# The CLIC Decelerator

Overview and beam physics

**International Workshop on Linear Colliders 2010** 

October 20, 2010

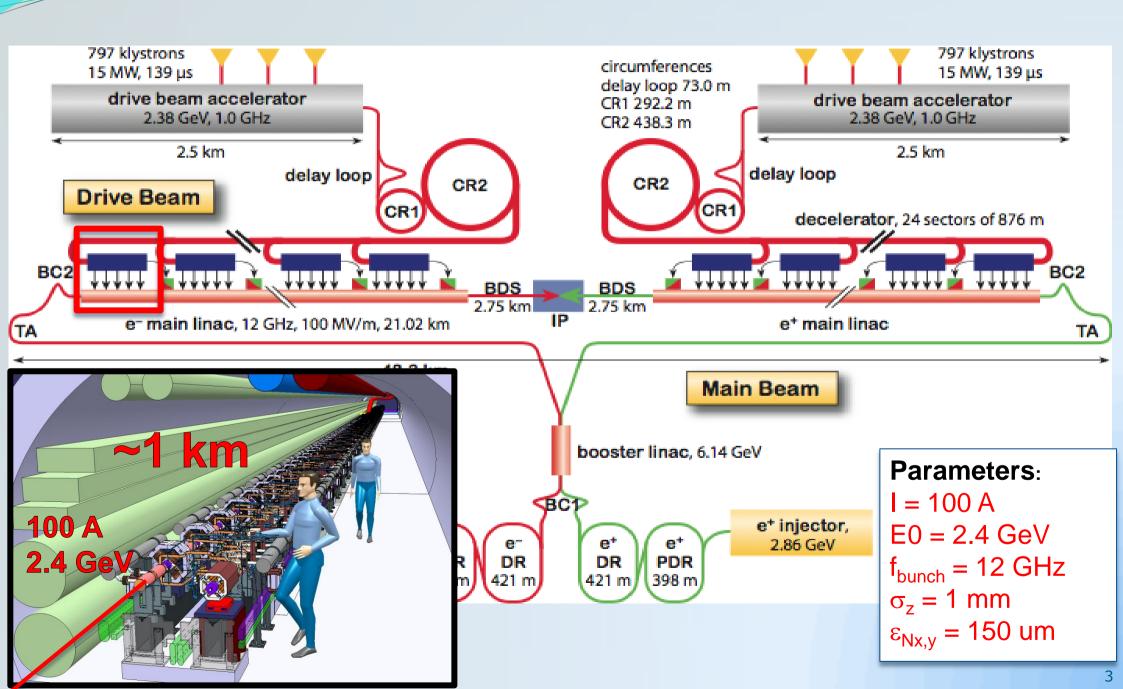
Erik Adli, Department of Physics, University of Oslo and CERN



- Requirements
- Beam physics
- Component specifications
- Test facilities
- Conclusions



### The CLIC Drive Beam decelerators





### Requirements: power production

Decelerator: power source for main linacs. Power production:

$$P=rac{1}{4}(R'/Q)rac{\omega_{
m rf}}{v_g}L_{
m PETS}^2\eta_{\Omega,
m PETS}^2I^2F^2(\lambda)F^2(\phi)$$

Requirements for 1% luminosity loss:  $\Delta E/E < 7 \times 10^{-4}$ . Converted to drive beam decelerator requirements:

$$\frac{\Delta \mathcal{L}}{\mathcal{L}} \approx 0.01 \left[ \left( \frac{\sigma_{\phi,coh}}{0.2^{\circ}} \right)^{2} + \left( \frac{\sigma_{\phi,inc}}{0.8^{\circ}} \right)^{2} + \left( \frac{\sigma_{I,coh}}{0.75 \times 10^{-3} I} \right)^{2} + \left( \frac{\sigma_{I,inc}}{2.2 \times 10^{-3} I} \right)^{2} + \left( \frac{\sigma_{\sigma_{z},coh}}{1.1 \times 10^{-2} \sigma_{z}} \right)^{2} + \left( \frac{\sigma_{\sigma_{z},inc}}{3.3 \times 10^{-2} \sigma_{z}} \right)^{2} \right]$$

D. Schulte, WG2,6,7,8, Wednesday 14:00

The drive beam generation is discussed separately:

WG6, Session5, Thursday 08:30

Here we focus on the consequences for the decelerator.



