


e^-e^- Physics for ILC

Y. Kurihara
(KEK)

ILC School @ Toyama

23/July/2013

Outline

- Introduction
 - Why e^-e^- ?
 - What is different from e^+e^- ?
 - Physics of e^-e^- collisions
 - Majorana Neutrino Search
 - s-electron Search/Study
 - And so on ...
 - Summary
- 

Introduction

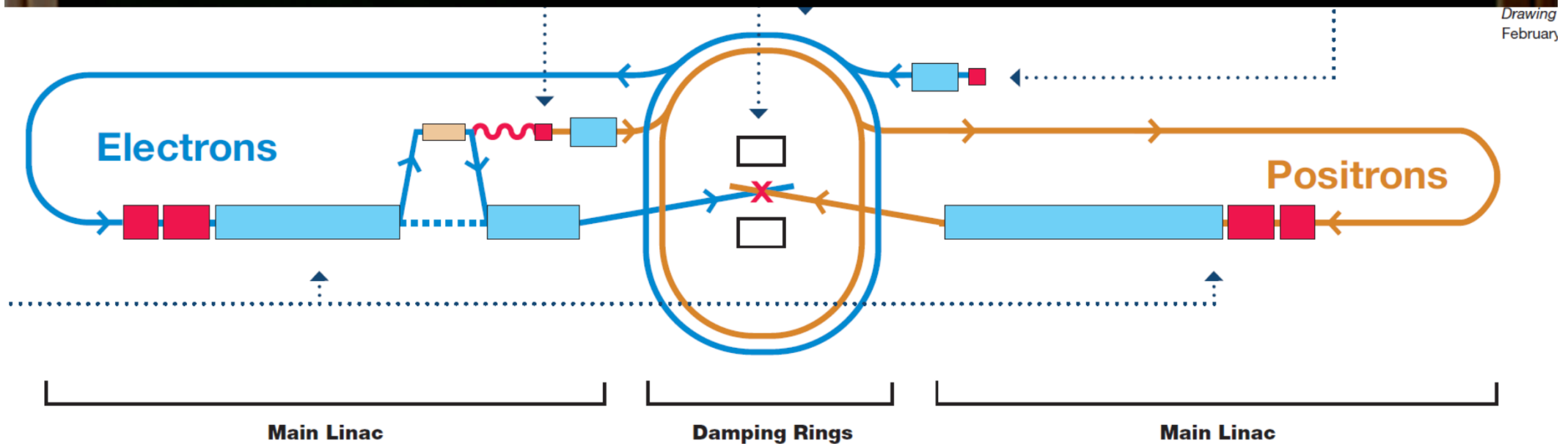


Introduction: Why $e^- e^-$?

- 電子・電子特有の物理があり、電子・陽電子実験と相補的な情報が得られる。
- (Higgsも作れる。)
- 加速器の変更は、ほとんどない。
- (電子源は安くて簡単。)
- 測定器/final focusの変更は、必要ない。

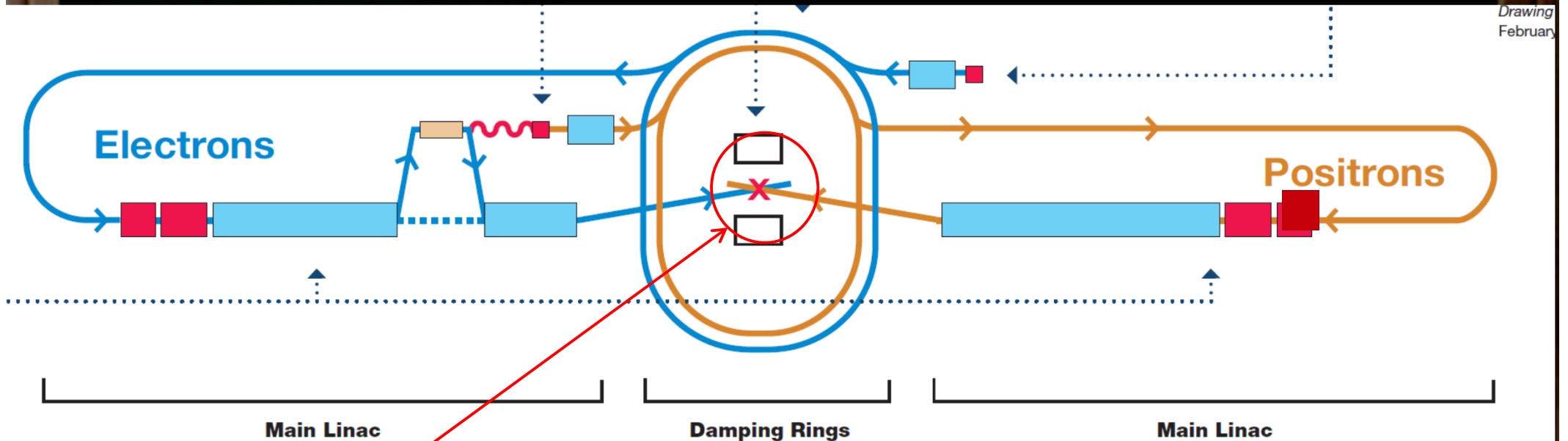
Introduction

• An e^-e^- option in ILC



Introduction

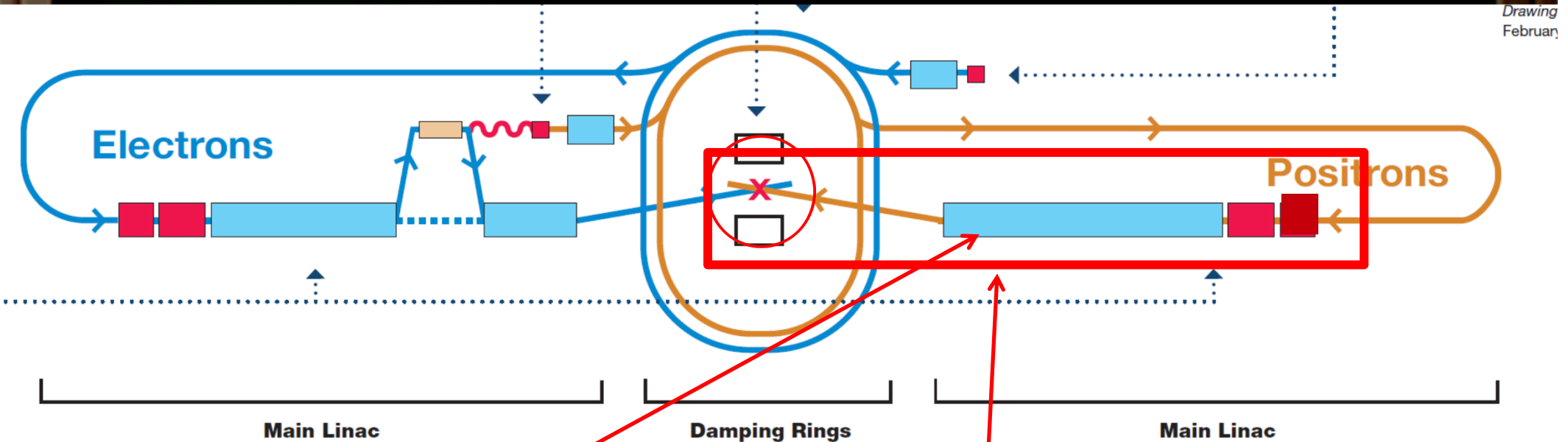
• An e^-e^- option in ILC



Final focus system (crossing angle) is the same as the e^+e^- collider
Detector can be common to the e^+e^- collider

Introduction

• An e^-e^- option in ILC



Acceleration Phase Change

Magnet Polarity Change

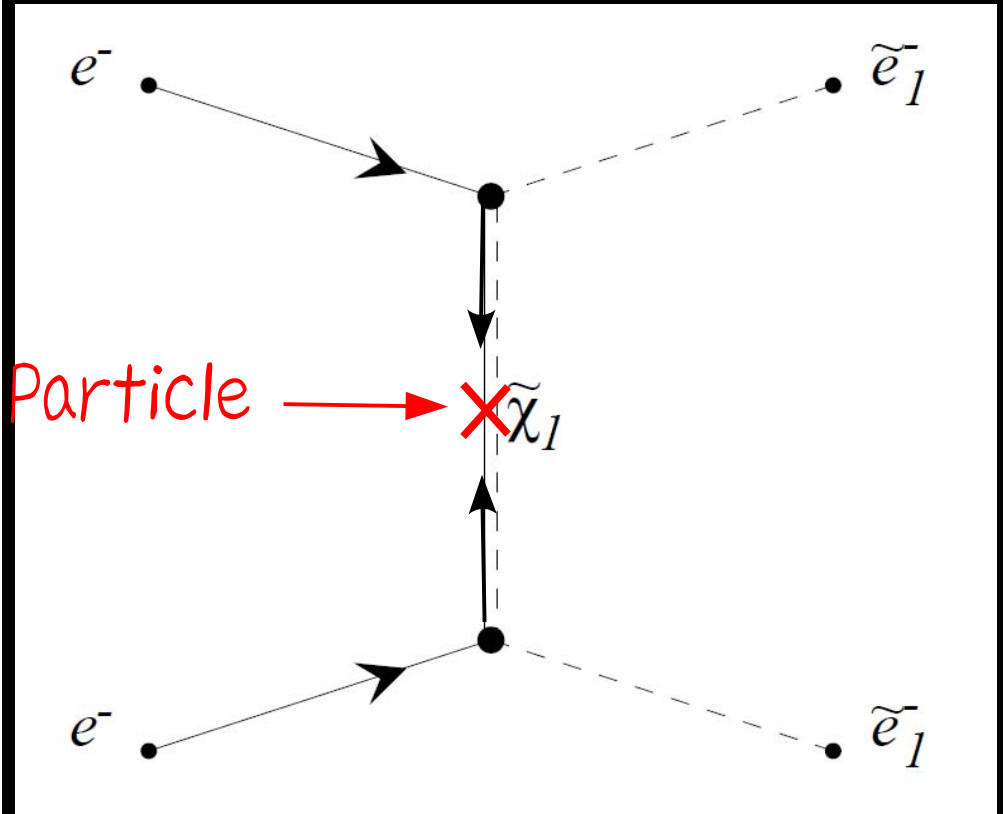
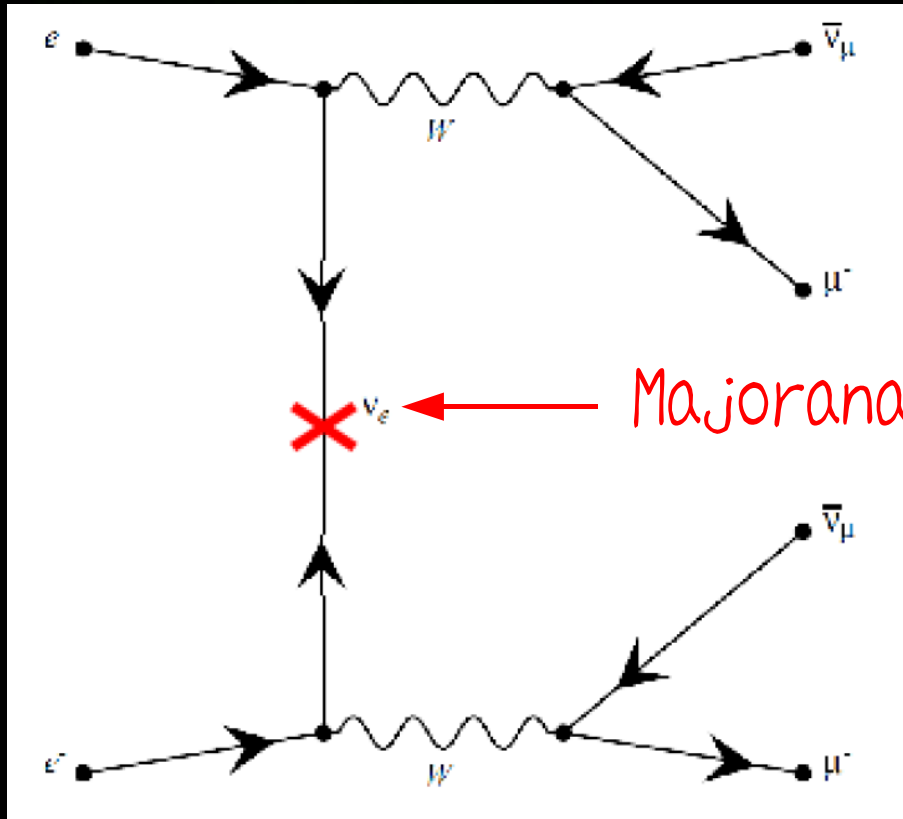
e^-e^-

