

Testing effective Yukawa couplings in Higgs studies at Linear Colliders

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in collaboration with B. Mele
talk based on [arxiv:1005.2498](https://arxiv.org/abs/1005.2498)

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What if Higgs boson is only responsible of M_W , M_Z but not of fermion masses ?

- fermion masses $m_f \longrightarrow$ ChSB
- in SM, ChSB and EWSB (M_W, M_Z) generated by the Higgs mechanism at same scale $\sim \langle H \rangle$

$$\mathcal{L} = Y_f \bar{\psi}_f \psi_f H \longrightarrow m_f = Y_f \langle H \rangle$$

- not (yet) any experimental evidence supporting tree-level Yukawa couplings Y_f
- maybe ChSB & EWSB have different mechanisms
 \rightarrow compositeness, extra-dimensions, technicolor...

■ Anyhow, there are indications that SM Higgs could be at the origin of M_W and M_Z

- EW precision tests favour a light Higgs
- not very sensitive to Y_f
- $H \rightarrow$ perturbative unitarity in $WW \rightarrow WW$

■ What if fermion masses put in by hand ?

- no tree-level $Y_f \rightarrow$ NO spontaneous ChSB
- SM becomes non-renormalizable, but can be considered as an effective field theory

Proposed Scenario

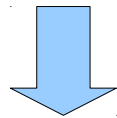
- **mf** generated by some new mechanism set at a high-energy scale $\Lambda \gg EW$ scale
- **NO tree-level** Yukawa couplings
- **ChSB mechanism** assumed to give a small contribution to EWSB
- **H** providing the main contribution to **M_w** and **M_z**
- **only SM degrees of freedom** propagating below Λ , with a light **H**

- Λ  scale where Y_f are vanishing
(i.e. scale of fermion mass generation)

- Y_f not protected against radiative corrections
due ChSB \rightarrow radiatively generated

- large logs $\underbrace{g_i^{2n} \log^n (\Lambda/m_H)}_{\text{can be summed up}}$
by solving Renormalization Group Equations

- SM RGE are not suitable due to tree-level Yukawas



RGE derived by keeping Y_f and m_f as
independent parameters

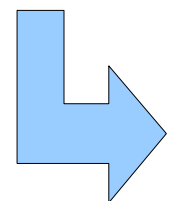
1-loop RGE for Yukawa couplings

$$\begin{aligned} \frac{d\mathbf{Y}_U}{dt} &= \frac{1}{16\pi^2} \left\{ 3\xi_H^2 (\mathbf{Y}_U - \mathbf{Y}_U^{\text{SM}}) - 3\mathbf{Y}_U^{\text{SM}}\mathbf{Y}_D^{\text{SM}} (\mathbf{Y}_D - \mathbf{Y}_D^{\text{SM}}) + \frac{3}{2}\mathbf{Y}_U (\mathbf{Y}_U\mathbf{Y}_U - \mathbf{Y}_D^{\text{SM}}\mathbf{Y}_D^{\text{SM}}) \right. \\ &\quad \left. - \mathbf{Y}_U \left(\frac{17}{20}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 - \text{Tr}(\mathbf{Y}) \right) \right\}, \end{aligned} \quad (1)$$

$$\begin{aligned} \frac{d\mathbf{Y}_D}{dt} &= \frac{1}{16\pi^2} \left\{ 3\xi_H^2 (\mathbf{Y}_D - \mathbf{Y}_D^{\text{SM}}) - 3\mathbf{Y}_D^{\text{SM}}\mathbf{Y}_U^{\text{SM}} (\mathbf{Y}_U - \mathbf{Y}_U^{\text{SM}}) + \frac{3}{2}\mathbf{Y}_D (\mathbf{Y}_D\mathbf{Y}_D - \mathbf{Y}_U^{\text{SM}}\mathbf{Y}_U^{\text{SM}}) \right. \\ &\quad \left. - \mathbf{Y}_D \left(\frac{1}{4}g_1^2 + \frac{9}{4}g_2^2 + 8g_3^2 - \text{Tr}(\mathbf{Y}) \right) \right\}, \end{aligned} \quad (2)$$

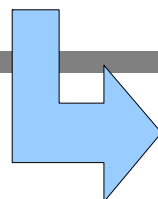
W(L) polarizations

$$\mathbf{Y} \equiv N_c \mathbf{Y}_U \mathbf{Y}_U + N_c \mathbf{Y}_D \mathbf{Y}_D + \mathbf{Y}_E \mathbf{Y}_E$$



$$\xi_H \equiv \frac{g_2 m_H}{2M_W},$$

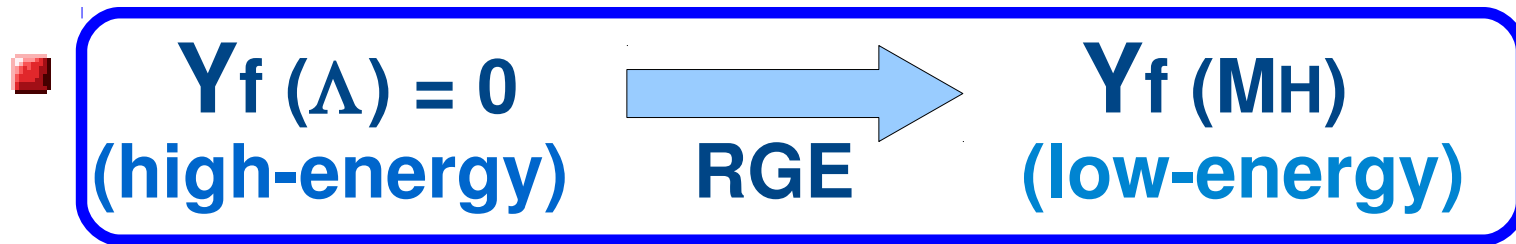
$$\mathbf{Y}_f^{\text{SM}} \equiv \frac{g_2}{\sqrt{2}M_W} \text{diag}[\mathbf{m}_f]$$



ChSB terms

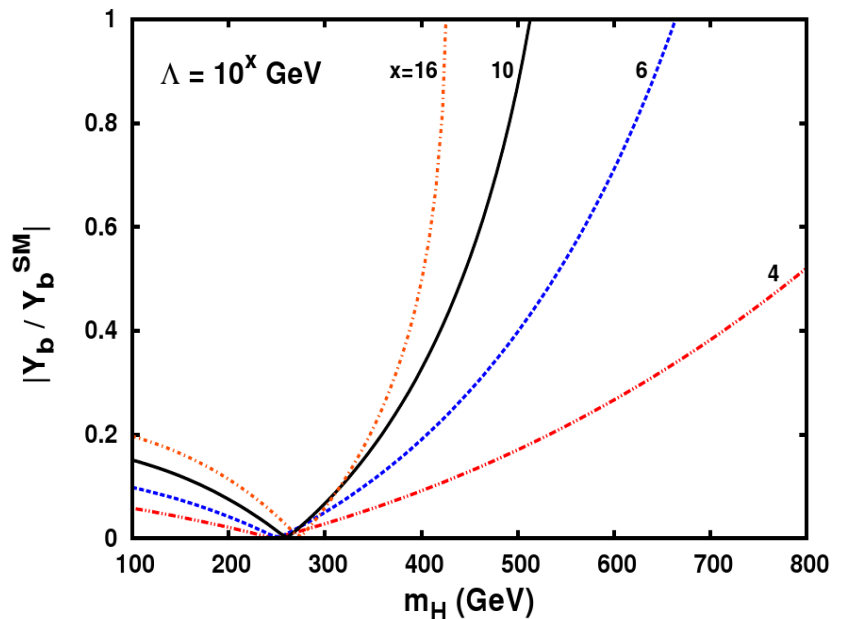
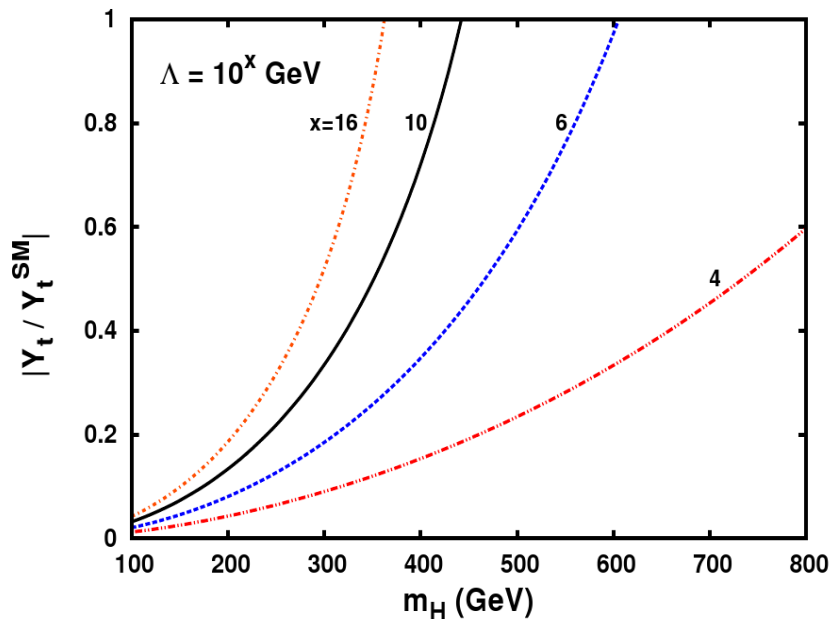
SM RGE recovered for $\mathbf{Y}_f^{\text{SM}} \rightarrow \mathbf{Y}_f$