# Decelerator requirements

The decelerator beam transport: robust performance of each of the 2 x 24 decelerator sectors - 42 km beam line. Will require a very large number of components.

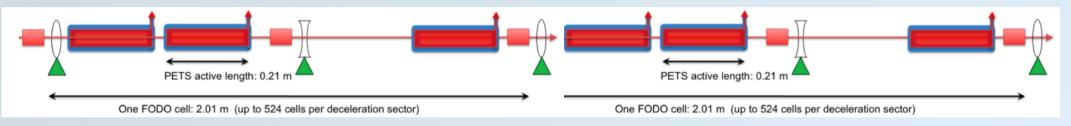
A number of studies have been performed to with the aim to optimize specifications, in order to contain cost and power consumption (this talk).

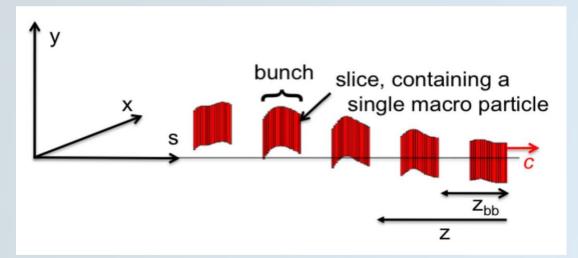
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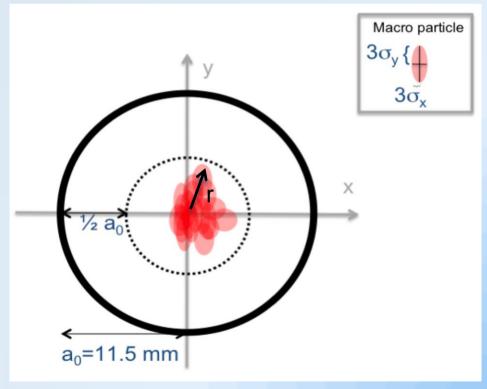
### Decelerator studies

Main tool: simulation studies (tracking code PLACET), with an element representing the Power Extraction and Transfer Structures, including fundamental and dipole modes wake field calculations, and both single and multi-bunch effects.





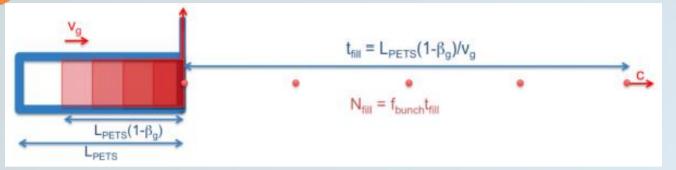
Macro-particle beam model, sliced beam with tracking of 2<sup>nd</sup> order moments



Simulation metric:  $r = 3\sigma$  of worst beam slice

# clc

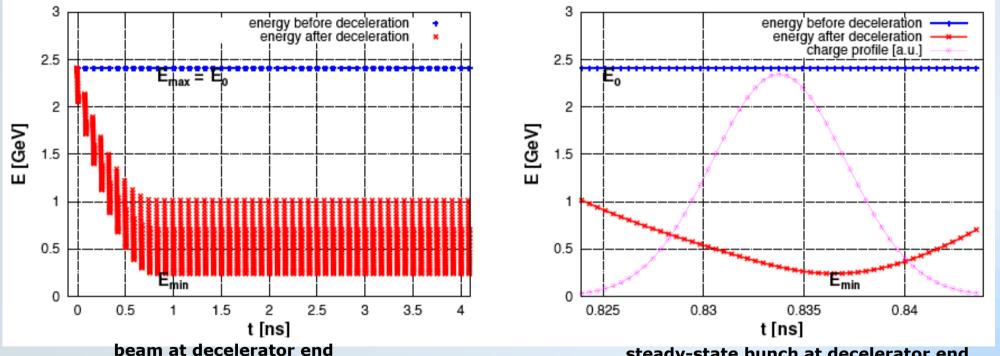
# PETS induced energy spread



(pilot beam, w/o beam loading compensation)

