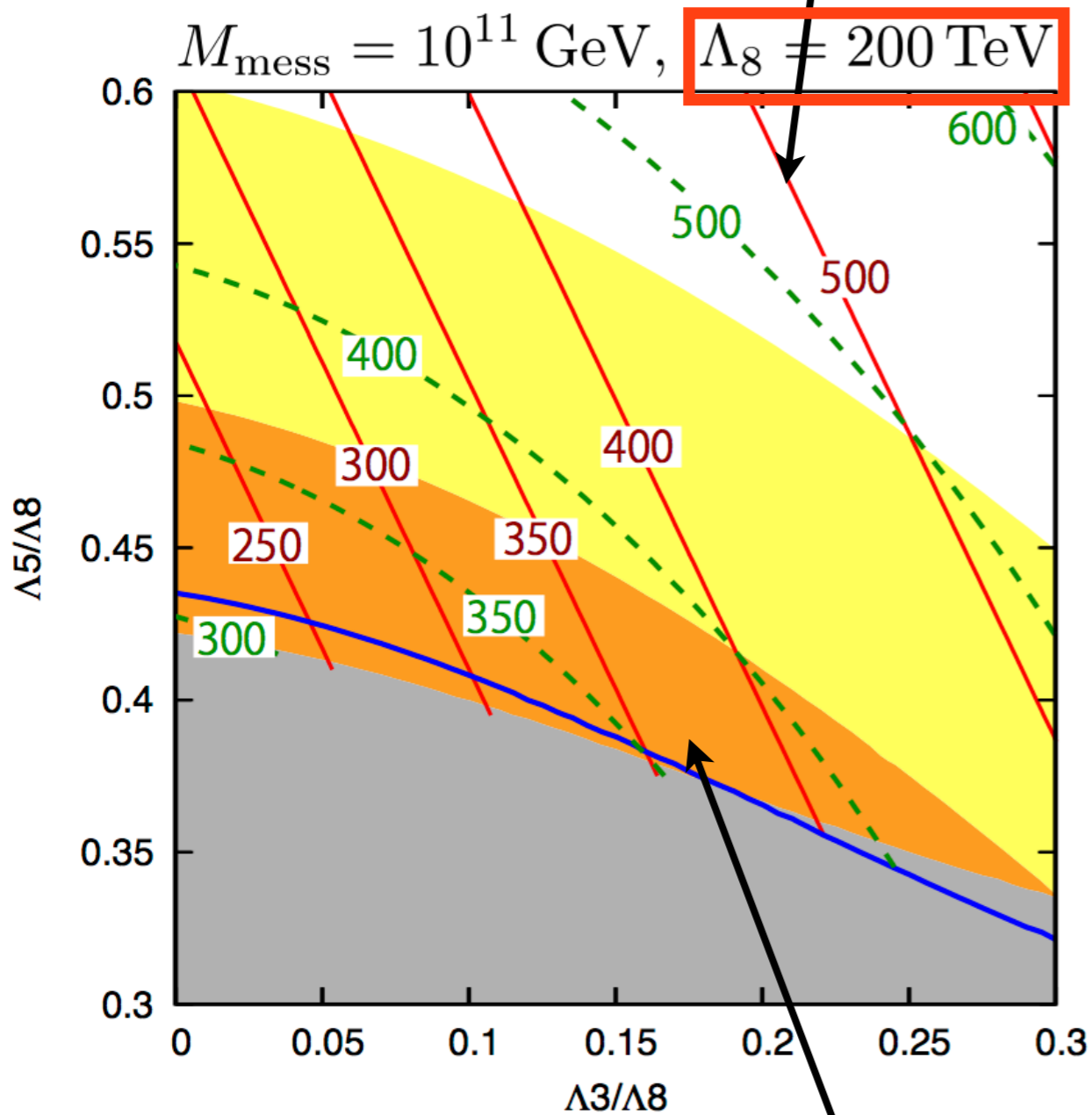
stau NLSP

$g-2 \ 2 \ \sigma$ [g-2: FeynHiggs]