Theoretical Framework

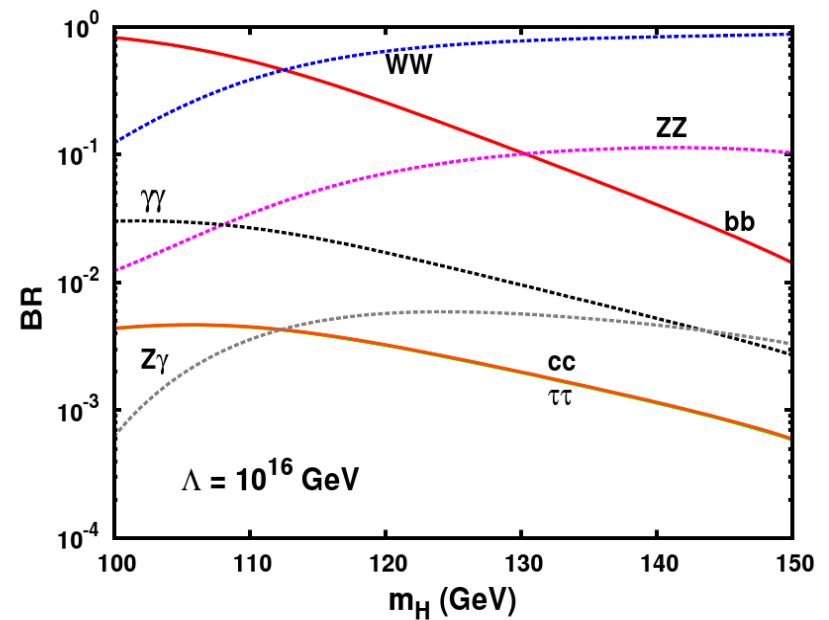
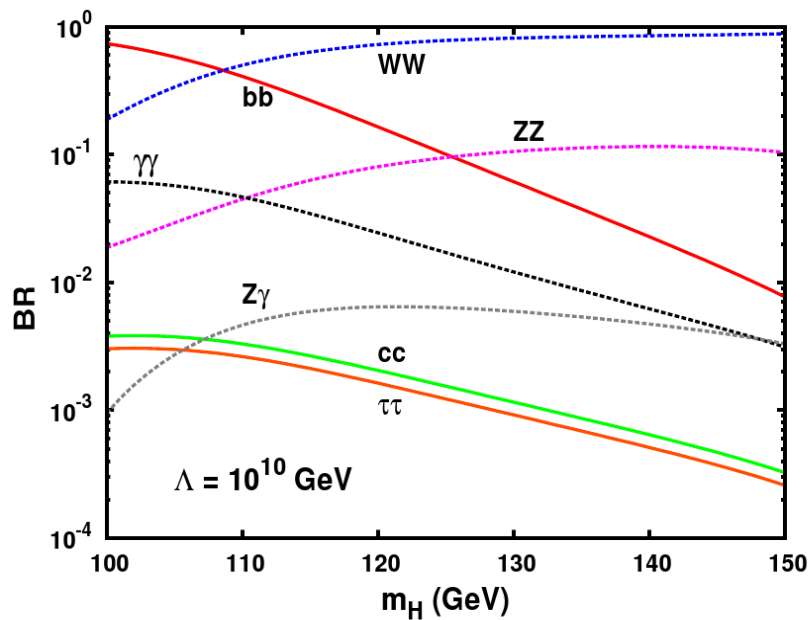
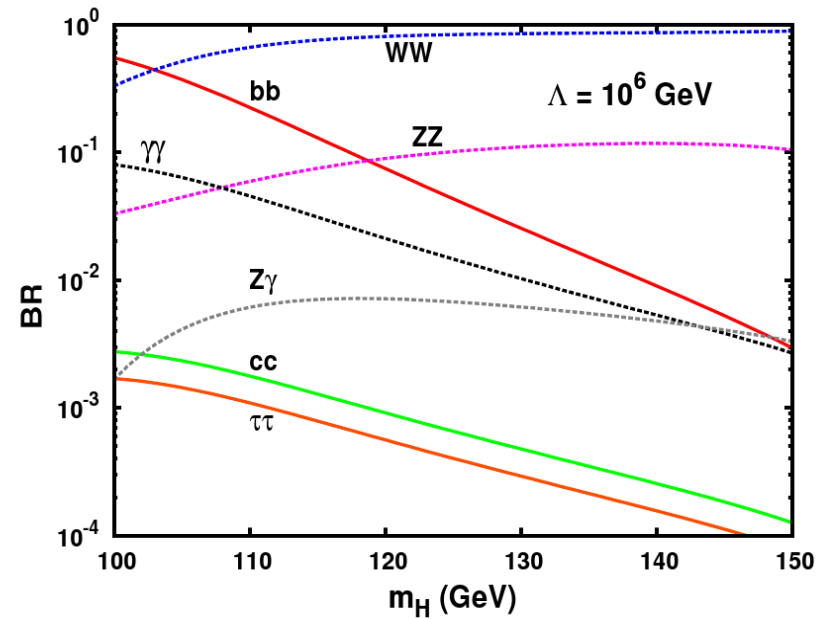
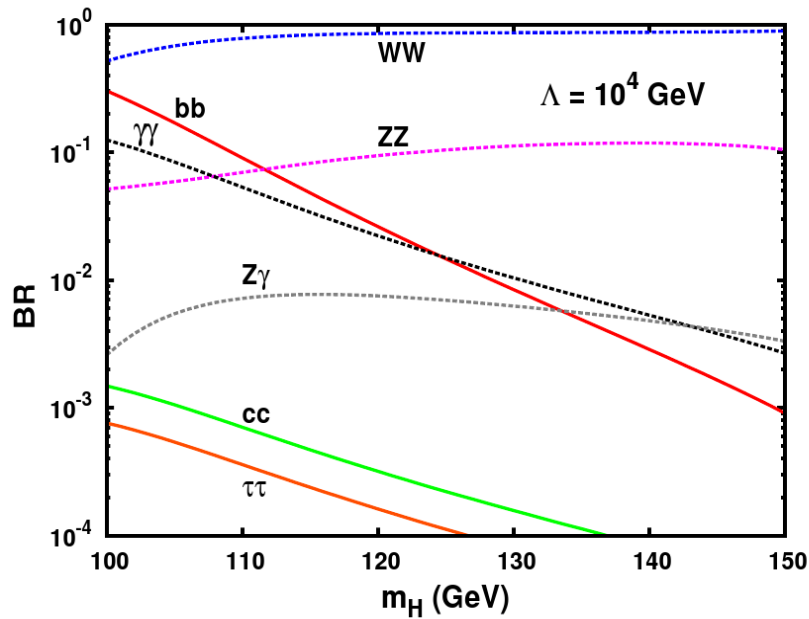