 $\eta_{\text{extr}} = 0.90$ 

$$E_{\text{most}}(n) = E_0(1 - \eta_{\text{extr}} \frac{n}{N_{\text{PETS}}})$$



steady-state bunch at decelerator end (pilot beam, w/o beam loading compensation)

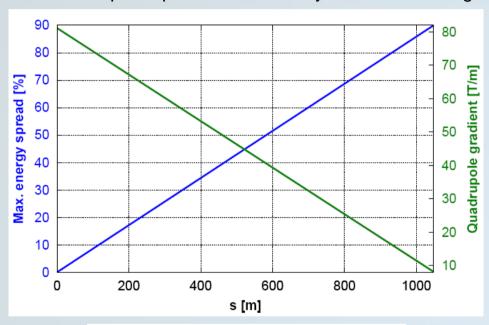
Decelerator beam: the high group velocity PETS will induce up to **90% energy spread** at the decelerator end, as well as significant intra-bunch energy spread. To ensure reliable rf power production it is of importance that electrons of all energies are robustly transported along the lattice.

Overall criterion for beam transport  $r < \frac{1}{2}$  radius (5.8 mm)

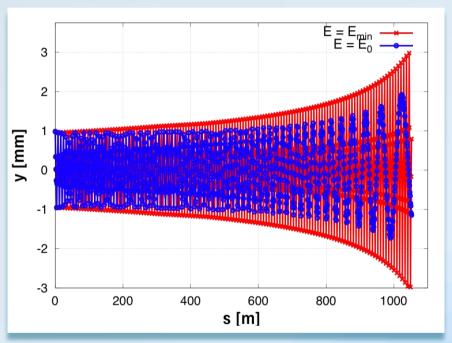
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### Focusing strategy

- Focusing strategy: lowest energy particles ideally see constant phase-advance μ≈90°
- Higher energy particles see phase-advance decreasing from μ≈90° to μ≈10°
  - Perfect machine and beam : high energy envelope contain in low energy envelope
  - Energy acceptance: -3% of E<sub>0</sub> at the entrance; but increasing along the lattice
  - · Each of quadrupoles should ideally have a different gradient



$$k_{E_0}(n) = k_0(1 - \eta_{\text{extr}} \frac{n}{N_{PETS}})$$



3-sigma particles in a perfectly aligned machine

Least decelerated particle has a tune of  $\mu \approx 70$ , and an increase of  $\beta$  of  $\beta_{Fmax}(E_0)/\beta_{F0}=4$ .

Most decelerated particle has a tune of  $\mu \approx 135$ , and an increase of action of  $J_{max}(E_{min})$  /  $J_0 = \gamma_0/\gamma_f = 10$ 

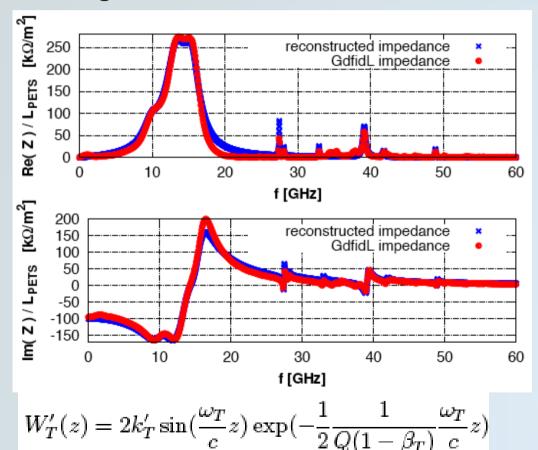
 $\sigma(s) = \sqrt{2J(s)\beta(s)}$ 

3-σ envelope for perfect machine: r<sub>ad</sub> = 3 mm



### Transport challenge: dipole wake

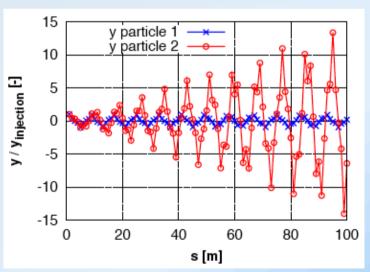
Sufficient mitigation of transverse electromagnetic forces, due to the **PETS high group velocity dipole wake**, has been a major challenge for the two-beam accelerator concept.



