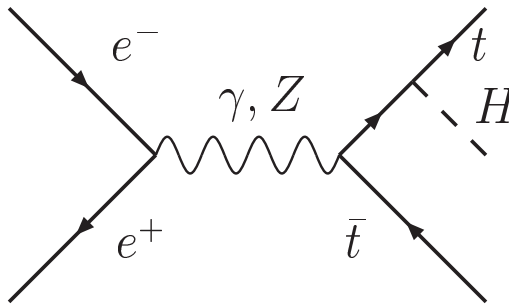




NLL QCD Corrections to $e^+e^- \rightarrow t\bar{t}H$ at 500 GeV



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in Collaboration with André Hoang

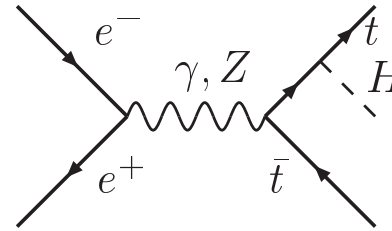
Phys.Rev.D72:014007, 2005 [hep-ph/0506253]

Outline

- Physical motivation
- Threshold effects of the top quarks, vNRQCD
- Results for $e^+e^- \rightarrow t\bar{t}H$ at NLL order
- Analysis at $\sqrt{s} = 500$ GeV
- Preliminary: Effects of e^+e^- polarization

Physical Motivation

Why $e^+e^- \rightarrow t\bar{t}H$?



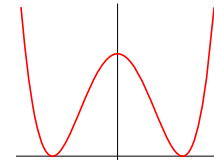
\Rightarrow Top-Yukawa coupling Y_t

Test of the EWSB mechanism:

Higgs mechanism predicts $m_t = Y_t v$

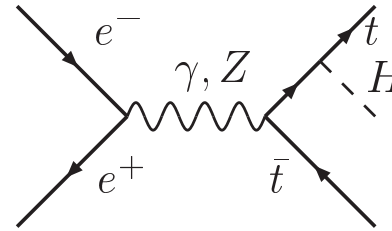
Experimental accuracy at the ILC: $\delta Y_t / Y_t \sim \%$

\Rightarrow Theoretical prediction: $\delta\sigma/\sigma \sim \%$



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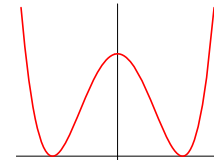
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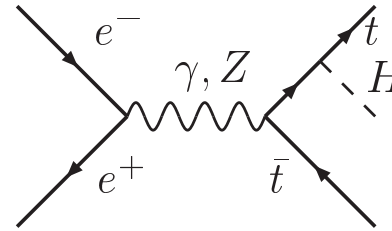
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- One-loop corrections [Dawson, Reina, Dittmaier, Belanger, Denner, Roth, Weber ...]

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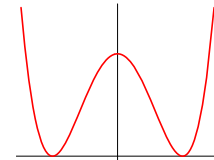
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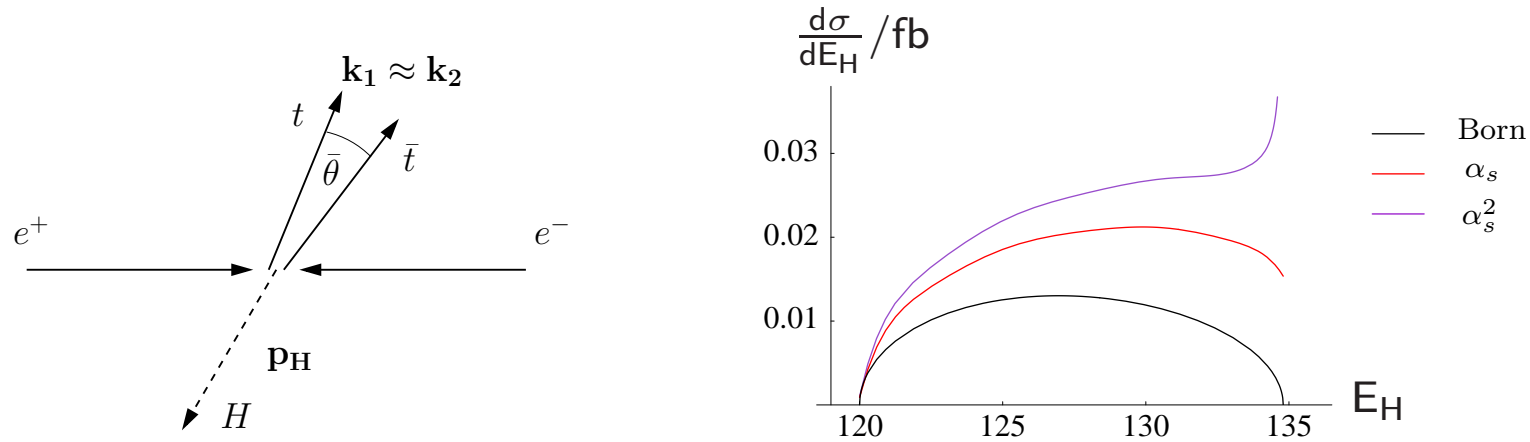
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But: Interesting kinematical region in $d\sigma/dE_H$

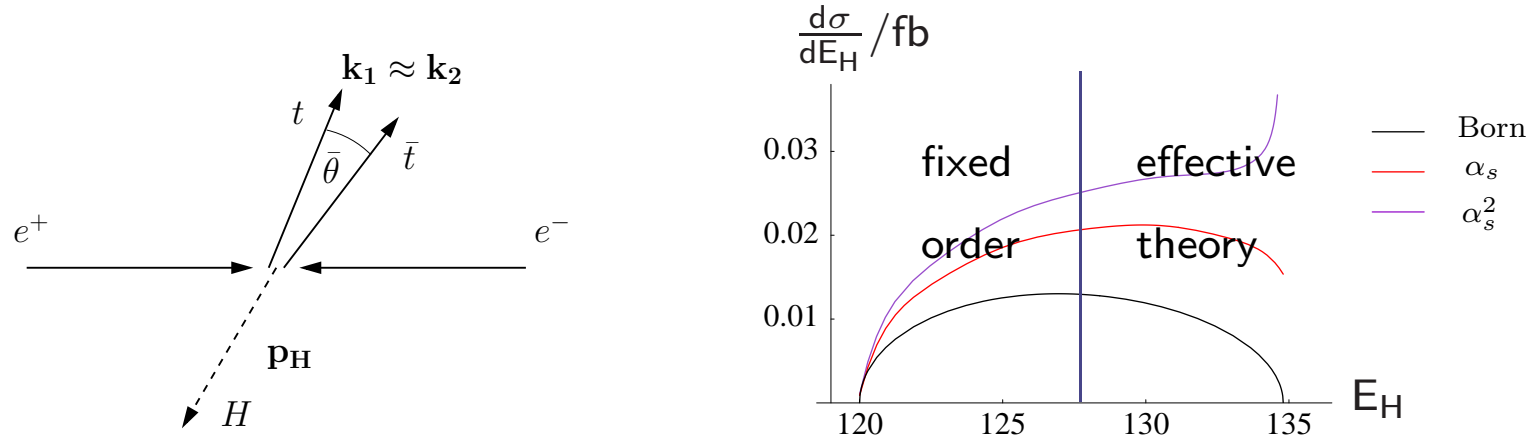
The Limit of Large Higgs Energies



The top quarks become collinear

- Coulomb Singularities $\sim \left(\frac{\alpha_s}{v}\right)^n$ and
- $(\alpha_s \ln v)^n$ singularities appear
- Fixed-order calculation breaks down

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\Rightarrow Effective Theory

$$\frac{d\sigma}{dE_H} \sim v \sum_n \left(\frac{\alpha_s}{v}\right)^n (\alpha_s \ln v)^n (1 \text{ (LL)} + \# \alpha_s \text{ (NLL)})$$

vNRQCD (velocity Non-Relativistic QCD)

[Luke, Manohar, Rothstein, Stewart, Hoang]

- Describes $t\bar{t}$ pairs near threshold
- Effective Theory with additional expansion parameter v
- Hierarchy:

$$m \gg mv \sim \vec{p} \gg mv^2 \sim E_{\text{kin}} \sim i\Gamma_t \gg \Lambda_{\text{QCD}}$$

