

# Top threshold at ILC

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LCWS2013 Tokyo

TTbarXSection: Bneke, YK, Schuller :

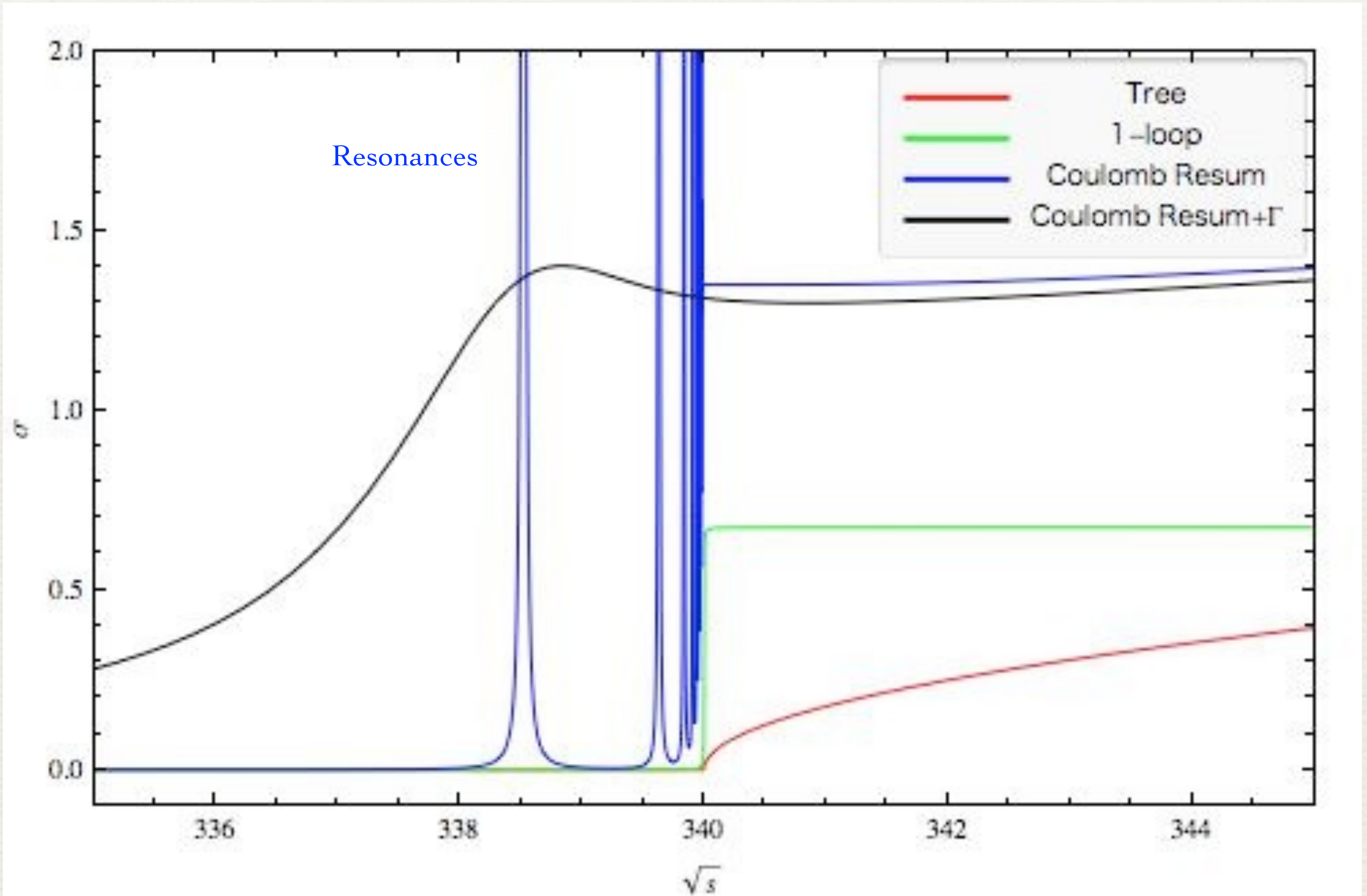
hep-ph.0501289, 0705.4518[hep-ph], 0801.3464[hep-ph], in preparation

Ultrasoft correction: Beneke, YK, Penin:

0804.4004[hep-ph], 0706.2733[hep-ph]