Tracking simulations approximate simulated PETS impedance by a number of discrete modes, each characterized by  $\{f,k_T,Q,\beta\}_i$ 



PETS 12 GHz prototype (TBTS 1 m)



Principal effect of dipole wake: resonant linear increase of betatron amplitude of driven particle.

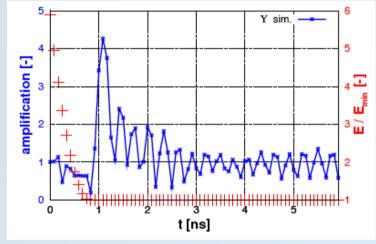
A. Cappelletti, WG6, Wednesday 10:30



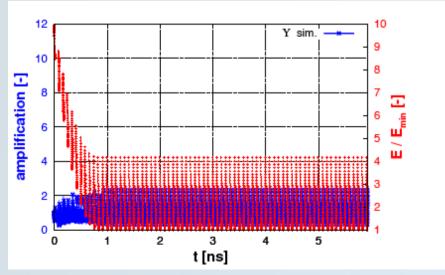
### Transport challenge: dipole wake

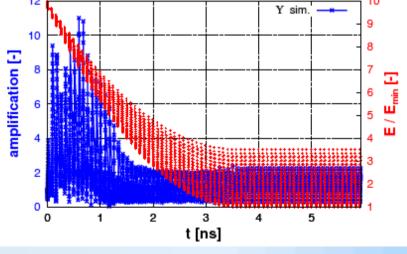
The multi-bunch amplification due to the dipole wake is large. The PETS induced energy spread mitigates the amplification, however, to a level depended on the PETS design. Here illustrated by calculating the amplification of action due to dipole wakes, at the decelerator end.

Rf power production is proportional with  $(R'/Q) / v_g = const.$  However, PETS with too low group velocity do not develop energy spread fast enough to decohere the wake build-up.



#### Point-like bunches.



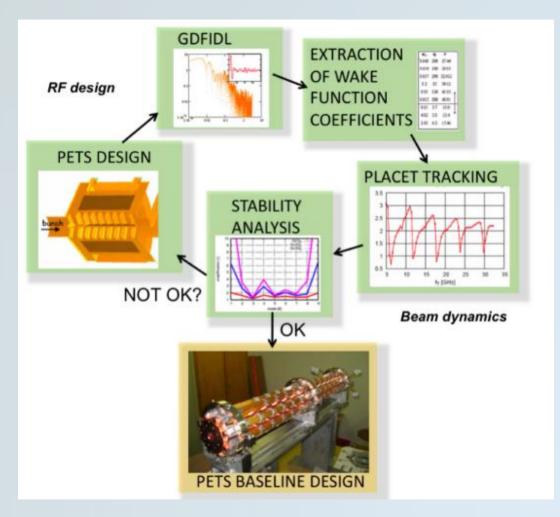


1 mm bunches, baseline PETS

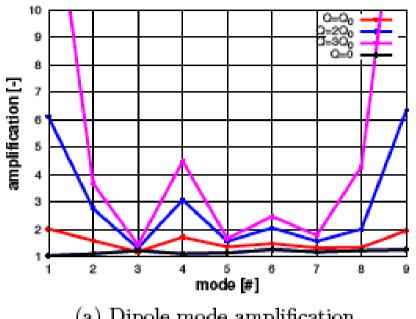
1 mm bunches, "slow" PETS design



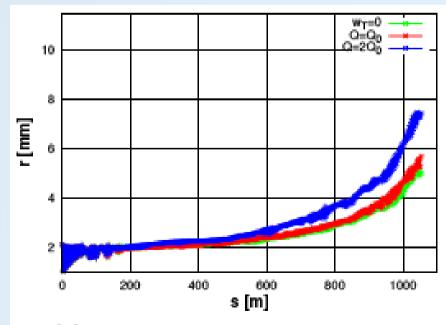
### Dipole wake status: PETS baseline design



Large series of potential PETS design variants have been examined for robust mitigation of the transverse wake, for all beam modes, and all errors sources. After several iterations rf and beam dynamics expertise, a PETS baseline indicating adequate wake mitigation, has been secured.