only SM degrees of freedoms assumed below Λ

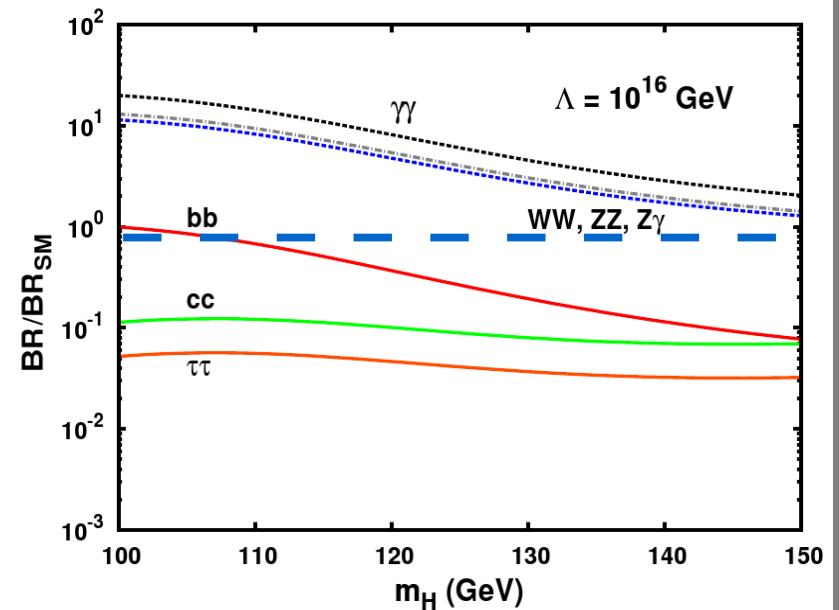
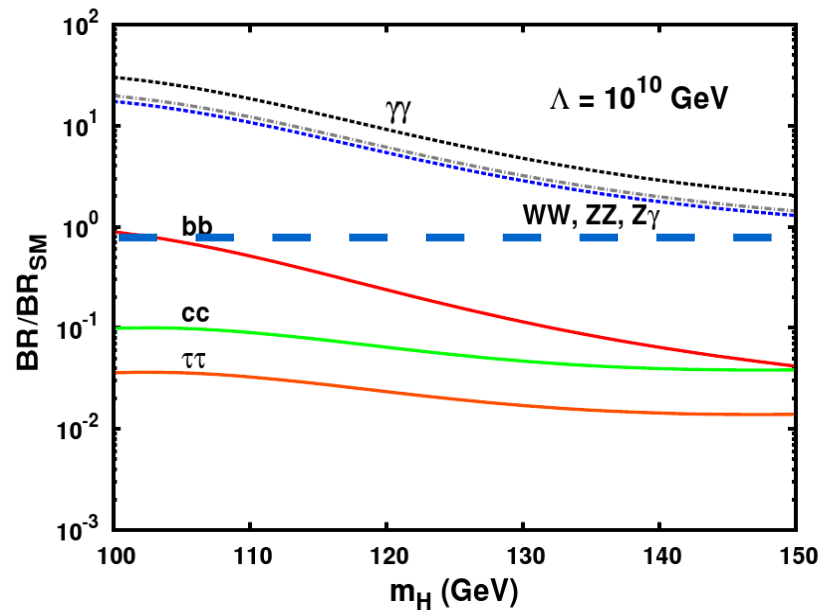
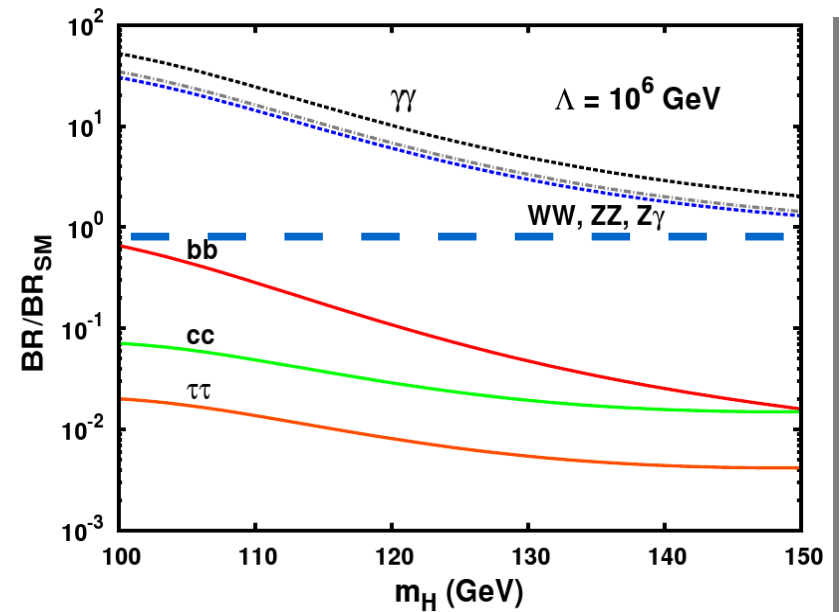
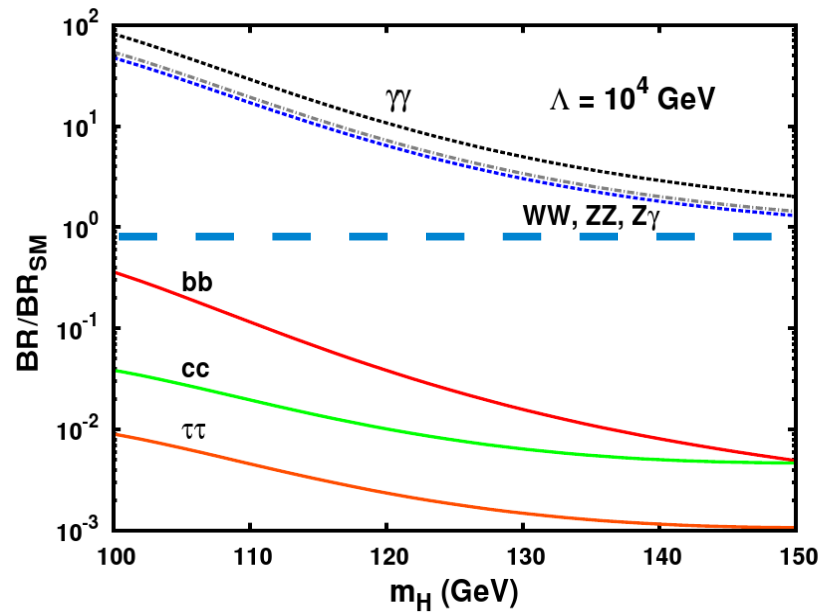
$Y_f(m_H)$ perturbative for $m_H < \langle H \rangle \sim 246 \text{ GeV}$



Higgs BRs dramatically affected



Higgs BRs normalized to SM ones



Higgs decay modes dramatically affected

- differs from **naïve fermiophobic scenarios** where Yukawa couplings are set to zero at EW scale
- main **H** decays modes \rightarrow **WW**, **bb**, **$\gamma\gamma$** , **ZZ**
- **BR(H \rightarrow bb) comparable or larger then enhanced $\gamma\gamma$**