- Width Γ_t as infrared regulator suppressing hadronization effects

- Lagrangian includes the fields of all resonant degrees of freedom
 → Schrödinger equation in the CMS
- Effective currents for the creation and annihilation of the $t\bar{t}$ pair:

$$|^{1S_0} = (\psi_{\vec{p}}^\dagger \tilde{\chi}_{-\vec{p}}) c_0(\mu) \quad (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$|^{3S_1} = (\psi_{\vec{p}}^\dagger \vec{\sigma} \tilde{\chi}_{-\vec{p}}) c_1(\mu) \quad (|\uparrow\uparrow\rangle, |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle, |\downarrow\downarrow\rangle)$$
- Wilson coefficients contain the non-resonant contributions
 - RGE running: $c(\mu) = c(m_t) \cdot u(m_t, \mu)$
- General structure of the $t\bar{t}$ production rate:

$$\sigma \sim \text{Im} \int e^{-iqx} d^4x \langle 0 | \mathbf{T} j(x) j^\dagger(0) | 0 \rangle \sim c^2(\mu) \text{Im}[G^0(\mu)]$$

Calculation

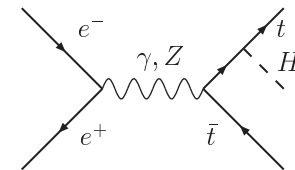
$$\frac{d\sigma}{dE_H} (E_H \approx E_H^{\max}) \sim \left[f_0^2 c_0^2(\mu, \sqrt{s}, m_t, m_H) + f_1^2 c_1^2(\mu, \sqrt{s}, m_t, m_H) \right] \text{Im } G_{\text{Coulomb}}^{\text{NLL}}(\mu, \sqrt{s}, m_t)$$

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$f_{0,1}^2$: electroweak information in the endpoint

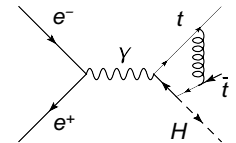
$$\sim \left(\frac{d\sigma_{0,1}}{dE_H} \right)_{\text{Born}} \text{ for } E_H \rightarrow E_{H,\max}$$



Calculation

$$\frac{d\sigma}{dE_H} (E_H \approx E_H^{\max}) \sim \left[f_0^2 c_0^2(\mu, \sqrt{s}, m_t, m_H) + f_1^2 c_1^2(\mu, \sqrt{s}, m_t, m_H) \right] \text{Im } G_{\text{Coulomb}}^{\text{NLL}}(\mu, \sqrt{s}, m_t)$$

$c_{0,1}^2$: hard QCD corrections
in the endpoint



$f_{0,1}^2$: electroweak information in the endpoint

Calculation

$$\frac{d\sigma}{dE_H} (E_H \approx E_H^{\max}) \sim \left[f_0^2 c_0^2(\mu, \sqrt{s}, m_t, m_H) + f_1^2 c_1^2(\mu, \sqrt{s}, m_t, m_H) \right] \text{Im } G_{\text{Coulomb}}^{\text{NLL}}(\mu, \sqrt{s}, m_t)$$

Known: - Renormalization group running of $c_{0,1}$

- $G_{\text{Coulomb}}^{\text{NLL}}$

New: - Matching conditions $f_{0,1}, c_{0,1}(\mu = m_t)$

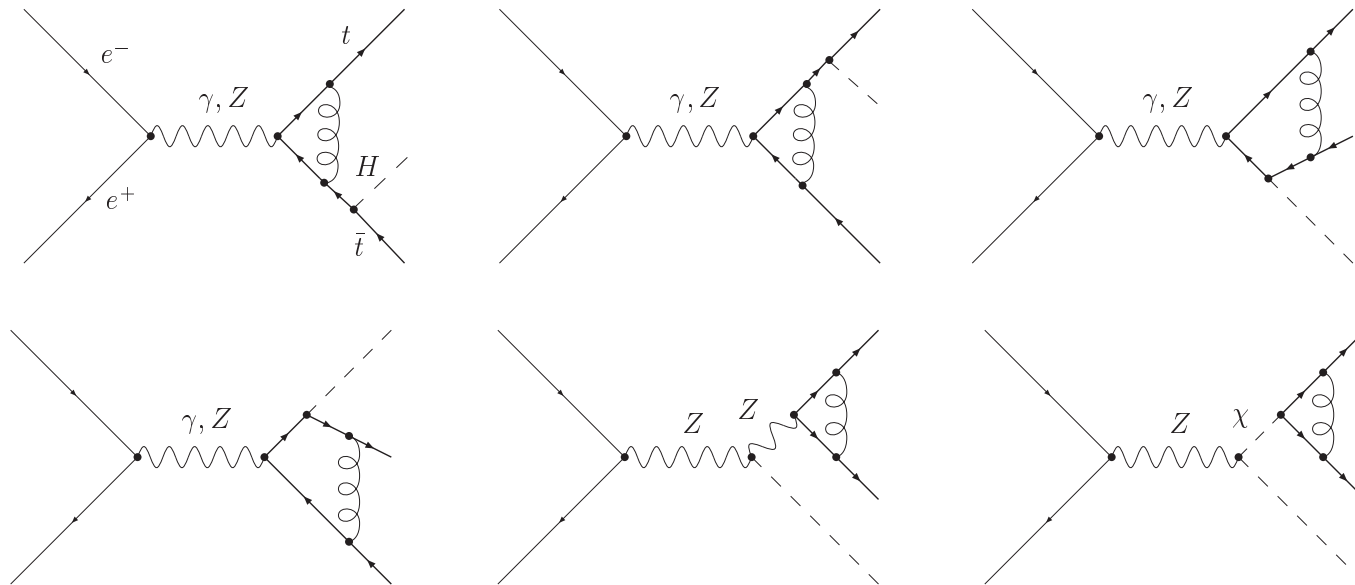
Ongoing: - Inclusion of e^+e^- polarization

Matching

$\mathcal{O}(\alpha_s)$ result in the limit $E_H \rightarrow E_H^{\max}$

$\stackrel{!}{=} \text{NLL result in the effective theory in } \mathcal{O}(\alpha_s)$

\Rightarrow Numerical Matching Conditions of the
Wilson coefficients $c_0(\mu = m_t)$, $c_1(\mu = m_t)$

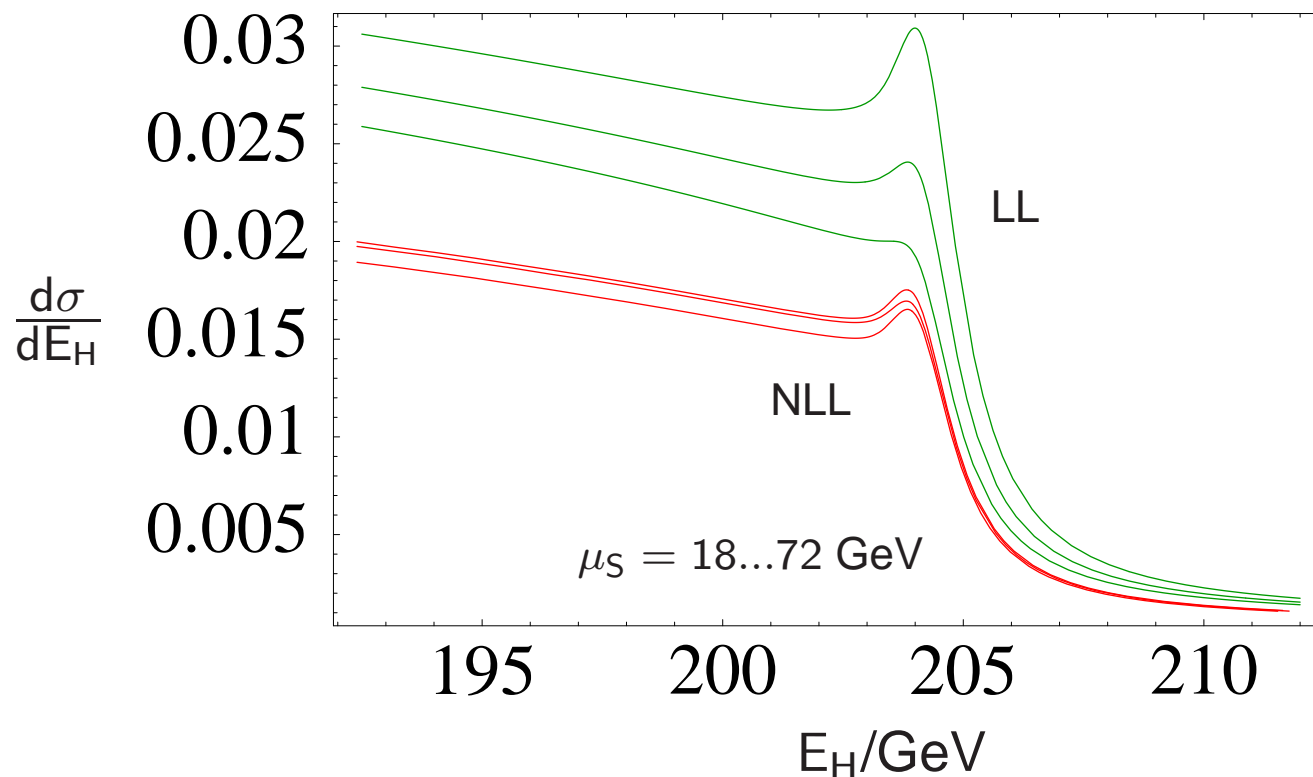


+ self-energy diagrams

[Denner, Dittmaier, Roth, Weber]