(b) Envelope growth along the lattice 12

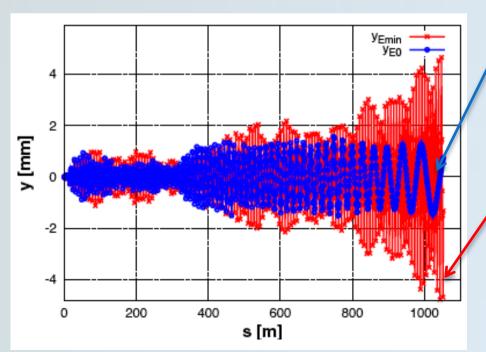


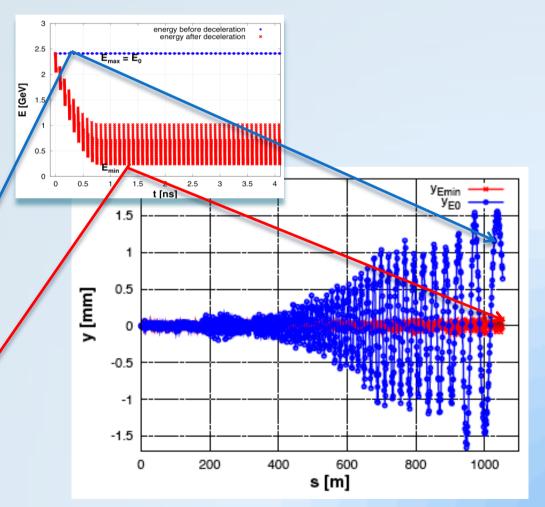
### Transport challenge: orbit correction

Kicks from misaligned quadrupoles might drive beam envelope out of vacuum chamber, even for pre-alignment of 20  $\mu$ m. Estimate for uncorrected machine sets scale :

$$\langle r_c \rangle_{\rm Emin} = \Sigma_{j=1}^N R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \text{With 1000 quadrupoles and 20 } \mu \text{m rms offset,} \\ \text{the expected centroid envelope is ca. 4 mm.} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm FODO}/2)} \sqrt{\frac{(1-\frac{1}{2}\eta_{\rm extr})}{(1-\eta_{\rm extr})}} \\ \frac{\langle r_c \rangle_{\rm Emin}}{\langle r_c \rangle_{\rm Emin}} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j=1}^{N} R_{Fi}(\xi_j/f) \sqrt{\frac{\gamma_j}{\gamma_f}} \approx 2\sqrt{N} \frac{\sigma_{\rm quad}}{\cos(\phi_{\rm Emin})} = \frac{1}{2} \sum_{j$$

90% energy spread of decelerator beam poses a challenge for beam transport: Dispersive trajectories of higher / lower energy particles: 1-to-1 correction does properly correct only the beam centroid.