H production at LHC changes

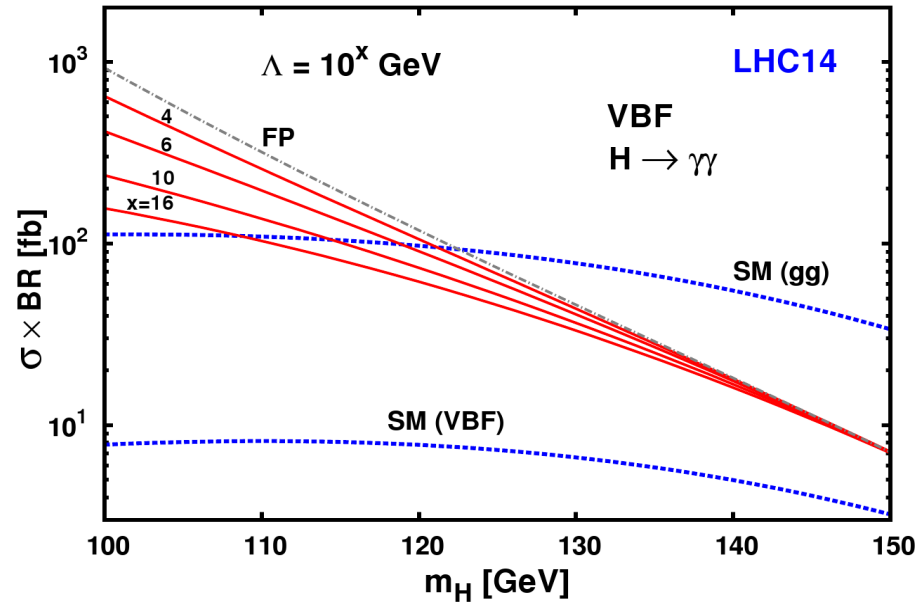
- **Y_t coupling radiatively induced (small)**
- **gluon-gluon fusion suppressed**
- **Vector boson fusion (VBF) becomes the dominant production mechanism at LHC**

LHC14 cross sections X BR [fb]

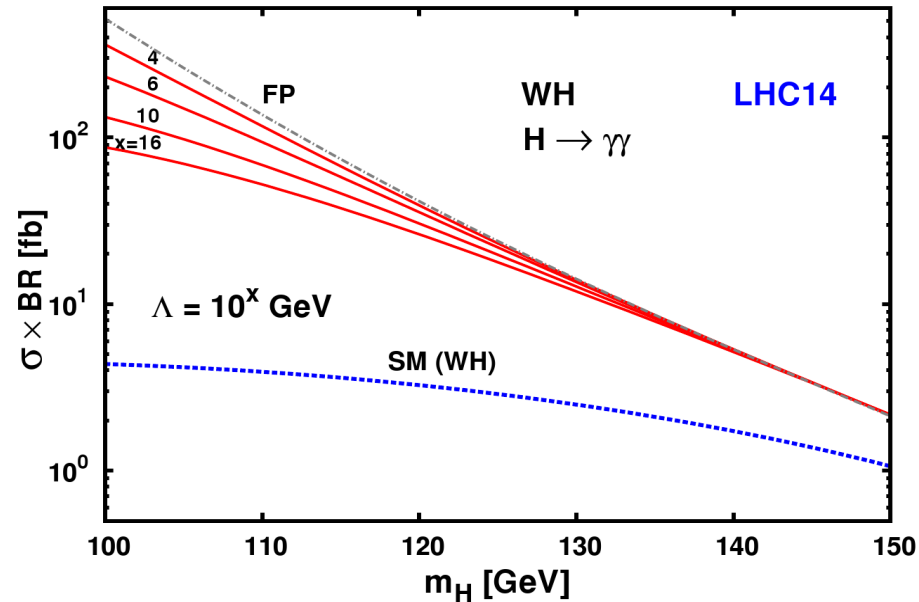
VBF

$\Lambda = 10^x \text{ GeV}$

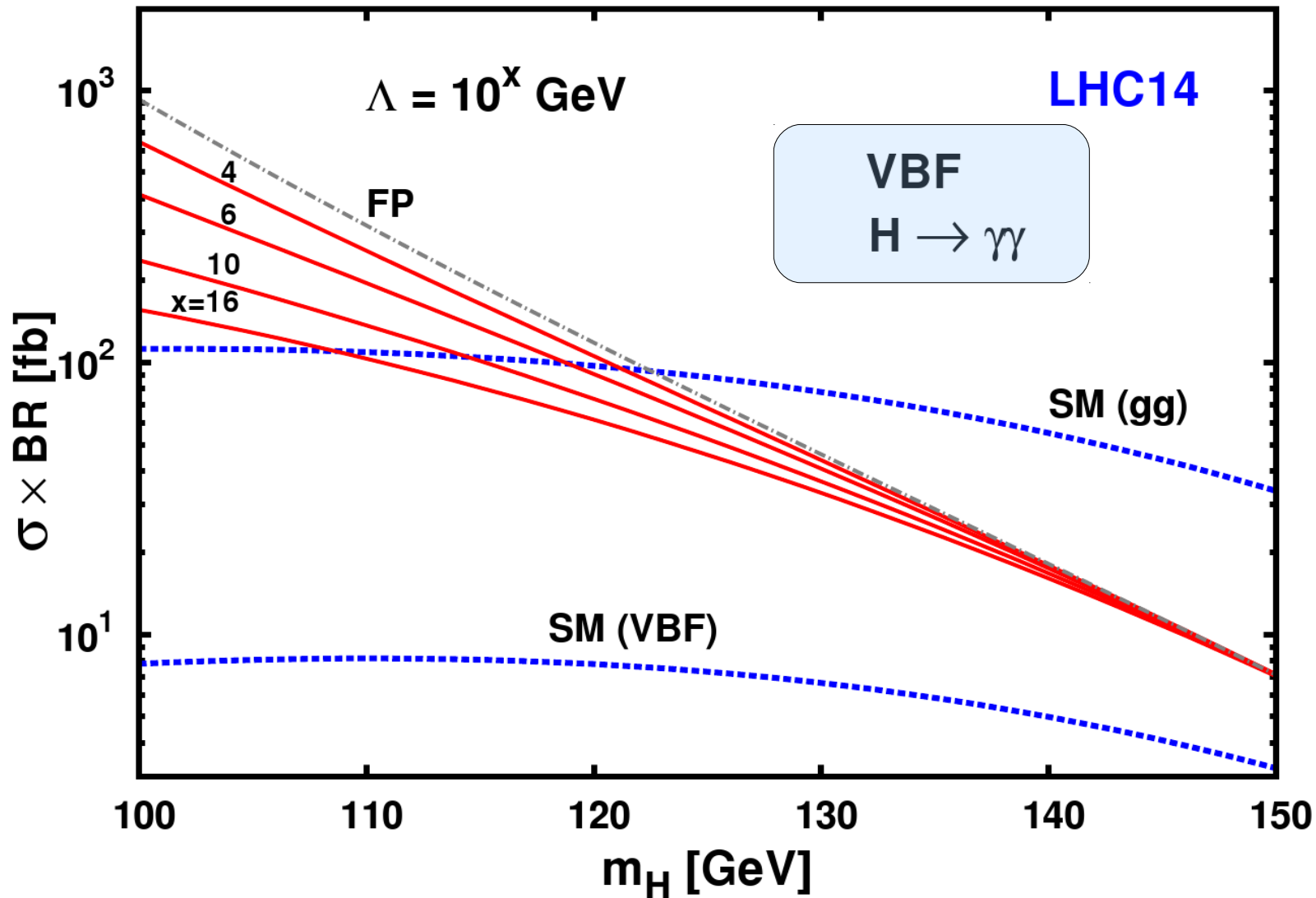
Fermiophobic
scenario \rightarrow **FP**



