

SLOOPs

An automatic program for full one-loop calculations in the SM/MSSM

Guillaume CHALONS
LAPTH, Annecy-le-Vieux

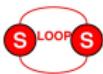
The **SloopS** team :

Fawzi BOUDJEMA, Guillaume DRIEU LA ROCHELLE, Sun Hao
(LAPTH, Annecy-le-Vieux)

Nans BARO
(ITTK, RWTH Aachen University)

Andreï SEMENOV
(JINR Dubna)

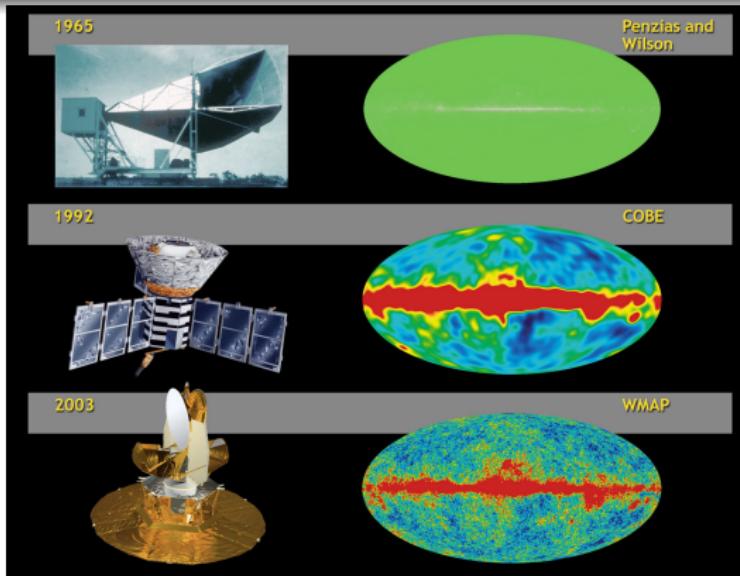
IWLC 2010



NEED FOR PRECISE THEORETICAL PREDICTIONS

RELIC DENSITY OF DARK MATTER

- WMAP : $0.0975 < \Omega_{DM} h^2 < 0.1223$ (10% precision)
- PLANCK : 2% precision



PRECISION MEASUREMENTS

RELIC DENSITY IN THE STANDARD SCENARIO

$$\Omega_{DM} h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 \text{s}^{-1}}{\langle \sigma(\chi\chi \rightarrow SM) v \rangle}$$

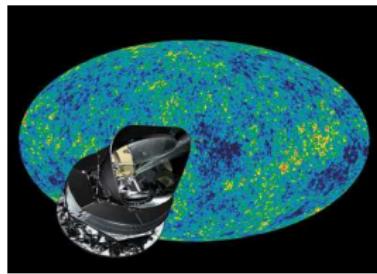
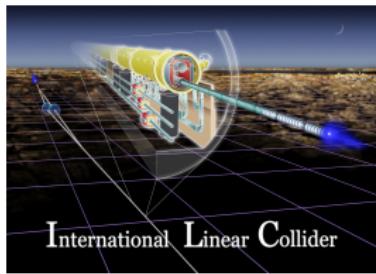


RELIC DENSITY IN THE STANDARD SCENARIO

$$\Omega_{DM} h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 \text{s}^{-1}}{\langle \sigma(\chi\chi \rightarrow SM)v \rangle}$$

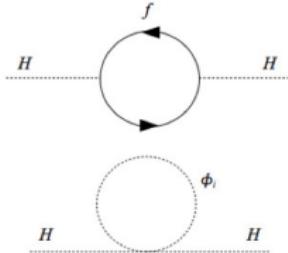
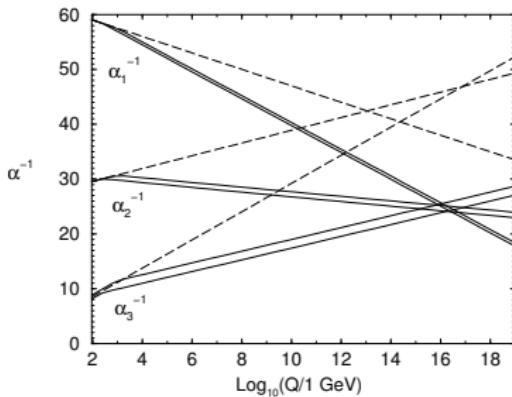
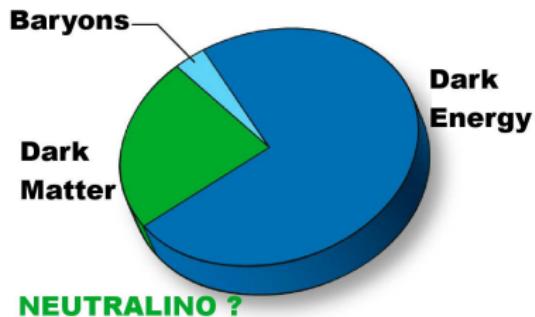
PRECISION