Physics in ILC



e^-e^- Physics

Majorana Fermion Productions

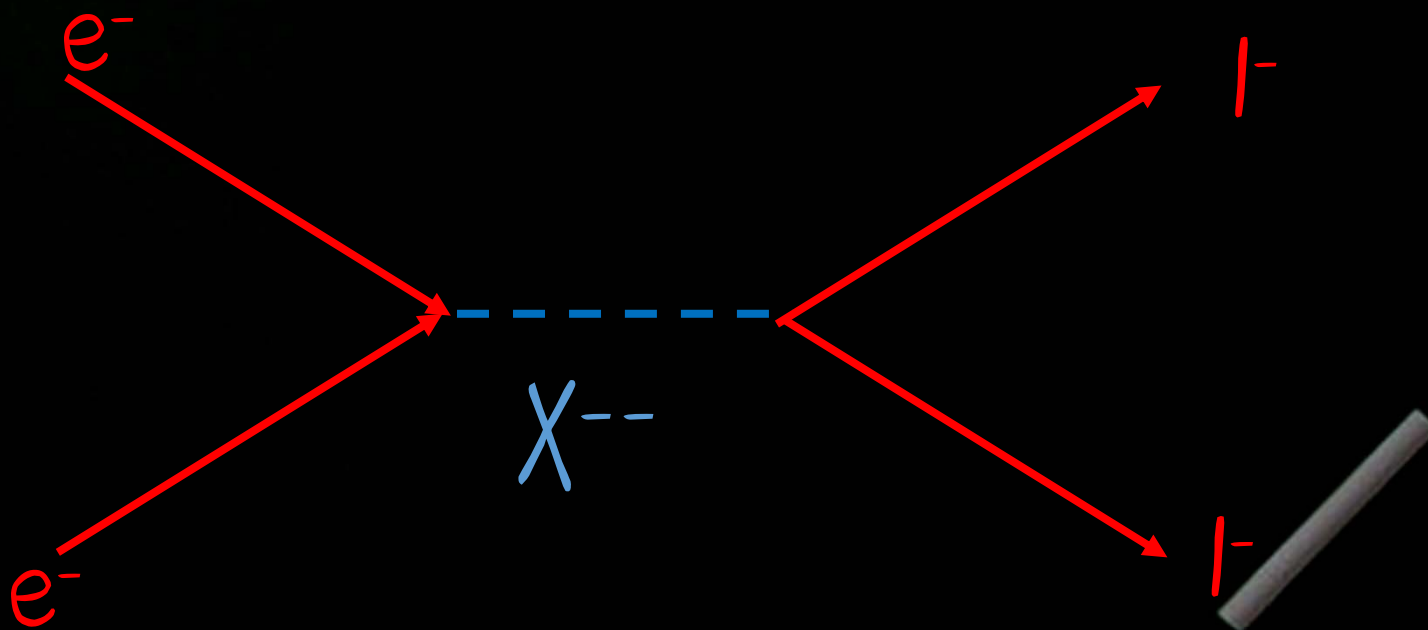


Majorana Neutrino Search

s-electron Search

e^-e^- Physics

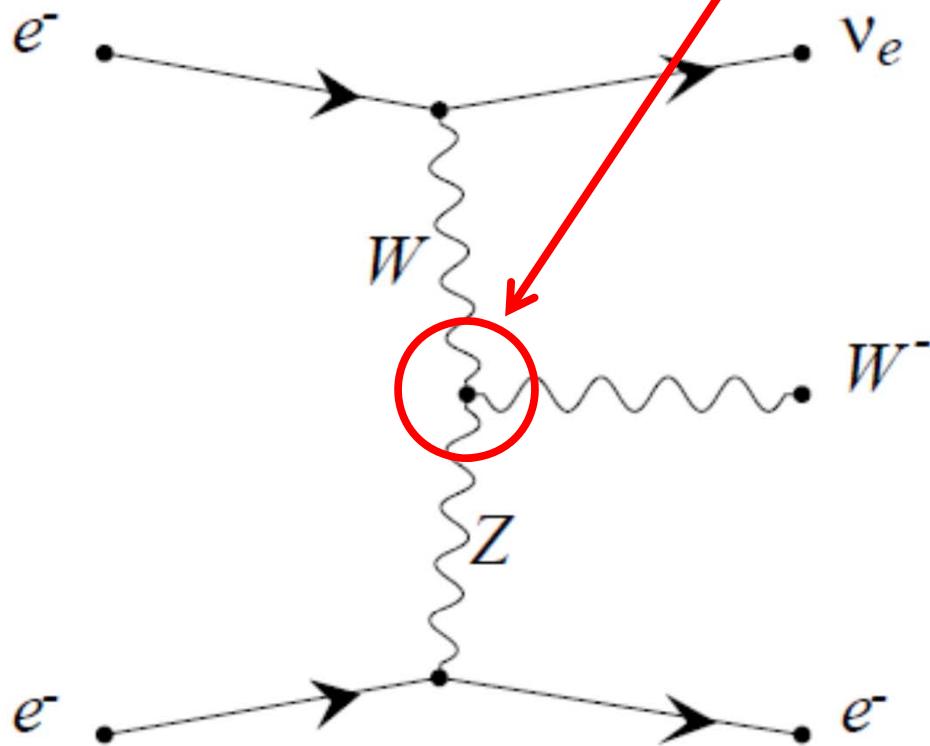
Doubly charged particle



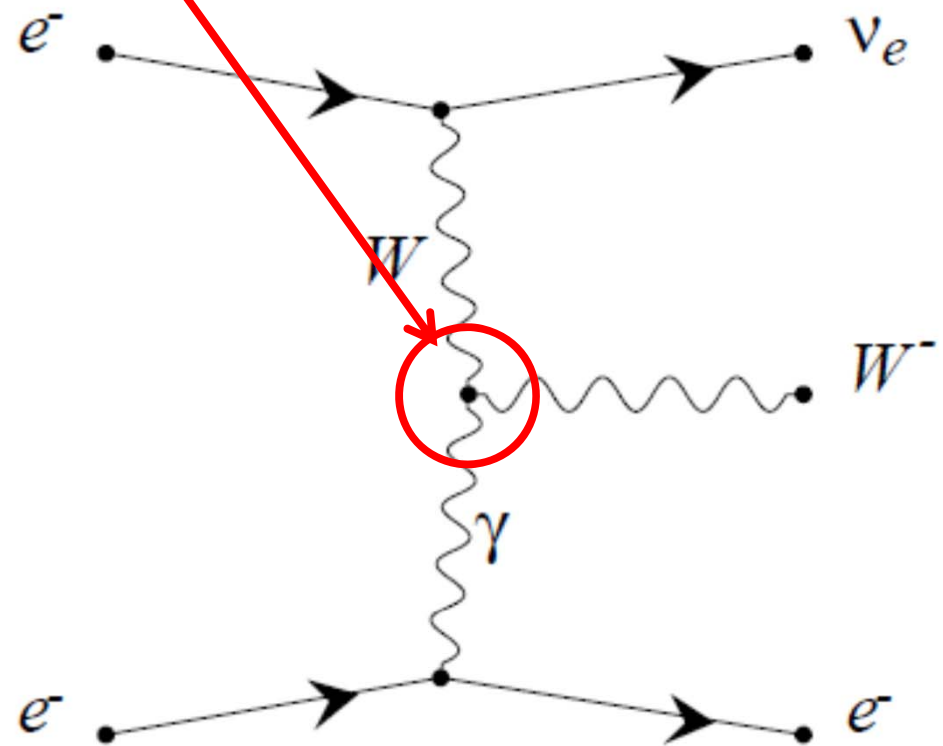
e^-e^- Physics

Gauge boson couplings

Graph 15

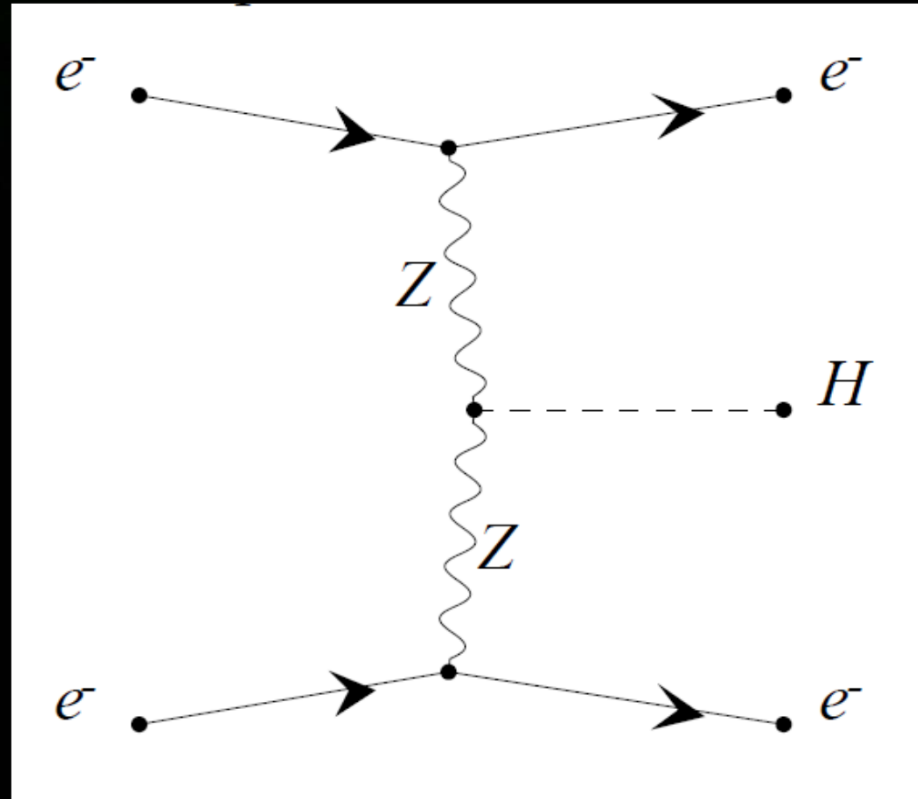


Graph 16




e^-e^- Physics

Higgs Production in e^-e^- collider



e^-e^- Physics:

