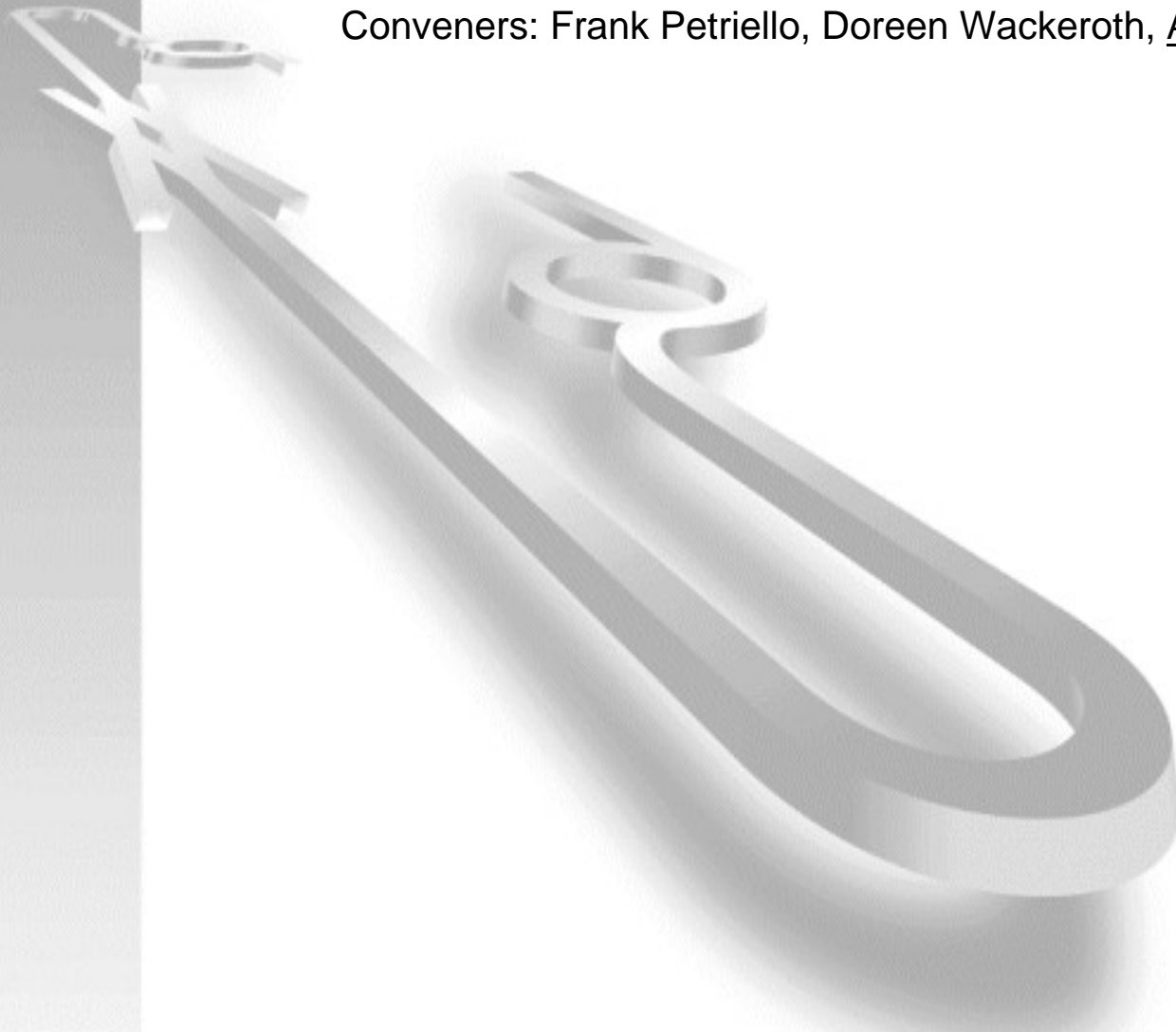


VLCW06, Vancouver, July 19-23, 2006

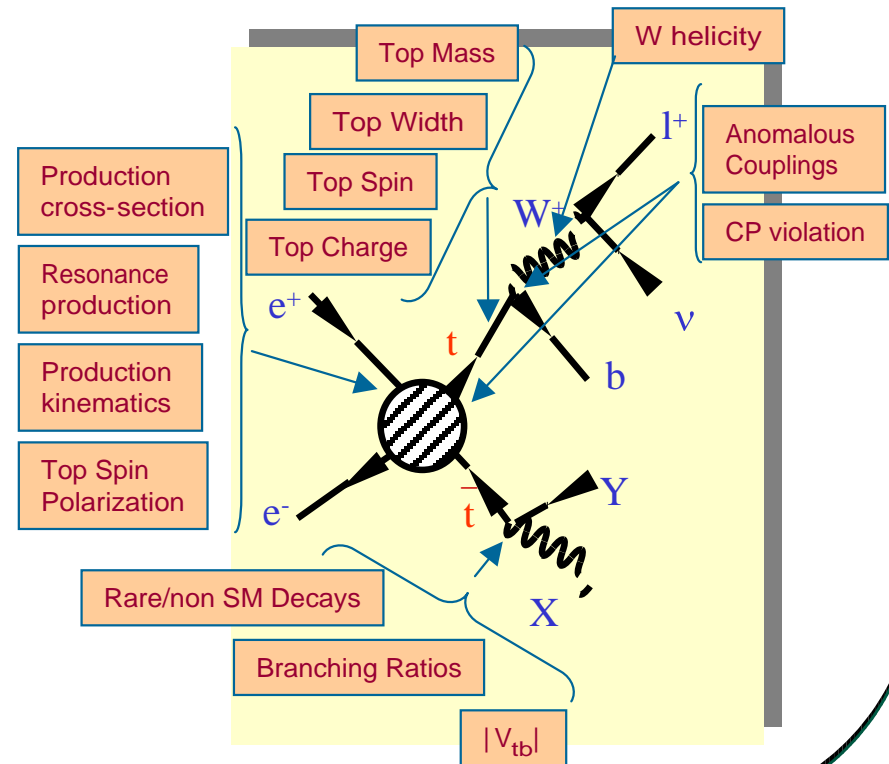
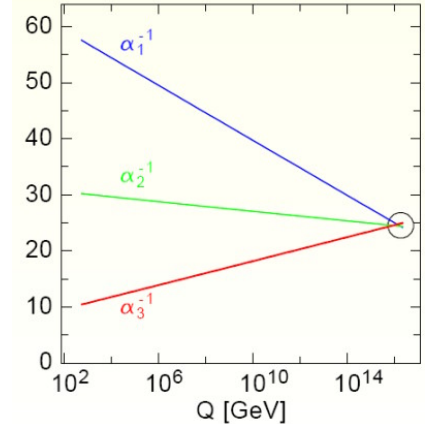
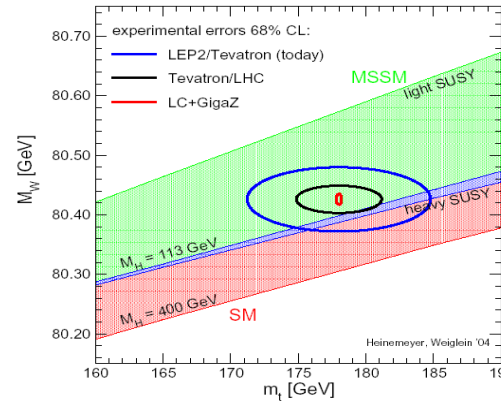
Precision Physics WG Summary