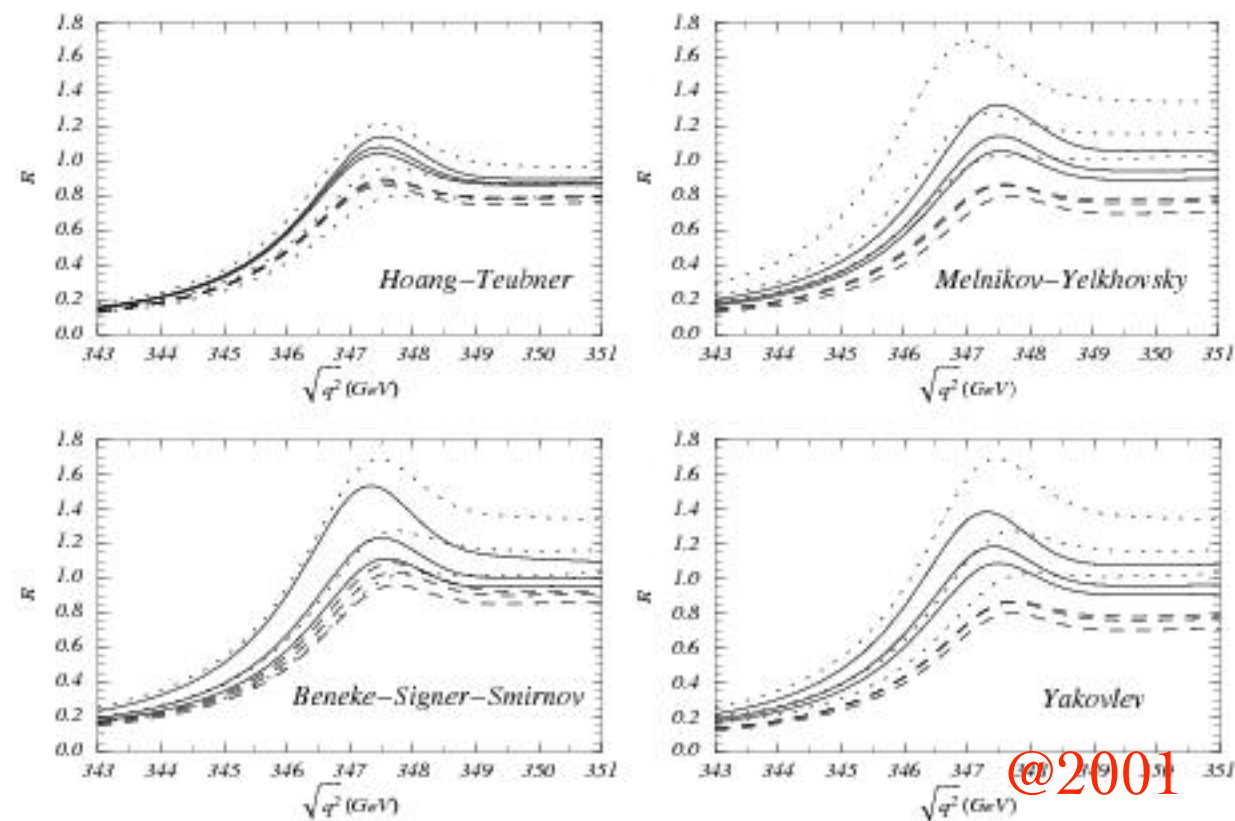


# $e^+e^- \rightarrow tt$ near threshold



# Top threshold

- ◆  $m_t$  and  $y_t$  measurement is important to test SM/BSM
- ◆  $(\Delta m_t)_{\text{exp}} \leq 50 \text{ MeV} \rightarrow$  theory goal  $\delta\sigma \leq 3\%$
- ◆ NNLO result @2001 **Scale uncertainty about 20 %**



- ◆ We will advance the theory calculation to N<sup>3</sup>LO



# Effective Field Theory

- Integrate out **Hard** (Caswell-Lepage('86))

$$\mathcal{L}_{\text{NRQCD}} = \psi^\dagger \left( iD_0 + \frac{\vec{D}^2}{2m_t} \right) \psi + [\psi \rightarrow \chi] + \dots$$

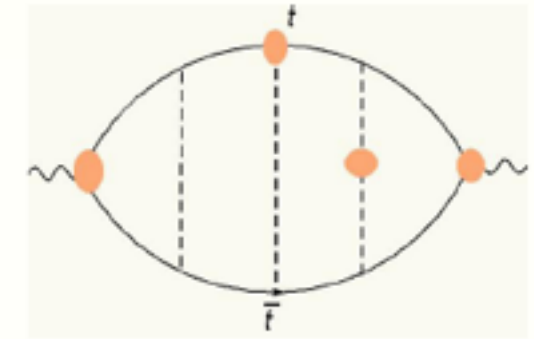
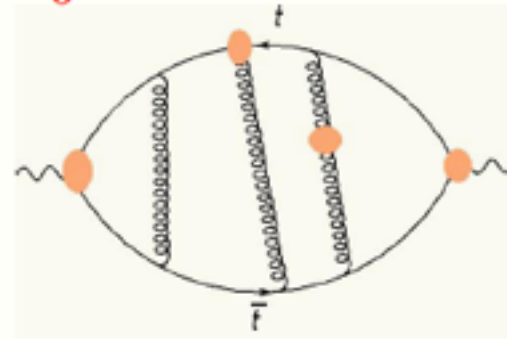
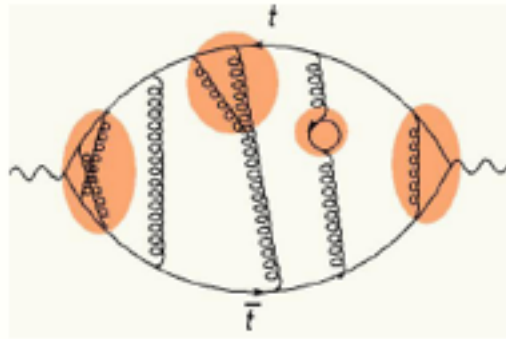
- Integrate out **Soft/Potential** gluons  
(Pineda-Soto('98); Luke-Manohar-Rothstein('99))

$$\begin{aligned} \mathcal{L}_{\text{pNRQCD}} = & \psi^\dagger \left( i\partial_0 + \frac{\vec{\nabla}^2}{2m_t} \right) \psi + \int d\vec{r} [\psi^\dagger \chi] V_{\text{pot}}(r) [\chi^\dagger \psi] \\ & + ig \psi^\dagger \left[ A_{0,us} + \frac{\nabla \cdot \vec{A}_{us}}{m} \right] \psi - \frac{1}{4} F_{us}^2 + \dots \end{aligned}$$

- Remaining Mode is **Ultra Soft** gluon:  $k \sim m(v^2, \vec{v}^2)$

# Threshold XS in EFT

Principal quantity is  $\Pi(q) = i \int d^4x e^{iqx} \langle 0 | J^\mu(x) J_\mu(0) | 0 \rangle$



- Integrating **hard** mode  $\rightarrow$  NRQCD (Caswel-Lepage '86):

$$J^i(x) = [\bar{t} \gamma^i t] \rightarrow c_v [\psi^\dagger \sigma^i \chi]$$

- Integrating **soft/potential** modes  $\rightarrow$  PNRQCD

(Pineda-Soto '97/Luke - Manohar-Rothstein'99):

$$\mathcal{L}_{\text{PNRQCD}} = \psi^\dagger \left[ i\partial_0 + \frac{\partial^2}{2m} + \frac{\partial^4}{8m^3} - g_s \mathbf{x} \mathbf{E}(t, \mathbf{0}) \right] \psi + (\psi \leftrightarrow \chi) \\ + \int d\mathbf{r} [\psi^\dagger \psi](x + \mathbf{r}) V_{\text{pot}}(\mathbf{r}) [\chi^\dagger \chi](x) + \dots$$

$$\Pi(q) = i \int d^4x e^{iEx} c_v^2 \langle 0 | [\psi^\dagger \sigma^i \chi](x) [\chi^\dagger \sigma_i \psi](0) | 0 \rangle$$

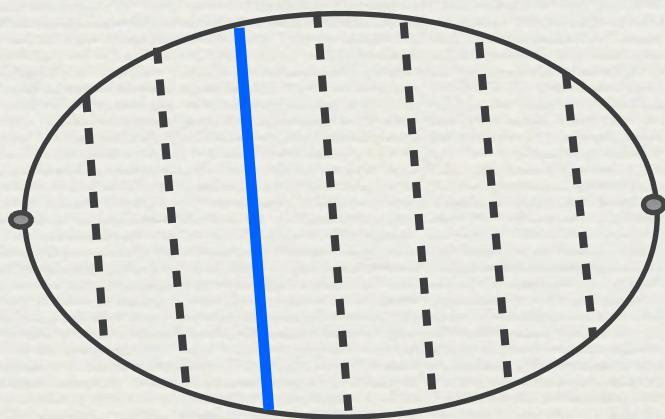


# Potential Perturbation

$$V_{pot} = -\frac{C_F \alpha_s}{r} + \frac{C_2}{r^2} + C_3 \delta(\mathbf{r}) + \dots$$