Majorana ν Search



Majorana Neutrino Search

Heavy Majorana-neutrino search
at future e^-e^- colliders

J. Fujimoto and Y. Kurihara

November 14, 2011

Majorana Neutrino Search

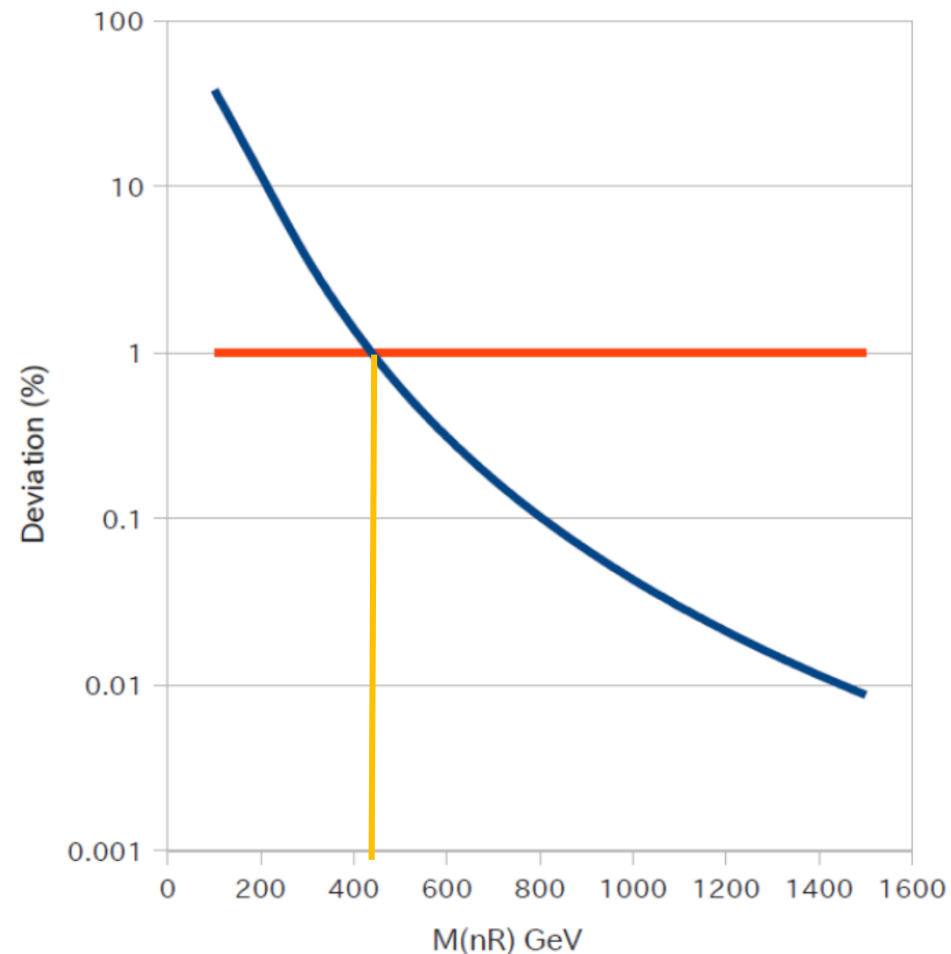
1 Model

- $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$
- Gauge boson masses on $SU(2)_L \times SU(2)_R$ are degenerated: $m_{W_L^\pm} = m_{W_R^\pm}$.
- Neutrinos are Majorana particles.
- A mass of a left-handed neutrino is light: $m_{\nu_L} \ll \mathcal{O}(\text{eV})$.
- A mass of a right-handed neutrino is heavy: $m_{\nu_R} \geq \mathcal{O}(\text{weak scale})$.
- Right-handed gauge-bosons, W_R^\pm , do not contribute to low energy experiments.

Majorana Neutrino Search

2 Experiment

- Heavy neutrino mass
- An electron-neutrino
- Total cross section of 1% level. A ν_{eR} mass diagram is less than allowed experiment



ts.

experiments with accurac
s mass is not heavy enough
 ν_{eR} in a t -channel exchange
t mass above 500 GeV is ti

Majorana Neutrino Search

2 Experimental limit for ν_R mass

- Heavy neutrino masses must be greater than $m_Z/2$ from LEP-I experiments.
- An electron-neutrino mass has additional constraint from LEP-II experiments.
- Total cross sections of a W -pair production have been measured by LEP-II experiments with accuracy of 1% level. A ν_{eR} can contribute the W -pair production cross-sections, if its mass is not heavy enough. When the ν_{eR} -mass is heavier than about 500 GeV, a contribution from ν_{eR} in a t -channel exchange diagram is less than 1% as shown in Figure 1. Then we assume a ν_{eR} with a mass above 500 GeV is allowed experimentally.

Majorana Neutrino Search

3 Parameters

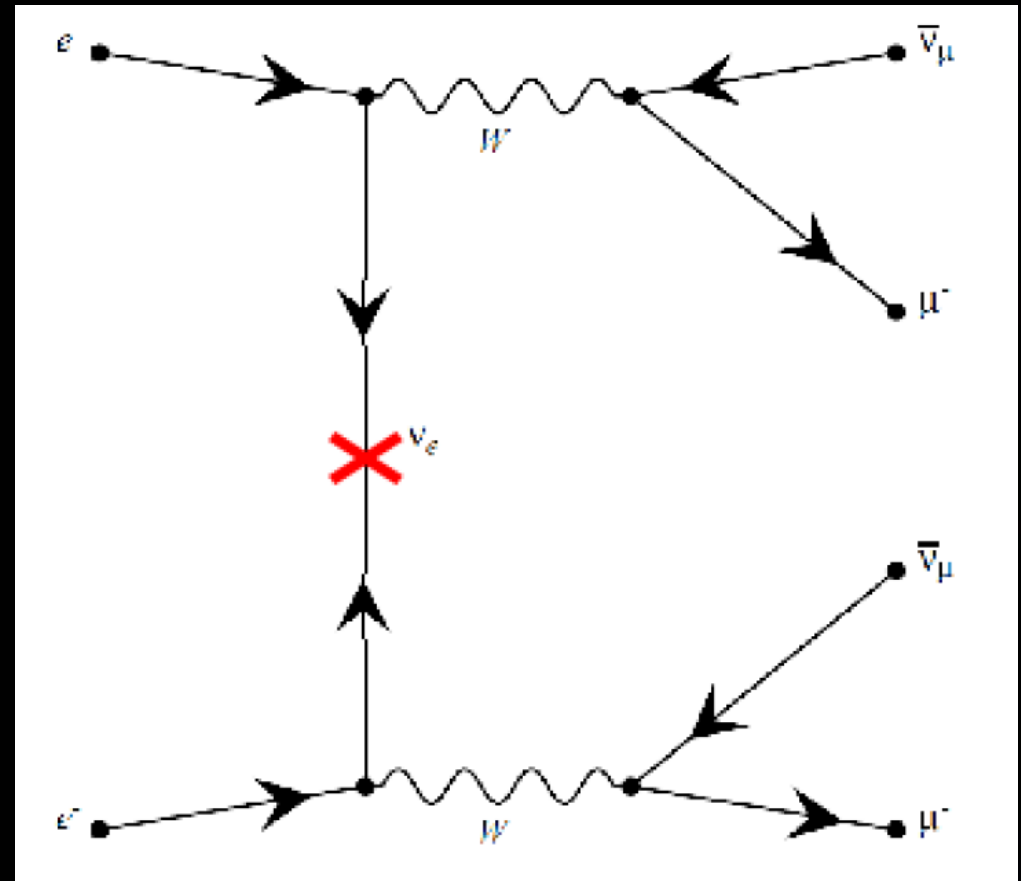
- Accelerator: a future e^-e^- linear collider experiment.
- $\alpha=1/137.036$
- $m_W=80.22$ GeV
- $m_Z=91.187$ GeV
- $\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$
- $m_e = 0.511 \times 10^{-3}$ GeV

Majorana Neutrino Search

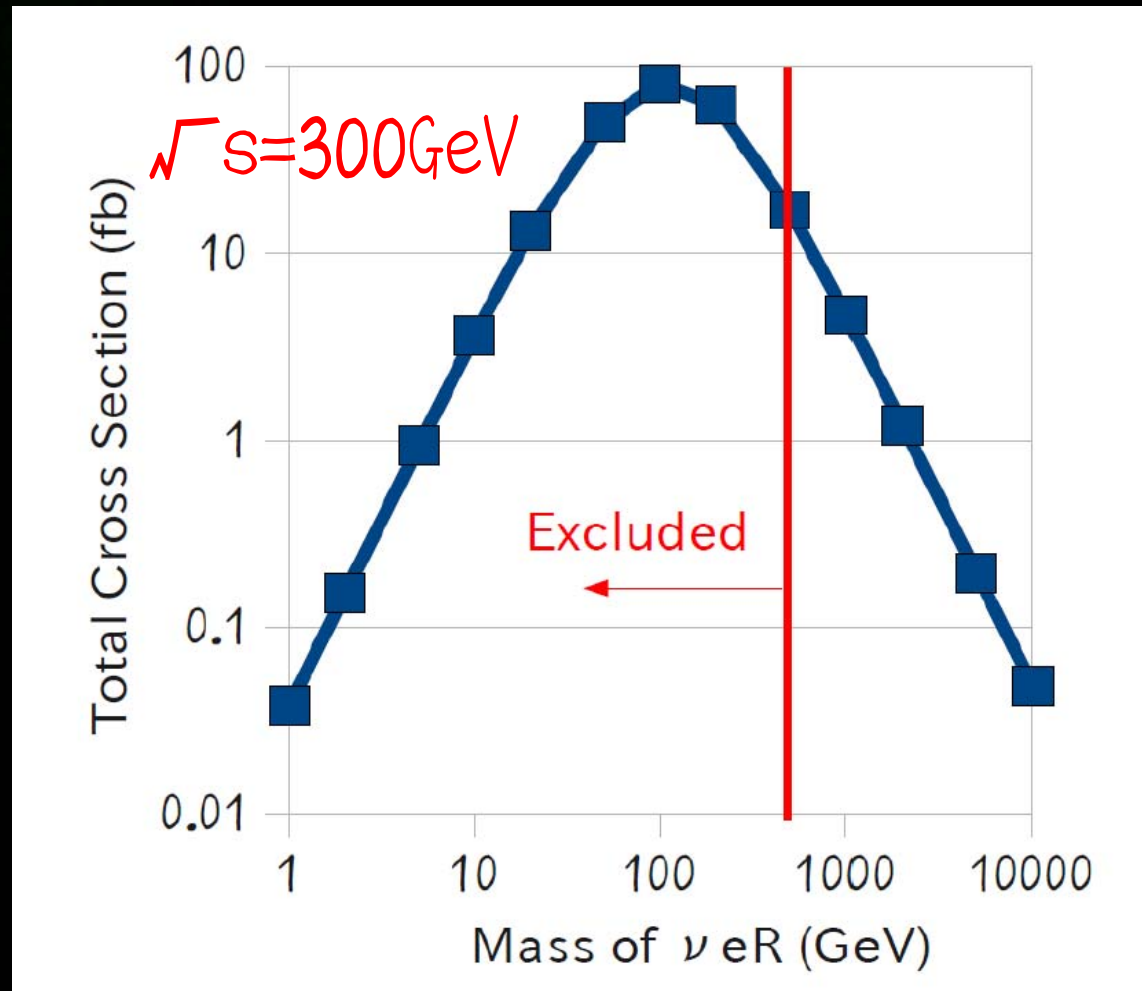
Signal:

$\mu^- \mu^-$ pair + missing

B.G.: $\nu_e \nu_e W^- W^-$

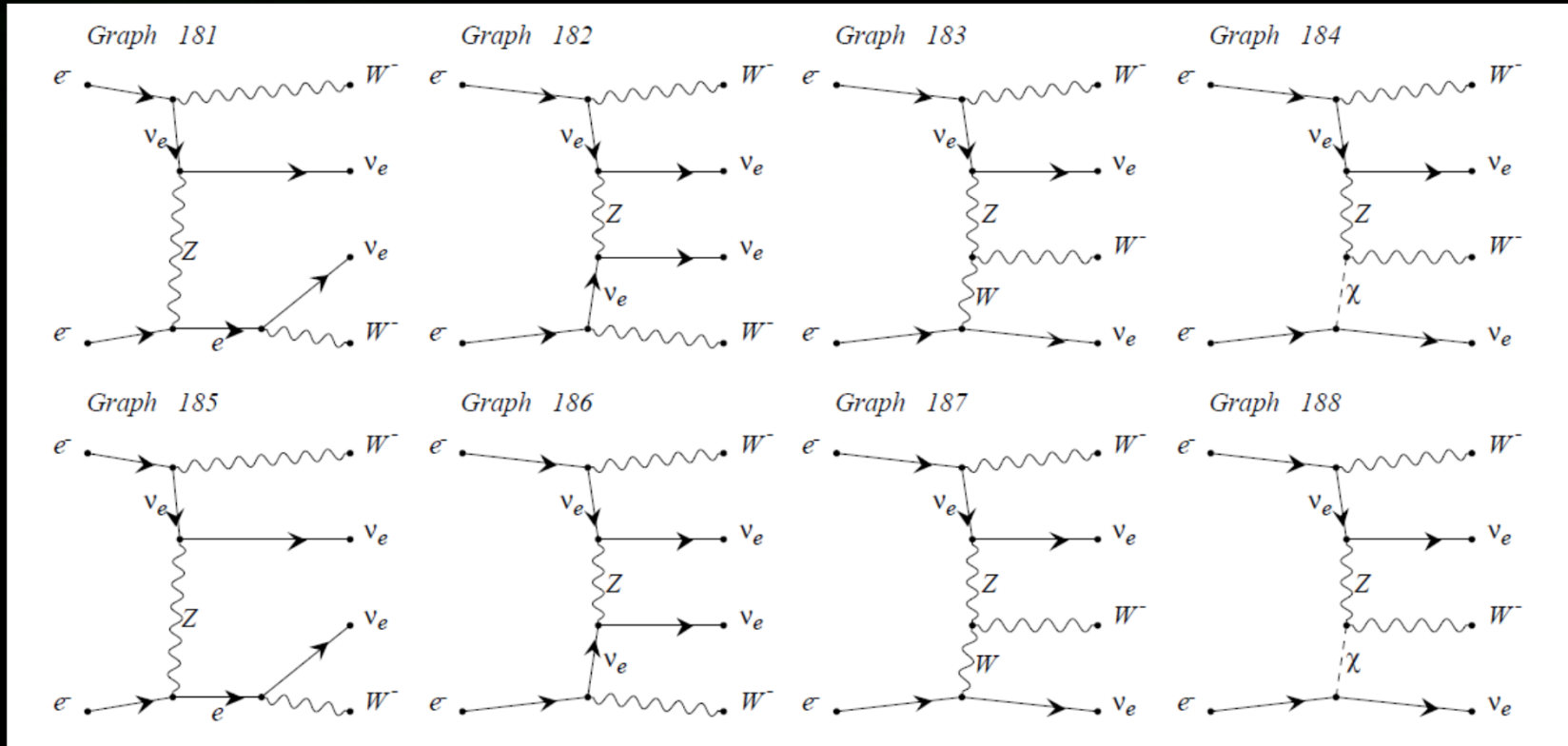


Majorana Neutrino Search



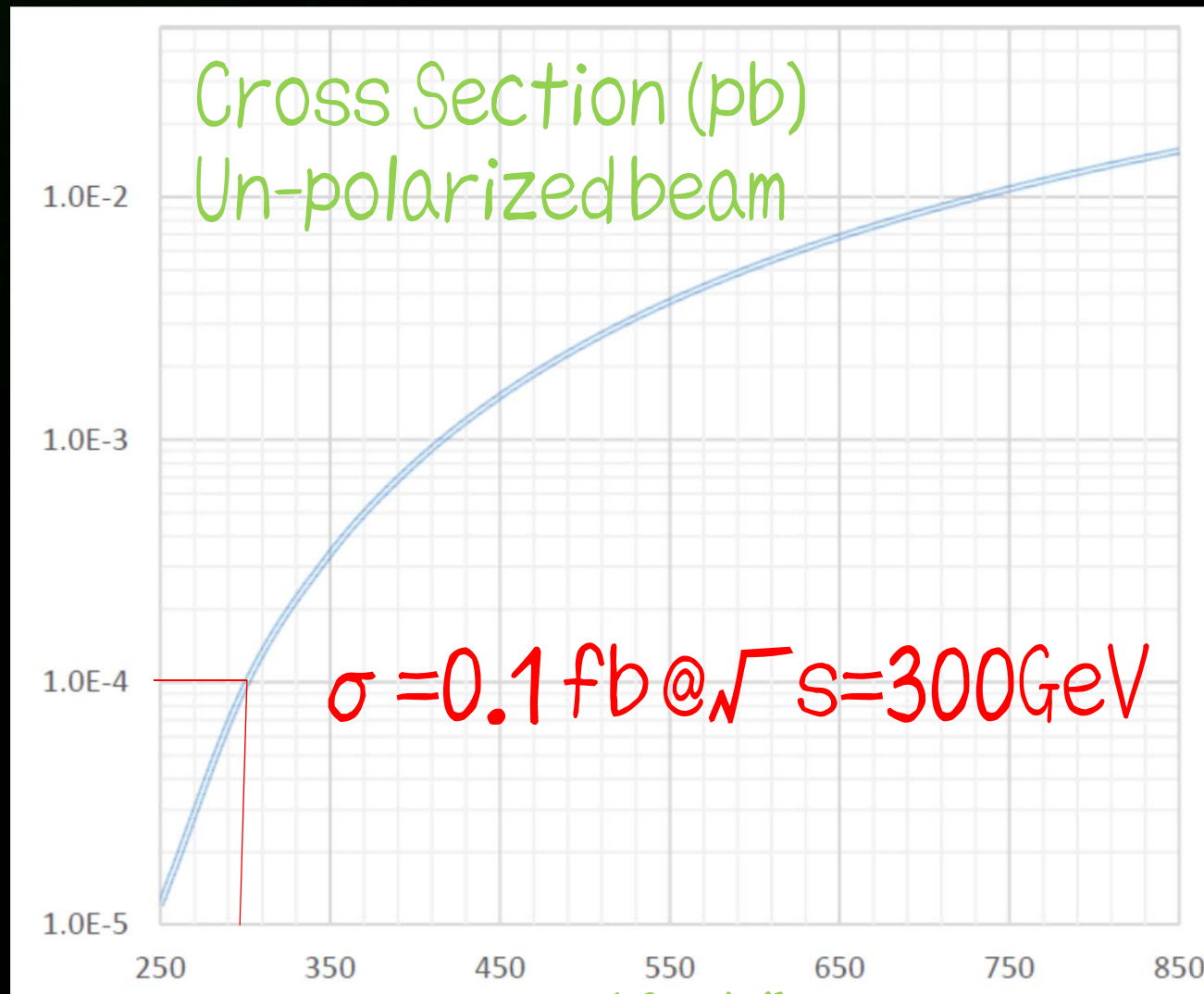
W-W-Scattering

Background



#Diagrams: 286

W-W-Scattering



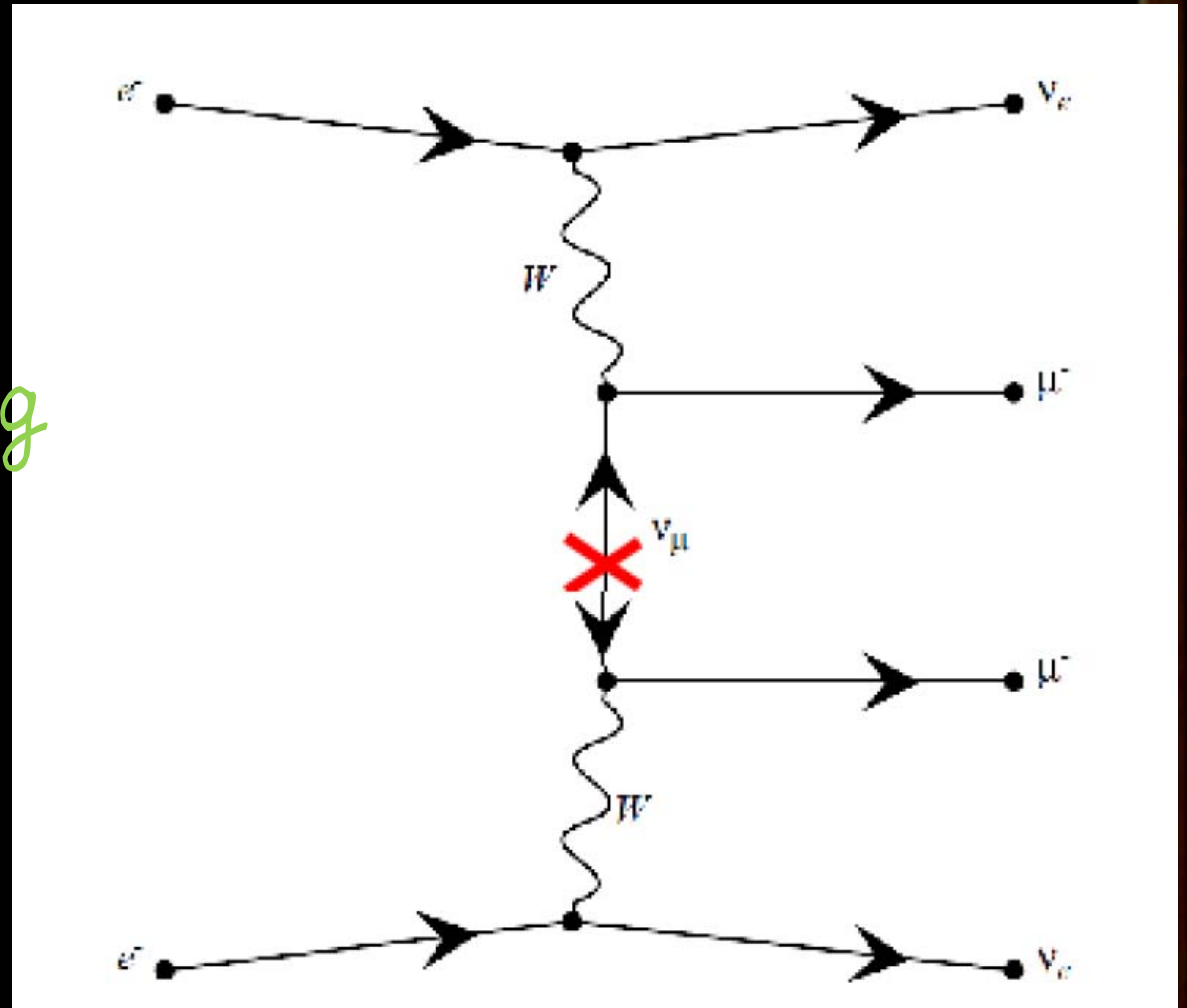
\sqrt{s} (GeV)