Conveners: Frank Petriello, Doreen Wackerath, Aurelio Juste

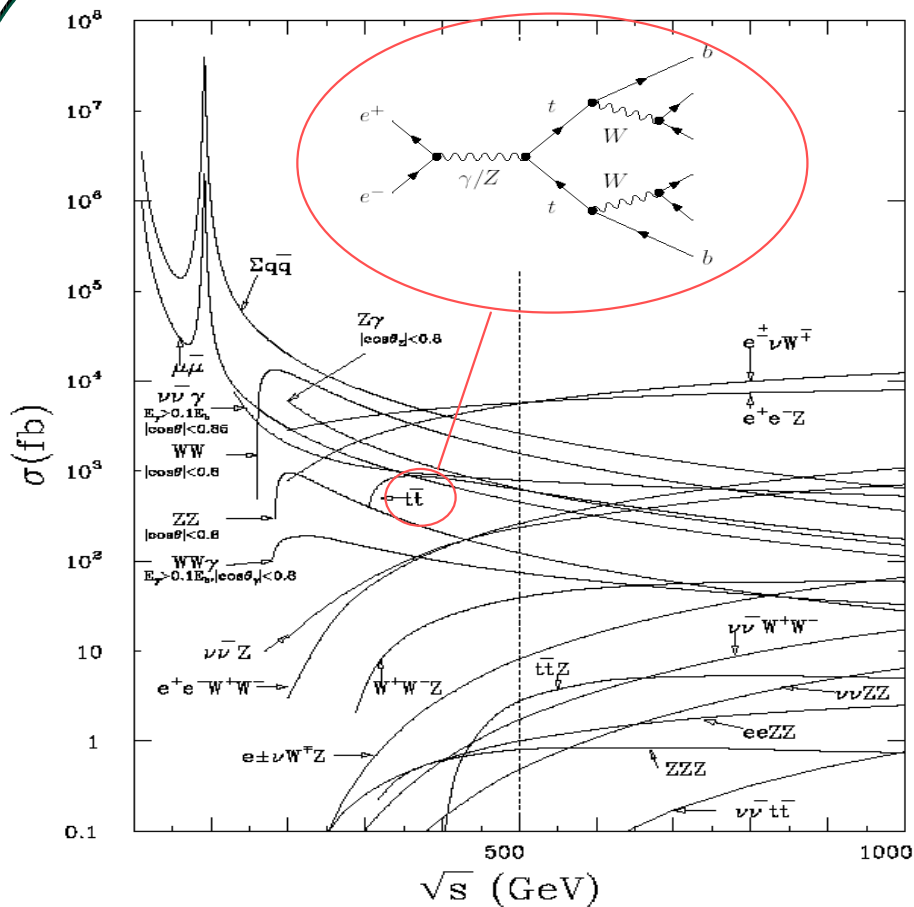


The Power of Precision Physics

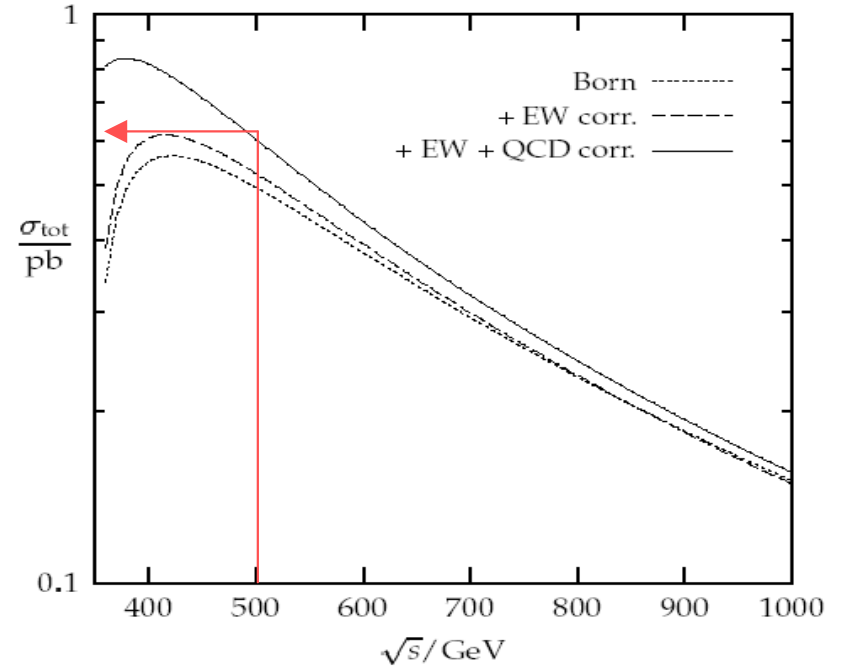
- The LHC will be probing the relevant energy scale and should definitely discover signs of the EWSB dynamics.
- The main strength of the ILC resides on its precision and model independence \Rightarrow will complement the LHC by providing essential information to interpret and exploit these discoveries.
- Here we cover Top, QCD and EW Physics, although at the ILC, precision measurements will extend well beyond these topics.
- Precise measurements of EW (e.g. M_W) and QCD (e.g. α_s) parameters essential to provide precise theoretical calculations, constrain models of New Physics, extrapolate to GUT scale, ...
- $m_t \sim 175 \text{ GeV} \Rightarrow \lambda_t = \sqrt{2} m_t/v \approx 1$
 \Rightarrow The top quark may either play a key role in EWSB, or serve as a window to New Physics related to EWSB.
 Fully outlining the top quark profile will be critical to unravel the secrets of EWSB.



Top Production in e^+e^- Collisions



- Top pair production via γ/Z exchange dominates



$\sigma_{tt} \sim 0.6 \text{ pb}$ at $\sqrt{s} = 500 \text{ GeV}$
 $\Rightarrow \sim 200\text{k}$ events/year ($L = 2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)

- The anticipated experimental accuracy must be matched with precise theoretical predictions

Available {

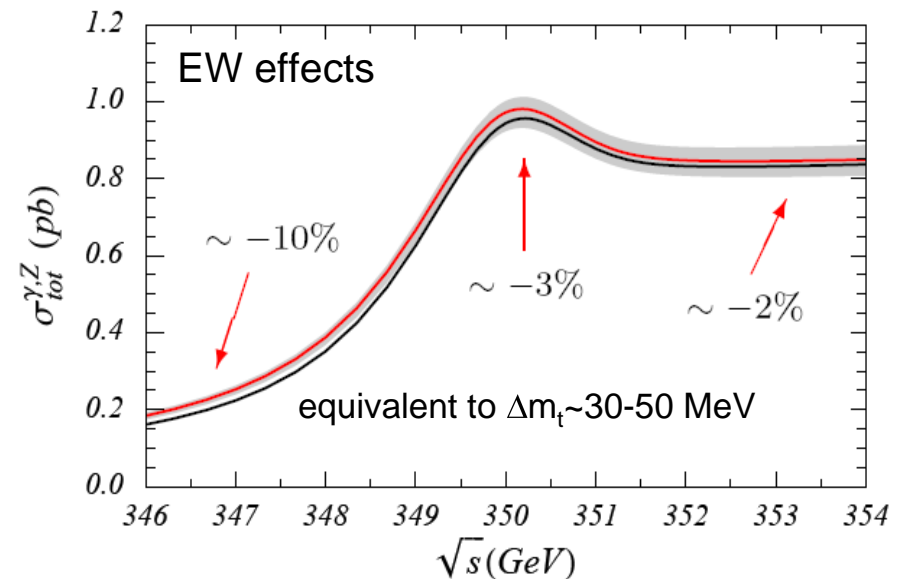
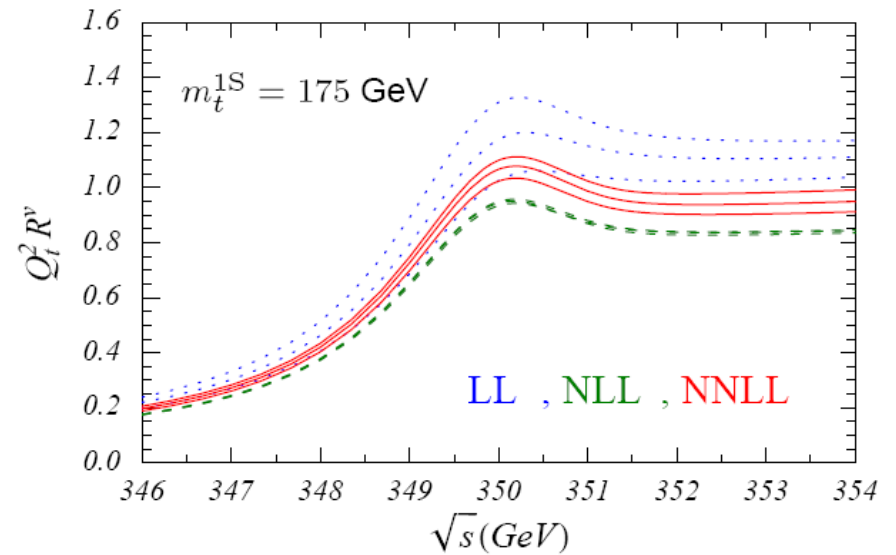
- Total cross section
- tt , threshold : NNLL QCD, N(LL) EW
- continuum: $O(\alpha_s^2)$, $O(\alpha_{EW})$, 2-loop Sudakovs
- ttH : $O(\alpha_s)$, $O(\alpha_{EW})$, tt threshold effects