The stau NLSP needs to be heavier than 340GeV

chargino mass

Left-handed slepton mass



g-2 1σ

g-2 2σ

[g-2: FeynHiggs]

The stau NLSP needs to be heavier than 340GeV

LHC results constrain on the electroweak production modes quite severely

sleptons

$$\tilde{l}_L^* \tilde{l}_L \rightarrow l^+ l^- 2\chi_1^0$$

$$\tilde{l}_R^* \tilde{l}_R \rightarrow l^+ l^- 2\chi_1^0$$

Wino

$$\chi_1^\pm \chi_2^0 \rightarrow (3l + \nu) 2\chi_1^0$$

(tau dominated)

Some of the parameter space is excluded already

$$M_{\text{wino}} \gtrsim 300 - 350 \text{ GeV}$$

(vacuum stability bound of the stau is OK!)

Λ_3/Λ_8	0.17	Λ_3/Λ_8	0.11
Λ_5/Λ_8	0.41	Λ_5/Λ_8	0.35
Λ_8	200 TeV		300 TeV
M_{mess}	10^{11} GeV		10^{11} GeV
$\tan \beta$	10		10
μ	2.4 TeV	μ	3.5 TeV
m_{stop}	3.6 TeV	m_{stop}	5.1 TeV
δa_μ	20.3×10^{-10}	δa_μ	18.6×10^{-10}
m_{gluino}	4.4 TeV	m_{gluino}	6.3 TeV
m_{squark}	4.1 TeV	m_{squark}	5.8 TeV
$m_{\tilde{e}_L} (m_{\tilde{\mu}_L})$	379 GeV	$m_{\tilde{e}_L} (m_{\tilde{\mu}_L})$	425 GeV
$m_{\tilde{e}_R} (m_{\tilde{\mu}_R})$	181 GeV	$m_{\tilde{e}_R} (m_{\tilde{\mu}_R})$	218 GeV
$m_{\tilde{\tau}_1}$	123 GeV	$m_{\tilde{\tau}_1}$	133 GeV
$m_{\chi_1^0}$	100 GeV	$m_{\chi_1^0}$	128 GeV
$m_{\chi_1^\pm}/m_{\chi_2^0}$	375 GeV	$m_{\chi_1^\pm}/m_{\chi_2^0}$	411 GeV

ILC!

Table 1: Some reference mass spectra and $(\delta a_\mu)_{\text{SUSY}}$.

Comments on Cosmology

Gravitino problem

Thermally produced NLSP decays into the gravitino and a SM particle

e.g.

bino \rightarrow photon or Z + gravitino

If the NLSP decays during/after the BBN era (~ 0.01 s - 100 s), the successful BBN prediction may be destroyed!

Life-time of NLSP

Bino:
$$\Gamma(\tilde{B} \rightarrow \psi_{3/2}\gamma) \simeq \frac{\cos^2 \theta_W M_{\tilde{B}}^5}{48\pi M_{\text{PL}}^2 m_{3/2}^2}$$