Majorana Neutrino Search

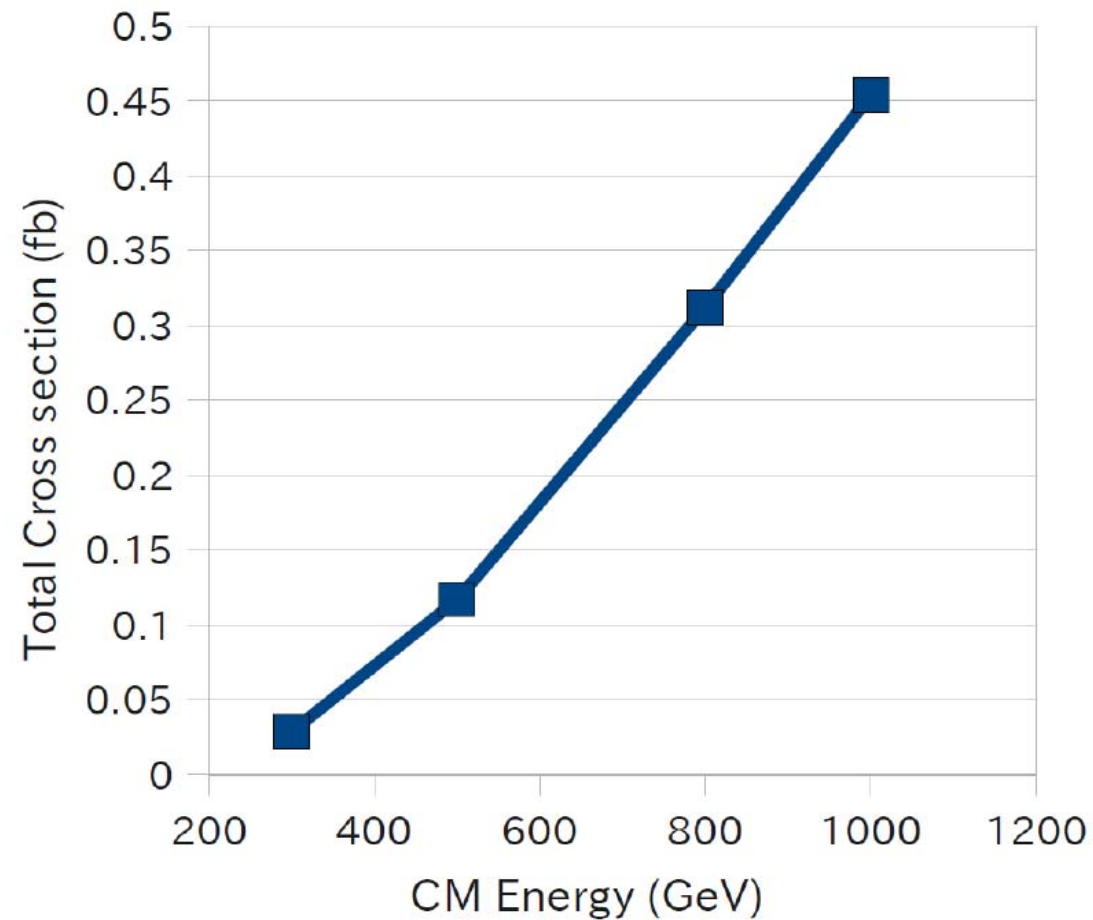
Signal:

$\mu^- \mu^-$ pair + missing

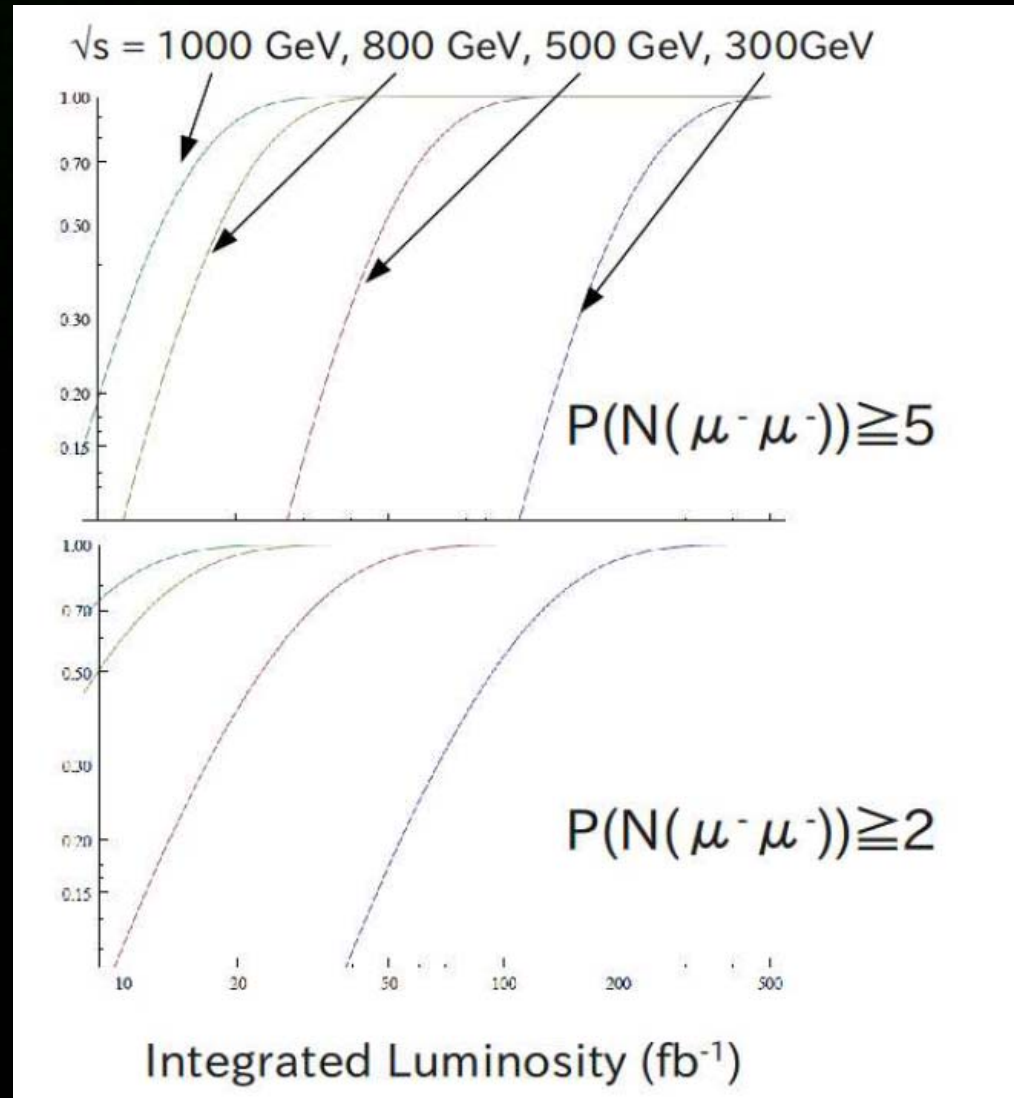
B.G.: $\nu_e \nu_e W^- W^-$



Majorana Neutrino Search



Majorana Neutrino Search



Majorana Neutrino Search

PHYSICAL REVIEW D **81**, 114001 (2010)

Inverse neutrinoless double beta decay revisited: Neutrinos, Higgs triplets, and a muon collider

Werner Rodejohann*

Max-Planck-Institut für Kernphysik, Postfach 103980, D-69029 Heidelberg, Germany

(Received 1 April 2010; published 1 June 2010)

We revisit the process of inverse neutrinoless double beta decay ($e^-e^- \rightarrow W^-W^-$) at future linear colliders. The cases of Majorana neutrino and Higgs triplet exchange are considered. We also discuss the processes $e^-\mu^- \rightarrow W^-W^-$ and $\mu^-\mu^- \rightarrow W^-W^-$, which are motivated by the possibility of muon colliders. For heavy neutrino exchange, we show that masses up to 10^6 (10^5) GeV could be probed for ee and $e\mu$ machines, respectively. The stringent limits for mixing of heavy neutrinos with muons render $\mu^-\mu^- \rightarrow W^-W^-$ less promising, even though this process is not constrained by limits from neutrinoless double beta decay. If Higgs triplets are responsible for inverse neutrinoless double beta decay, observable signals are only possible if a very narrow resonance is met. We also consider unitarity aspects of the process in case both Higgs triplets and neutrinos are exchanged. An exact seesaw relation connecting low energy data with heavy neutrino and triplet parameters is found.

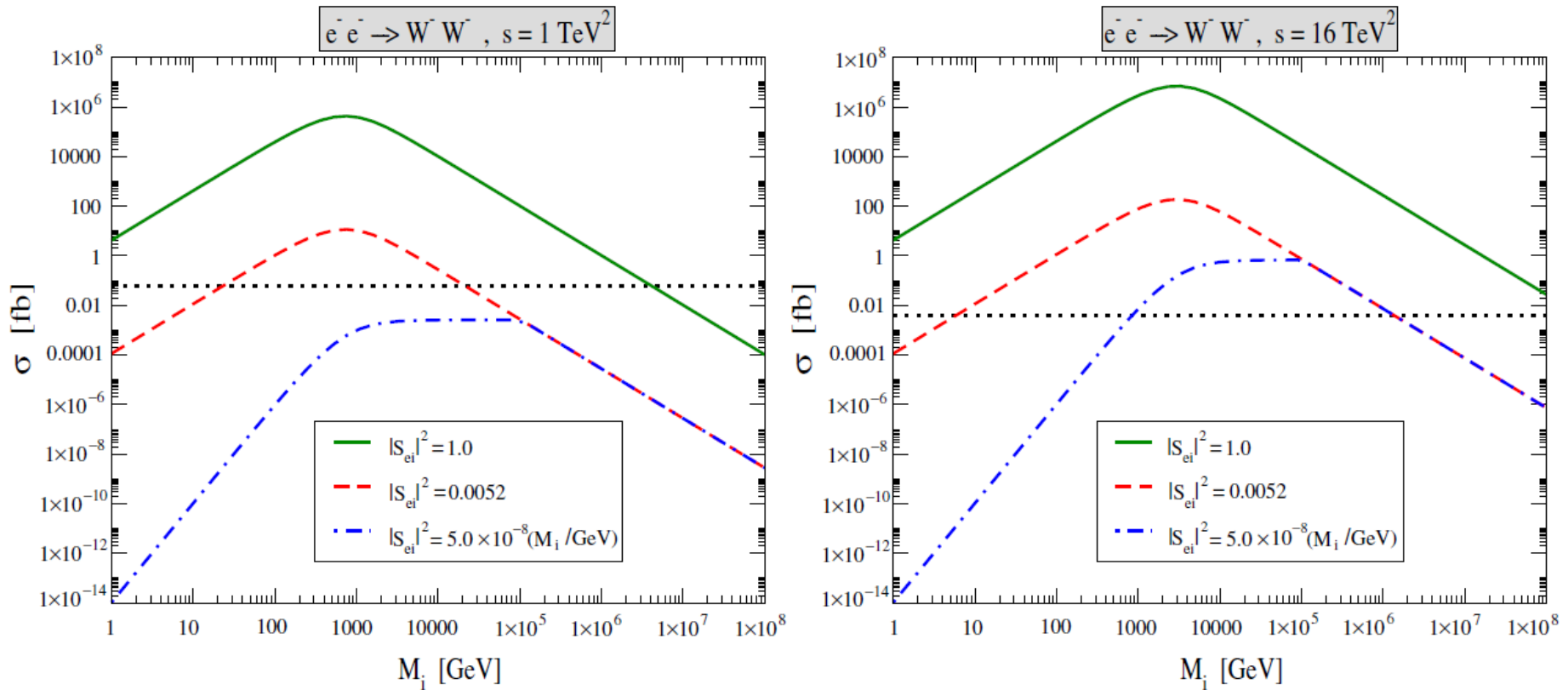
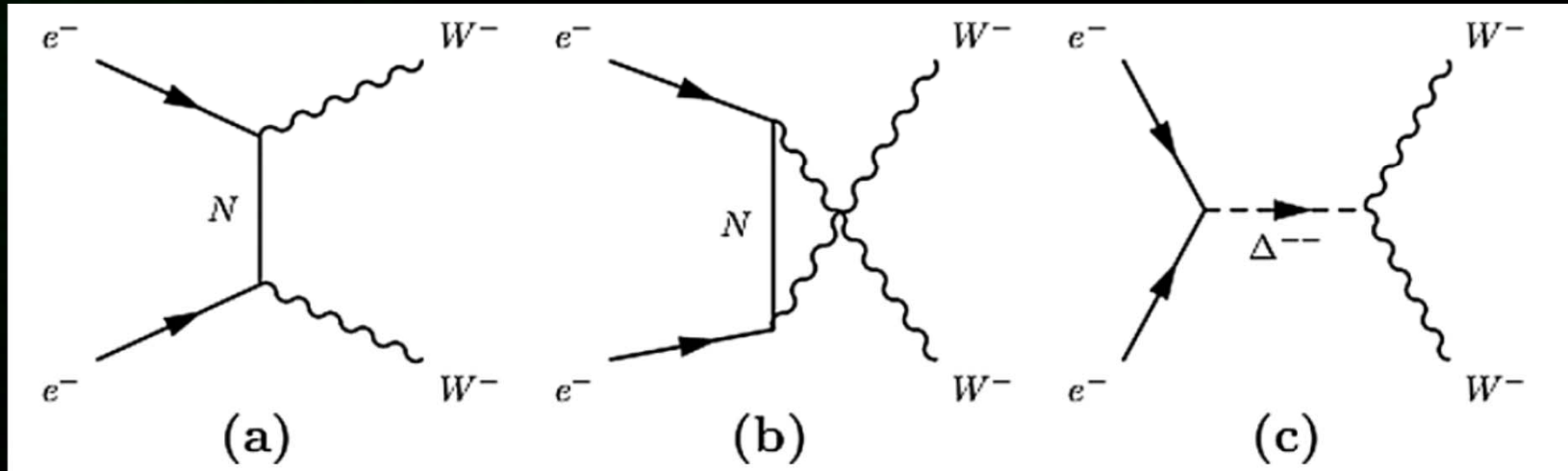


FIG. 2 (color online). Cross section for $e^-e^- \rightarrow W^-W^-$ with $\sqrt{s} = 1$ TeV (left panel) and $\sqrt{s} = 4$ TeV (right panel) and three limits for the mixing parameter $|S_{ei}|^2$. The dotted line corresponds to five events for an assumed luminosity of 80 (s/TeV^2) fb^{-1} .