- Higher order corr to the potential (case of Coulomb pot)

$$\tilde{V}_C = -\frac{4\pi C_F \alpha_s(\mathbf{q})}{\mathbf{q}^2} \times \left[ 1 + \frac{\alpha_s(\mathbf{q})}{4\pi} a_1 + \left( \frac{\alpha_s(\mathbf{q})}{4\pi} \right)^2 a_2 + \left( \frac{\alpha_s(\mathbf{q})}{4\pi} \right)^3 \left[ a_3 + 8\pi^2 C_A^3 \left( \frac{1}{3\epsilon} + \ln \frac{\mu_{US}^2}{\mathbf{q}^2} \right) \right] \right]$$



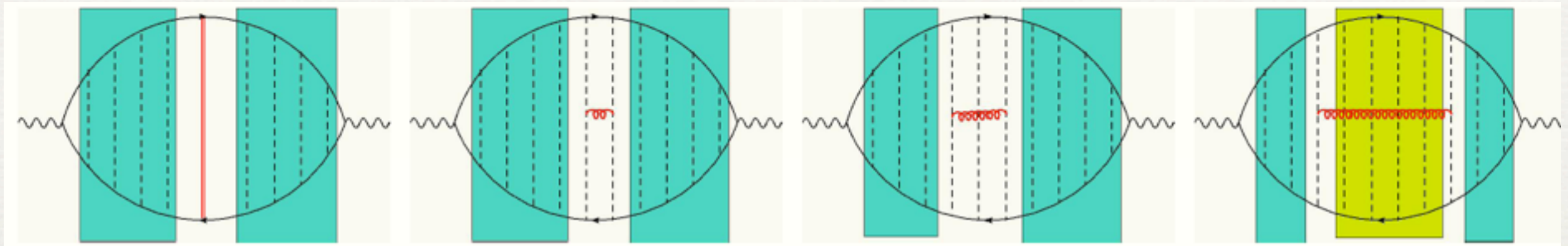
All order Coulomb ladder

Single insertion of  $a_1$ ,  $a_2$ ,  $a_3$  is NLO, NNLO and NNNLO, respectively

We treat higher order potentials as perturbation to the all order resummation of LO Coulomb exchange.

→ We call Potential Insertion

# Ultrasoft Correction

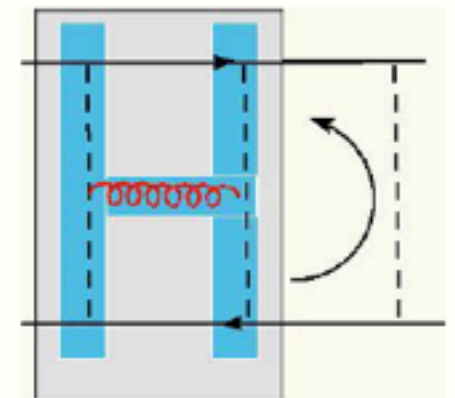


- quark-gluon vertex is  $1/m$  suppressed;  $\psi^\dagger (iD_0 + \frac{D^2}{2m}) \psi$
- $n_g$ , number of potential exchange  $\sim \Delta t$ ;  $n_g > 1 \Leftrightarrow$  UV finite
- **ADM**  $1/\epsilon$  of the Coulomb pot is a counter term for us corr

$$\delta V_C^{\text{ADM}} = \frac{1}{\epsilon} \frac{C_A^3 \alpha_s^4}{q^2}$$

$$H_0 = -\frac{\alpha_s^3 C_A^2}{\epsilon} \left( \frac{1}{mq} + \frac{2(p^2/m - E)}{q^3} + \epsilon L_{\text{Bethe}} \right) + O(\epsilon)$$

$$H_1 = -\frac{1}{\epsilon} \frac{C_A^2 (C_A - 2C_F) \alpha_s^4}{q^2} + L_{\text{Bethe}} + O(\epsilon)$$



- UV cancelation happens tricky way:

$$\frac{(p^2/m - E)}{q^3} \Rightarrow \frac{C_F \alpha_s}{q^2} \text{ (Eq. of motion)}$$

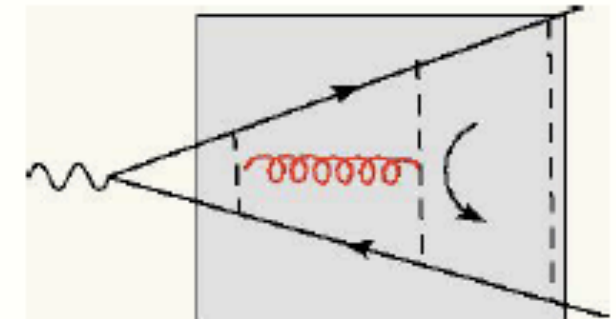
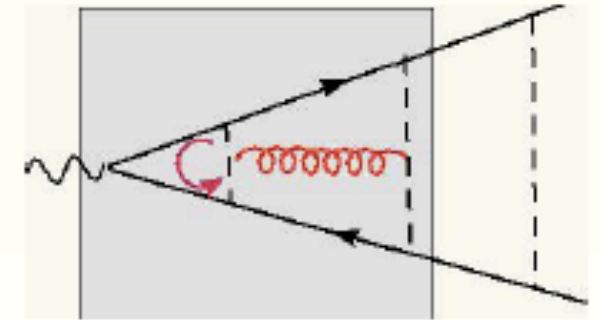


# US renormalization II: Vertex

$$\delta V_C^{\text{ADM}} = \frac{1}{\epsilon} \frac{C_A^3 \alpha_s^4}{q^2}$$

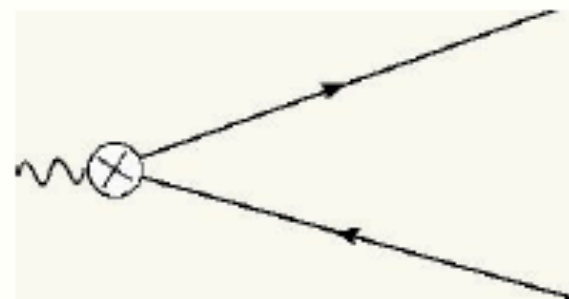
$$H_0 = -\frac{\alpha_s^3 C_A^2}{\epsilon} \left( \frac{1}{mq} + \frac{2(p^2/m - E)}{q^3} + \epsilon L_{\text{Bethe}} \right) + O(\epsilon)$$

$$H_1 = -\frac{1}{\epsilon} \frac{C_A^2 (C_A - 2C_F) \alpha_s^4}{q^2} + L_{\text{Bethe}} + O(\epsilon)$$



- Loop near photon vertices are more singular  
 $\Leftrightarrow$  Vertex Renormalization
- $1/\epsilon$  cancelation does not happen exactly anymore
  - due to vertex divergence, (3+1) dimensional Eq. of motion is invalid
  - diagrams with different loop order get mis-match of Eq. of motion