Scale Dependence

$$\sqrt{s} = 600 \text{ GeV}, m_H = 120 \text{ GeV}, m_t = 180 \text{ GeV}$$



$$\text{LL} \sim \left(\frac{\alpha_s}{v}\right)^n$$

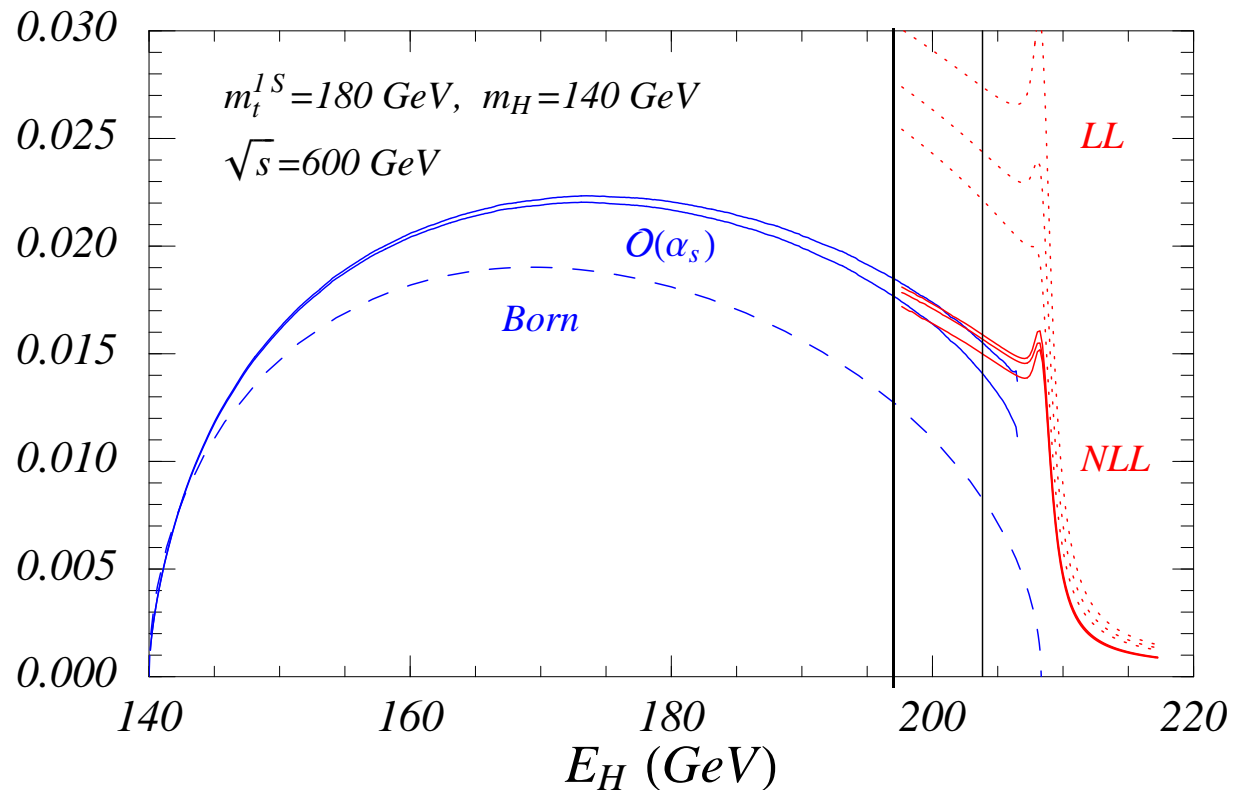
$$\text{NLL} \sim \alpha_s \left(\frac{\alpha_s}{v}\right)^n$$