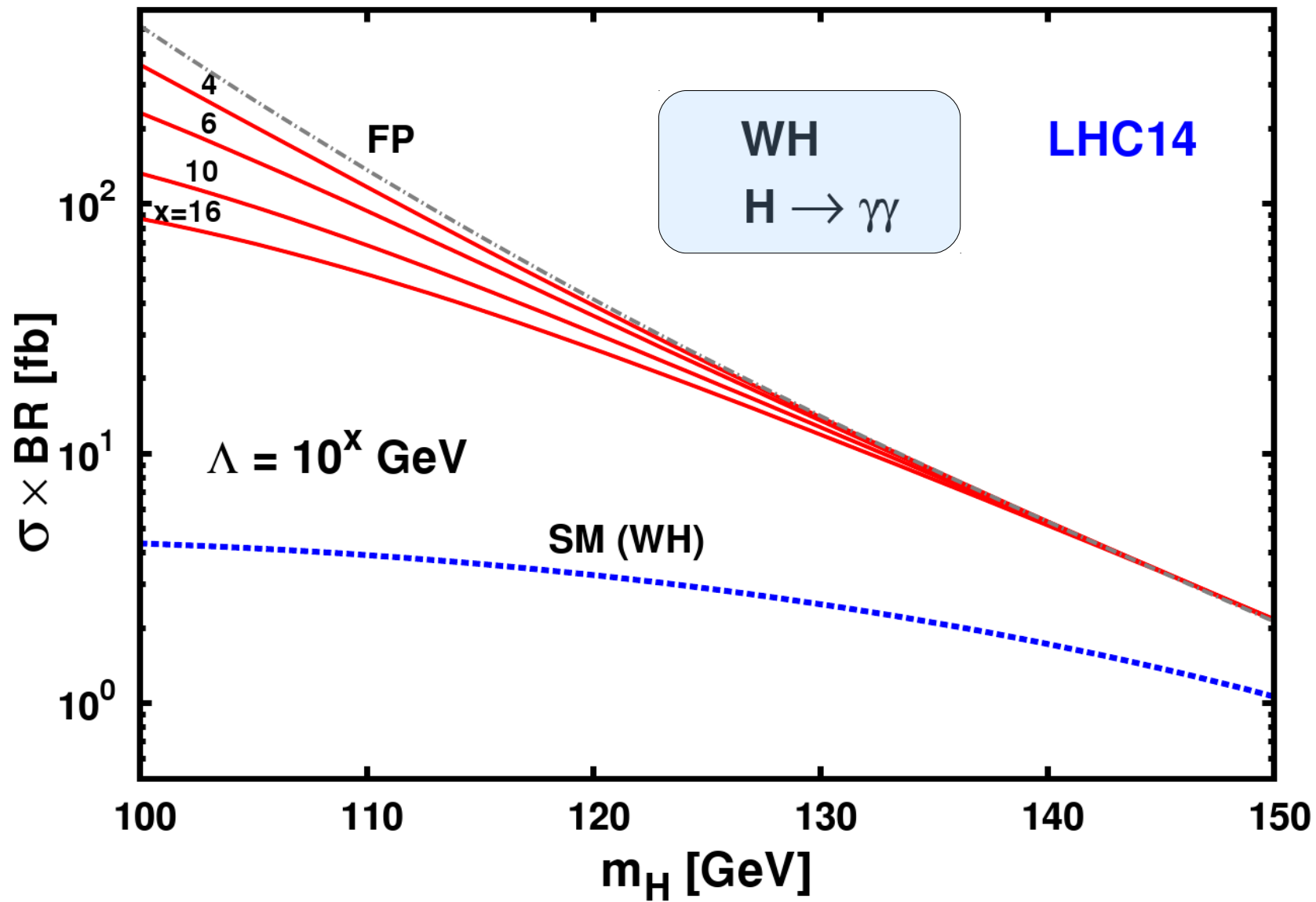
$H \rightarrow \gamma\gamma$



$H \rightarrow \gamma\gamma$



- Comparable rate to the SM (gg)
but VBF/BCKG much cleaner

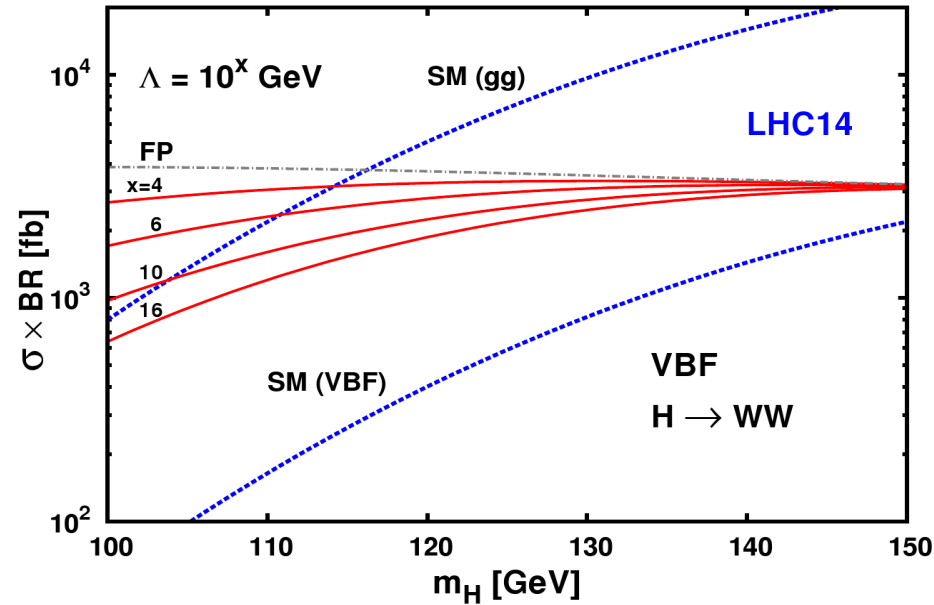


LHC14 cross sections X BR [fb]

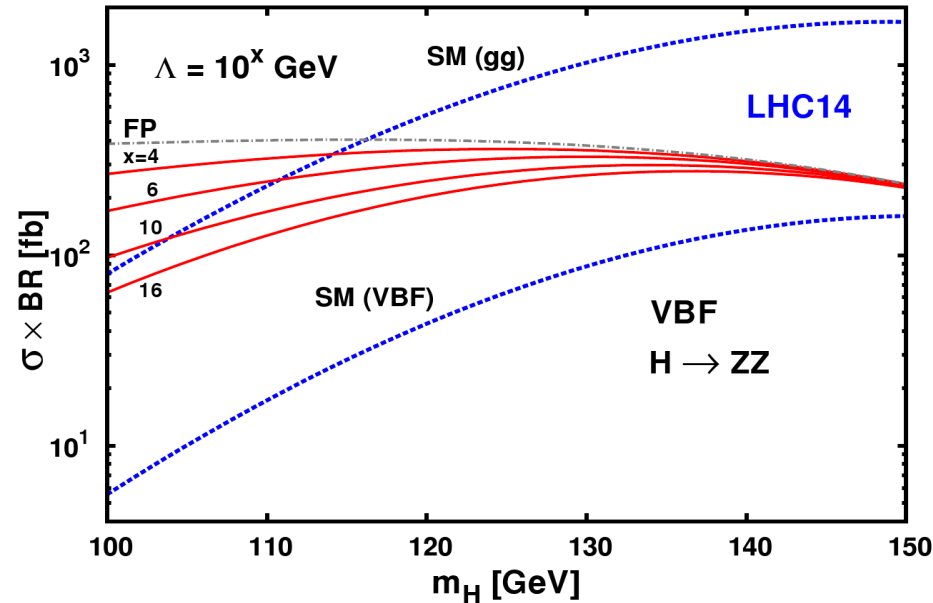
VBF

$\Lambda = 10^x \text{ GeV}$

Fermiophobic
scenario \rightarrow **FP**



H \rightarrow WW



H \rightarrow ZZ

Tevatron: present bounds

- **CDF & D0** analyzed various fermiophobic scenarios with enhanced $H \rightarrow \gamma\gamma$
- **$m_H < 110$ GeV excluded for pure FP**
- a dedicated analysis is necessary here due to non-trivial depletion of $BR(H \rightarrow b\bar{b})$
- **it could probe at best $m_H < 110$ GeV**

Future searches: LHC

- **Excellent probe of this scenario**
- $\sigma \times BR$ much larger than SM ($\gamma\gamma$, WW , ZZ , $Z\gamma$)
- better S/B ratio compared to SM (VBF, VH)
- better theoretical accuracy (VBF, VH)