- Need for precise theoretical predictions w.r.t experimental measurements.
- Precision needed at the level of $\sigma \Rightarrow$ One-loop calculations (at least).
- Reconstruction of fundamental underlying parameters at LHC and LC.
- Radiative corrections must be under control to be able to constrain the cosmological underlying scenario.
- What precision required at colliders and theory to constrain cosmology ?

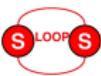


ATTRACTIVE FEATURES

- Stabilization of the scalar sector
- Better Unification of coupling constants
- Dark Matter Candidate(s)
- ...



$$\Delta M_H^2 \sim \frac{\lambda_f^2}{4\pi^2} [(m_f^2 - m_S^2) \log(\frac{\Lambda}{m_s})]$$



COMPLICATIONS

- Not observed yet, neither the Higgs boson...
- \mathcal{L}_{soft} unknown.
- Lots of free parameters ($\simeq 105$).
- Calculations become extremely tedious and involved.

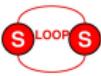


COMPLICATIONS

- Not observed yet, neither the Higgs boson...
- $\mathcal{L}_{\text{soft}}$ unknown.
- Lots of free parameters ($\simeq 105$).
- Calculations become extremely tedious and involved.

BEYOND LEADING ORDER IN SUSY

- At LO : $m_h < m_Z$ but no Higgs found !
- LEP Bound on Higgs mass $m_h > 114 \text{ GeV}$
- At higher orders : Higgs mass can get large corrections.
- Generically SUSY processes get large radiative corrections.
- Calculations become even more complicated...



COMPLICATIONS

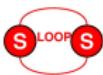
- Not observed yet, neither the Higgs boson...
- $\mathcal{L}_{\text{soft}}$ unknown.
- Lots of free parameters ($\simeq 105$).
- Calculations become extremely tedious and involved.

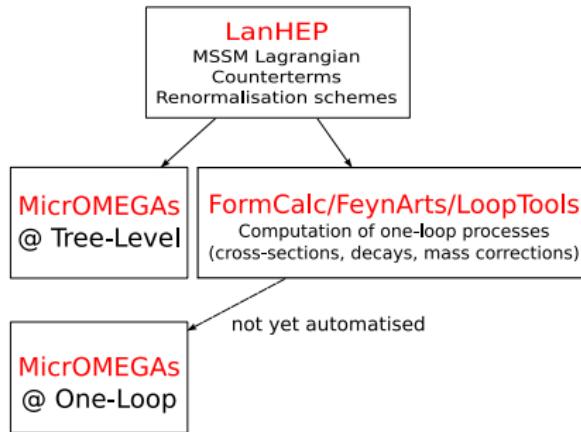
BEYOND LEADING ORDER IN SUSY

- At LO : $m_h < m_Z$ but no Higgs found !
- LEP Bound on Higgs mass $m_h > 114 \text{ GeV}$
- At higher orders : Higgs mass can get large corrections.
- Generically SUSY processes get large radiative corrections.
- Calculations become even more complicated...

RADIATIVE CORRECTIONS ARE
IMPORTANT

AUTOMATION NEEDED





SLOOPs

A code for calculation of **loops** diagrams in the MSSM with application to **colliders**, **astrophysics** and **cosmology**.

- Evaluation of one-loop diagrams including a **complete** and **coherent** renormalisation of **each sector** of the MSSM with an **OS** scheme.
- Modularity between different renormalisation schemes.
- **Non-linear** gauge fixing.
- Checks : results **UV**,**IR** finite and **gauge** independent.