Needs external current renormalization





- ♦ all the logarithmic parts obtained analytically
- ♦ there are uncanceled  $i\Gamma/\epsilon$ , which in turn produces scale dependence  $\rightarrow$  unstable top EFT

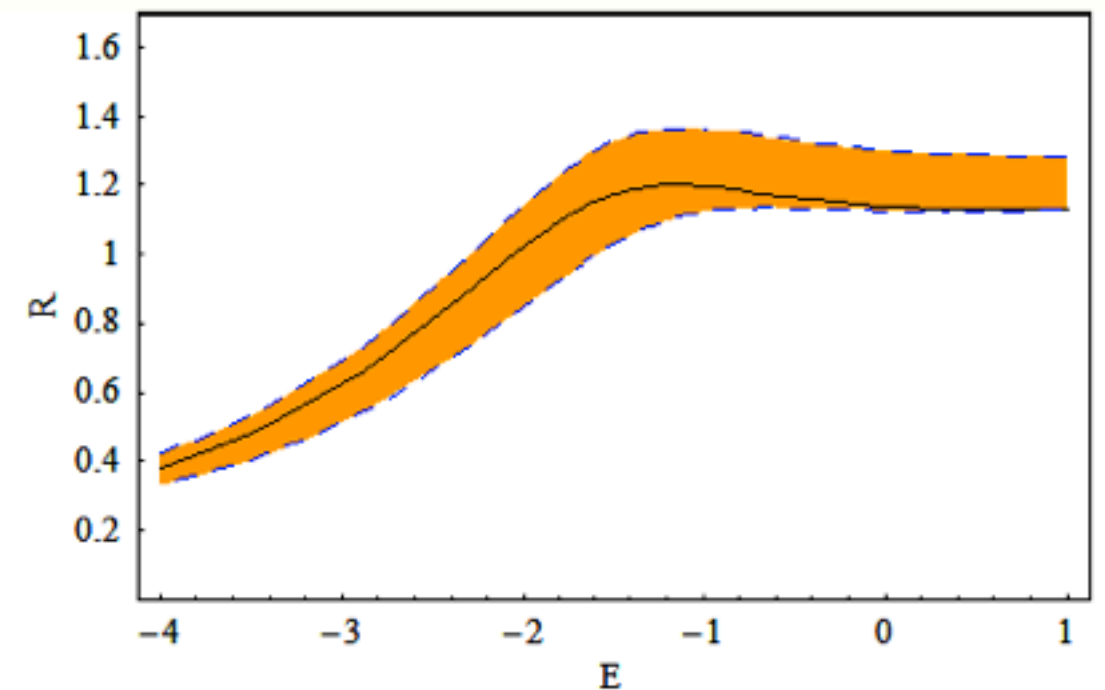
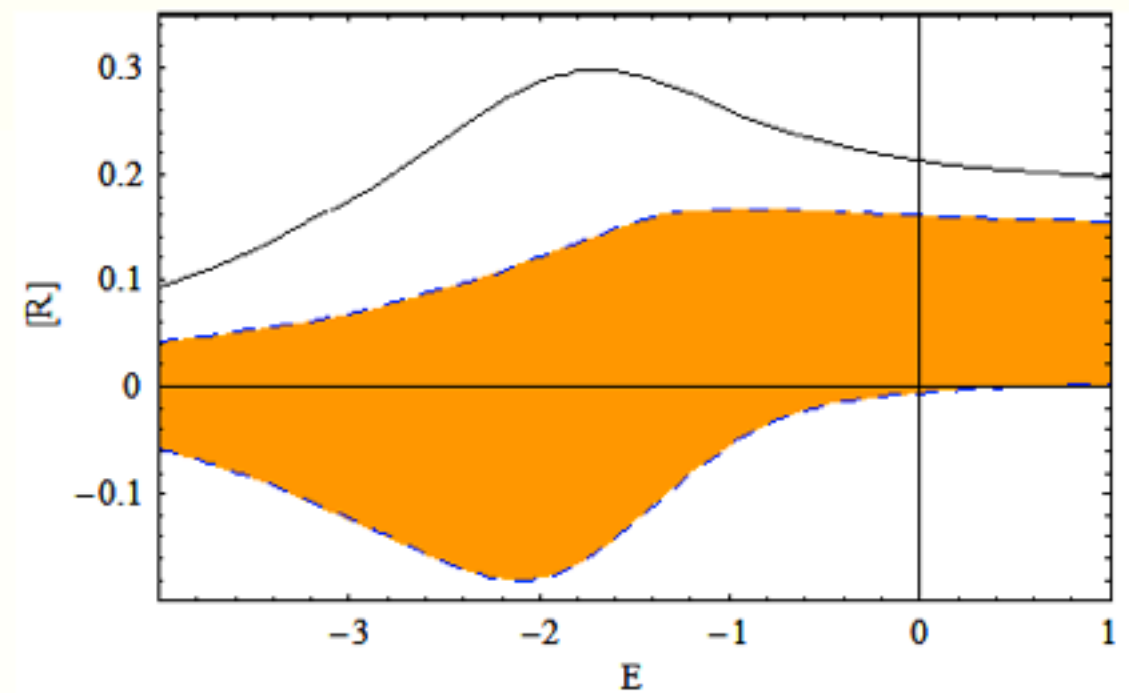
$$\delta^{us} G(E) = \frac{2m^2 \alpha_s^4}{9\pi^2} \left\{ \left[ \frac{17 i\hat{\Gamma}_t}{24} + \frac{527 \hat{G}_C}{72} \right] \frac{1}{\epsilon^2} + \left[ \frac{17 i\hat{\Gamma}_t}{12} + \frac{221 \hat{G}_C}{36} \right] \frac{L_\mu}{\epsilon} + \left[ \left( \frac{19}{12} \ln 2 - \frac{91}{72} \right) i\hat{\Gamma}_t \right. \right. \\ \left. \left. + \left( -\frac{119}{12} \ln 2 + \frac{2059}{108} \right) \hat{G}_C \right] \frac{1}{\epsilon} + \left[ -\frac{34 i\hat{\Gamma}_t}{3} - \frac{595 \hat{G}_C}{9} \right] L_{\alpha_s}^2 + \left[ -\frac{17 i\hat{\Gamma}_t}{12} - \frac{833 \hat{G}_C}{36} \right] L_\mu^2 \right. \\ \left. + \left[ \frac{34 i\hat{\Gamma}_t}{3} + \frac{748 \hat{G}_C}{9} \right] L_{\alpha_s} L_\mu + \left[ \frac{2380 \mathcal{P}^2}{27} + \left( \frac{272 \ln 2}{9} - \frac{23483}{162} + \frac{2380}{27\lambda} + \frac{272}{27\lambda^2} \right) \mathcal{P} \right. \right. \\ \left. \left. + \left( \frac{27\lambda}{2} - \frac{16}{3\lambda} \right) \psi' + \frac{64}{27\lambda^3} + \frac{4(-1331 + 306 \ln 2)}{81\lambda} + \frac{4(-199 + 114 \ln 2)}{81\lambda^2} \right] L_{\alpha_s} \right. \\ \left. + \left[ -\frac{1496 \mathcal{P}^2}{27} + \left( -\frac{34 \ln 2}{3} + \frac{5065}{72} - \frac{1496}{27\lambda} - \frac{136}{27\lambda^2} \right) \mathcal{P} + \left( \frac{8}{3\lambda} - \frac{81\lambda}{8} \right) \psi' \right. \right. \\ \left. \left. - \frac{32}{27\lambda^3} + \frac{163 - 114 \ln 2}{27\lambda^2} + \frac{271 - 51 \ln 2}{9\lambda} \right] L_\mu + \delta^{us}(\hat{E}) \right\},$$