Physics Letters B

Volume 689, Issue 1, 17 May 2010, Pages 28–35



Probing the Majorana nature of TeV-scale radiative seesaw models at collider experiments

Mayumi Aoki^a, , Shinya Kanemura^b, , 

^a Department of Physics, Tohoku University, Aramaki, Aoba, Sendai, Miyagi 980-8578, Japan

^b Department of Physics, University of Toyama, 3190 Gofuku, Toyama 930-8555, Japan

<http://dx.doi.org/10.1016/j.physletb.2010.04.024>, [How to Cite or Link Using DOI](#)

[Permissions & Reprints](#)

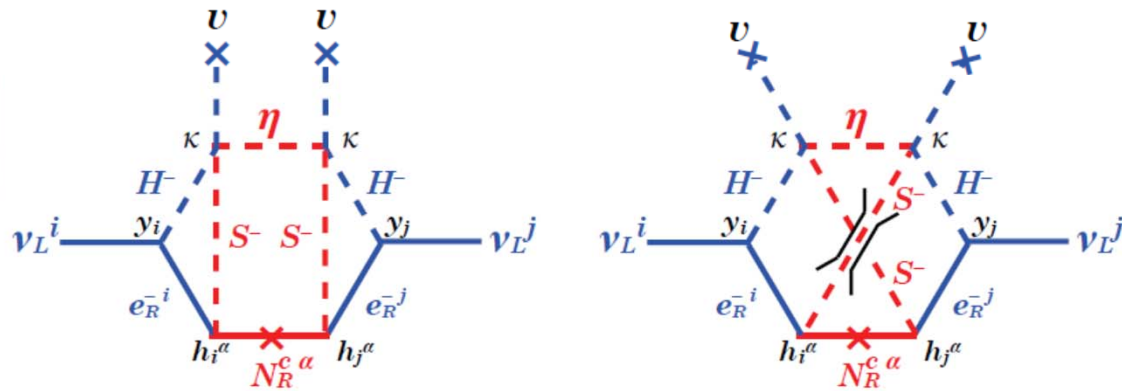
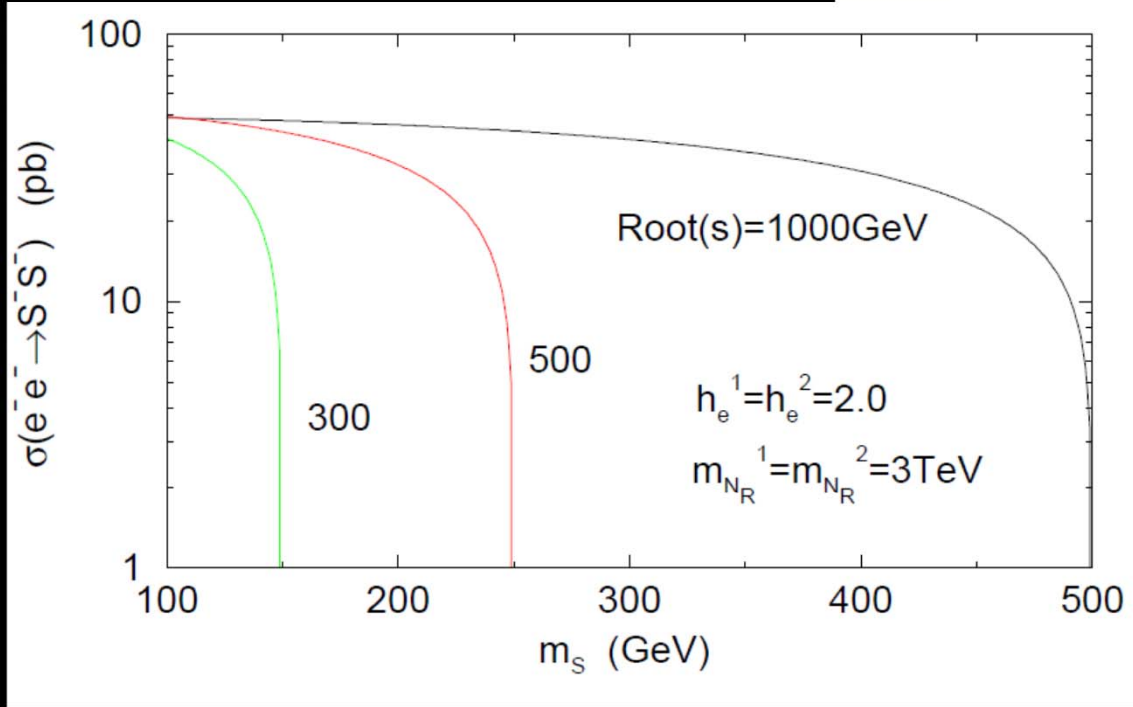
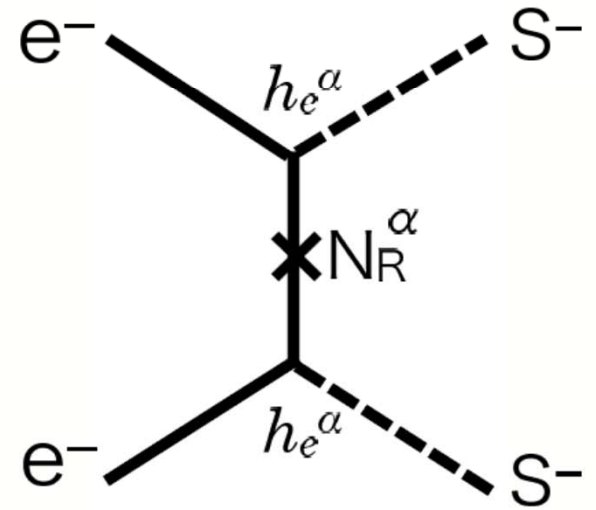
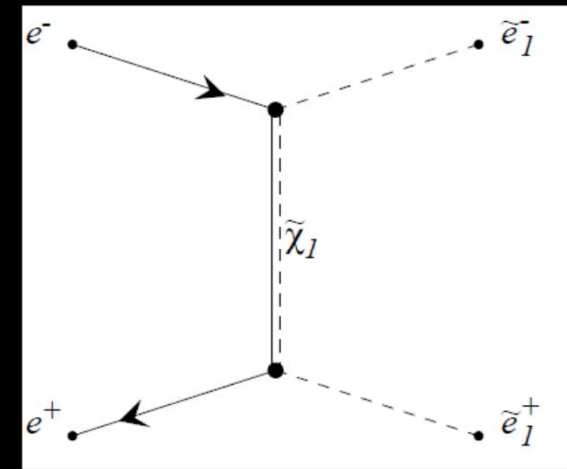
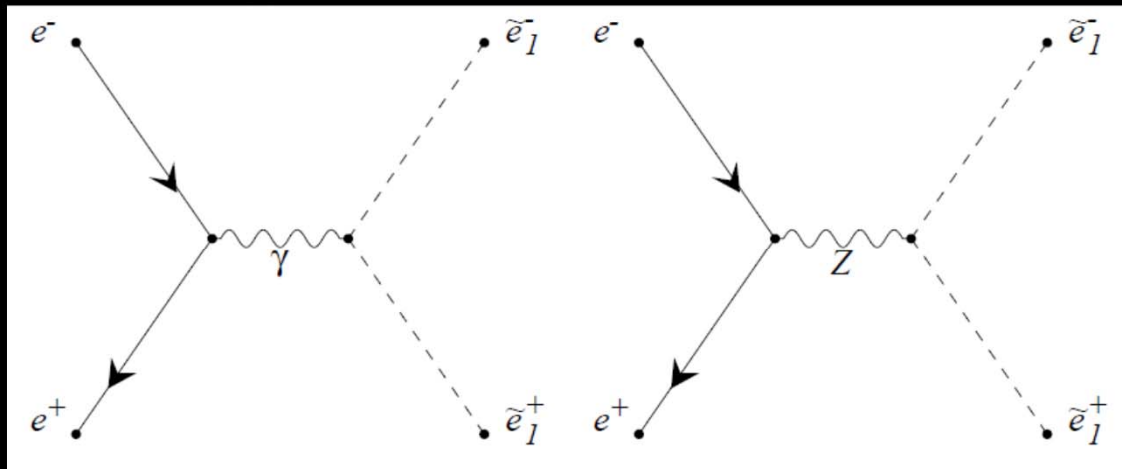
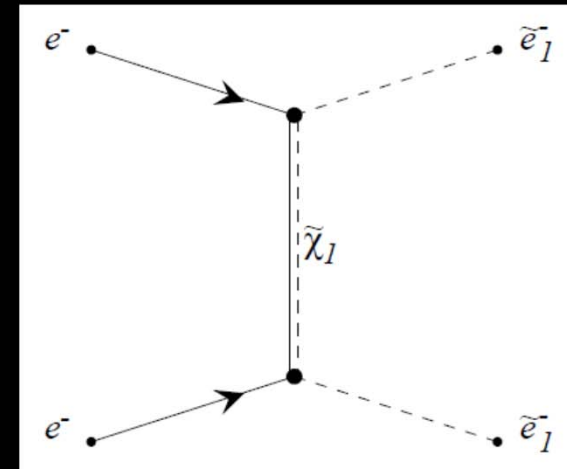


Figure 1: The diagrams for generating tiny neutrino masses.