Event generators
 $e^+e^- \rightarrow (tt) \rightarrow WbWb: O(\alpha_s)$

Will be needed: $e^+e^- \rightarrow 6f$ (lusifer) and $e^+e^- \rightarrow 8f$ to $O(\alpha_s)$
 consistent treatment of unstable particles, non-factorizable corrections, ...

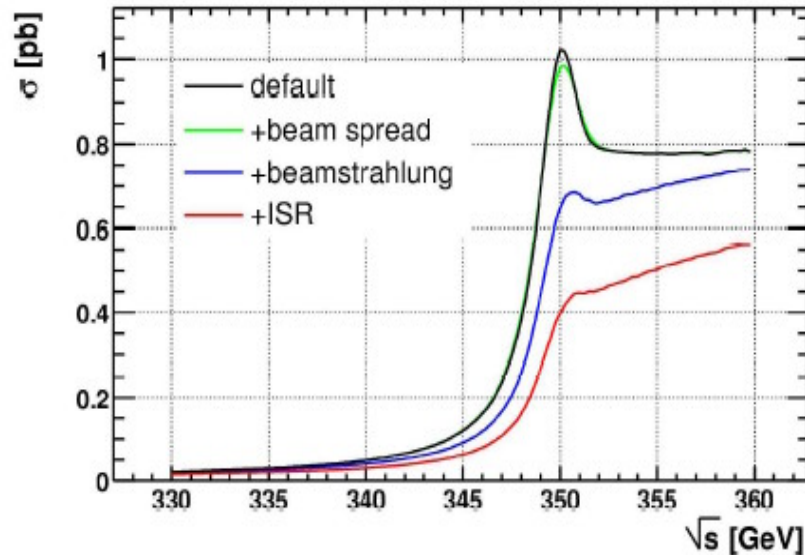
Top Pair Production at Threshold (I)

- Large Γ_t : cutoff for non-perturbative QCD effects
 - Top decays before top-flavored hadrons or $t\bar{t}$ -quarkonium bound states can form.
 - Use non-relativistic pQCD to compute $\sigma_{t\bar{t}}$ near threshold.
- Remnants of toponium S-wave resonances induce a fast rise of $\sigma_{t\bar{t}}$ near threshold.
 Basic parameters: $\sigma_{t\bar{t}}(m_t, \alpha_s, \Gamma_t)$
 \Rightarrow high precision expected (color singlet system, counting experiment,...)
- Convergence of calculation sensitive to m_t definition used: **pole mass is not IR-safe**
 $\Rightarrow \sigma_{t\bar{t}}^{\text{peak}}$ not stable vs \sqrt{s}
 Solution is to use threshold masses: e.g. 1S mass (=1/2 the mass of the lowest $t\bar{t}$ bound state in the limit $\Gamma_t \rightarrow 0$).
 High accuracy in absolute normalization requires velocity resummation.
 State of the art (NNLL): $(\Delta\sigma_{t\bar{t}})_{\text{QCD}} \sim 6\%$
- **Goal: 3%** \Rightarrow important to take into account previously neglected %-level effects: EW corrections (Γ_t + non-resonant W^+bW^-b background, QED), non-factorizable QCD corrections,... \Rightarrow **a lot of work ahead!**



\Rightarrow Talk by Andre Hoang

Top Pair Production at Threshold (II)



- Lineshape significantly distorted:

$$\sigma^{\text{obs}}(\sqrt{s}) = \frac{1}{L_0} \int_0^1 L(x) \sigma(x\sqrt{s}) dx$$

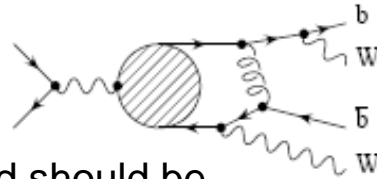
- Beam energy spread+Beamstrahlung: must be measured (acollinearity in Bhabha events) \Rightarrow detector/theory precision.
- Bremsstrahlung (ISR): can a-priori be calculated precisely (enough?)
- Precise determination of $dL/d\sqrt{s}$ and $\langle\sqrt{s}\rangle$ critical.

Not only σ_{tt} but also differential observables are important!

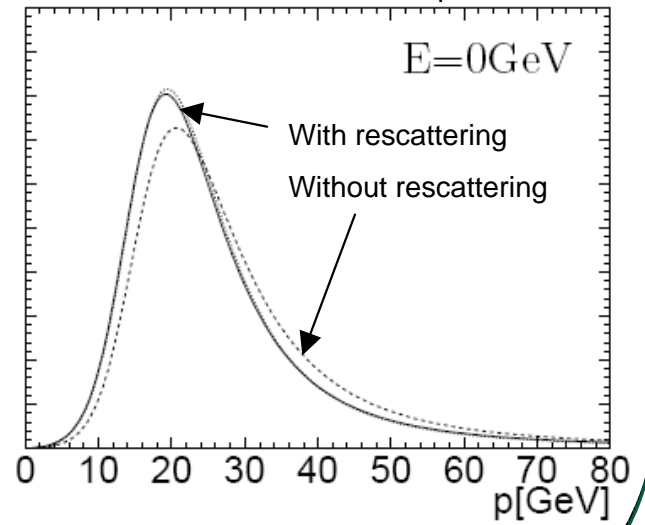
- To implement effect of experimental cuts (MC event generators)
- Exploit additional experimental information from A_{FB} , $d\sigma/dp_t$, s_t, \dots
 - Additional sensitivity to m_t , α_s and Γ_t
 - Reduce correlations

\Rightarrow Simultaneous determination of parameters possible when using all threshold observables.

- Non-factorizable QCD (aka “rescattering”) corrections important.
- N(N)LO QCD corrections are available and should be implemented in a MC event generator for more realistic experimental studies.