Effective Y_f at Linear Colliders

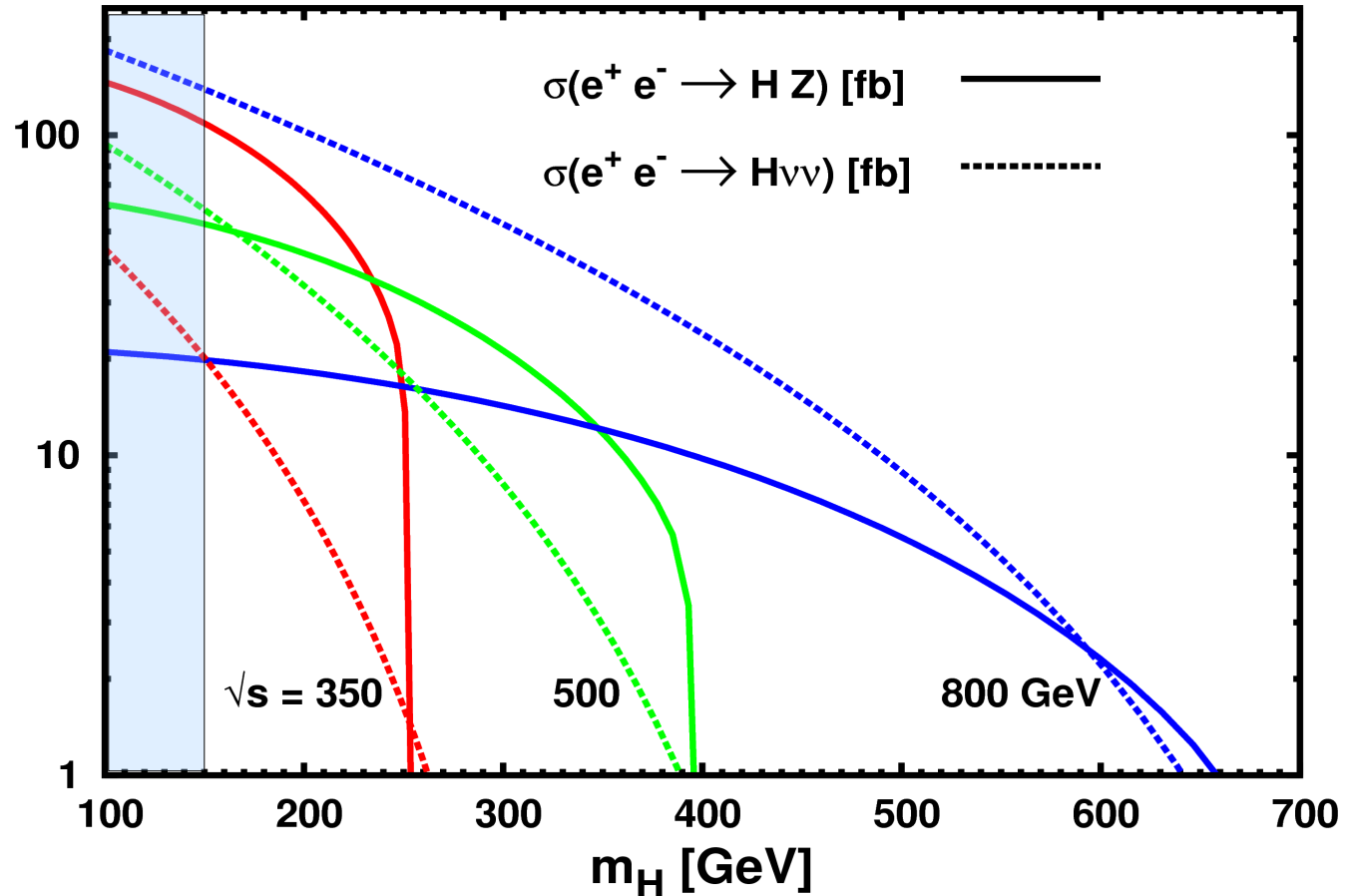
- no new degrees of freedom below TeV except an almost fermiophobic light Higgs
- **very good option** →

$$\sqrt{S} = 350 \text{ GeV}$$

advantages

- optimized Higgs production rate via ZH
- $\sigma(\text{ZH}) > \sigma(\text{H}\nu\nu)$ **almost monochromatic H**
- **good potential to measure radiative Yukawa couplings**
- **could constrain the scale Λ of ChSB**

Higgs production cross sections

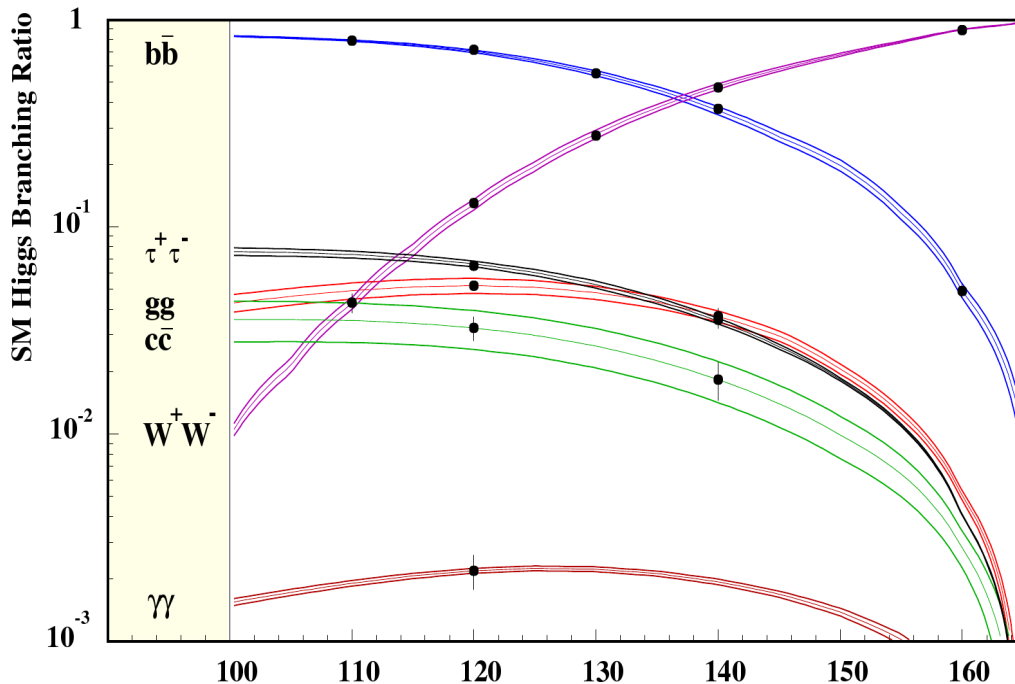


■ we will show only $\sigma(ZH) \times \text{BR}$ at 350 GeV →
for interesting channels

■ note that:

$\sigma(H\nu\nu) + \sigma(HZ)$ @500GeV \sim $\sigma(HZ)$ @350GeV for $M_H < 150$ GeV

Accuracy on BRS for SM Higgs boson



[hep-ph/0106315]
Tesla, Technical Design
Report (part III)