<http://code.sloops.free.fr/>



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model $m_f, \alpha(0), M_W, M_Z$



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model $m_F, \alpha(0), M_W, M_Z$

HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model $m_F, \alpha(0), M_W, M_Z$

HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence
- MH : δt_β is defined from the measurement of the mass m_H



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model

 $m_f, \alpha(0), M_W, M_Z$

HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence
- MH : δt_β is defined from the measurement of the mass m_H
- $A^0\tau\tau$: δt_β is defined from the decay $A^0 \rightarrow \tau^+\tau^-$ ($vertex \propto m_\tau t_\beta$)



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model $m_f, \alpha(0), M_W, M_Z$

HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence
- MH : δt_β is defined from the measurement of the mass m_H
- $A^0\tau\tau$: δt_β is defined from the decay $A^0 \rightarrow \tau^+\tau^-$ ($vertex \propto m_\tau t_\beta$)



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model $m_f, \alpha(0), M_W, M_Z$

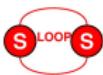
HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence
- MH : δt_β is defined from the measurement of the mass m_H
- $A^0\tau\tau$: δt_β is defined from the decay $A^0 \rightarrow \tau^+\tau^-$ ($vertex \propto m_\tau t_\beta$)

SFERMIONS SECTOR

Input parameters : 3 sfermions masses $m_{\tilde{d}_1}, m_{\tilde{d}_2}, m_{\tilde{u}_1}$ and 2 conditions for $A_{u,d}$



FERMION + GAUGE SECTOR

Input parameters as in the Standard Model

 $m_f, \alpha(0), M_W, M_Z$

HIGGS SECTOR

Input parameters : $M_{A^0}, t_\beta = v_2/v_1$. Several definitions for δt_β :

- \overline{DR} : δt_β is a pure divergence
- MH : δt_β is defined from the measurement of the mass m_H
- $A^0\tau\tau$: δt_β is defined from the decay $A^0 \rightarrow \tau^+\tau^-$ ($vertex \propto m_\tau t_\beta$)

SFERMIONS SECTOR

Input parameters : 3 sfermions masses $m_{\tilde{d}_1}, m_{\tilde{d}_2}, m_{\tilde{u}_1}$ and 2 conditions for $A_{u,d}$

NEUTRALINOS/CHARGINOS SECTOR

Input parameters : 2 charginos $m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_2^\pm}$ and 1 neutralino $\tilde{\chi}_1^0$ 

Linear gauge fixing

$$\begin{aligned}\mathcal{L}_{GF} = & -\frac{1}{\xi_W} |\partial_\mu W^{\mu+} + i\xi_W \frac{g}{2} v G^+|^2 \\ & -\frac{1}{2\xi_Z} (\partial_\mu Z^\mu + \xi_Z \frac{g}{2c_w} v G^0)^2 \\ & -\frac{1}{2\xi_A} (\partial_\mu A^\mu)^2\end{aligned}$$

$$\Gamma^{VV} = \frac{-i}{q^2 - M_V^2 + i\epsilon} \left[g_{\mu\nu} + (\xi_V - 1) \frac{q_\mu q_\nu}{q^2 - \xi_V M_V^2} \right]$$



Linear gauge fixing

$$\begin{aligned}\mathcal{L}_{GF} = & -\frac{1}{\xi_W} |\partial_\mu W^\mu + i\xi_W \frac{g}{2} v G^+|^2 \\ & -\frac{1}{2\xi_Z} (\partial_\mu Z^\mu + \xi_Z \frac{g}{2c_w} v G^0)^2 \\ & -\frac{1}{2\xi_A} (\partial_\mu A^\mu)^2\end{aligned}$$

$$\Gamma^{VV} = \frac{-i}{q^2 - M_V^2 + i\epsilon} \left[g_{\mu\nu} + (\xi_V - 1) \frac{q_\mu q_\nu}{q^2 - \xi_V M_V^2} \right]$$