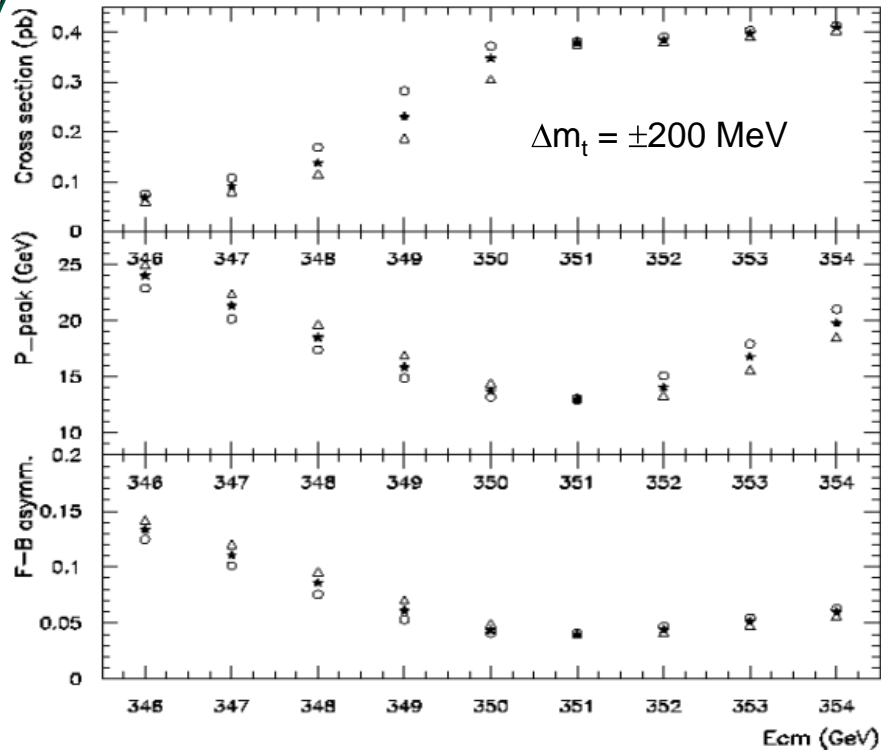


Top momentum spectrum hep-ex/9604328



Experimental Threshold Scan (I)

hep-ph/0207315



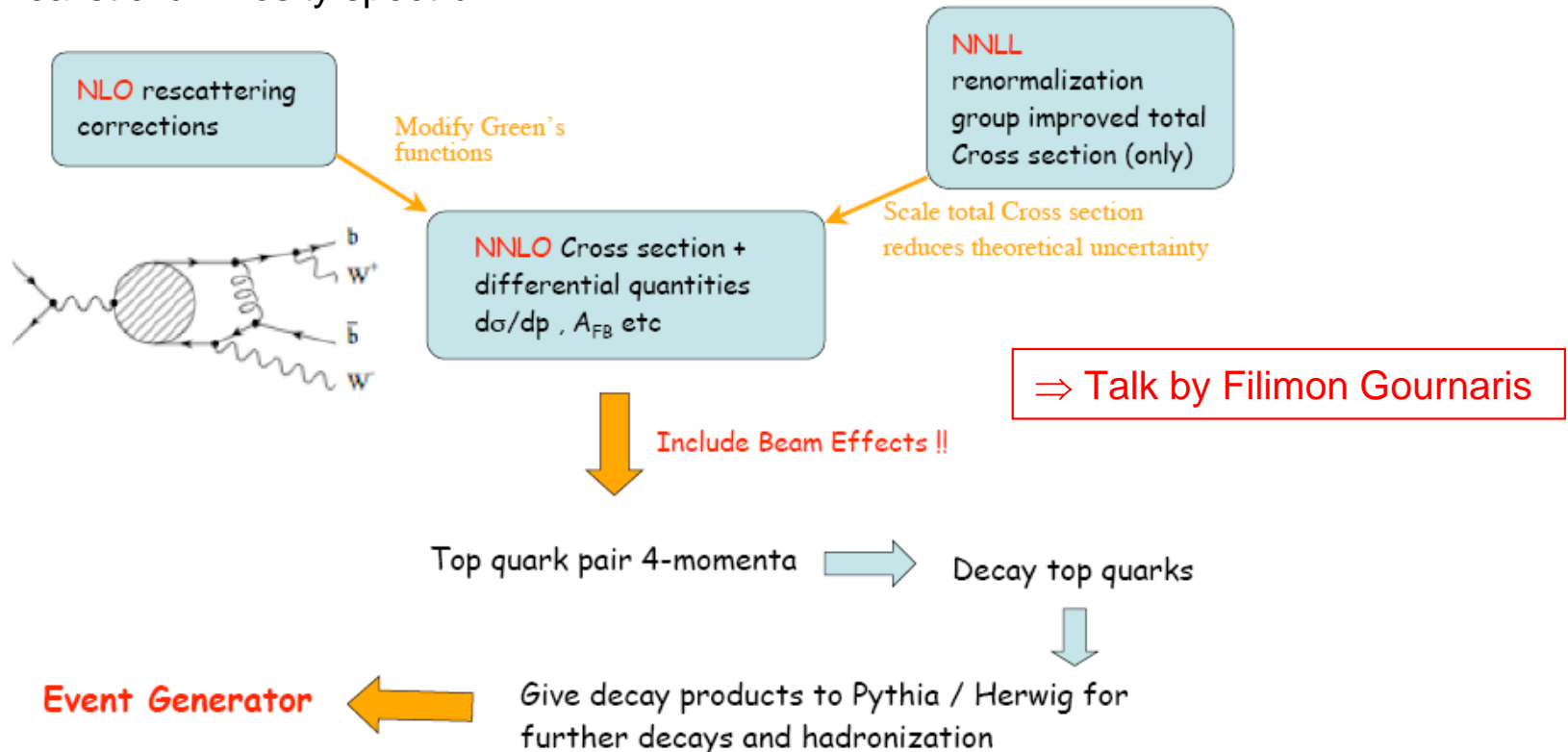
Example of attainable precision:

- 9+1 point scan with $30 \text{ fb}^{-1}/\text{point}$.
- Lepton+jets and alljets final states.
- **Assumptions:**
 - **Theoretical uncertainties:**
 $\Delta\sigma_{tt}=3\%$, $\Delta p_t^{\text{peak}}=\Delta A_{\text{FB}}=0$
 - Perfect knowledge of luminosity spectrum
 - SM top couplings
- Simultaneous determination of m_t , α_s and Γ_t (experimental uncertainties only):
 $\Delta m_t(1S)=19 \text{ MeV}$, $\Delta\alpha_s=0.0012$, $\Delta\Gamma_t=32 \text{ MeV}$, $\rho_{ij}<0.5$
- **Theoretical uncertainty on $m_t \sim O(100 \text{ MeV})$.** More work needed for a precise estimate.

- **Experimental analysis should be repeated with increased level of realism:**
 - Run parameter optimization (beamstrahlung, integrated luminosity/point, beam polarization, ...)
 - Current detector concepts: event reconstruction and selection
 - Realistic energy and luminosity spectrum uncertainties
 - Improved Monte Carlo event generator (including rescattering corrections)
 - Realistic theoretical uncertainties
 - Simultaneous determination of top anomalous couplings and/or implementation of external constraints for a model-independent determination of threshold parameters
 - ...

