

Global Fits Beyond the SM

Martin Wiebusch

in collaboration with

Ulrich Nierste, Otto Eberhardt, Alexander Lenz,
Heiko Lacker, Geoffrey Herbert, Andreas Menzel

based on [PRL 109 (2012) 241802], [JHEP07 (2013) 118]

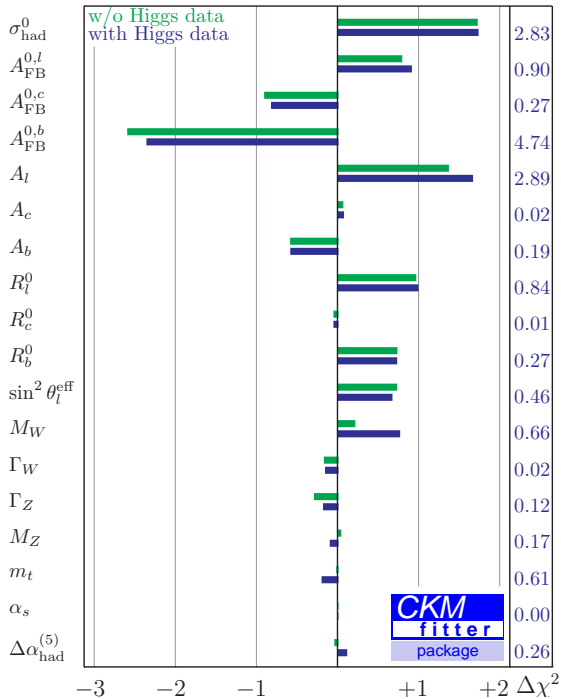


LCWS13, November 2013

SM Electroweak Fit

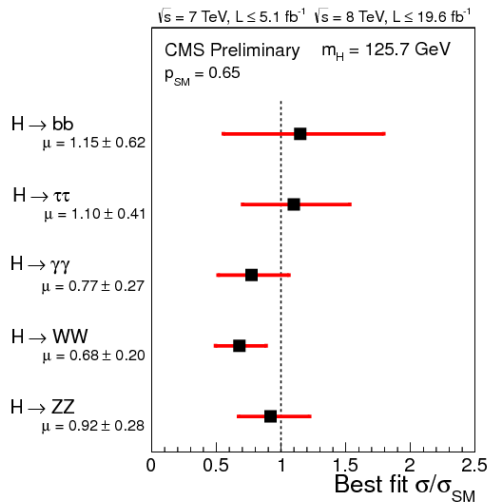
- Fits were used in the LEP era to determine the free parameters of the SM.
- With the Higgs discovery the last free parameter of the SM has been measured directly.
- The SM electroweak fit determines these Parameters from the electroweak precision observables (EWPOs). \rightarrow
- The influence of this measurement on the SM electroweak fit is relatively small.

[PRL 109 (2012) 241802]



Higgs Signal Strengths

- The EWPOs also severely constrain many models beyond the SM.
- Constraints from the observed Higgs signal strengths can already be equally powerful.
- Combined fits of both types of data can (and should) be used to study models beyond the SM, but...
- the statistical interpretation of the fits can be (conceptually and technically) less trivial for models beyond the SM.



Chi-squares and p -values

- The usual analytic relation between $\Delta\chi^2$ and the p -value (Wilks' theorem) requires **nested models**.

- The additional fermions of the SM4 **do not decouple**.

⇒ You cannot obtain the SM3 as a limiting case of the SM4.

⇒ The computation of the p -value requires a **very expensive numerical simulation**, which is **unfeasible without special simulation methods**.

- These methods were implemented in the public code **myFitter** (<http://myfitter.hepforge.org>) and documented in [CPC 184 (2013) 2438].

- The SM4 is excluded at **5.3 standard deviations**. (Wilks' theorem gives 3.5 standard deviations.)