⇒ NLL: scale dependence of the order of 3%

⇒ Error estimate of 5%

Differential Cross Section

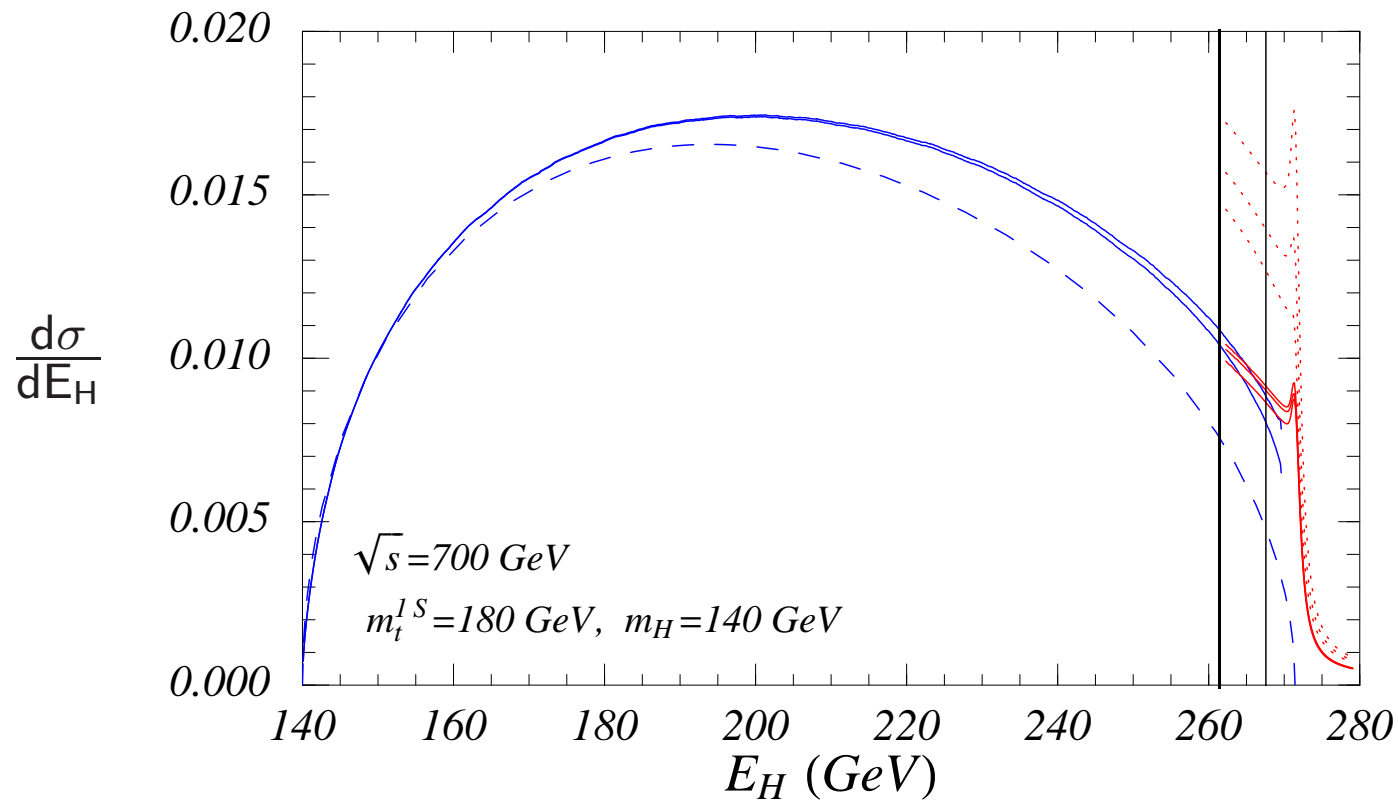
$$\sqrt{s} = 600 \text{ GeV}$$



- In subleading order: Good transition around $\beta = 0.2$

Differential Cross Section

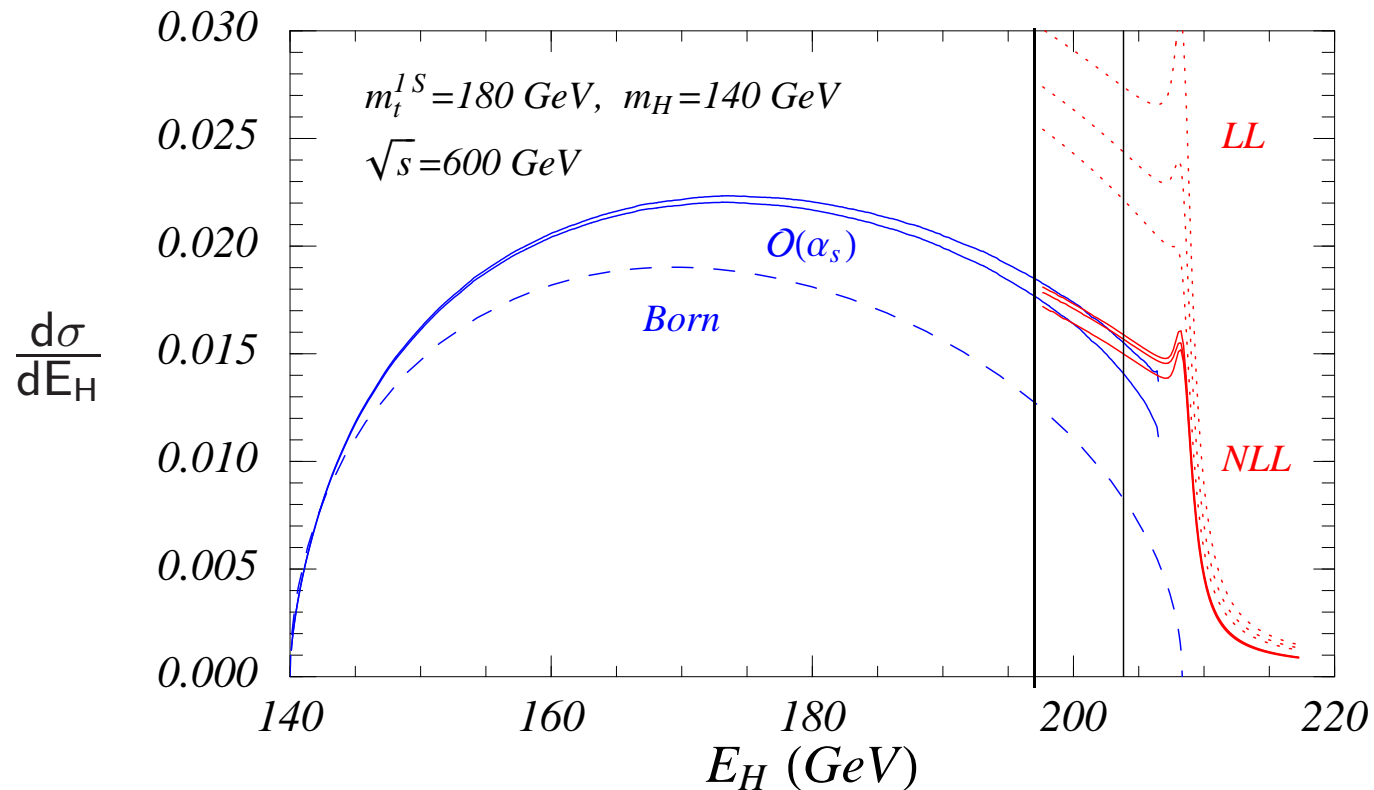
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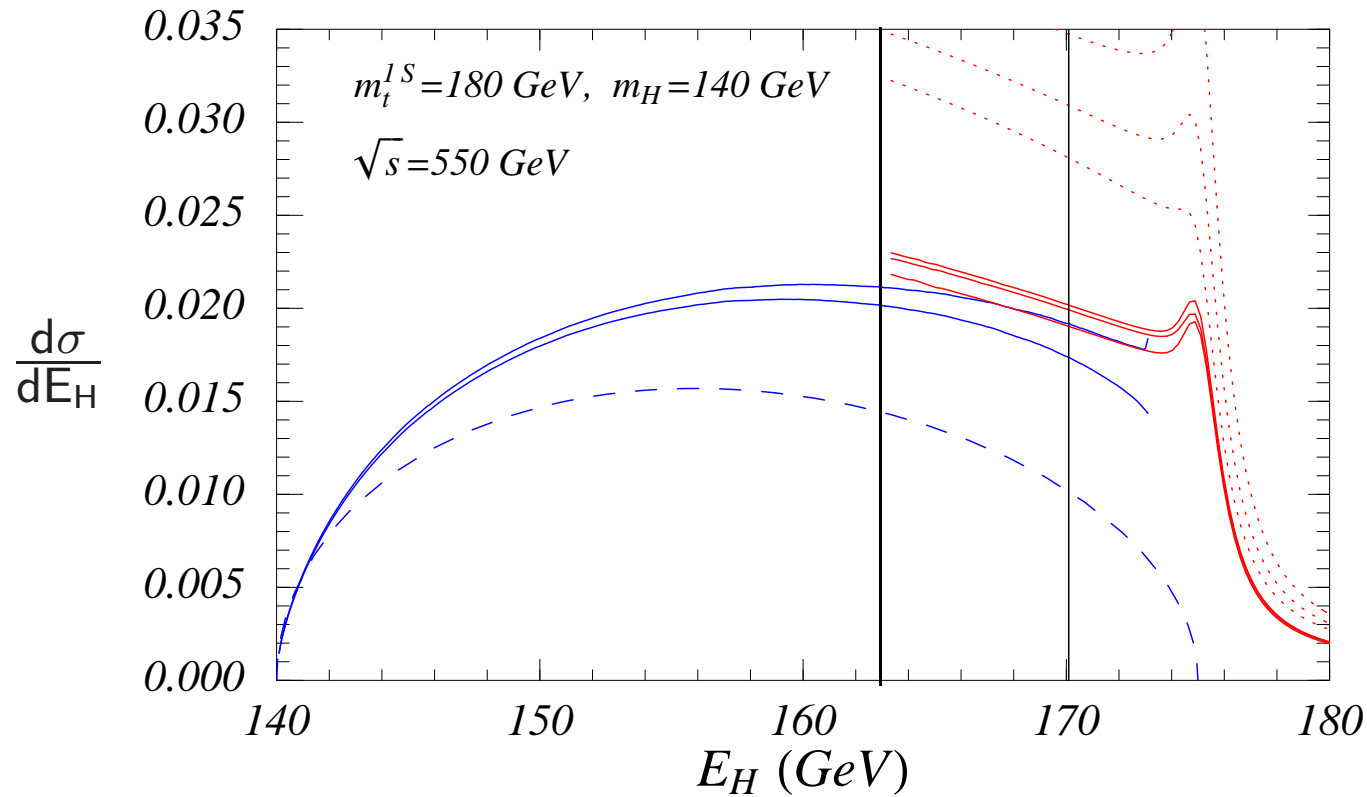
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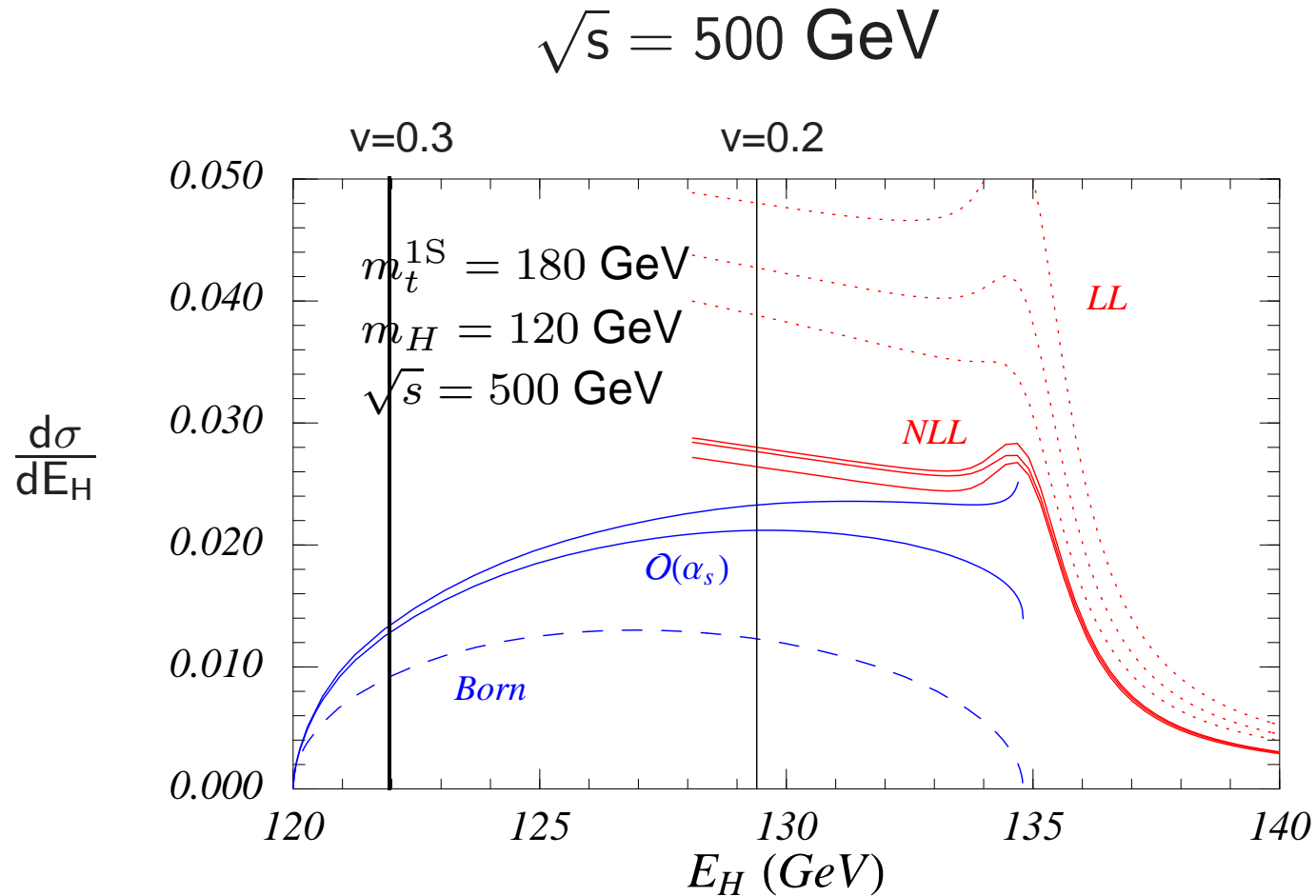
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ILC at 500 GeV