$$\xi_{W,Z,A} = 1 \text{ (Feynman gauge)}$$



Non-Linear gauge fixing

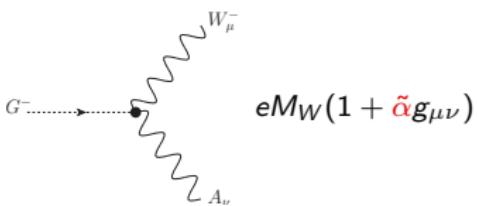
$$\begin{aligned}
 \mathcal{L}_{GF} = & -\frac{1}{\xi_W} |(\partial_\mu - ie\tilde{\alpha}A_\mu - igc_w\tilde{\beta}Z_\mu)W^\mu + \\
 & + i\xi_W \frac{g}{2}(v + \tilde{\delta}h^0 + \tilde{\omega}H^0 + i\tilde{\kappa}G^0 + i\tilde{\rho}A^0)G^+|^2 \\
 & - \frac{1}{2\xi_Z} (\partial_\mu Z^\mu + \xi_Z \frac{g}{2c_w}(v + \tilde{\epsilon}h^0 + \tilde{\gamma}_H^0)G^0)^2 \\
 & - \frac{1}{2\xi_A} (\partial_\mu A^\mu)^2
 \end{aligned}$$

$\xi_{W,Z,A} = 1$ (Feynman gauge)



Non-Linear gauge fixing

$$\begin{aligned}\mathcal{L}_{GF} = & -\frac{1}{\xi_W} |(\partial_\mu - ie\tilde{\alpha}A_\mu - igc_w\tilde{\beta}Z_\mu)W^\mu + \\ & + i\xi_W \frac{g}{2}(v + \tilde{\delta}h^0 + \tilde{\omega}H^0 + i\tilde{\kappa}G^0 + i\tilde{\rho}A^0)G^+|^2 \\ & - \frac{1}{2\xi_Z} (\partial_\mu Z^\mu + \xi_Z \frac{g}{2c_w}(v + \tilde{\epsilon}h^0 + \tilde{\gamma}_H^0)G^0)^2 \\ & - \frac{1}{2\xi_A} (\partial_\mu A^\mu)^2\end{aligned}$$



$$\xi_{W,Z,A} = 1 \text{ (Feynman gauge)}$$

- Gauge parameter dependence in gauge/Goldstone/ghost vertices.
- No "unphysical" threshold, no higher rank tensor.



FIRST CHECKS ON THE CODE

TREE LEVEL CALCULATIONS

Comparison with public codes : Grace and CompHEP

Nans Baro PhD Thesis

Cross-section [pb]	SloopS	CompHEP	Grace
$h^0 h^0 \rightarrow h^0 h^0$	3.932×10^{-2}	3.932×10^{-2}	3.929×10^{-2}
$W^+ W^- \rightarrow \tilde{t}_1 \bar{\tilde{t}}_1$	7.082×10^{-1}	7.082×10^{-1}	7.083×10^{-1}
$e^+ e^- \rightarrow \tilde{\tau}_1 \bar{\tilde{\tau}}_2$	2.854×10^{-3}	2.854×10^{-3}	2.854×10^{-3}
$H^+ H^- \rightarrow W^+ W^-$	6.643×10^{-1}	6.643×10^{-1}	6.644×10^{-1}
Decay [GeV]		 # 200 processes checked
$A^0 \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-$	1.137×10^0	1.137×10^0	1.137×10^0
$\tilde{\chi}_1^+ \rightarrow t \bar{b}_1$	5.428×10^0	5.428×10^0	5.428×10^0
$H^0 \rightarrow \tilde{\tau}_1 \bar{\tilde{\tau}}_1$	7.579×10^{-3}	7.579×10^{-3}	7.579×10^{-3}
$H^+ \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^0$	1.113×10^{-1}	1.113×10^{-1}	1.113×10^{-1}

... # 200 processes checked

ONE-LOOP PROCESSES THAT DO NOT NEED RENORMALISATION

Comparison with public codes : PLATON and DarkSUSY

Implementation of a special routine for loop integrals at $v = 0$

Boudjema, Semenov, Temes, Phys. Rev. D72, 055024 (2005)

- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow \gamma\gamma$
- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow gg$
- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow \gamma Z^0$

APPLICATIONS IN THE HIGGS SECTOR

N. Baro., F. Boudjema, A. Semenov, *Phys. Lett.* **B660** (2008) 550, 0710.1821 [hep-ph]

- One-loop corrections to Higgs masses H^+, h^0 Freitas, Stockinger, *Phys. Rev.* **D66** (2002) 095014, hep-ph/0205281