$$L_\mu = \ln \frac{\mu}{m_t}, \quad L_{\alpha_s} = \ln \alpha_s, \quad \lambda = \frac{C_F}{2\sqrt{-\hat{E}}}, \quad \mathcal{P} = \ln \left( \frac{C_F}{\lambda} \right) + \gamma_E + \psi(1 - \lambda),$$

# US corr

- ◆ Ultrasoft correction alone; constant (solid), log + const. orange band
- ◆  $R_{LO} + R_{US}$

Ultrasoft correction is scale dependent and not physical





# Bulding blocks

$$J^i = c_v \psi^\dagger \sigma^i \chi + d_v \frac{1}{6m_t^2} \psi^\dagger \sigma^i \mathbf{D}^2 \chi$$

$c^{(2)}$ : Beneke-Signer-Smirnov('97),  
Czanecki-Melnikov('97)

$c_{n_f}^{(3)}$ : Marquard-Piclum- Seidel-Steinhauser(06)

$d_v^{(1)}$ : Luke-Savage('97)

$$\mathcal{L}_{\text{QCD}} \Leftrightarrow \mathcal{L}_{\text{PNRQCD}}$$

$a_2$ : Schröder('98)

$a_{3,\text{pade}}$ : Chishtie-Elias (01)  
(New:  $a_{3, n_f}$ )

Smirnov-Smirnov-Steinhauser (Sep.08)

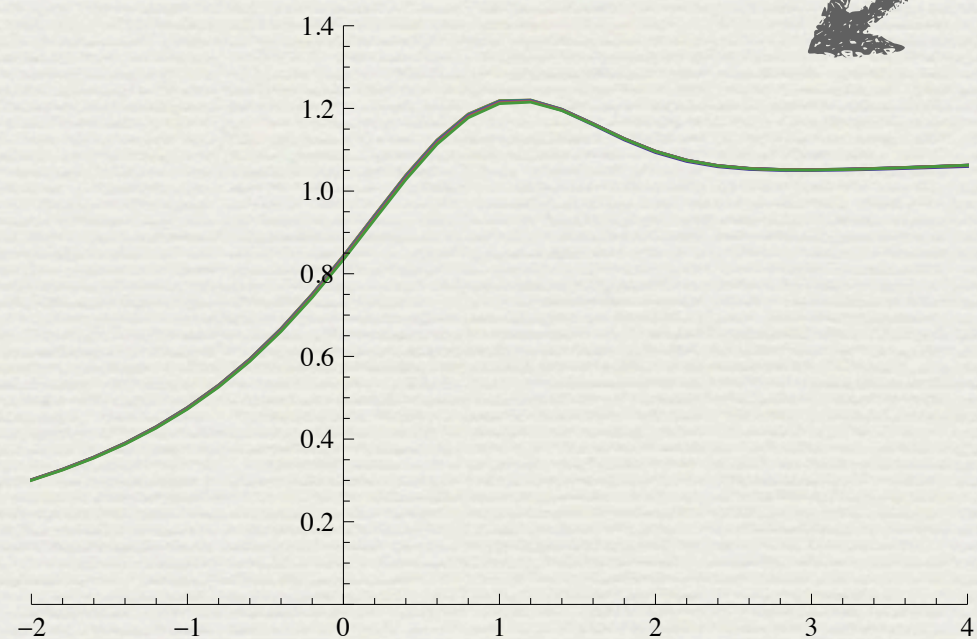
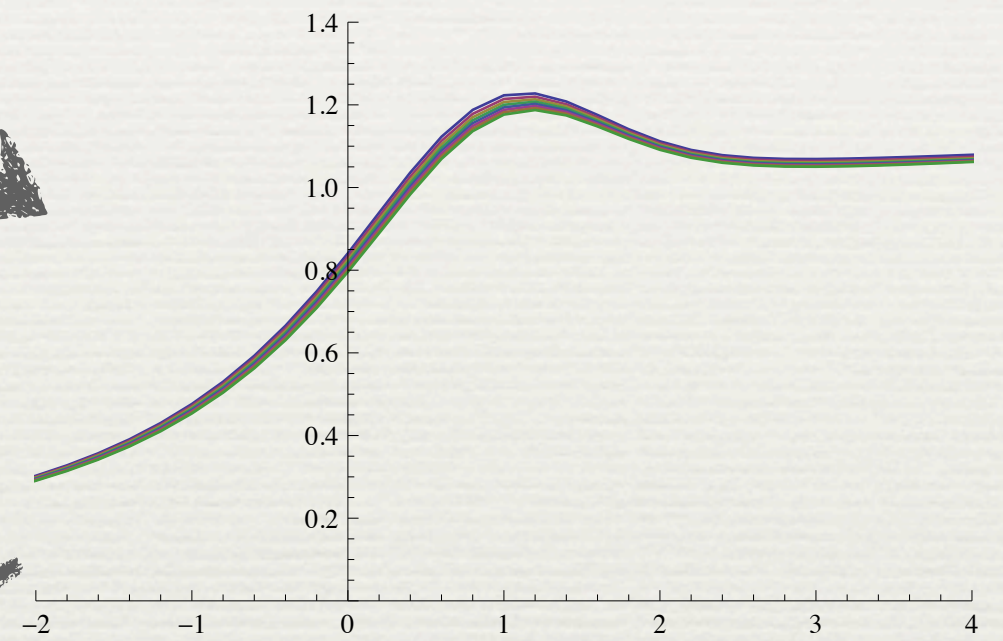
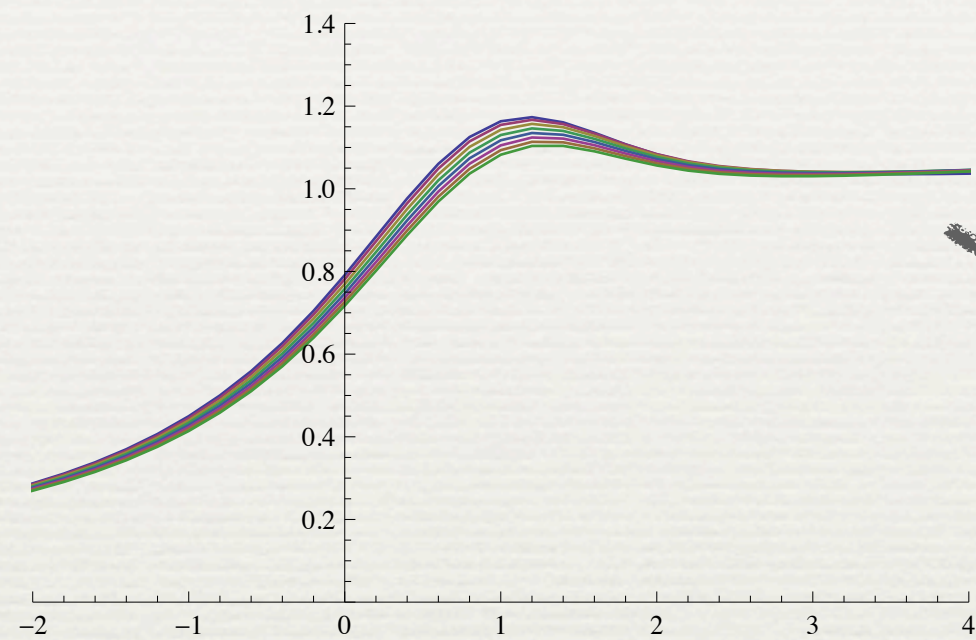
$\delta\mathcal{L}^{(1)}$ : Manohar('97),  
Beneke-Signer-Smirnov('99),  
Wüster-Schuller('03)