Experimental Threshold Scan (II)

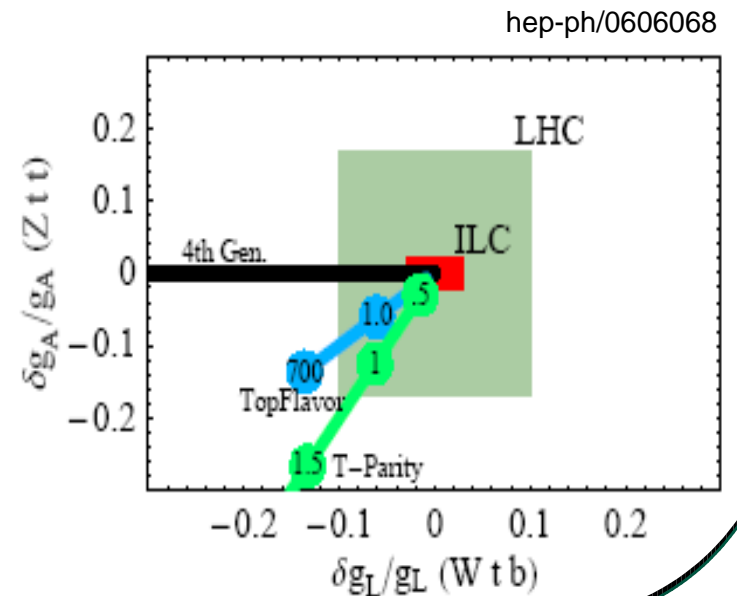
- Ongoing work to build a MC event generator for precise $t\bar{t}$ threshold studies, including:
 - State of the art theoretical calculations for inclusive cross section and differential distributions
 - Top polarization
 - Realistic luminosity spectrum



- Important requirement: speed (NNLO calculation in TOPPIK takes $>1.5s/event$)
 \Rightarrow pre-compute Green functions with TOPPIK and perform fast 4D interpolation in $(m_t, \Gamma_t, \alpha_s, \sqrt{s})$
- Powerful tool for the next generation of threshold studies:** top spin physics at threshold, optimization of scan strategy, additional observables, meaningful systematic uncertainties,...

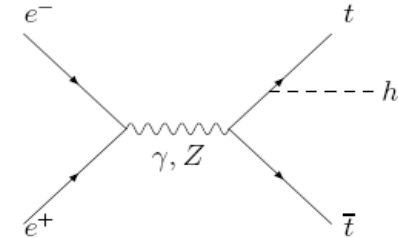
Top Quark and New Physics

- The top quark plays a prominent role in many models of New Physics.
 - Radiative effects in the top sector drive EWSB (e.g. Little Higgs)
 - New particles preferentially coupled to the top quark:
 - Vector gauge bosons (e.g. g_t in Topcolor, V_{KK} in Randall-Sundrum extra-dimensions,..)
 - Charged scalars (e.g. H^\pm in generic 2HDMs)
 - Neutral scalars (e.g. η_T in Technicolor)
 - Mixing with the top quark (e.g. TopFlavor)
 - ...
- In many instances the new states can be sufficiently heavy that they are not produced directly at the ILC (or even the LHC).
 - ⇒ The first indication for New Physics might be in the form of modified top quark interactions.
- Anomalous top couplings can manifest themselves affecting many observables:
 - total cross-sections,
 - $t\bar{t}$ invariant mass distribution,
 - angular distributions of decay products,
 - rare decays,
 - ...
- Several operators can contribute to a given observable. Must disentangle the effect of different operators. Beam polarization will be a crucial tool.
- High precision and model independence critical to be able to rule/figure out specific models of New Physics.



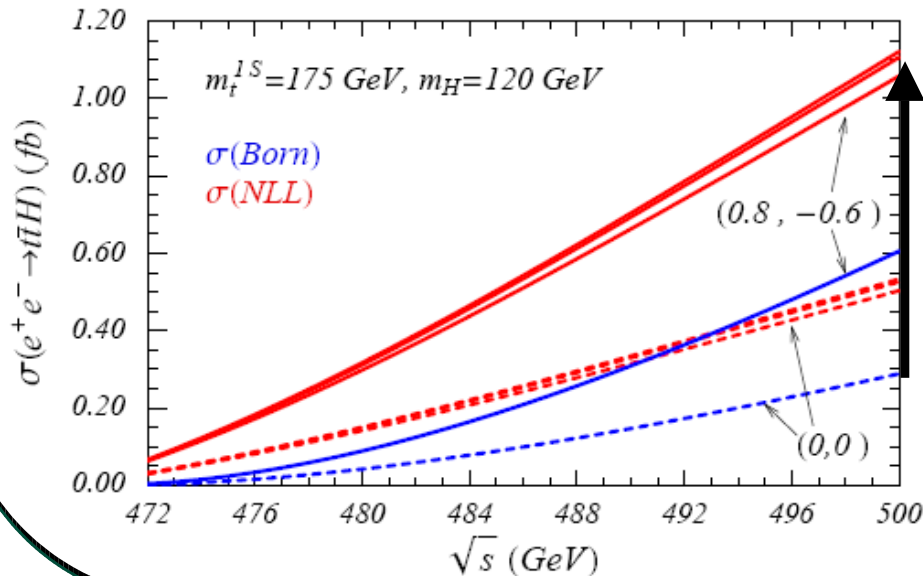
Top-Higgs Yukawa Coupling

- The top-Higgs Yukawa coupling is the largest coupling of the Higgs boson to fermions. **Precise measurement important since the top quark is the only “natural” fermion from the EWSB standpoint.**
- Can be determined via cross section measurement: $\sigma_{t\bar{t}h} \propto g_{t\bar{t}h}^2$
- $\sigma_{t\bar{t}h}(\text{Born}) \sim 0.2(2.5) \text{ fb}$ at $\sqrt{s}=500(800) \text{ GeV}$ for $m_h=120 \text{ GeV}$
- Previous study:
 $\sqrt{s}=800 \text{ GeV}$, $L=1 \text{ ab}^{-1}$, $\Delta g_{t\bar{t}h}/g_{t\bar{t}h} \sim 6(10)\%$ for $m_H=120(190) \text{ GeV}$
 \Rightarrow **What are the prospects at $\sqrt{s}=500 \text{ GeV}$?**
 First estimate: $\Delta g_{t\bar{t}h}/g_{t\bar{t}h} \sim 23\%$ for $m_H=120 \text{ GeV}$, $L=1 \text{ ab}^{-1}$ [AJ, 2002]
- However, at $\sqrt{s}=500 \text{ GeV}$ the $t\bar{t}$ dynamics is non-relativistic
 \Rightarrow **must use vNRQCD as in the $t\bar{t}$ threshold**



hep-ph/9910301
 hep-ph/0604034

\Rightarrow **Talk by Andre Hoang**



Considering $\sigma_{t\bar{t}h}$ enhancement due to:

- Large QCD resummation effect:
 $\sim x2.4$ for $m_h=120 \text{ GeV}$
- Use of beam polarization:
 $\sim x2.1$ for $(P(e^-), P(e^+)) = (-0.8, +0.6)$

Anticipate:

$(\Delta g_{t\bar{t}h}/g_{t\bar{t}h})_{\text{stat}} \sim 10\%$ for $m_H=120 \text{ GeV}$, $L=1 \text{ ab}^{-1}$

Top Quark and Extra-Dimensions (I)

- Solving the hierarchy problem is one of the main focus of models of New Physics (e.g. SUSY, Technicolor, extra-dimensions,..)