Maximum cms energy $\sqrt{s} = 500$ GeV

Little phase space ($2m_t + m_H \gtrsim 465$ GeV)

Dedicated analysis using the Born cross section:

$$\delta Y_t / Y_t \approx 25\% \text{ (A. Juste, 2002)}$$

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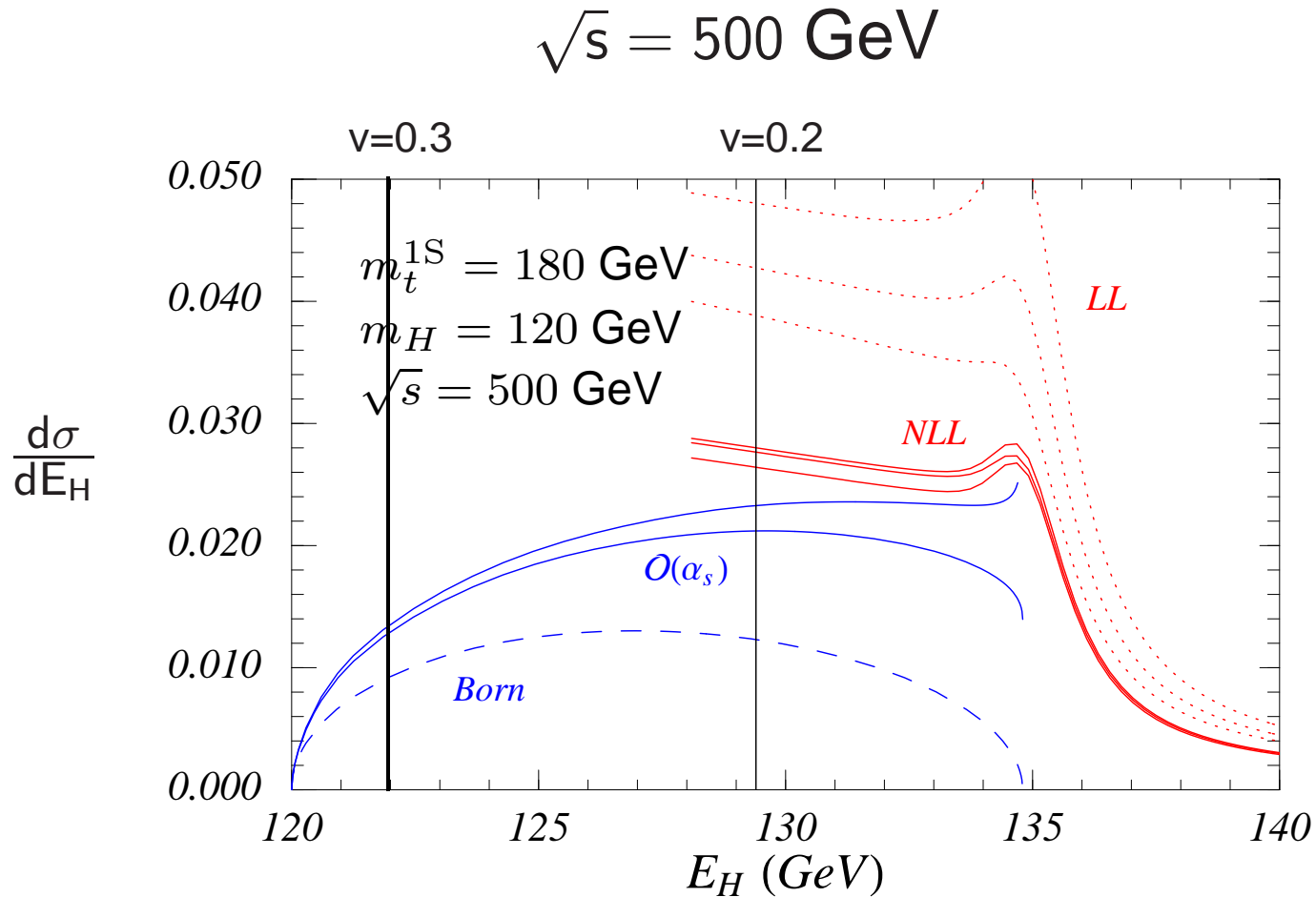
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$$\left(\frac{d\sigma}{dE_H} \right)_{E_H \approx E_H^{\max}} \sim [f_0^2 c_0^2(\mu, \sqrt{s}, m_t, m_H) + f_1^2 c_1^2(\mu, \sqrt{s}, m_t, m_H)] \text{Im } G_{\text{Coulomb}}^{\text{NLL}}(\mu, \sqrt{s}, m_t, m_H)$$

$$\left(\frac{d\sigma_{\text{Born}}}{dE_H} \right)_{E_H \rightarrow E_{H, \max}} \longrightarrow \frac{d\sigma_{\text{Born}}}{dE_H}$$

\Rightarrow Inclusion of the kinematical information in the Born CS

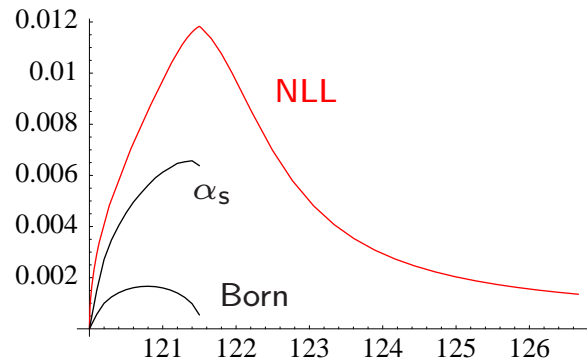
Differential Cross Section



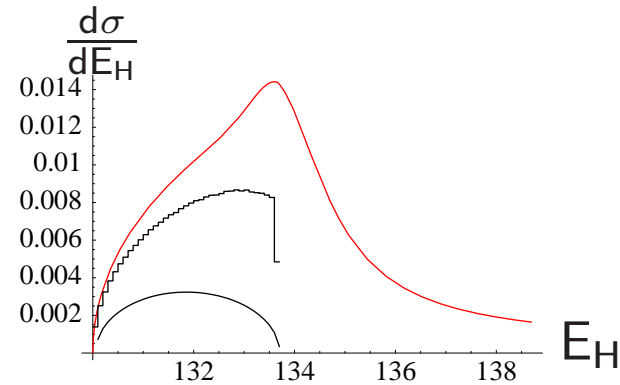
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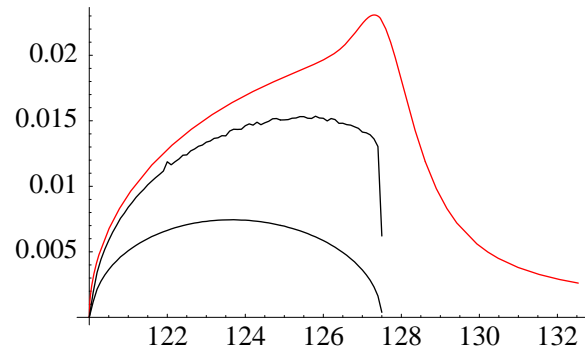
$\sqrt{s} = 482 \text{ GeV}, m_H = 120 \text{ GeV}$