2HDM Fits

The 2HDM of type II

- Two scalar SU(2) doublets.
- A **softly broken \mathbb{Z}_2 symmetry** which forbids FCNCs.
- No Higgs-sector CP violation.
- Scalar particle content: h, H, A, H^\pm .
- Independent sets of **real parameters** are
 - $v_2/v_1 \equiv \tan\beta, m_{12}^2, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5$
 - $\tan\beta, \beta - \alpha, m_{12}^2, m_h, m_H, m_A, m_{H^\pm}$
- This time we have a **decoupling limit**:
 $\beta - \alpha = \pi/2, m_H, m_A, m_{H^\pm} \gg m_h$.
- The map between the two parametrisations **near the decoupling limit** is **not very smooth**.

Theoretical Constraints

- The Higgs potential must be **bounded from below**

$$\lambda_1 > 0 \quad , \quad \lambda_2 > 0 \quad , \quad \lambda_3 > -\sqrt{\lambda_1 \lambda_2} \quad , \\ |\lambda_5| < \lambda_3 + \lambda_4 + \sqrt{\lambda_1 \lambda_2} \quad .$$

- ‘Our’ minimum of the Higgs potential must be the **global minimum**

$$m_h^2, m_H^2, m_A^2, m_{H^\pm}^2 \geq 0 \quad , \\ m_{12}^2(m_{11}^2 - m_{22}^2\sqrt{\lambda_1\lambda_2})(\tan\beta - (\lambda_1/\lambda_2)^{1/4}) > 0 \quad .$$

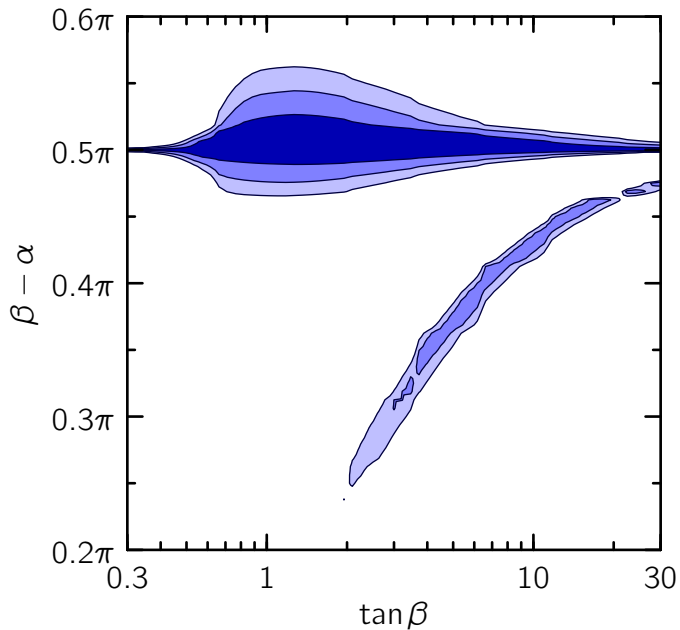
- The Higgs self-couplings must be **perturbative**.
(Implemented by requiring **tree-level unitarity**.)
- ⇒ Fit requires **optimisation under non-linear constraints**.

Experimental Constraints

- Full set of **electroweak precision observables** (no S , T , U).
- **Signal strengths** of the light Higgs boson (including correlations between different production mechanisms).
- Limits on heavy $H \rightarrow WW, ZZ$ and $H \rightarrow \tau\tau$ resonances.
- **Flavour observables** relevant for the low $\tan\beta$ region:
 Δm_{B_s} and $\text{Br}(\bar{B} \rightarrow X_s \gamma)$.