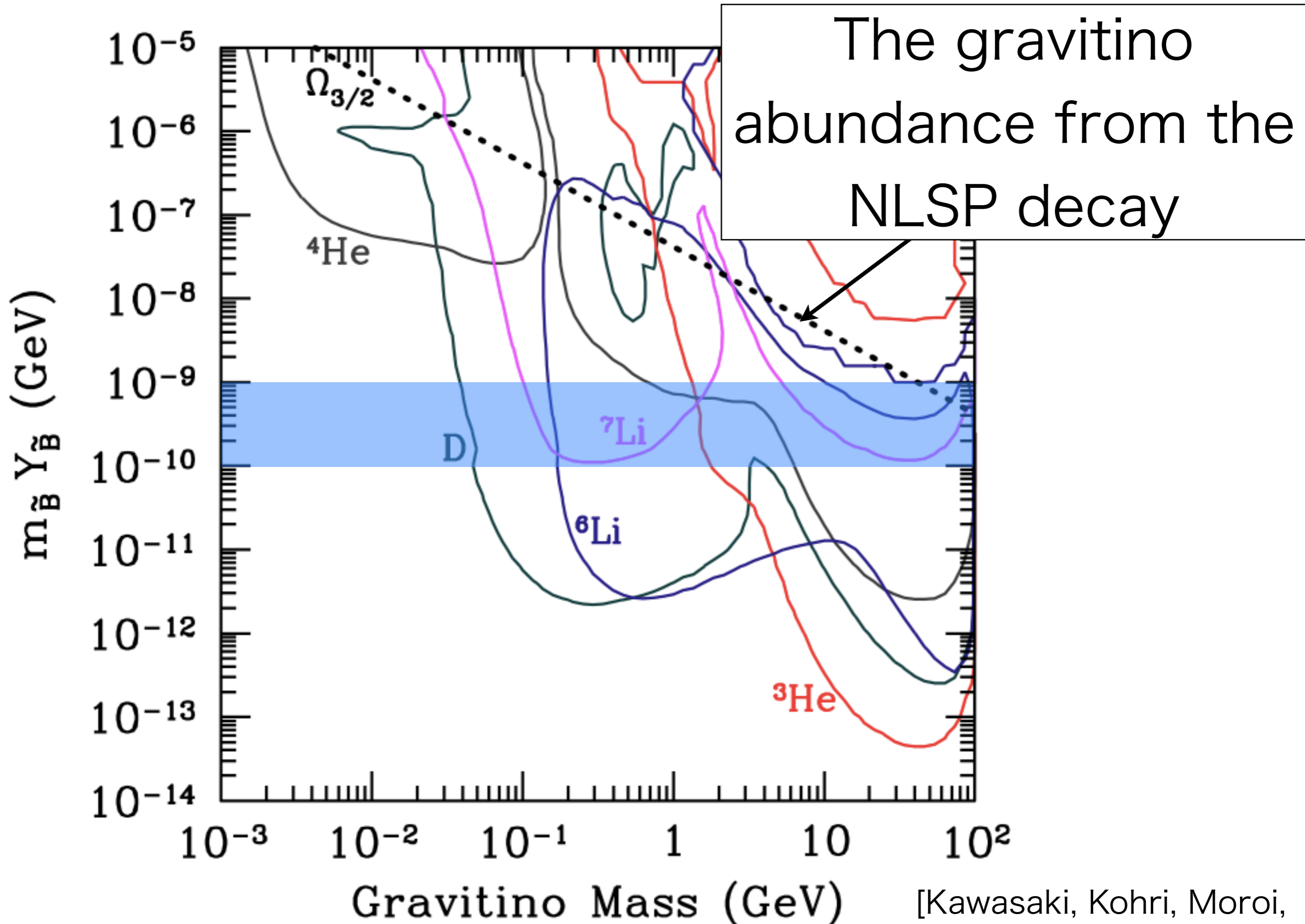
$$\Gamma(\tilde{B} \rightarrow \psi_{3/2}Z) \simeq \frac{\sin^2 \theta_W M_{\tilde{B}}^5}{48\pi M_{\text{PL}}^2 m_{3/2}^2}$$