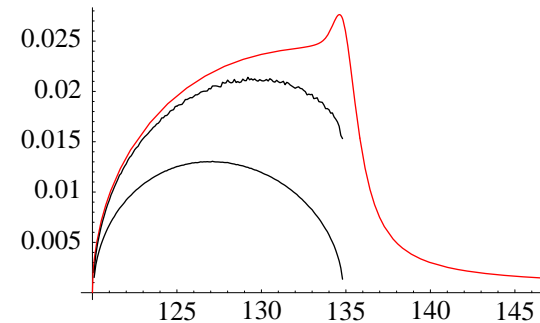
$\sqrt{s} = 500 \text{ GeV}, m_H = 130 \text{ GeV}$



$\sqrt{s} = 490 \text{ GeV}, m_H = 120 \text{ GeV}$



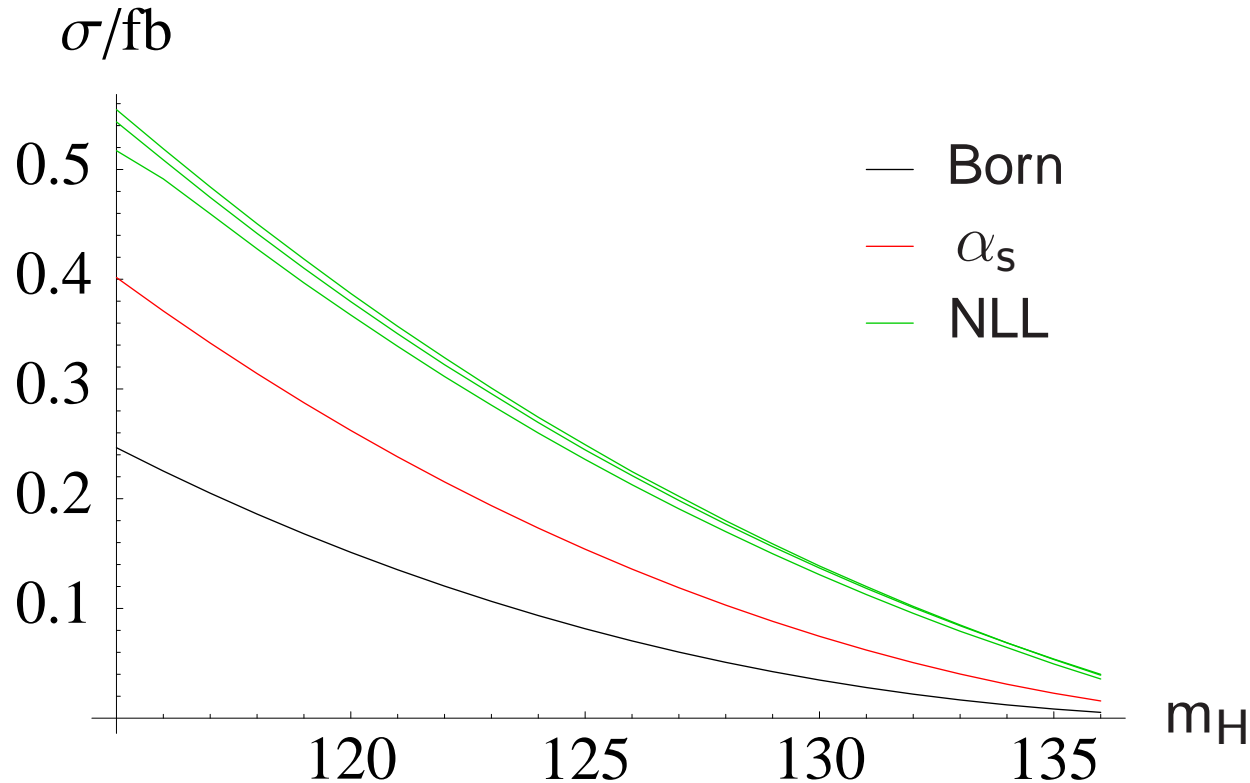
$\sqrt{s} = 500 \text{ GeV}, m_H = 120 \text{ GeV}$



⇒ Behavior far from threshold is well reproduced
⇒ Increase of total cross section

ILC at 500 GeV

$$\sqrt{s} = 500 \text{ GeV}, m_t = 180 \text{ GeV}$$

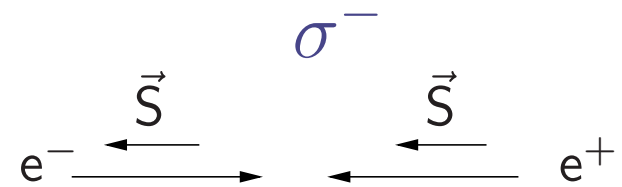
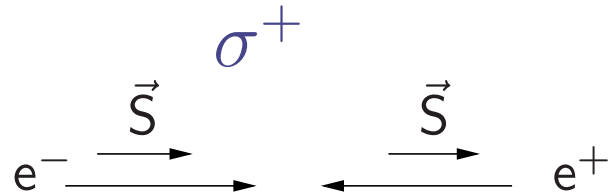


Cross section: NLL 2-4 times larger than Born

\Rightarrow Measurement of Y_t is possible
for all accessible Higgs masses!

e^+e^- Polarization (preliminary)

Cross section depends on the helicity of e^+ and e^- :



Total cross section:

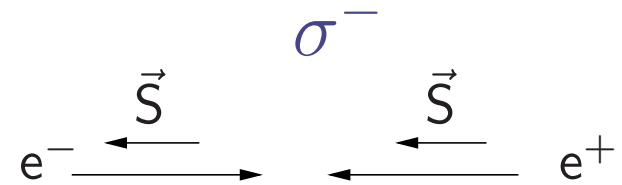
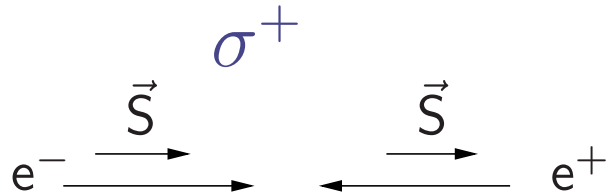
$$\begin{aligned}\sigma_{\text{pol}} &= \frac{1}{4}(1 + P_+)(1 - P_-) \sigma^+ + \frac{1}{4}(1 - P_-)(1 + P_+) \sigma^- \\ &= \sigma_{\text{unpol}} [1 - P_- P_+ - A_{\text{LR}}(P_+ - P_-)]\end{aligned}$$

P_{\pm} : degree of e^{\pm} polarization

left-right asymmetry: $A_{\text{LR}} = \frac{\sigma^- - \sigma^+}{\sigma^- + \sigma^+}$

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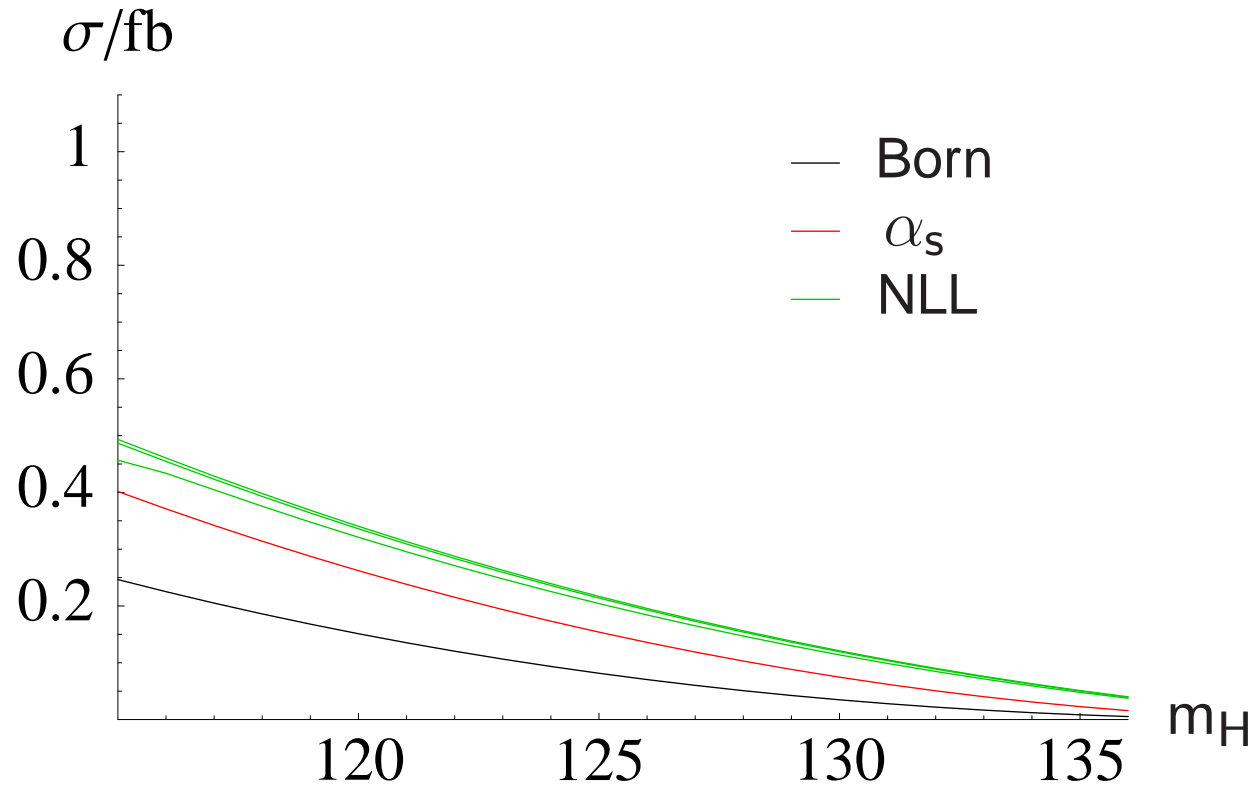
left-right asymmetry: $A_{\text{LR}} = \frac{\sigma^- - \sigma^+}{\sigma^- + \sigma^+}$