$t_\beta = 3$	$mhmax$	$large \mu$	$nomix$
Tree Level	72.51	72.51	72.51
DCPR	134.28	97.57	112.26
MH	140.25	86.68	117.37
$A\tau\tau$	134.25	97.59	112.27
$\overline{DR} \bar{\mu} = m_{A^0}$	134.87	98.10	112.86
Light Higgs mass m_{h^0}			

- $A^0 \rightarrow \tau^+ \tau^-$, $A^0 \rightarrow Z^0 h^0$, $H^0 \rightarrow Z^0 Z^0$, $H^0 \rightarrow \tau^+ \tau^-$

$t_\beta = 3$	$mhmax$	$large \mu$	$nomix$
Tree Level	9.35×10^{-3}	9.35×10^{-3}	9.35×10^{-3}
DCPR	-1.09×10^{-4}	-7.96×10^{-5}	-1.09×10^{-4}
MH	$+6.28 \times 10^{-3}$	-7.91×10^{-3}	$+4.47 \times 10^{-3}$
$A\tau\tau$	-1.45×10^{-4}	-7.09×10^{-5}	-1.01×10^{-4}
$\overline{DR} \bar{\mu} = m_{A^0}$	$+5.08 \times 10^{-4}$	$+3.24 \times 10^{-4}$	$+4.17 \times 10^{-4}$
$H^0 \rightarrow \tau^+ \tau^-$ at one-loop with no QED			

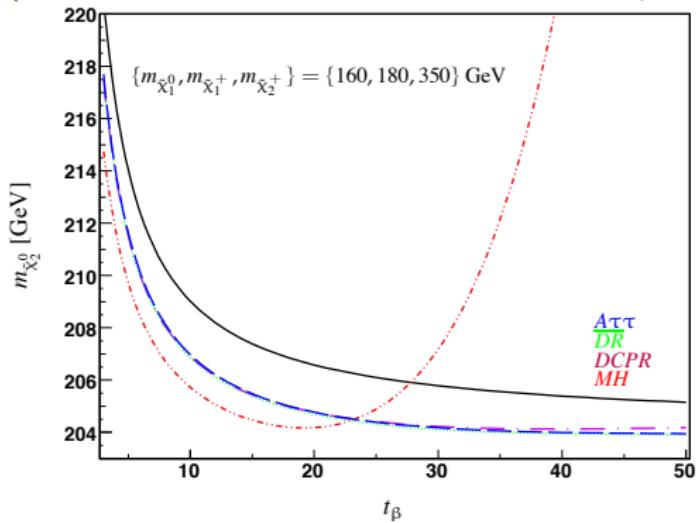
- Theoretical issue due to non-linear gauge fixing and modified Ward-Slavnov-Taylor Identity in the Higgs sector :

$$m_{A^0}^2 \times A^0 \dashrightarrow \textcircled{O} \dashrightarrow Z^0 + m_{Z^0} \times A^0 \dashrightarrow \textcircled{O} \dashrightarrow G^0 = (m_{A^0}^2 - m_{Z^0}^2) \frac{ie}{s_{2W}} [\bar{\epsilon} \times \textcircled{O}_{h^0}^{G^0} \dashrightarrow A^0 + \bar{\gamma} \times \textcircled{O}_{H^0}^{G^0} - \textcircled{S}^{A^0 \text{ loop}}]$$

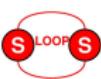
APPLICATIONS IN THE CHARGINO/NEUTRALINO SECTOR

N. Baro, F. Boudjema, *Phys. Rev. D* **80** (2009) 076010, arXiv :0906.1665[hep-ph].

- One-loop corrections to neutralino masses $\tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$
(see also T. Fritzsch, W. Hollik, *Eur. Phys. J.* **C24** (2002) 619, hep-ph/0203159.)