s-electron Search/Study

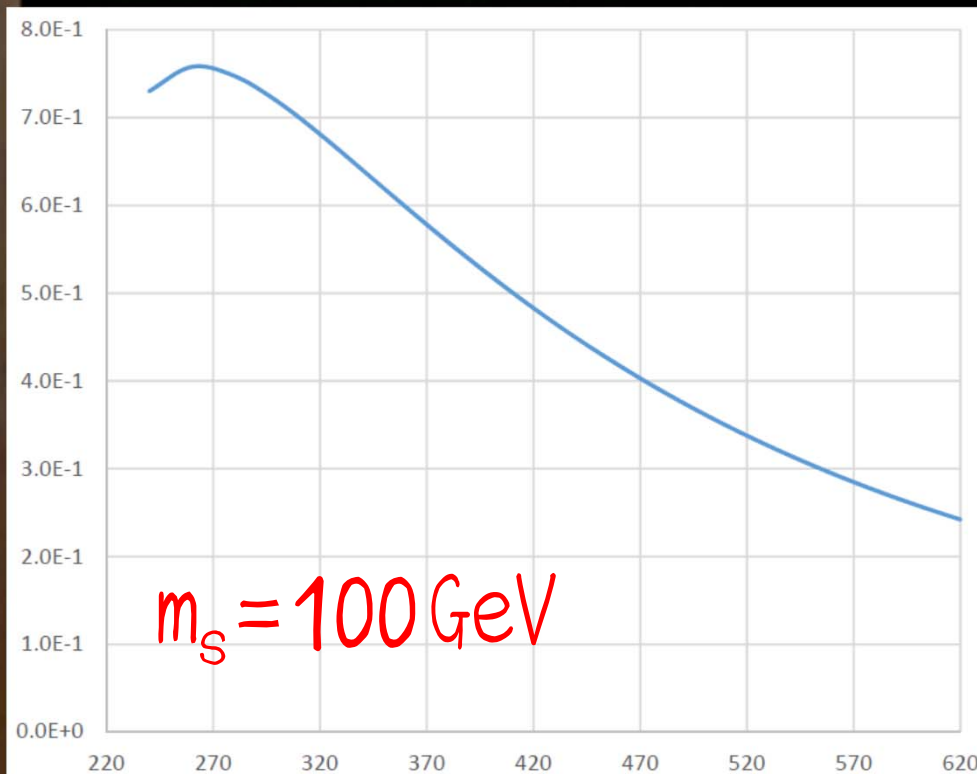
$$e^-e^\pm \rightarrow \tilde{e}_{1,2}^-\tilde{e}_{1,2}^\pm$$



s-electron Search/Study

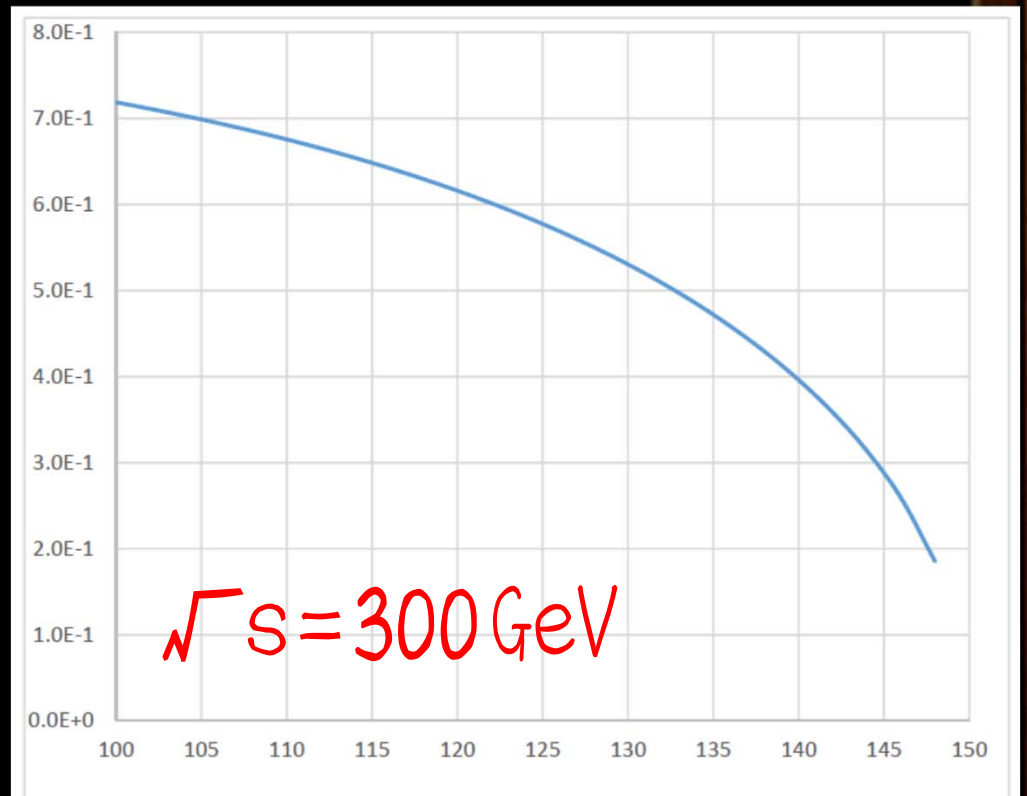
Cross Section (pb)
Un-polarized beam

$m_{\chi_1} = 145.46 \text{ GeV}$
 $m_{\chi_2} = 294.09 \text{ GeV}$
 $m_{\chi_3} = 605.15 \text{ GeV}$
 $m_{\chi_4} = 611.65 \text{ GeV}$



$m_s = 100 \text{ GeV}$

\sqrt{s} (GeV)



$\sqrt{s} = 300 \text{ GeV}$

m_s (GeV)