$\delta\mathcal{L}^{(2)}$ : Kniehl-Penin- Smirnov-Steinhauser(02)  
( $\delta\mathcal{L}^{(2)} = \mathcal{O}(\epsilon)$  not known)

$\delta\mathcal{L}^{(us)}$ : Brambilla-Pineda-Soto-Vairo('99),  
Kniehl-Penin- Smirnov-Steinhauser(02)

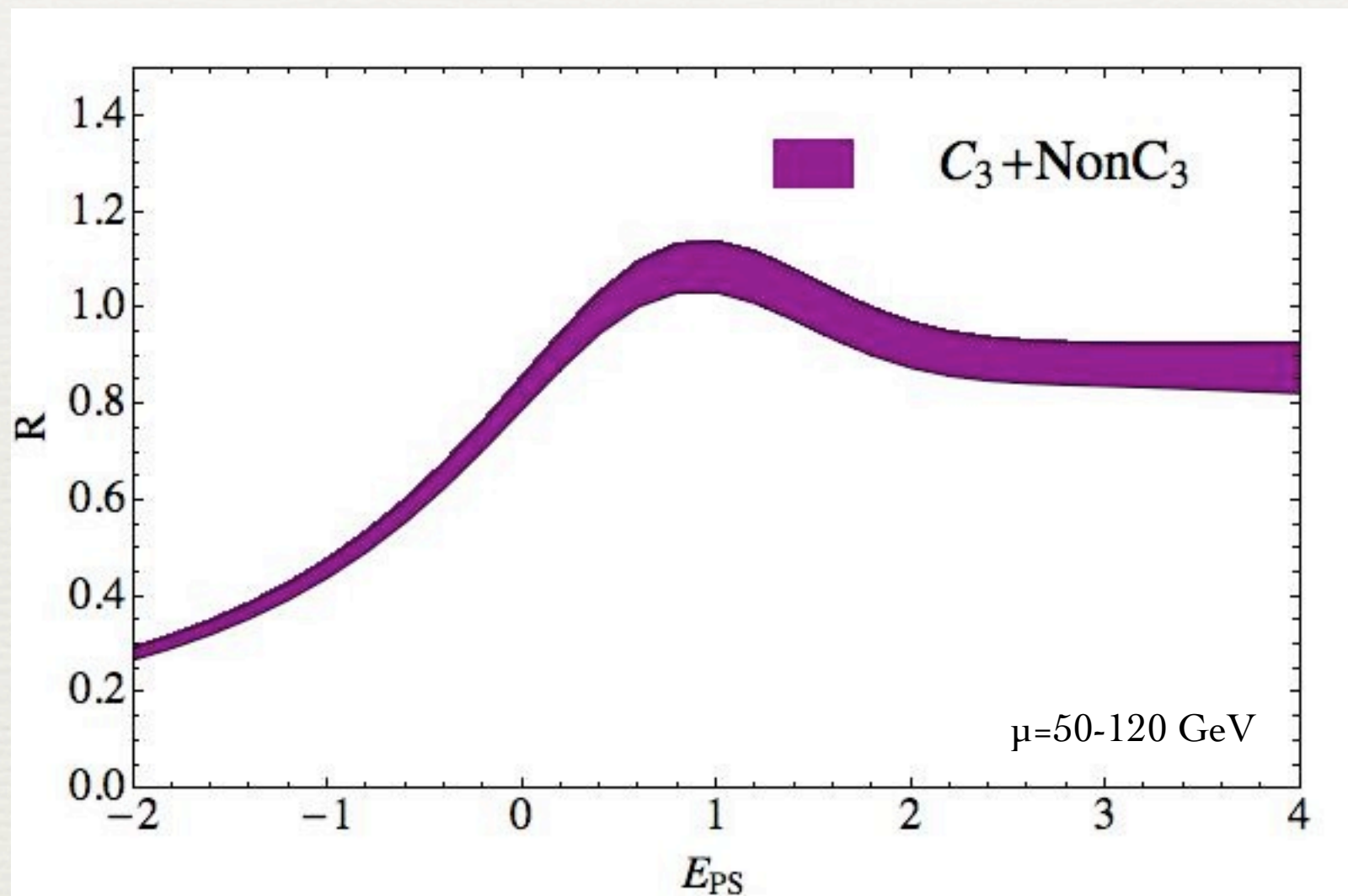
Anzai-YK-Sumino,  
Smirnov-Smirnov-  
Steinhauser (2010)