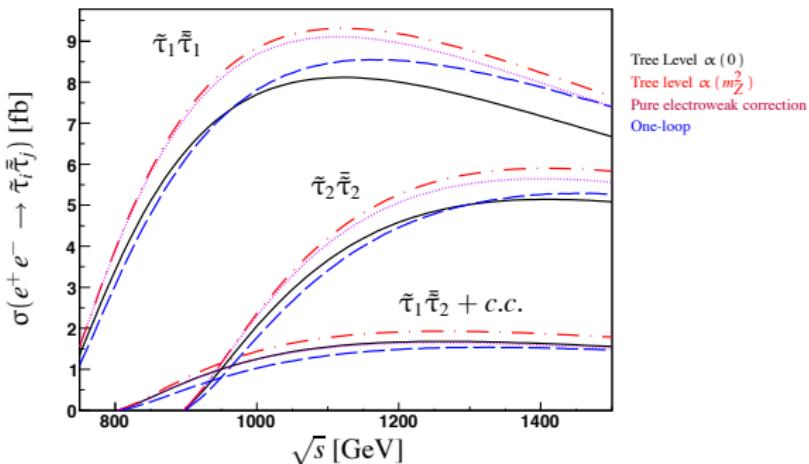


- Chargino decays at one-loop
(comparison with J. Fujimoto et al., *Phys. Rev. D* **75** (2007) 113002, hep-ph/0701200.)



APPLICATIONS TO COLLIDER PHYSICS

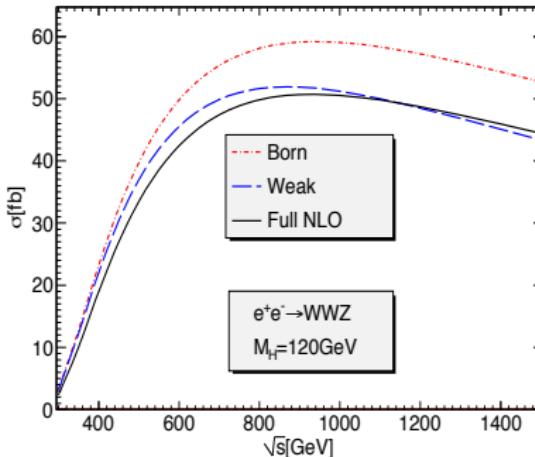
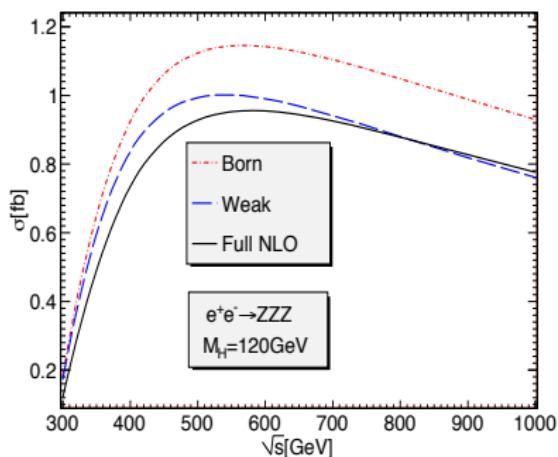
- $e^+ e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-$ J. Fujimoto et al., Phys. Rev. D75 (2007) 113002, hep-ph/0701200.
- $e^+ e^- \rightarrow \tilde{\tau}_i \tilde{\tau}_j^*$ K. Kovarik et al., Phys. Rev. D72 (2005) 053010, hep-ph/0506021.



APPLICATIONS TO PURE STANDARD MODEL PROCESSES

F. Boudjema, Le Duc Ninh, Sun Hao, M. M. Weber, *Phys. Rev.* **D81** 073007 (2010)

- $e^+e^- \rightarrow W^+W^-Z^0$
- $e^+e^- \rightarrow Z^0Z^0Z^0$
- Important processes to test the quartic gauge couplings and Higgs mechanism

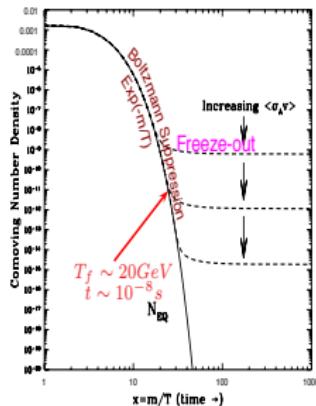


See also Su Ji-Juan et al. *Phys.Rev* **D78** 016007, Sun Wei et al. *Phys.Lett* **B680**, 321



THERMAL RELIC

- $\Omega_\chi h^2 \propto 1/(\sigma(\chi\chi \rightarrow \text{SM}))$
- Relic density calculated through the interface of **SloopS** with **micrOMEGAs** (Bélanger et al.)