s-electron Search/Study

Beam Polarization

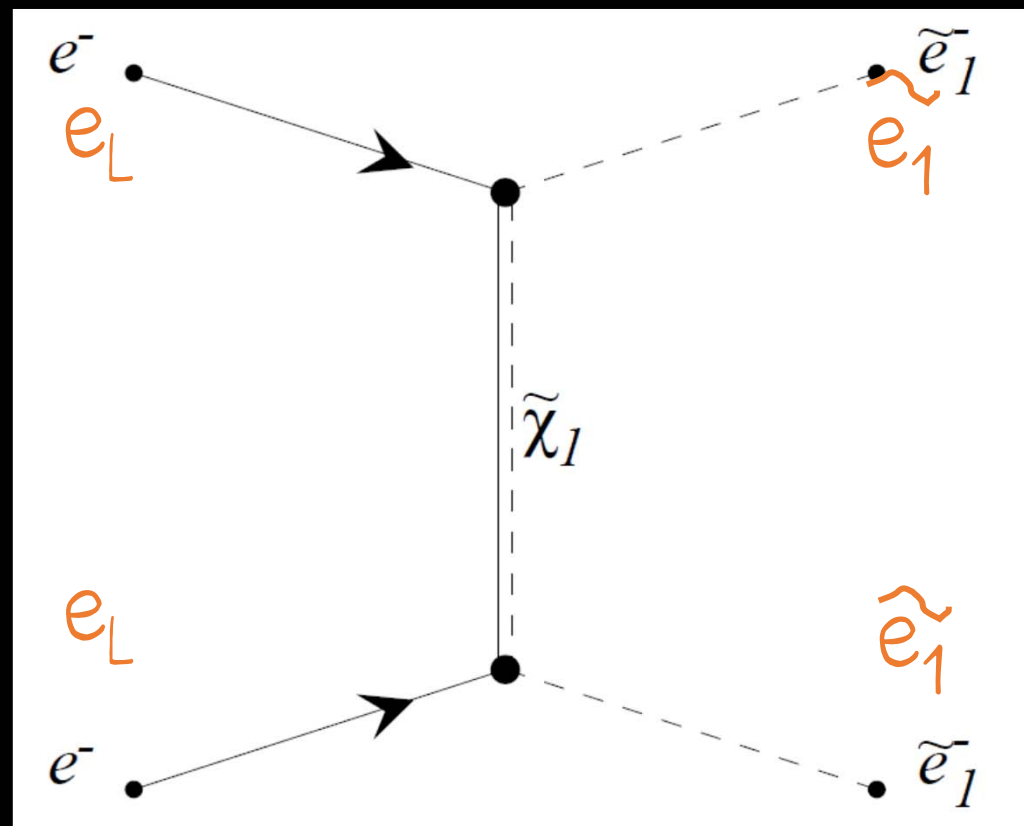
$e^- : 80\%$ $e^+ : 30\%$

Small Mixing

$$e_L^\pm \rightarrow \tilde{e}_1^\pm$$

$$e_R^\pm \rightarrow \tilde{e}_2^\pm$$

$$m_1 \doteq m_2$$



s-electron Search/Study

Beam Polarization

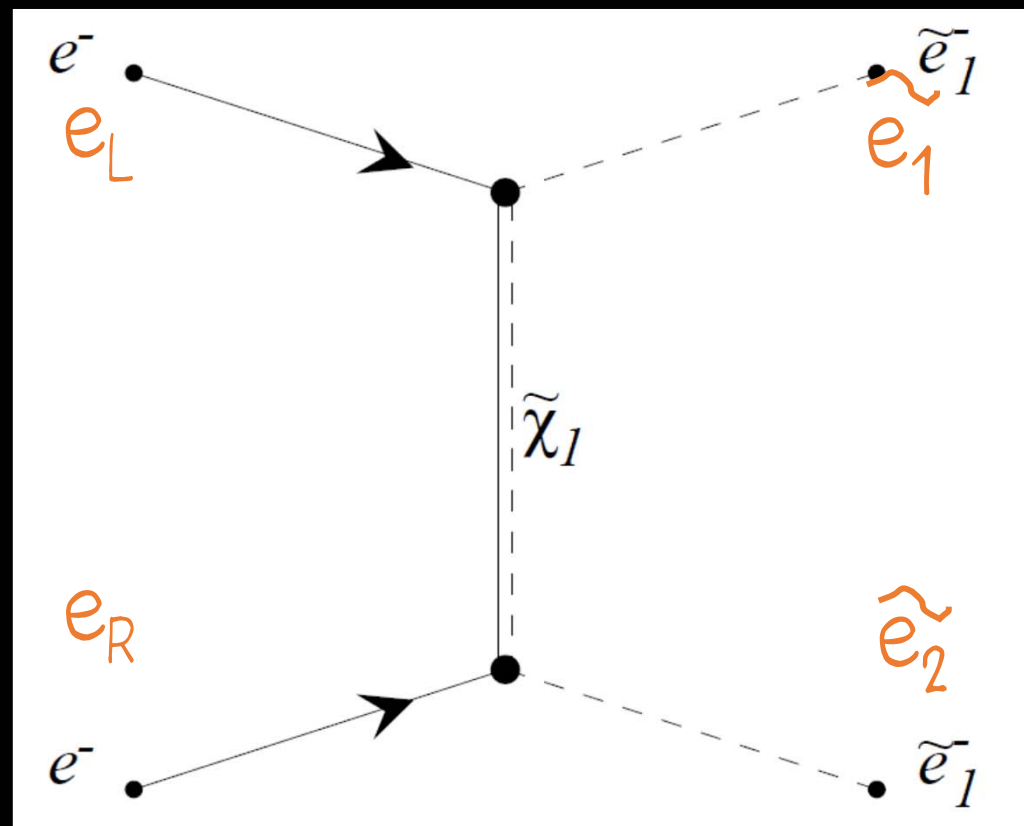
e^- : 80% e^+ : 30%

Small Mixing

$$e_L^\pm \rightarrow \tilde{e}_1^\pm$$

$$e_R^\pm \rightarrow \tilde{e}_2^\pm$$

$$m_1 \doteq m_2$$



s-electron Search/Study

Beam Polarization

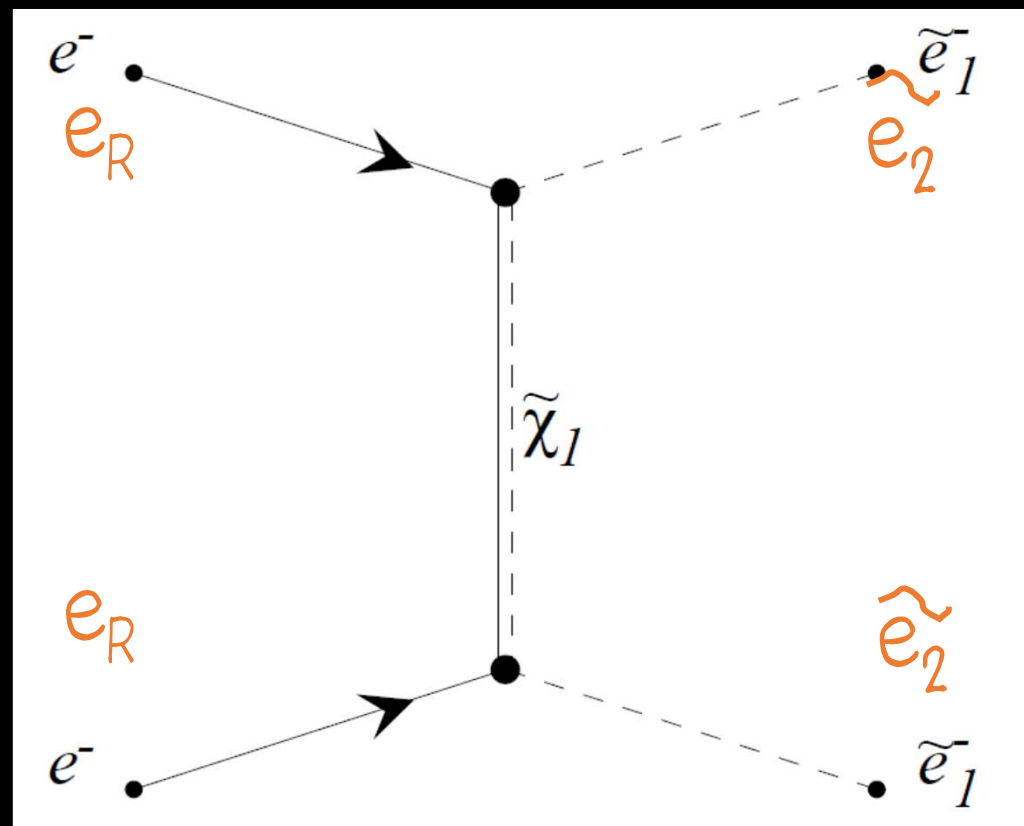
$e^- : 80\%$ $e^+ : 30\%$

Small Mixing

$$e_L^\pm \rightarrow \tilde{e}_1^\pm$$

$$e_R^\pm \rightarrow \tilde{e}_2^\pm$$

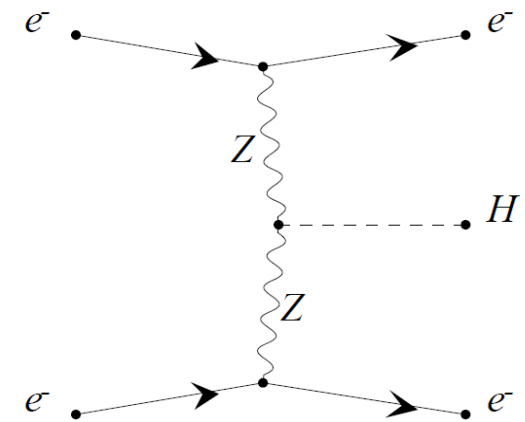
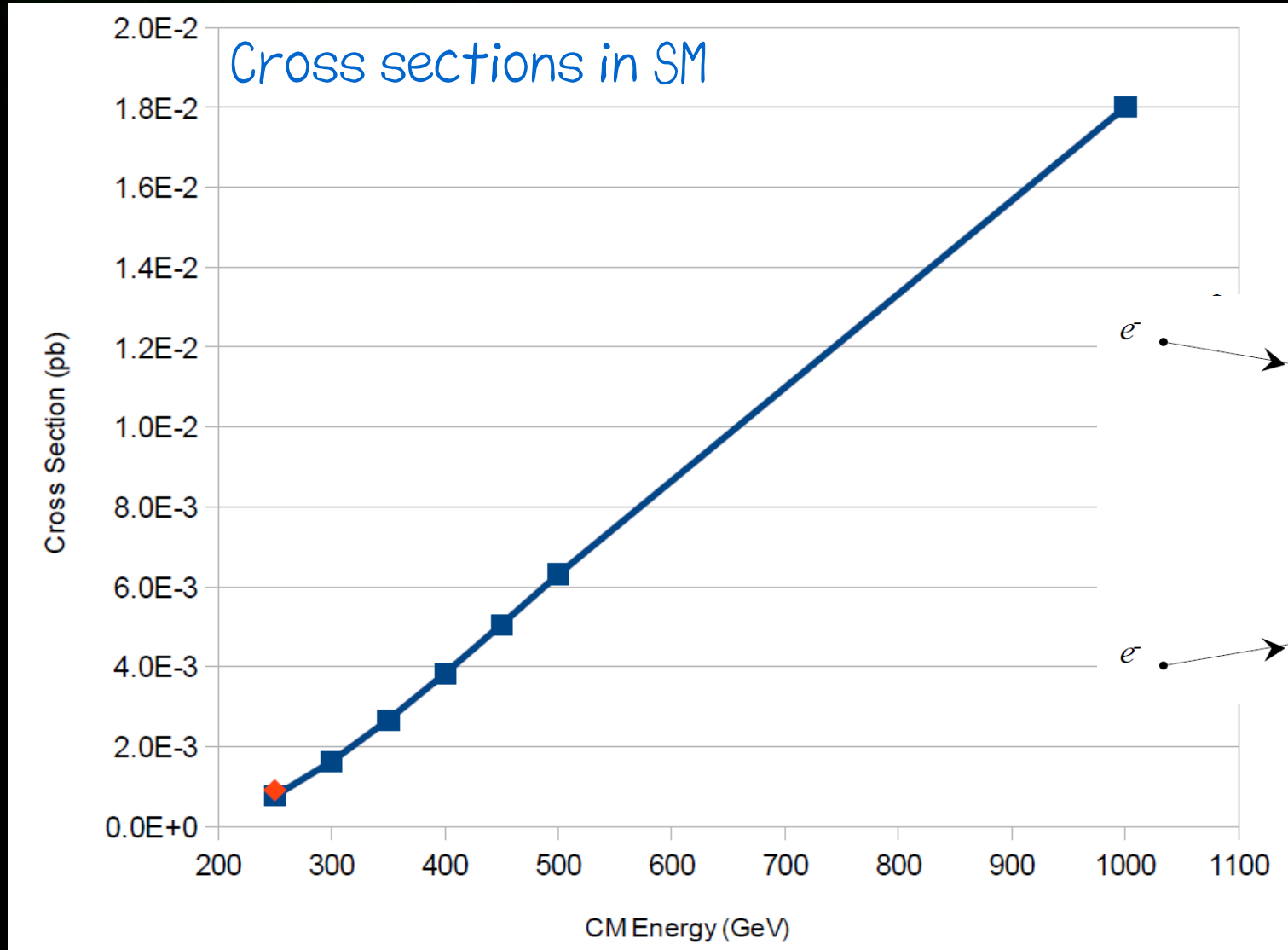
$$m_1 \doteq m_2$$



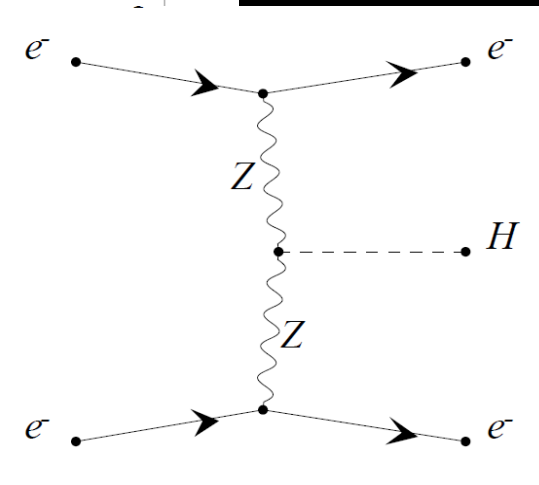
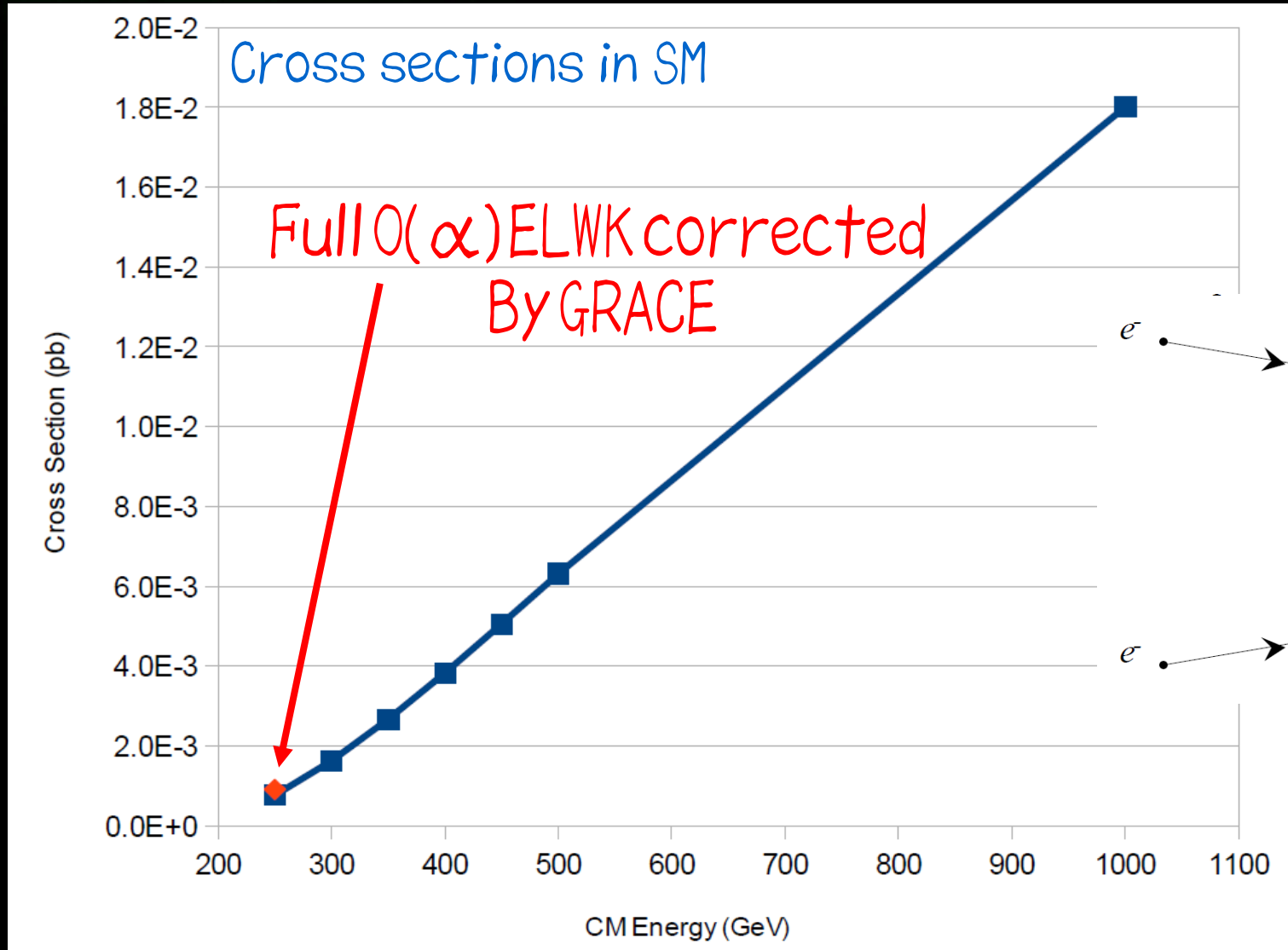
e^-e^- Physics:
Higgs Production



BSM Signal Search in Higgs couplings



BSM Signal Search in Higgs couplings



Summary



Summary

e^-e^- physics in ILC

- Majorana fermion search
 - Majorana neutrino
 - s -electron (beam pol.)
- X^{--} Search
- Higgs Physics
- An e^-e^- option of ILC is interesting!!