$\tan\beta$ vs. $\beta - \alpha$



SM Fits

The Fourth
Generation

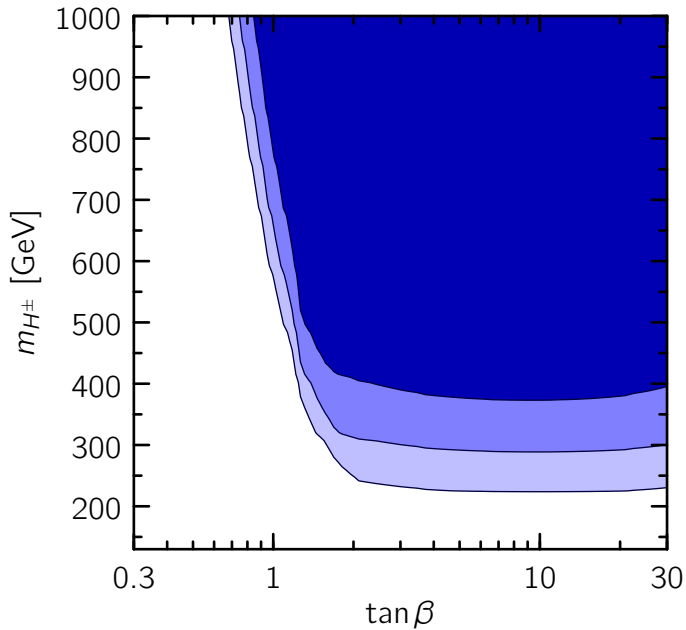
2HDM Fits

Constraints

Fit results

Conclusions

$\tan\beta$ vs. m_{H^\pm}



SM Fits

The Fourth
Generation

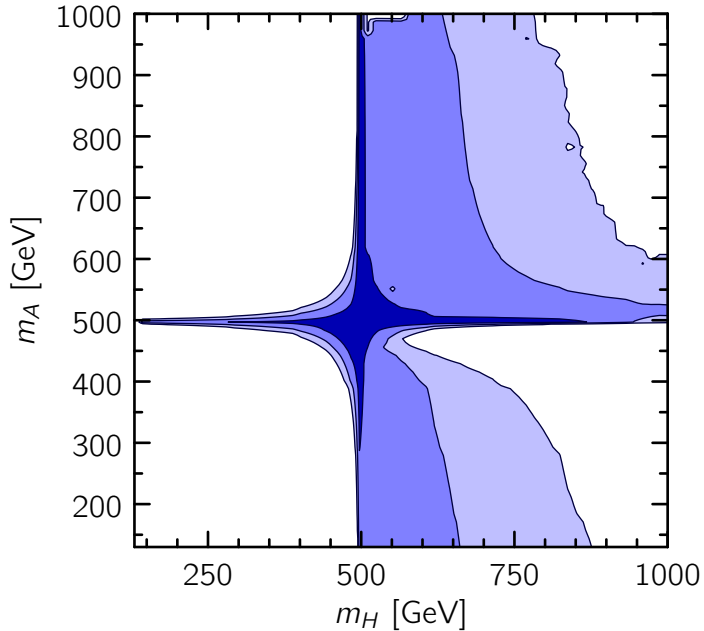
2HDM Fits

Constraints

Fit results

Conclusions

m_H vs. m_A , $m_{H^\pm} = 500$ GeV fixed



SM Fits

The Fourth
Generation

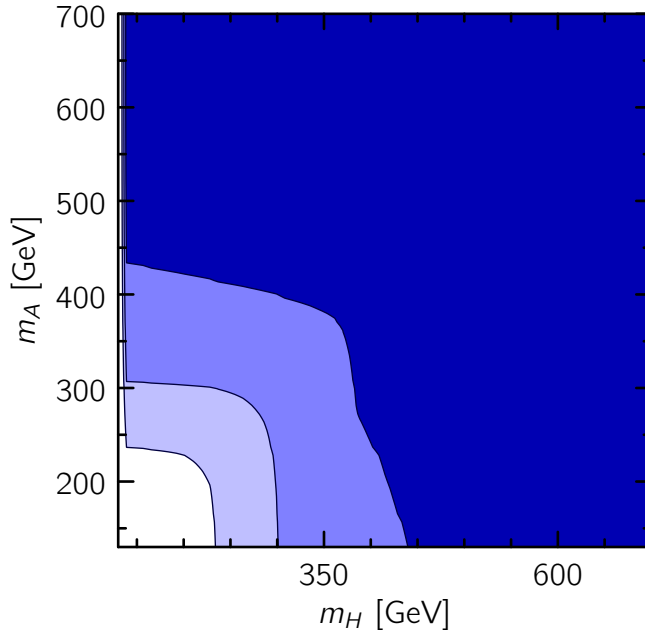
2HDM Fits

Constraints

Fit results

Conclusions

m_H vs. m_A , m_{H^\pm} free



SM Fits

The Fourth
Generation

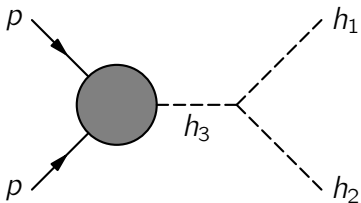
2HDM Fits

Constraints

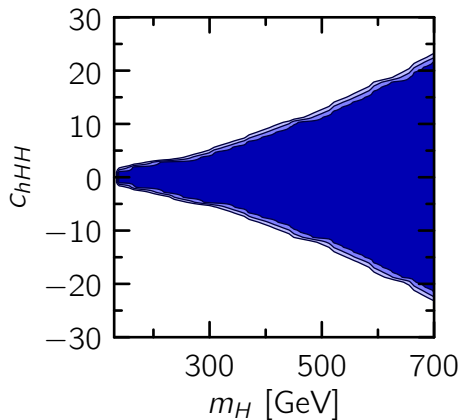
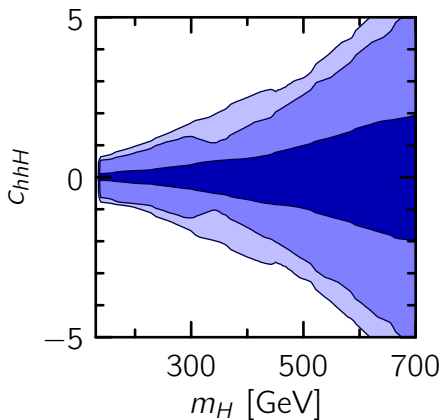
[Fit results](#)

Conclusions

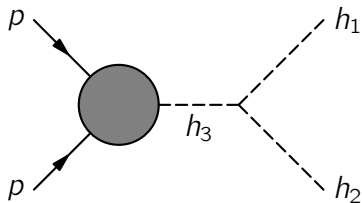