BUNCH OF FULL ONE-LOOP PROCESSES CALCULATED

Baro, Boudjema, Semenov, *Phys. Lett.* **B660**

Baro, Boudjema, G.C, Sun Hao, *Phys. Rev.* **D81** 015005 (2010) (2008) 550

- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow f\bar{f}$ (**bino**)
- $\tilde{\chi}_1^0 \tilde{\tau}_1^+ \rightarrow \tau^+ \gamma (Z^0)$ (**bino**)
- $\tilde{\tau}_1^+ \tilde{\tau}_1^+ \rightarrow \tau^+ \tau^+$ (**bino**)
- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow W^+ W^-, Z^0 Z^0$ (**bino-wino**, **bino-higgsino**, **higgsino**, **higgsino-bino**, **wino**)
- $\tilde{\chi}_1^0 \tilde{\chi}_1^+ \rightarrow u\bar{d}, t\bar{b}$ (**bino-wino**, **higgsino**, **higgsino-bino**, **wino**)
- $\tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow W^+ W^-, Z^0 Z^0$ (**wino**)

- What is required from **collider** data to get a **precise** prediction of $\Omega_\chi h^2$?
(see Allanach et al. JHEP 0412:020, 2004.)
- What are the relevant observables to **control** the **uncertainty** on the predicted $\Omega_\chi h^2$?
- What is the **required accuracy** in order to **achieve** PLANCK precision?

Basic example as an illustration

- Typical mSUGRA scenario : LSP is $\tilde{\chi}_1^0$ in the bulk region
- $\tilde{\chi}_1^0 \tilde{\chi}_1^0 \rightarrow \ell\bar{\ell}$ through R-sleptons for $\Omega_\chi h^2$
- At tree-level $m_{\tilde{\chi}_1^0}$, $m_{\tilde{\ell}}(m_{\tilde{\tau}_1})$ + mixing matrix of neutralinos and in the $\tilde{\tau}_1$ sector
needed to reconstruct $\Omega_\chi h^2$.

At one-loop level ?

- Sensitivity to **parameters** entering in **loops**? (non decoupling corrections, thresholds...)
- Sensitivity to **renormalisation schemes**?

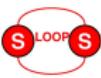


SENSITIVITY TO SQUARKS MASSES

	M_1	M_2	μ	t_β	$M_{\tilde{e}_R}$	$M_{\tilde{e}_L}$	M_3	M_{A^0}
Masses (GeV)	90	200	-600	5	110	250	800	500

$$250 \text{ GeV} \leq M_{\tilde{Q}} \leq 800 \text{ GeV}$$

- At tree-level $\Omega_\chi h^2$ sensitive to $M_{\tilde{Q}}$ when $M_{\tilde{Q}} \simeq 300$ GeV (new channels open).
- Effects of squarks relevant in loops for $M_{\tilde{Q}} \geq 300$ GeV ?

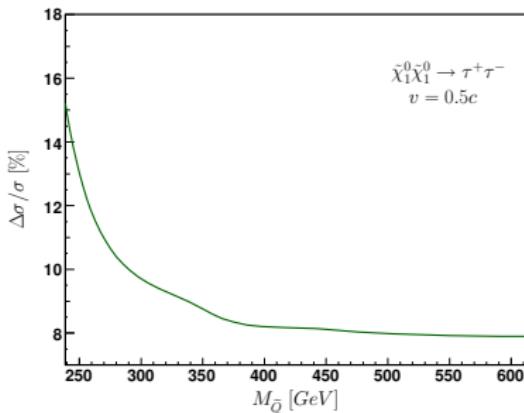


SENSITIVITY TO SQUARKS MASSES

	M_1	M_2	μ	t_β	$M_{\tilde{e}_R}$	$M_{\tilde{e}_L}$	M_3	M_{A^0}
Masses (GeV)	90	200	-600	5	110	250	800	500

$$250 \text{ GeV} \leq M_{\tilde{Q}} \leq 800 \text{ GeV}$$

- At tree-level $\Omega_\chi h^2$ sensitive to $M_{\tilde{Q}}$ when $M_{\tilde{Q}} \simeq 300$ GeV (new channels open).
- Effects of squarks relevant in loops for $M_{\tilde{Q}} \geq 300$ GeV ?