Beam transport for ideal injection into a misaligned machine

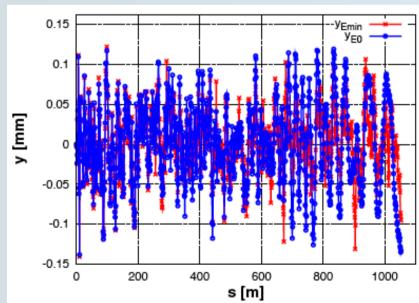


#### Transport challenge: orbit correction

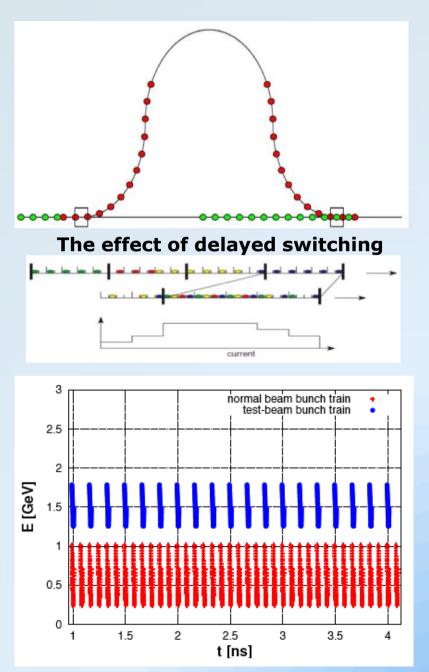
We seek to improve the situation by imposing that particles of different energies shall follow same trajectory, i.e. minimizing the energy dependence of the trajectories; a dispersionfree correction.

We propose a scheme based on drive beam bunch-manipulation and exploiting PETS beam loading, to generate a test-beam.

Use of "delayed switching" in the drive beam generation. The test beam can have almost any energy leverage. One-pulse correction.



Beam transport for ideal inj. into a dispersion-free steered machine

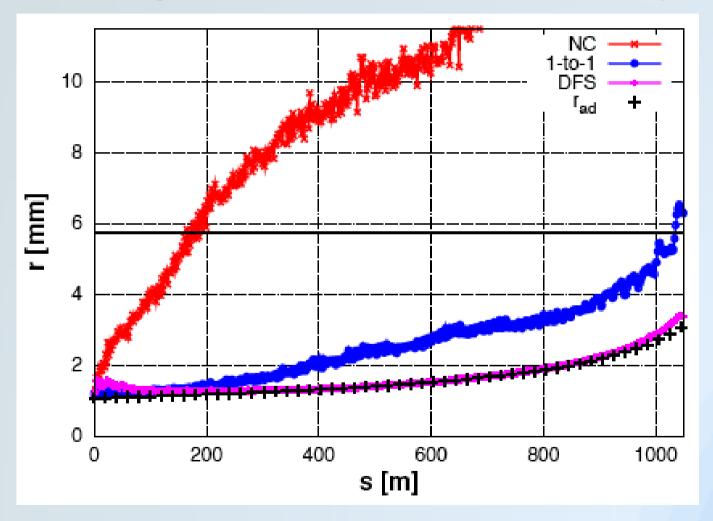


Energy profile of main beam and example test-beam

# clc

### Beam transport with dispersion-free correction

 Results of simulation including the combined effects of wake fields and misalignment, for the CLIC base line parameters :

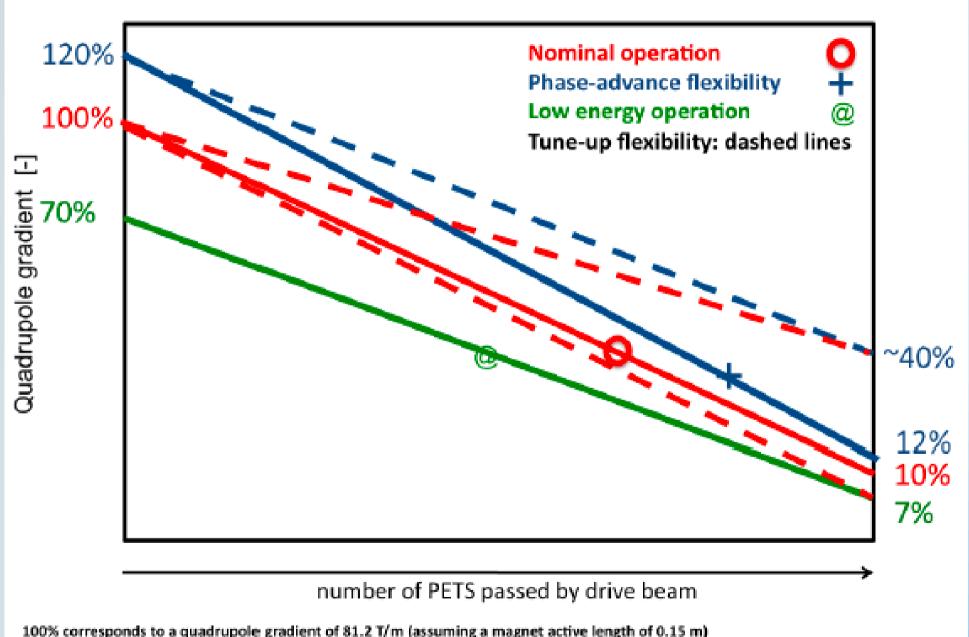


3-σ envelope of 500 simulated machines (worst case r)