Triple Higgs Couplings



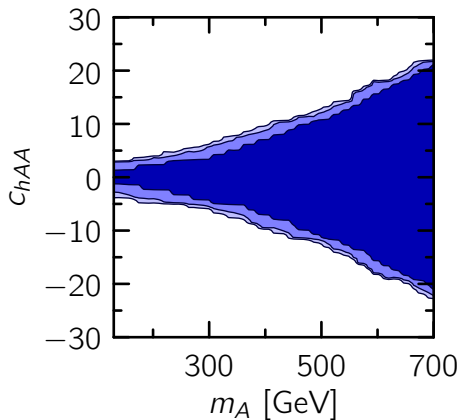
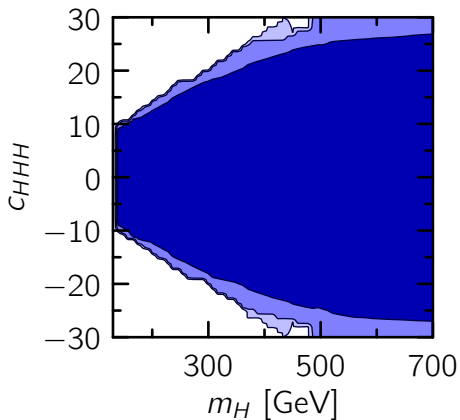
$$C_{h_1 h_2 h_3} = \frac{g_{h_1 h_2 h_3}^{2\text{HDM}}}{g_{HHH}^{\text{SM}}|_{m_H=126 \text{ GeV}}}$$



Triple Higgs Couplings



$$C_{h_1 h_2 h_3} = \frac{g_{h_1 h_2 h_3}^{2\text{HDM}}}{g_{HHH}^{\text{SM}}|_{m_H=126 \text{ GeV}}}$$



Conclusions (SM4)

- SM with a sequential **fourth generation** is **ruled out** by a combination of **Higgs and electroweak precision data**.
- Computation of **p -values** in **non-decoupling models** is **non-trivial** and requires numerical simulations which become unfeasible for small p -values.
- **Importance sampling** techniques as implemented in ***myFitter*** can **speed things up considerably**.