Expected experimental uncertainties ■ error bars → exp. sensitivities

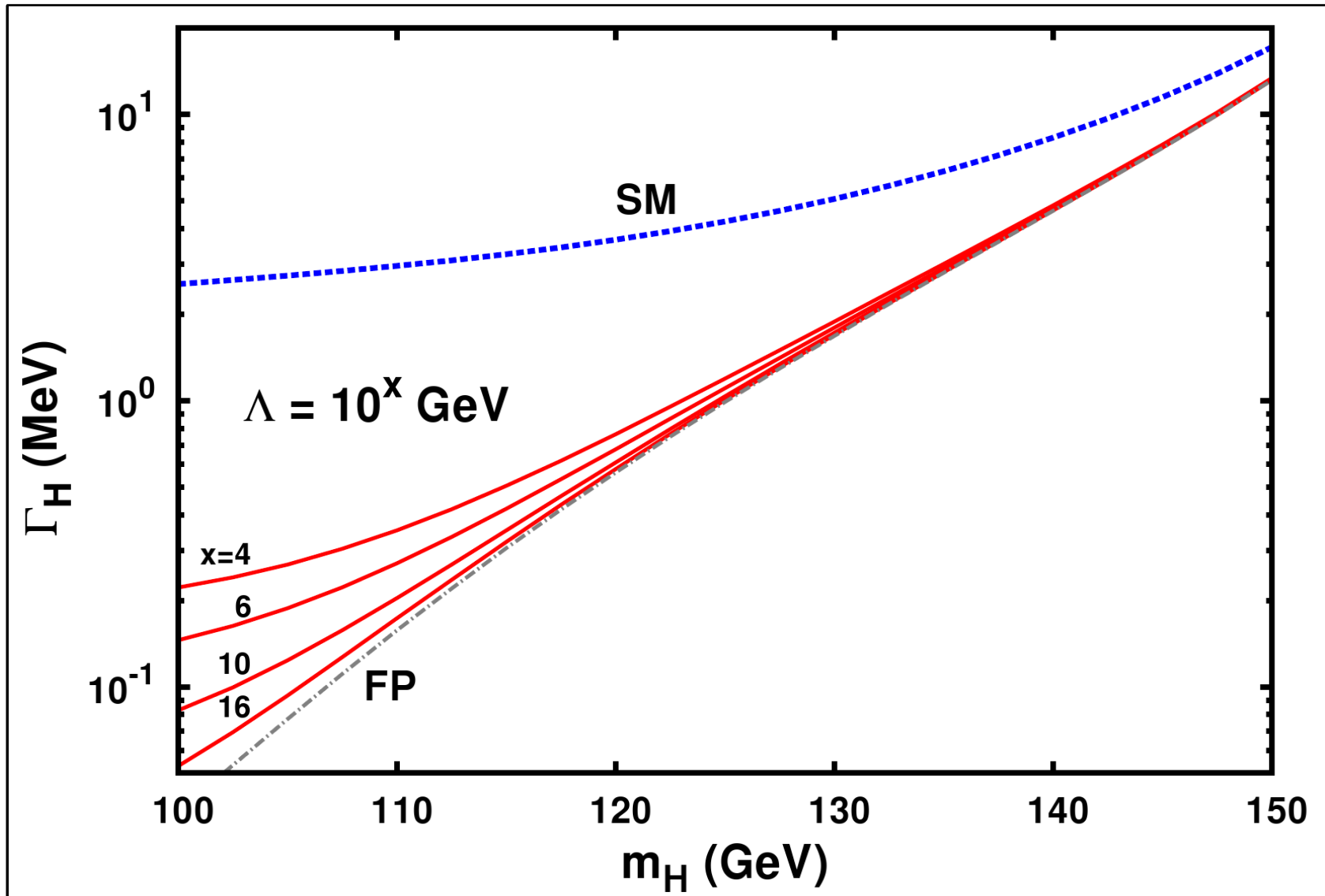
“Phase 1”: 500 fb⁻¹ at 350 GeV, no beam

SM Higgs branching ratio uncertainties

	$m_H = 120$ GeV	140 GeV
BR($b\bar{b}$)	2.4%	2.6%
BR($c\bar{c}$)	8.3%	19.0%
BR($\tau\tau$)	5.0%	8.0%
BR(WW)	5.1%	2.5%
BR(gg)	5.5%	14.0%

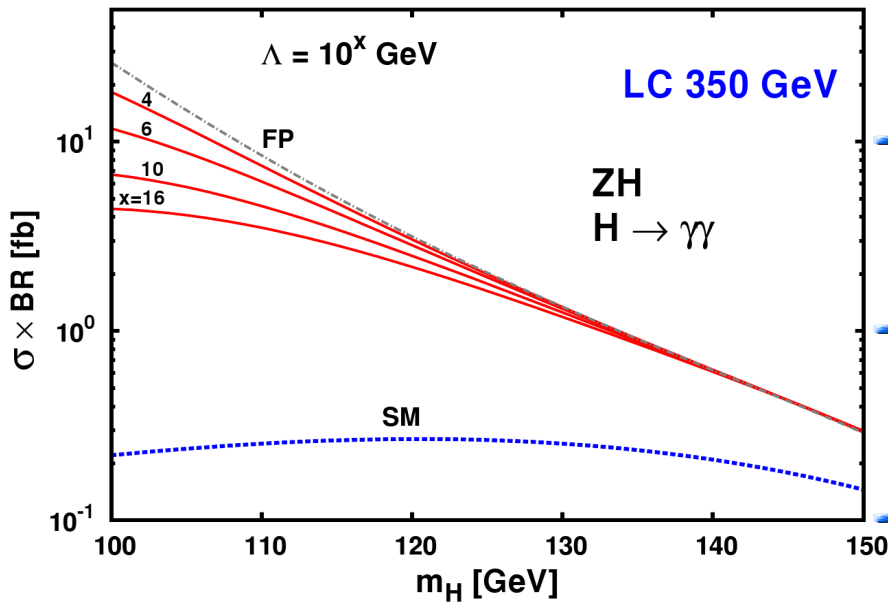
**New analysis needed
to establish $\Delta\text{BR}(\text{H} \rightarrow \text{ff})$**

Total Higgs decay width



$\sigma \times \text{BR}$ [fb]

$N_{\text{ev}}/500 \text{ fb}^{-1}$



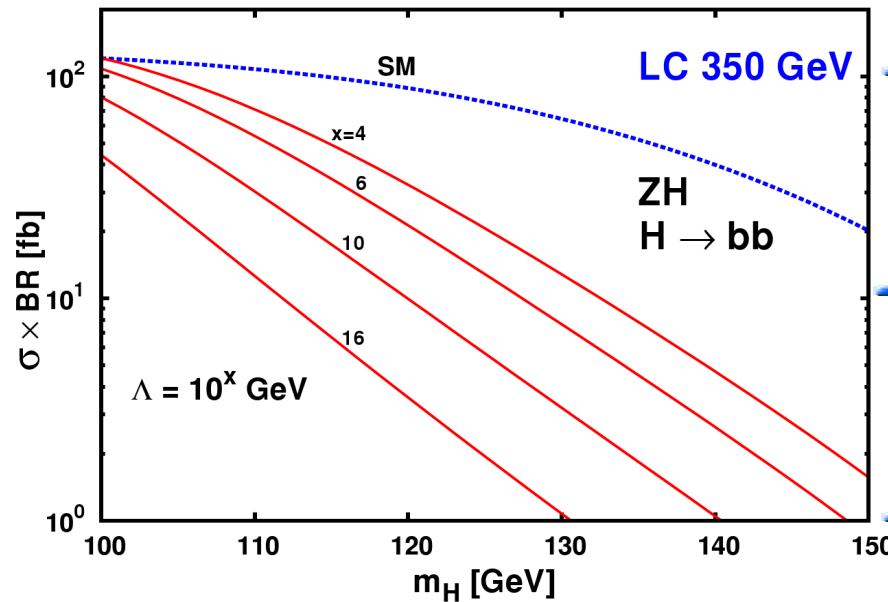
5000

LC 350

500

$H \rightarrow \gamma\gamma$

50



50000

LC 350

5000

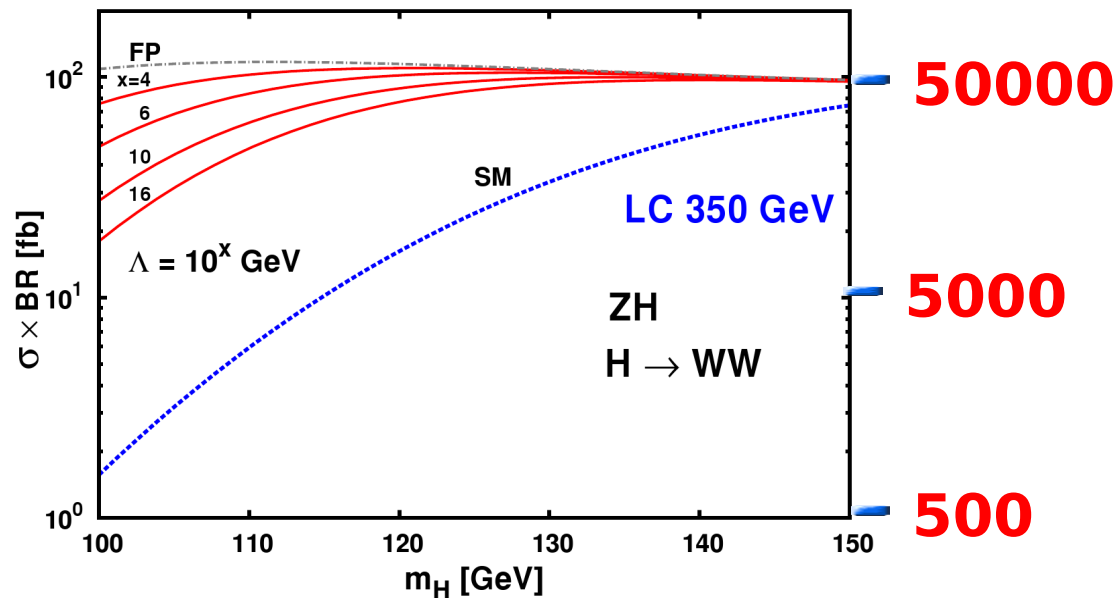
$H \rightarrow b\bar{b}$

500

Excellent
sensitivity to Λ

$\sigma \times \text{BR}$ [fb]