~100s for 150GeV Bino and
0.1 GeV gravitino

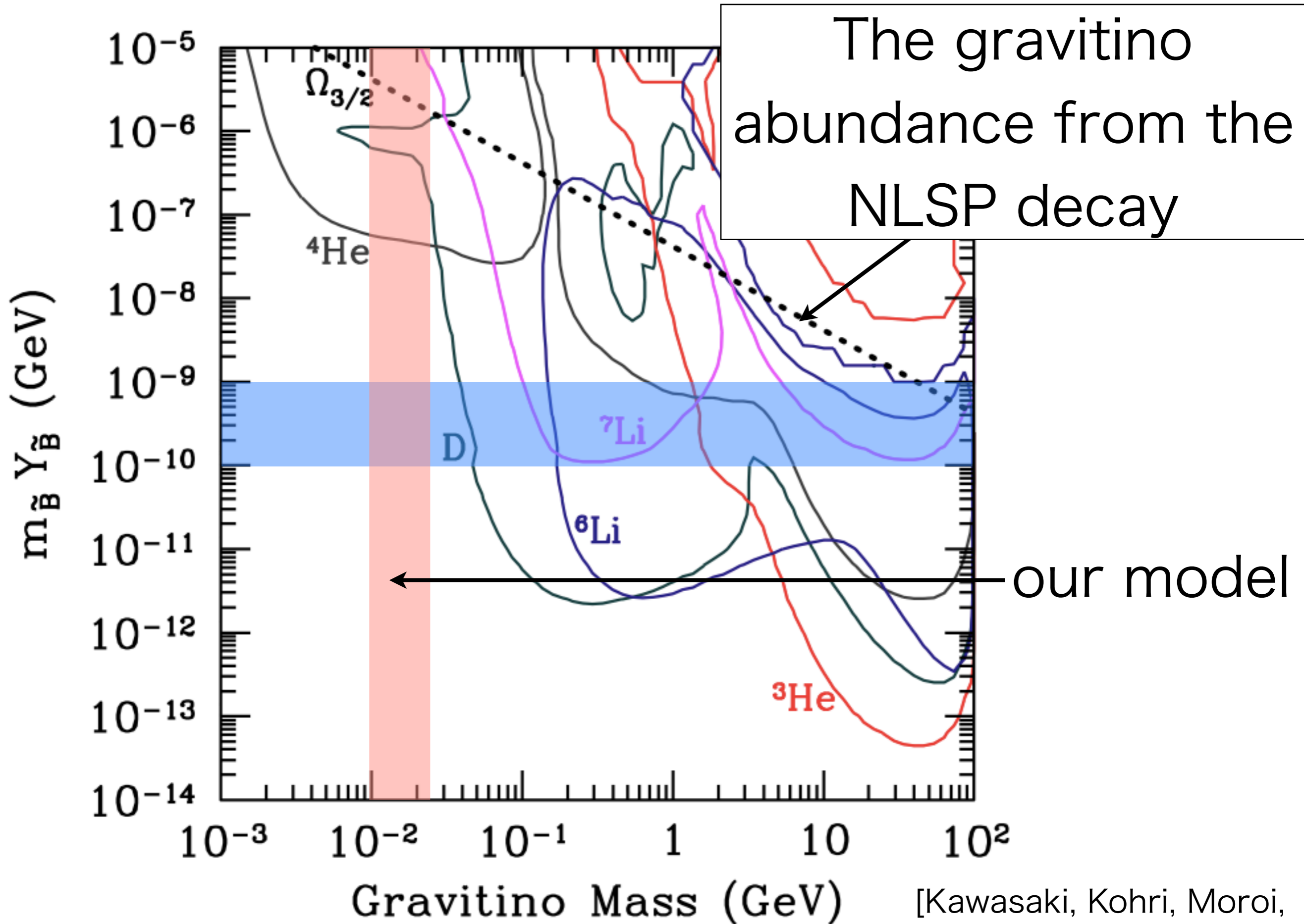
$c\tau \sim 10^{10}$ m Stable at the collider scale

The BBN prediction constrains the abundance of the NLSP and the gravitino mass

Roughly, if the gravitino mass is larger than about 0.1 GeV, then it becomes problematic.



[Kawasaki, Kohri, Moroi, Yotsuyanagi, 2008]



[Kawasaki, Kohri, Moroi, Yotsuyanagi, 2008]

Summary

Only few viable models can explain

Higgs boson mass, muon $g-2$
anomaly, LHC results

Consistent with cosmological
gravitino problem

**Light sleptons can be observed at
the ILC!**