- Randall-Sundrum extra-dimensions:**

- 5D spacetime w/ one dimension (y) compactified on a S_1/Z_2 orbifold (AdS₅ geometry)
- Bounded by Planck and TeV branes

$$M_P^* = M_P e^{-\pi k R} \sim TeV \text{ (if } kR \sim 12)$$

⇒ solution to hierarchy problem

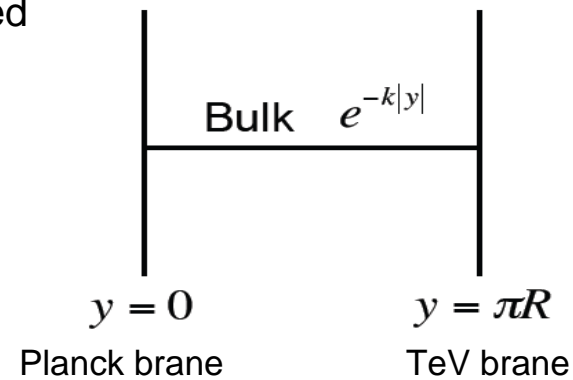
- Higgs on TeV brane
- Gauge bosons in 5D bulk ⇒ KK excitations

- CASE 1: fermions on TeV brane**

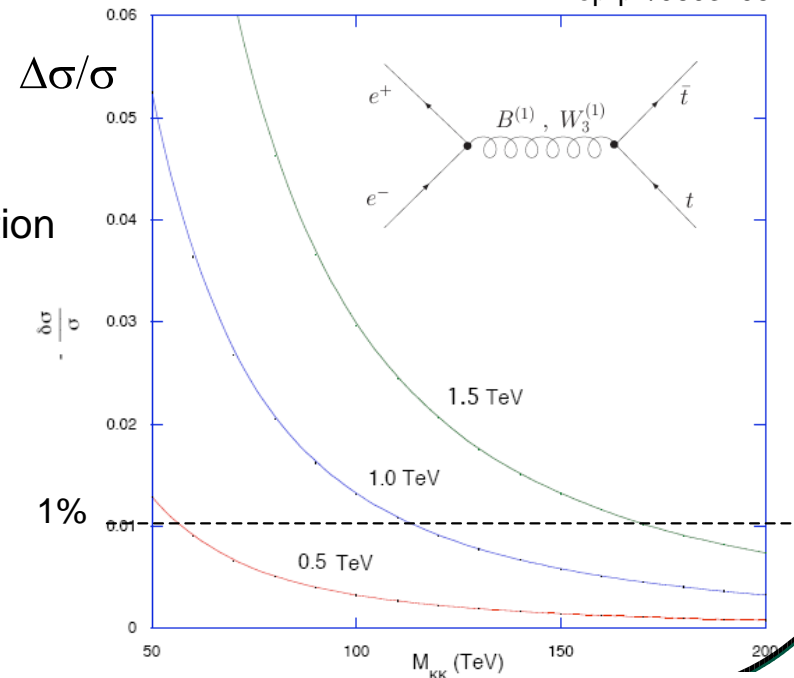
- Enhanced coupling to KK-gauge bosons
⇒ significant corrections to fermion pair production (here the top quark is not special)
- Severe constraints from precision EW data: $M_{KK} > 20$ TeV (inaccessible to LHC)
- Precise measurements at the ILC sensitive to very high mass.

⇒ Talk by Erin De Pree

$$ds^2 = e^{-2\sigma(y)} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$$



hep-ph/0603105

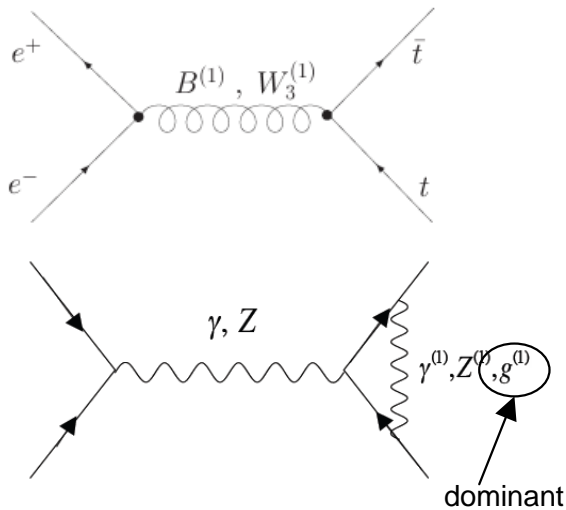


$M_{KK} \text{ (TeV)}$

Top Quark and Extra-Dimensions (II)

CASE 2: fermions in the bulk

- Precision EW constraints less severe: $M_{KK} > 10$ TeV
- Fermion mass hierarchy explained by “geography”
- Top quark is special: stronger coupling to KK-gauge bosons than the rest of fermions
- Different observables would be affected:



$$\frac{\delta\sigma}{\sigma} = (0.24\alpha_L + 0.14\alpha_R) \frac{s}{M_{KK}^2}$$

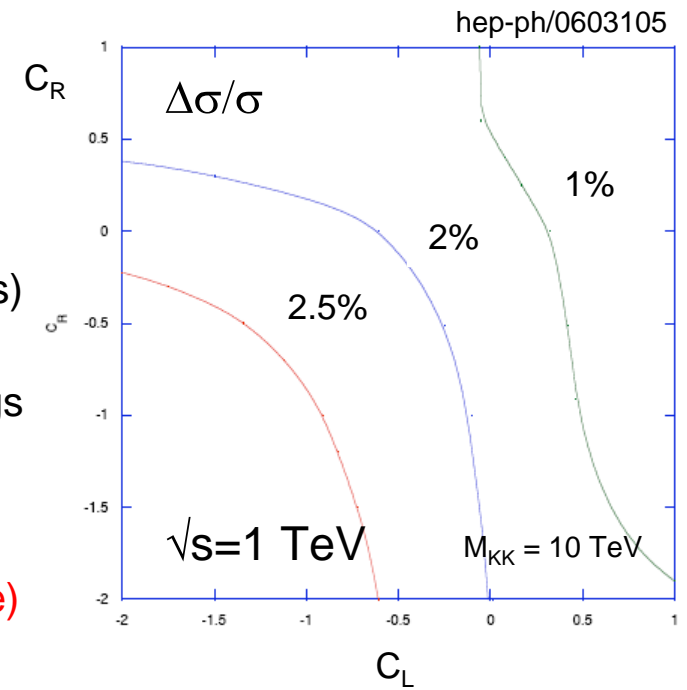
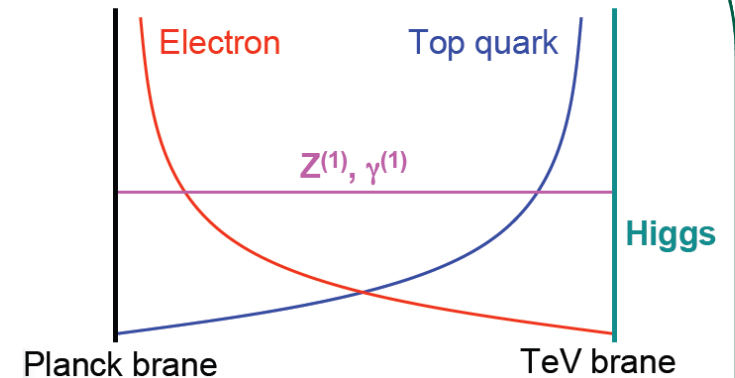
$$\delta A_{LR} = (0.26\alpha_L - 0.19\alpha_R) \frac{s}{M_{KK}^2}$$

Negligible effect on A_{FB}
(distinct from other extra-Z models)

Anomalous t-t- γ and t-t-Z couplings
(also possibly due to top/KK-top mixing effects)

- In summary, probe $M_{KK} > 50-150$ TeV (fermions on the brane) or up to $M_{KK} \sim 10$ TeV in large regions of parameter space (fermions off the brane).
- Future improvements: $P(e^+)$, top polarization observables,...