# Coulomb Corr

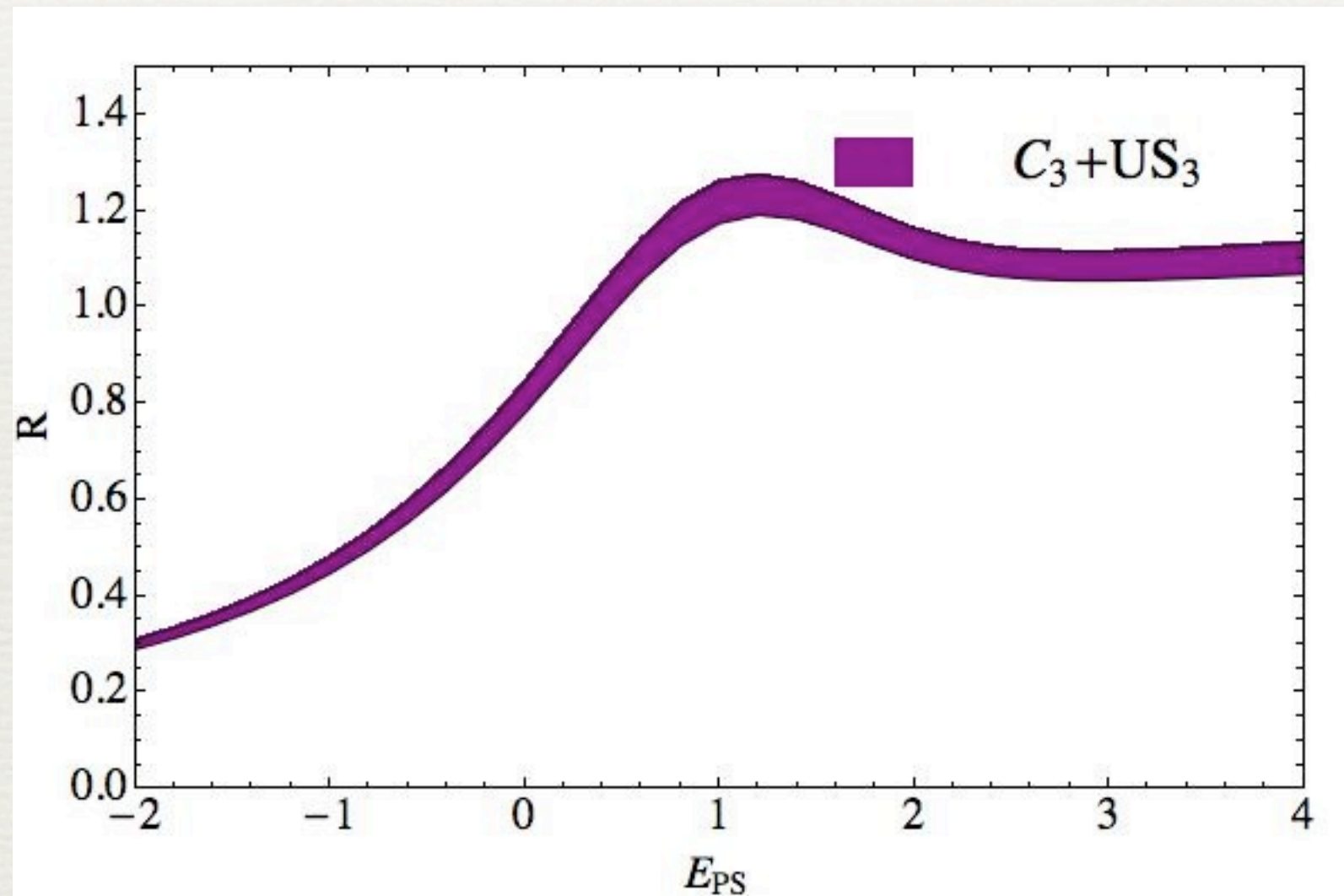




# NonCoulomb Corr

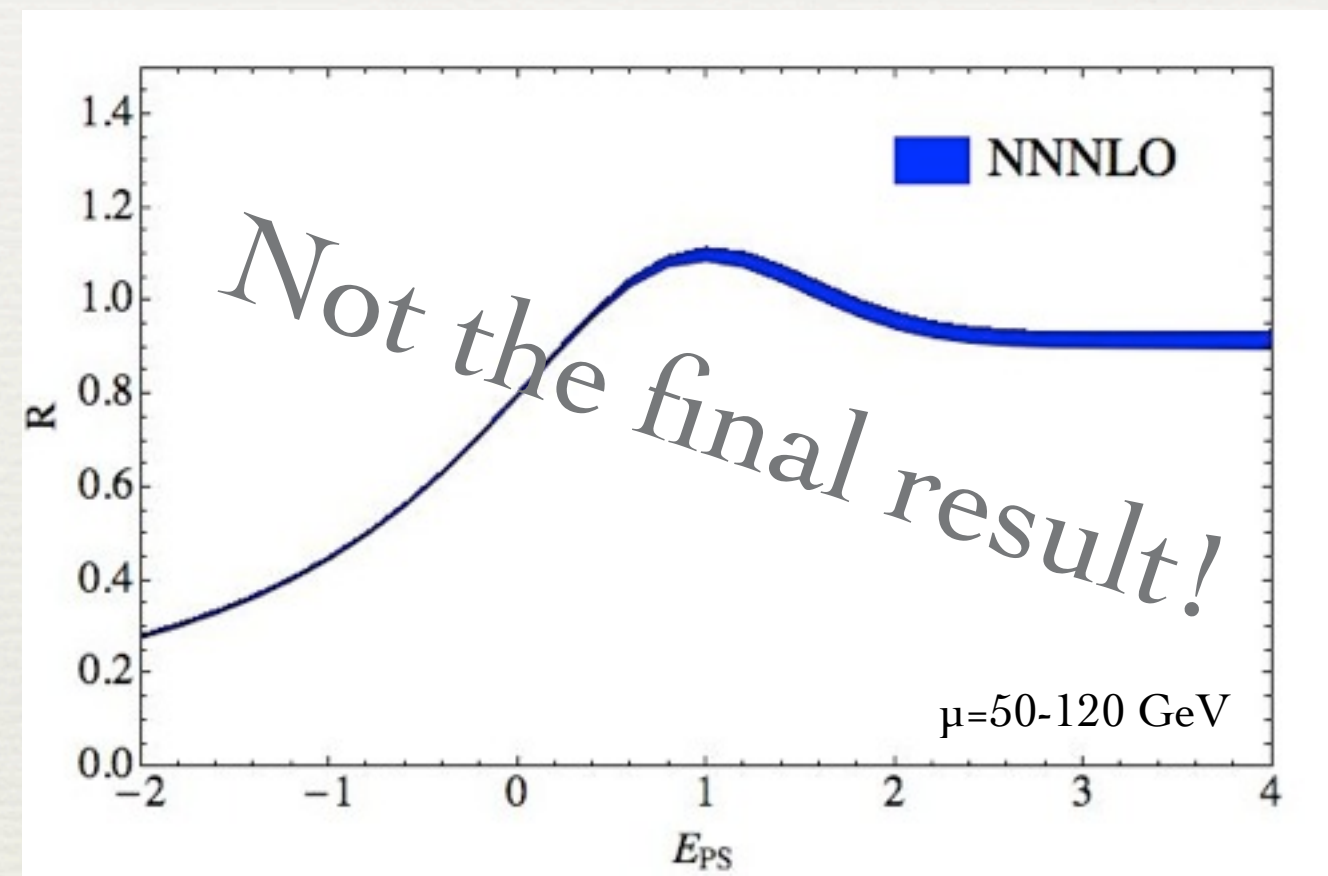
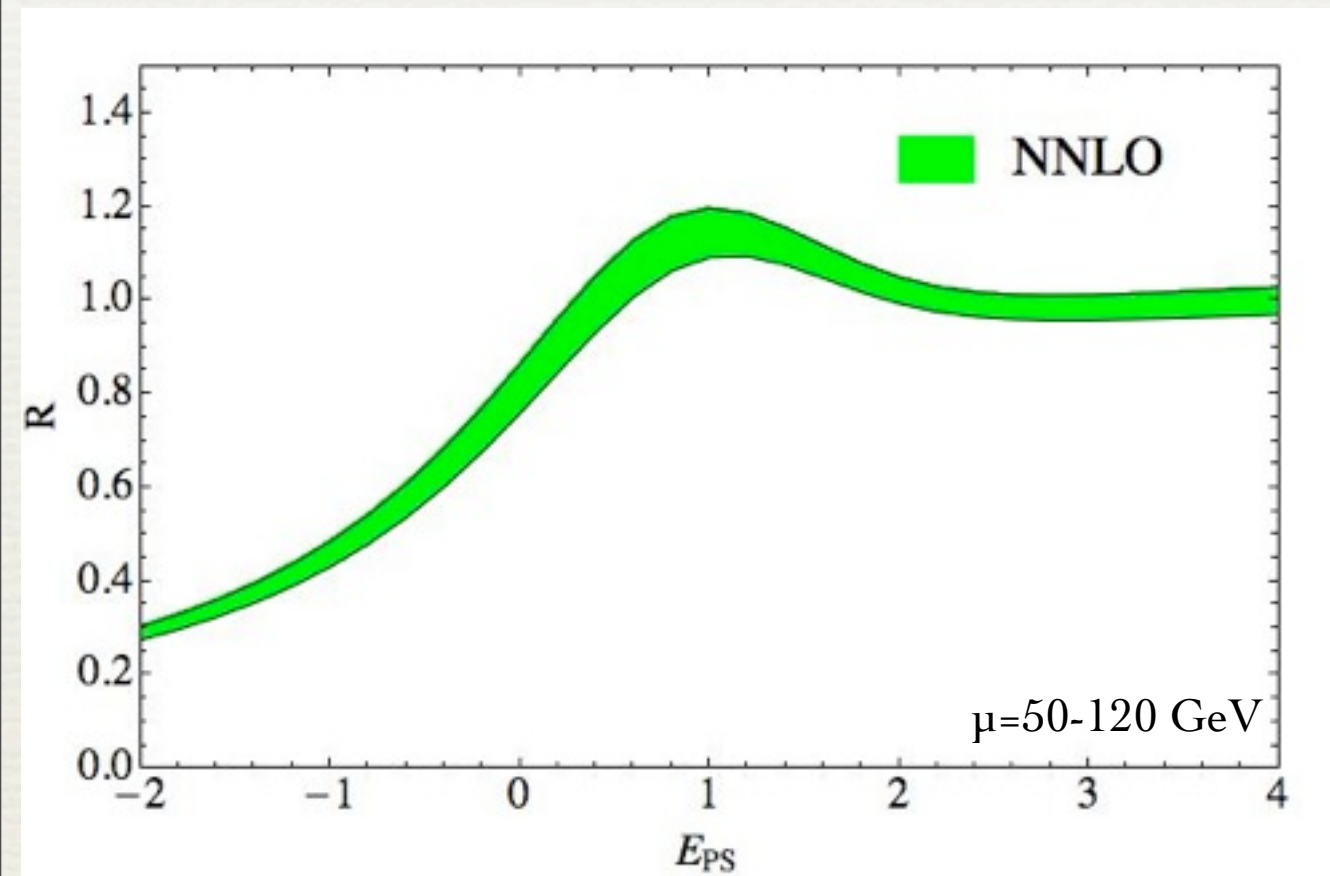


# Ultrasoft Corr





# $N^2$ LO and $N^3$ 'LO



Beneke-YK-Schuller in progress

Inputs:  
 $m_{PS}(20\text{GeV})=170$  GeV,  
 $\Gamma_t=1.4$  GeV,  
 $\alpha_s(M_Z)=0.1180$



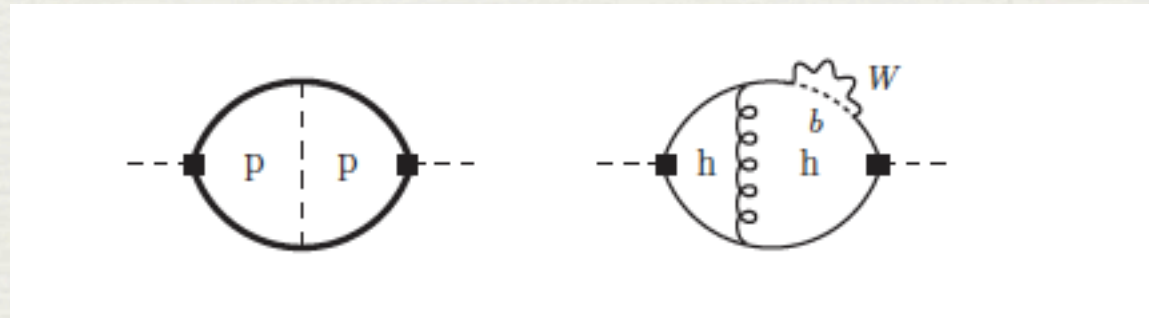
# Issue of unstable top

- ◆ There are several different effects concerning to unstable nature of top quarks. → talk by A. Hoang, and P. Ruiz
- ◆ We have taken into account one of unstable top effect by  $E \rightarrow E + i\Gamma_t$ , We have introduced some uncanceled scale dependence.

$$\delta^{(2)}G_{div} = \frac{(2\pi)^2}{3\epsilon\alpha_s} C_F(2C_F + 3C_A) G_C(E) + \frac{\Gamma}{\epsilon} \frac{2\pi C_F m}{\alpha_s}$$

multiplicative Z-factor

need to renormalized counter term by unstable top EFT



It was started to seriously think this issue few years ago  
 [ Hoang-Reisser-Ruiz, Beneke-Falgari-Schwinn-Signer-Zanderighi, Jatzten-Ruiz, Penin-Piclum,...].  
 We have started to implement a prescription in our code, to fix the scale dependence, but not yet completed.



# Summary

- ◆  $e^+e^- \rightarrow t\bar{t}$  cross section near threshold computed at N<sup>3</sup>LO  $\rightarrow$  sizable correction, reduction of scale dep.
- ◆ Still few missing parts  $\rightarrow$  c3,...
- ◆ Implementation of some of unstable effects started