\Rightarrow Increase of σ_{tot} by polarization possible

e^+e^- Polarization (preliminary)

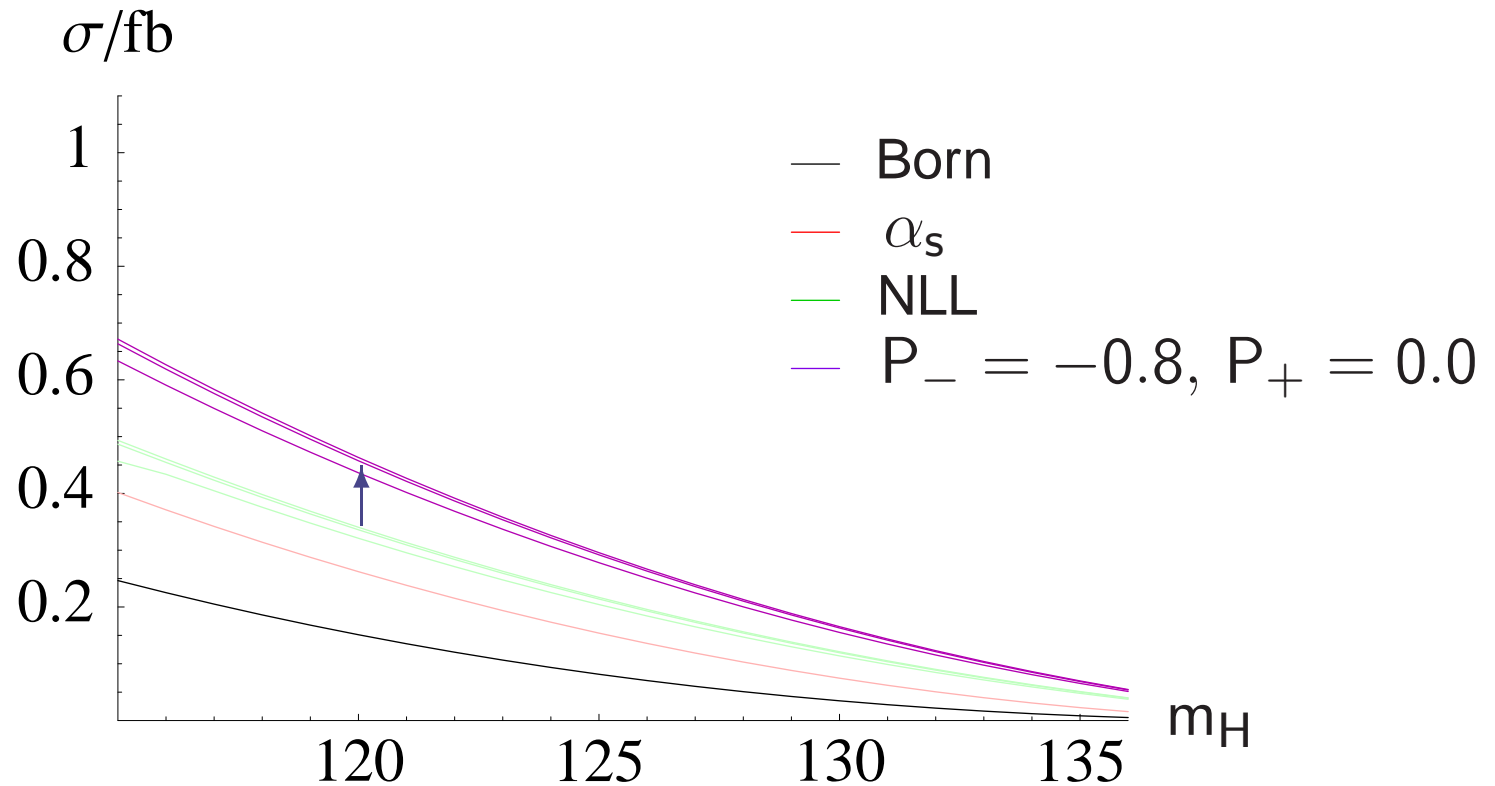
Unpolarized e^+e^- -beams:

$$\sqrt{s} = 500 \text{ GeV}, m_t = 180 \text{ GeV}$$



Polarized electron beam:

$$\sqrt{s} = 500 \text{ GeV}, m_t = 180 \text{ GeV}$$

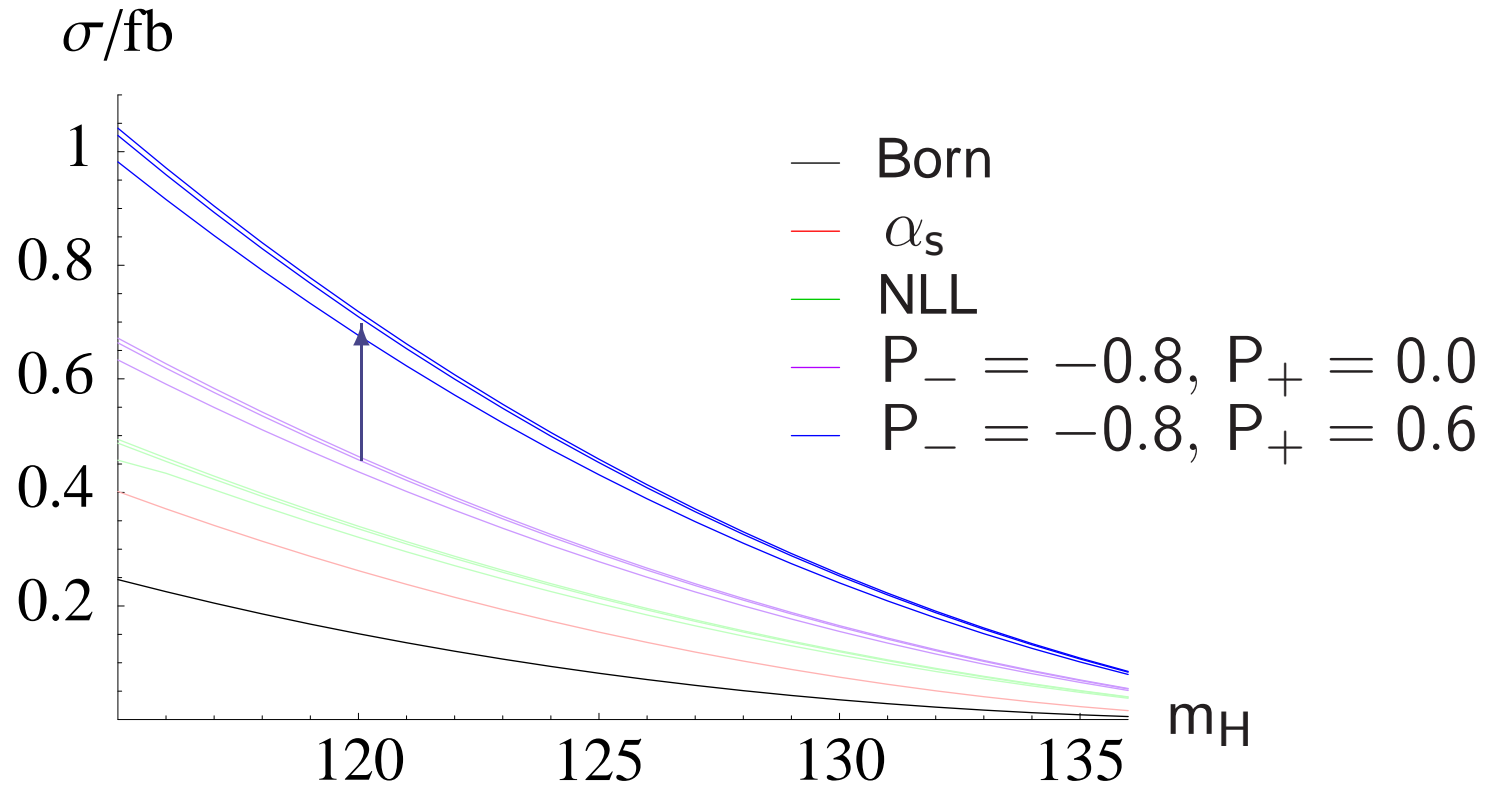


- Increase of σ_{tot} by $\sim 40\%$

e^+e^- Polarization (preliminary)