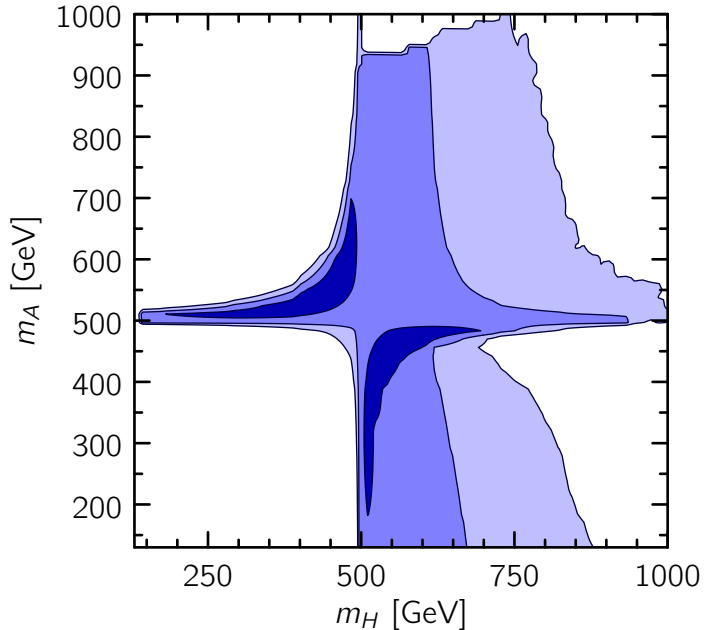
Conclusions (2HDM)

- Best-fit scenario of the type-II 2HDM is the decoupling limit.
- Scenarios with non-SM-like h couplings are allowed by Higgs data but disfavoured (at 1σ) by flavour observables.
- Scenarios with m_H and m_A below 300 GeV are ruled out at 2σ .
- With our inputs, H can still be light and have a factor 10 enhancement of its triple coupling.
- Be careful with purely scan-based analyses. They don't necessarily give you the full picture.

Backup Slides

m_H vs. m_A , $m_{H^\pm} = 500$ GeV fixed

Using S , T , U and R_b instead of EWPOs:



SM Fits

The Fourth
Generation

2HDM Fits

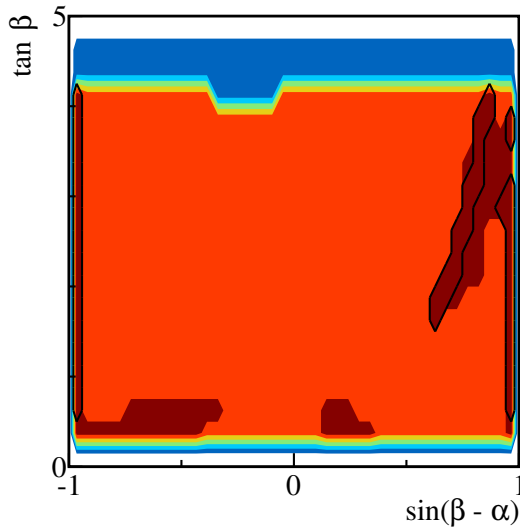
Constraints

Fit results

Conclusions

Scan-based analyses

From [B. Coleppa, F. Kling, and S. Su, arXiv:1305.0002]:



... take them with a grain of salt.

SM Fits

The Fourth
Generation

2HDM Fits

Constraints

Fit results

Conclusions