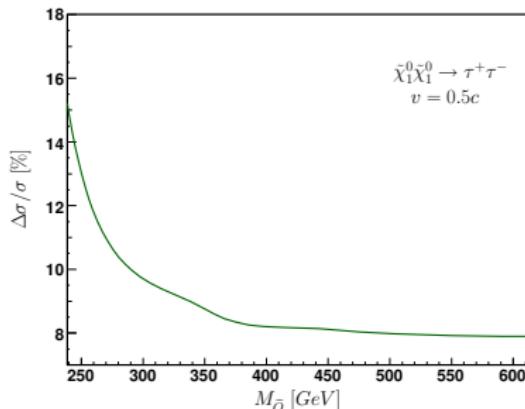


SENSITIVITY TO SQUARKS MASSES

	M_1	M_2	μ	t_β	$M_{\tilde{e}_R}$	$M_{\tilde{e}_L}$	M_3	M_{A^0}
Masses (GeV)	90	200	-600	5	110	250	800	500

$$250 \text{ GeV} \leq M_{\tilde{Q}} \leq 800 \text{ GeV}$$

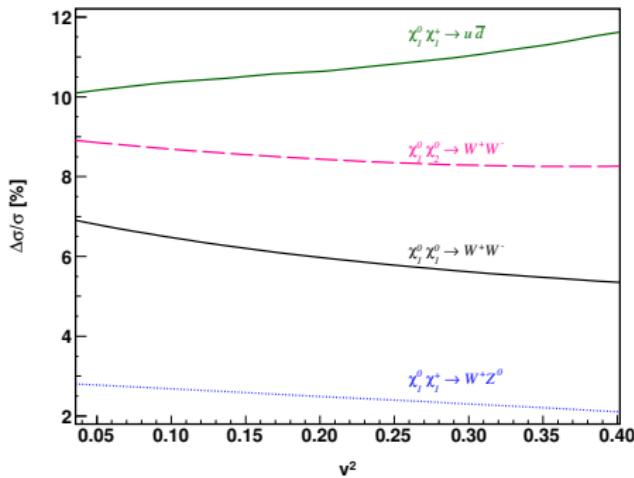
- At tree-level $\Omega_\chi h^2$ sensitive to $M_{\tilde{Q}}$ when $M_{\tilde{Q}} \simeq 300$ GeV (new channels open).
- Effects of squarks relevant in loops for $M_{\tilde{Q}} \geq 300$ GeV ?



More or less the same conclusion at one-loop level. This may not be the case for other set of parameters.

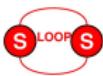
SENSITIVITY TO RENORMALISATION SCHEME

Parameter	M_1	M_2	μ	t_β	M_3	$M_{\tilde{L}, \tilde{Q}}$	A_i	M_{A^0}
Value	110	134.5	-245	10	600	600	0	600
$\tilde{\chi}_1^0 = 0.94\tilde{B} - 0.20\tilde{W} - 0.27\tilde{H}_1^0 - 0.10\tilde{H}_2^0$								



- Bulk of corrections to the **s-wave** coefficient
- Large δt_β scheme dependence
- QCD corrections to $u\bar{d} \simeq 2.5\%$

	$A_{\tau\tau}$	\overline{DR}	MH
$\Omega_\chi h^2$	0.105	0.102	0.097
$\delta \Omega h^2 / \Omega h^2$	-2.8%	-5.6%	-10.2%



CONCLUSIONS AND PERSPECTIVES

- Complete EW renormalisation of the MSSM and modularity with different schemes.
- One-loop corrections to masses, decays, cross sections at colliders.
- One-loop corrections to neutralino annihilation relevant for relic density and indirect detection.
- First steps done for the connection with micrOMEGAs .
- In any case for $\Omega_\chi h^2$ @ 1-2% \Rightarrow one-loop corrections mandatory.
- Then at one-loop level more input is needed for an efficient reconstruction of parameters compared to the tree level case.
- Gather all available data to construct efficient renormalisation schemes.
- Implementation of QCD renormalisation in SloopS ongoing.