Polarized e^+e^- - beams:

$$\sqrt{s} = 500 \text{ GeV}, m_t = 180 \text{ GeV}$$

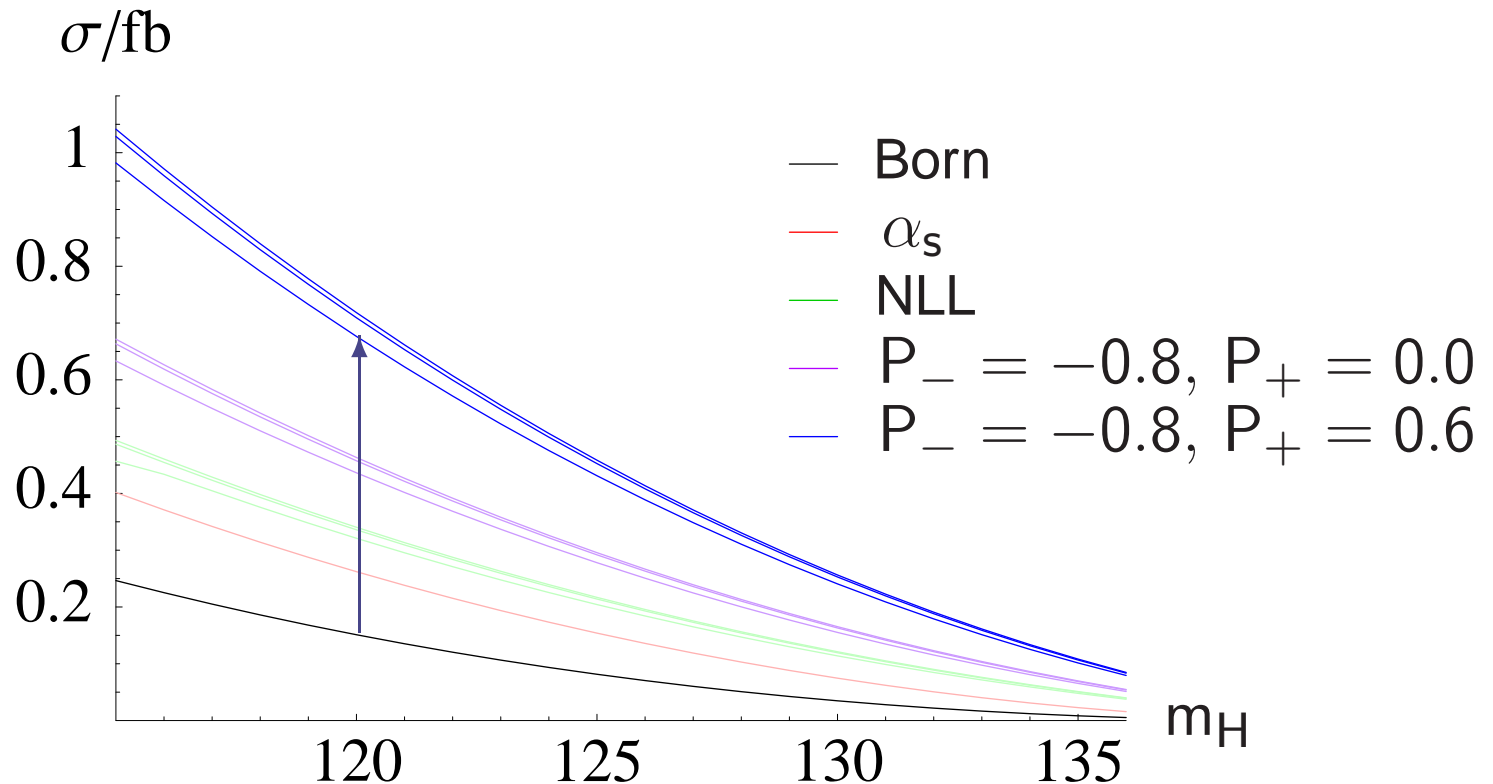


- Increase of σ_{tot} by $\sim 100\%$ vs. $P_+ = 0$

e^+e^- Polarization (preliminary)

Polarized e^+e^- - beams:

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- Increase of σ_{tot} by $\sim 400\%$ vs. Born Cross section
 \Rightarrow Decrease of statistical error by 50% to $\sim 12\%$

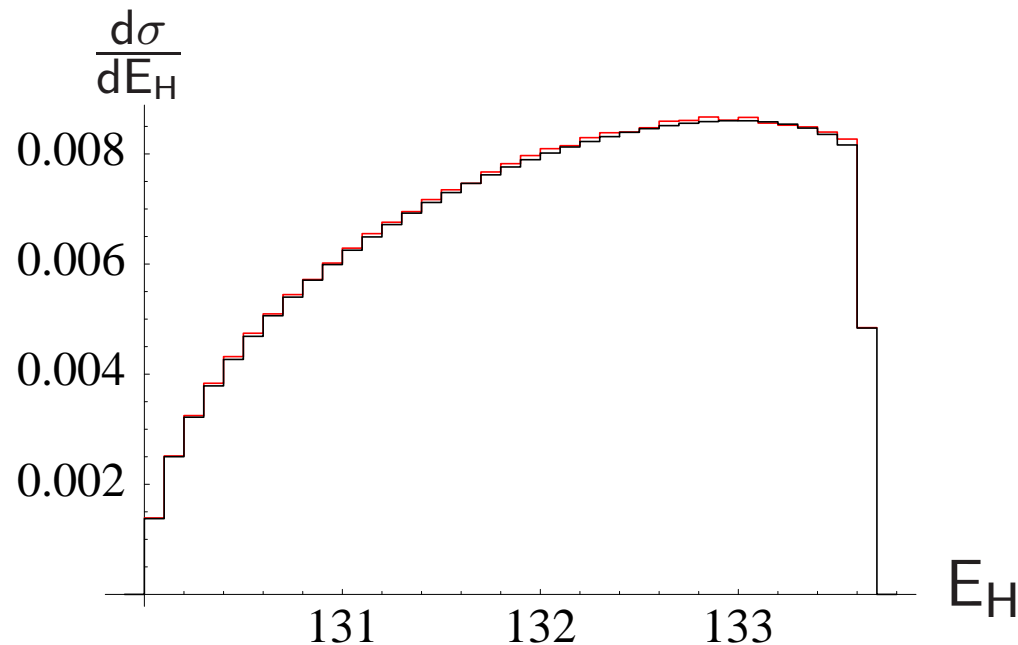
Conclusion

- Measurement of the top-Yukawa coupling for the verification of the Higgs mechanism
- First phase of the ILC
 - At energies of 500 GeV resummation must be included
⇒ vNRQCD
- Calculation at NLL order
 - Scale dependence of the order of 3%
 - Large enhancement of the cross section
- Outlook: Inclusion of e^+e^- polarization
 - Increase of the total cross section
 - Y_t might be measurable to 10-15% at 500 GeV

Reproduction of the $\mathcal{O}(\alpha_s)$ Cross Section

Comparison: - Full $\mathcal{O}(\alpha_s)$ result
- NLL result at $\mathcal{O}(\alpha_s^1 v^1)$

$\sqrt{s} = 495 \text{ GeV}$, $m_t = 180 \text{ GeV}$, $m_H = 125 \text{ GeV}$



Difference: \Rightarrow At threshold: \sim Per mille

\Rightarrow Up to $\sqrt{s} = 700 \text{ GeV}$: maximally $\sim 1.5\%$