5 dimensional wavefunctions



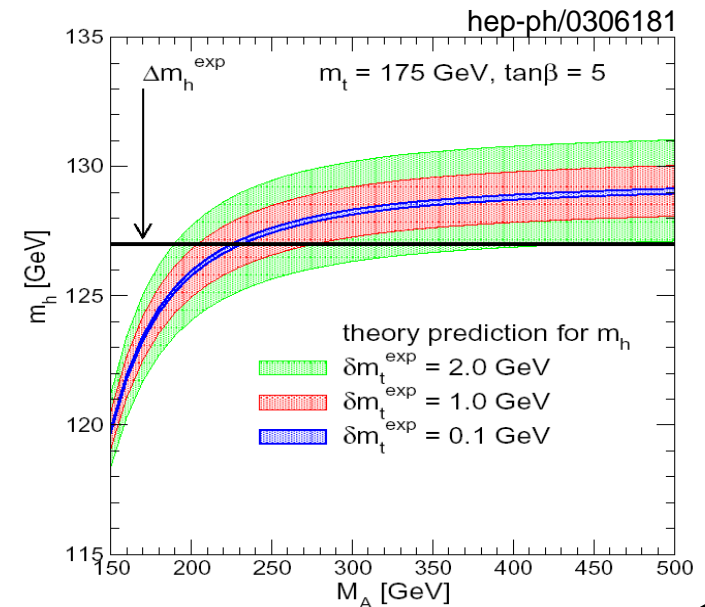
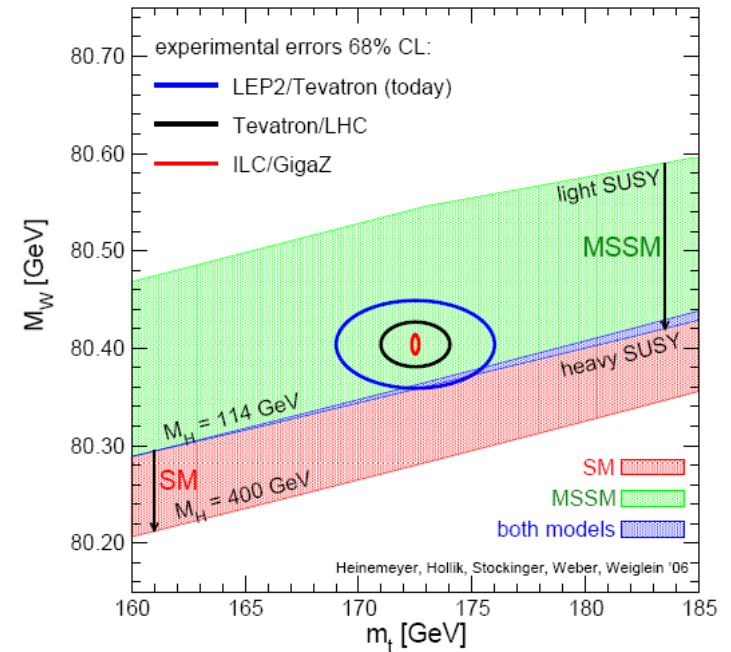
⇒ Talk by Erin De Pree

The Role of Precision Observables

- Important ingredient for EW precision analyses at the quantum level.
 - ⇒ incisive consistency checks
 - ⇒ constrain/rule out models of New Physics
 - ⇒ provide valuable information on the parameters of the Lagrangian
- Two prominent examples: M_W and m_t
- Experimental uncertainties:

	Today	Tevatron/LHC	ILC	“GigaZ”
ΔM_W [MeV]	32	20/15	10	7
Δm_t [GeV]	2.3	1.5	0.1	

- Theoretical uncertainties (within the SM):
 - Intrinsic: $\Delta M_W^{\text{intr}} \sim 4$ MeV (2 MeV future?)
 - Parametric (due to Δm_t):
 - Today: $\Delta m_t = 2.3$ GeV $\Rightarrow \Delta M_W^{\text{para,mt}} = 14$ MeV
 - LHC : $= 1.5$ GeV $\Rightarrow \Delta M_W^{\text{para,mt}} = 9$ MeV
 - ILC : $= 0.1$ GeV $\Rightarrow \Delta M_W^{\text{para,mt}} = 1$ MeV
- Need ILC precision on m_t :
 - Top match future accuracy on M_W (and $\sin^2\theta_{\text{eff}}$)
 - To exploit LHC (and ILC) precision on Higgs measurements
 - ...



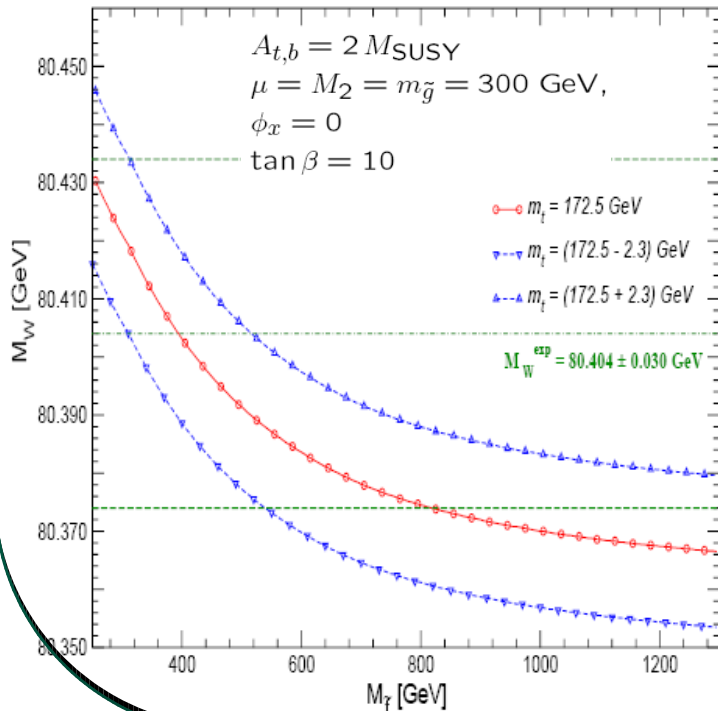
Precise M_W Prediction in the MSSM

- Precise theoretical predictions for electroweak observables required to fully exploit anticipated experimental accuracy.
- M_W prediction in terms of M_Z , α , G_μ and Δr :

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2} \right) = \frac{\pi \alpha}{\sqrt{2} G_\mu} (1 + \Delta r) \leftarrow \text{Loop corrections}$$

- Current intrinsic theoretical uncertainty within the SM ($\Delta M_W^{\text{intr}} \sim 4$ MeV): ~ok
- **What about within the MSSM?** Obtain best prediction possible by:
 - Using best available SM result
 - Implementing all available MSSM corrections

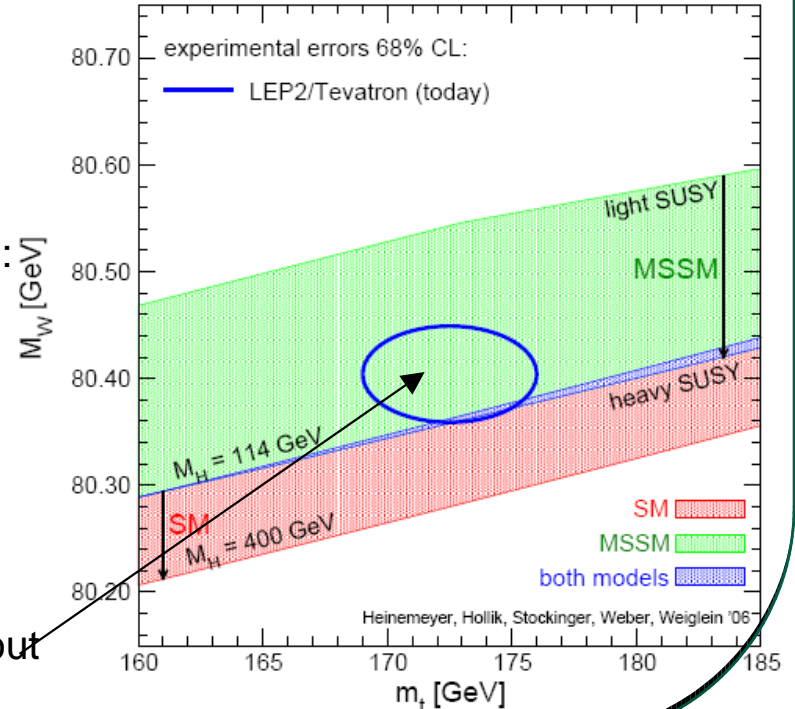
⇒ Talk by Sven Heinemeyer



Result:
 $\Delta M_W^{\text{intr}} \sim 5-11 \text{ MeV}$
 (must be improved)

Effect of different sectors:
 \tilde{t} / \tilde{b} : $\Delta M_W > 20 \text{ MeV}$
 $\tilde{\chi}^0 / \tilde{\chi}^\pm$: $\Delta M_W > 20 \text{ MeV}$
 ...