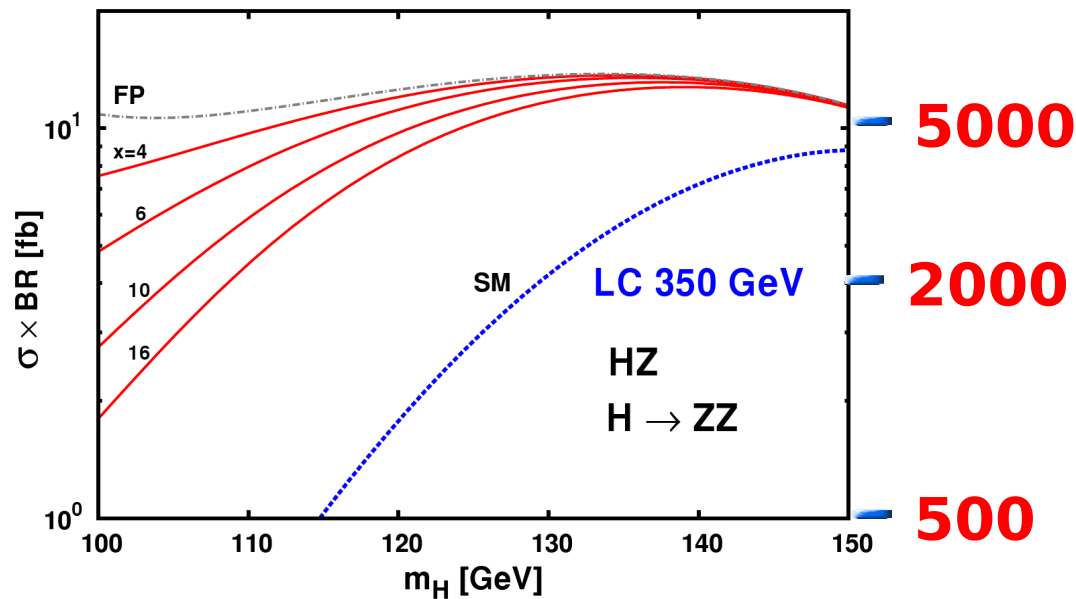
$N_{\text{ev}}/500 \text{ fb}^{-1}$



LC 350

$H \rightarrow WW$

$\Lambda = 10^x \text{ GeV}$

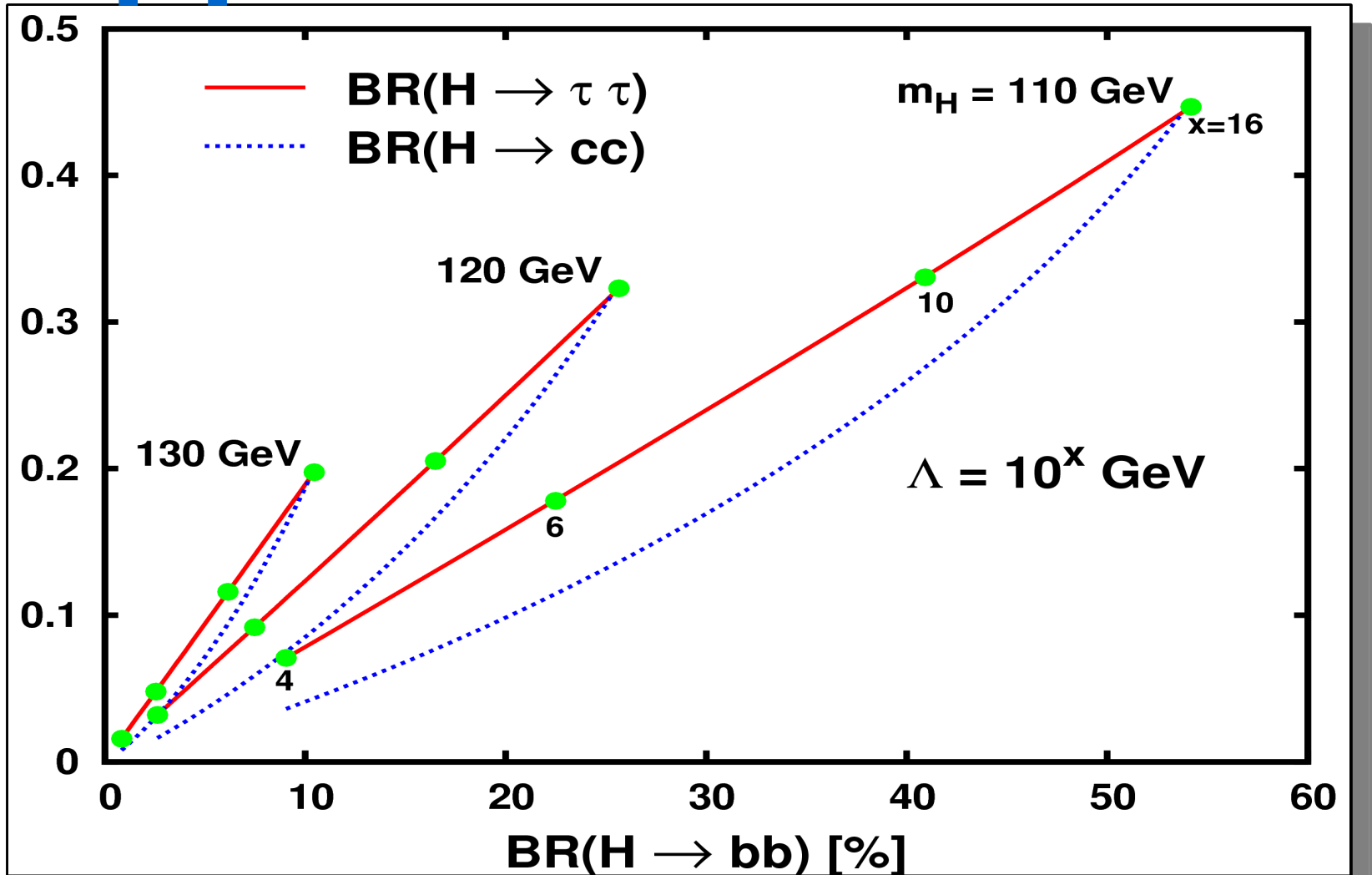


LC 350

$H \rightarrow ZZ$

Correlations of $\text{BR}(H \rightarrow ff)$

BR[%]



$\text{BR}[\tau\tau] \sim 1\text{-}2\% \text{ BR}[b\bar{b}]$

Outlook

- New predictive scenario where **mf's** arise at large scale Λ and induce **Yf** radiatively → **only 1 free param.**
- **big impact on Higgs phenomenology** → “improved” FP scenario: effective **Yf** and Higgs production via W/Z radiation (no $gg \rightarrow H$ at LHC!)
- **Enhanced BR($\gamma\gamma$, WW, ZZ, γZ) for $M_H < 140$ GeV** and possibly large radiative **BR($H \rightarrow bb$)**
- **Excellent LHC potential for testing this scenario,** but accurate study of **$H \rightarrow bb, cc, \tau\tau$** , requires LC
- Rates for **$e^+e^- \rightarrow ZH \rightarrow Zbb$** remarkably sensitive to the scale Λ !
- **New analysis of LC sensitivity to Higgs BR in the new framework needed**

Backup

Tevatron cross sections X BR (fb)