What can be learned about SUSY parameter space?

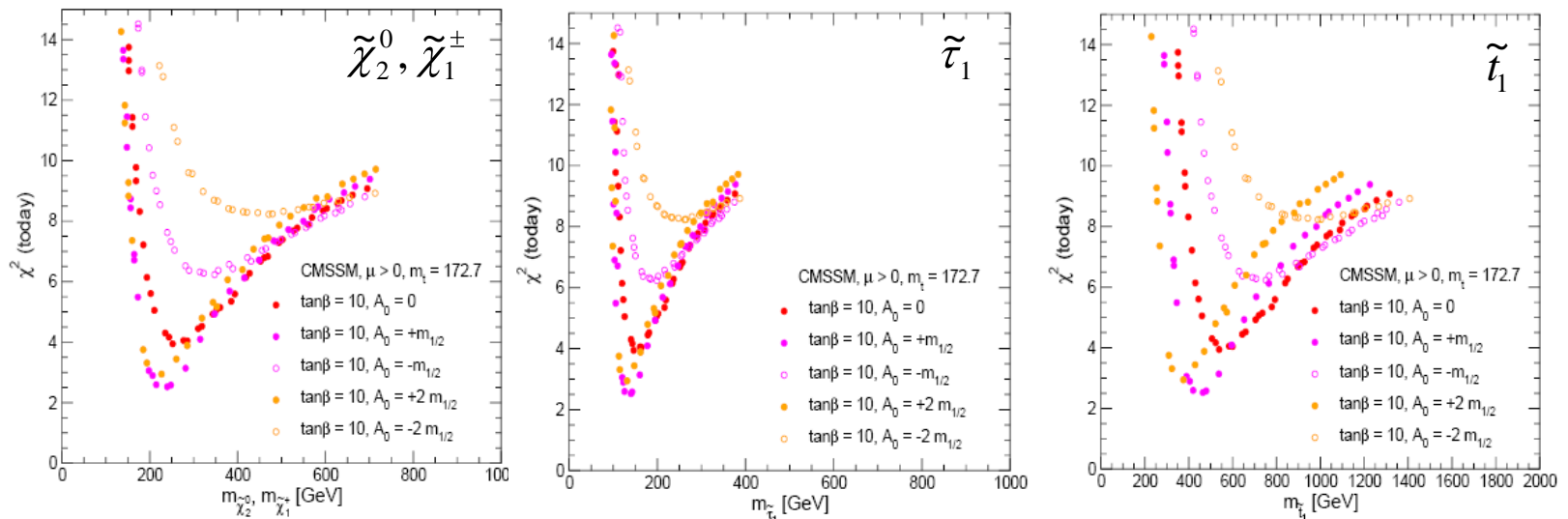


Prospects for SUSY at the ILC

- **Idea:** combine different precision observables and look for favored regions of parameter space.
- **Observables:** M_W , $\sin^2\theta_{\text{eff}}$, $B(b \rightarrow s\gamma)$, $(g-2)_\mu$, M_H
 - Build χ^2 taking into account experimental and theoretical (parametric+intrinsic) uncertainties
 - Obtain best fit for masses.
- Too many free parameters in MSSM.

⇒ Talk by Sven Heinemeyer
- ⇒ Consider a number of scenarios with reduced # free parameters:
 - CMSSM (mSUGRA): 5 parameters (m_0 , $m_{1/2}$, A_0 , $\tan\beta$, $\text{sign}(\mu)$)
- Impose hard constraint: LSP gives right amount of Cold Dark Matter.
- Rather good fit to current data ($\chi^2 \sim \mathcal{O}(2)$).

$\tan\beta = 10$



- **Similar conclusion regarding ILC reach in different scenarios:**
 - $\tan\beta = 10$: sleptons, charginos, neutralinos (partially) in reach. Some chance for light stop.
 - $\tan\beta = 50$: sleptons, charginos, neutralinos (partially) in reach. Hardly any chance for light stop or gluinos.
- ⇒ Rather good prospects for the ILC

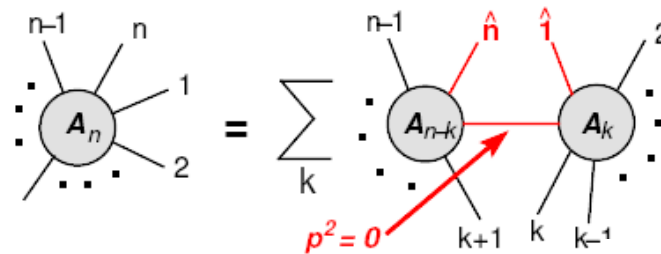
New Approaches to Higher Order Calculations

- Many physics processes at the ILC involve a large number of final state particles.
- Feynman diagram method is not optimized for multi-leg processes.
- **More efficient methods have been developed that exploit symmetries in QCD leading to “recursion relations”**

⇒ Talk by Carola Berger

• *“Trees are recycled into trees”*

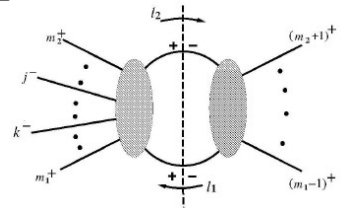
- 6 gluon amplitude in agreement w/ previous calculations [Giele et al]
- Many applications: massive particles, SUSY,...



• Method also works at one-loop level and remains “sustainable”:

- Some results for parts of Higgs+gluons @ NNLO
- A-priori no obstacles towards phenomenology

“Trees are recycled into loops”
“Loops are recycled into loops”



- In certain kinematic regions, precise theoretical calculations involve resummation of large logs.
- **Exploring application of Soft Collinear Effective Theory (SCET; developed for b-physics) for resummation in collider physics.**
- Application: large momentum transfer Q^2 , but hadronic final states with small M_X ($\alpha_s^n \ln^{2n}(M_X/Q) \uparrow$)
- Traditionally, resummation is performed in moment space.
Caveats: Landau poles, Mellin inversion only numerically
- **Solving Renormalization Group equations in SCET, one obtains resummed expressions directly in momentum space.**

- Clear scale separation. No Landau pole ambiguities.
- Analytic expressions for resummed rates.
- Simple connection with fixed order expressions.

⇒ Talk by Thomas Becher

Conclusions

- After the first signs of New Physics at the LHC, precision measurements at the ILC will likely catalyze the next revolution in Physics by helping answer truly fundamental questions.
- **Important to keep a dynamic and active Precision Physics Working Group:**
 - The physics case needs to keep being sharpened all the time;
 - An increased involvement of experimentalists crucial to establish a feedback loop between the detector and its physics (an existing gap that needs to be bridged);
 - Theorists must find enough incentives to work on the precise theoretical predictions/tools that will be required to carry out the ILC physics program (in many cases involve multi-year projects).
- Finally, many thanks to our speakers: **Thomas Becher, Carola Berger, Erin De Pree, Filimon Gounaris, Sven Heinemeyer, Andre Hoang,** and the rest of participants, for contributing to very productive and stimulating sessions.

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