- Requirements
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### Decelerator operational scenarios



Tune-up beams: low current, negligible beam-loading. Not needed (nor optimal) to go to full phase-advance

# ccc

### Quad specifications: baseline parameters

 Specification: one quadrupole per meter gives beta function (for most decelerated particles) of

$$<\beta> = 1.25 \text{ m}$$

- Deemed necessary for robust mitigation of dipole wake
- Gives r = 3.3 mm (out of  $a_0 = 11.5 \text{ mm}$ ) for ideal beam

A. Vorozhtsov, WG6, Wednesday 09:30

- Results in ~42'000 decelerator quads
- Powered magnets is the baseline
  - failure tolerant serial powering scheme a necessity
- Tuneable permanent magnets option investigated
  - must cover all operational scenarios

D. Siemaszko, S. Pittet, WG6, Wednesday 09:50

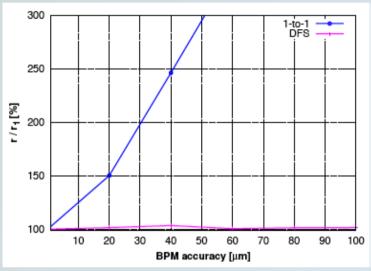
> J. Clarke, WG6, Wednesday 11:40

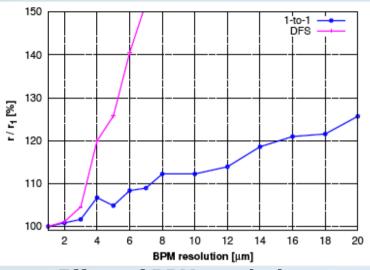
Quadrupole specifications			
Total nb. of quads	$N_{ m tot}$	~42000	-
Inner radius of vacuum chamber	$a_0$	11.5	mm
Max. integrated gradient	$\hat{G}$	12.2	T
Tunabiliy		See operational scenarios	
Magnet design accuracy rms	$\sigma(Gl)/(Gl)$	$1 \times 10^{-3}$	-
Resulting magnet design tolerance	$\Delta(Gl)/(Gl)$	$\sqrt{3} \times 10^{-3}$	-
Max. dodecapole component at 11 mm	Int B6 / Int B2	$3 \times 10^{-4}$	-
Power supply accuracy rms	$\sigma(I)/I$	$5 \times 10^{-4}$	-

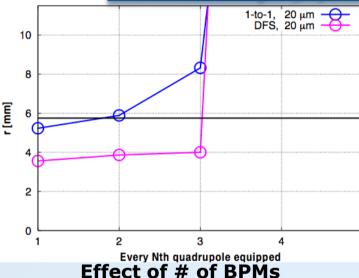
### BPM specifications: baseline parameters

Beam-based correction performance drive the BPM specifications. Target: negligible envelope growth due to quadrupole kicks.

S. Smith, WG8, Wednesday 09:40







**Effect of BPM accuracy** 

10

[mm]

**Effect of BPM resolution** 

0 200 400 600 800 1000 s [m]

**Results with baseline parameters** 

Quantity	Symbol	Value	Unit
# of quadrupoles per BPMs	N	~2/3	-
Total number of BPMS	$N_{tot}$	~28000	-
Production beam			
BPM accuracy	$\sigma_{ m acc}$	20	$\mu\mathrm{m}$
BPM resolution	$\sigma_{ m res}$	2	$\mu\mathrm{m}$
Time resolution	$t_{ m res}$	60	ns
Pilot beam			
BPM accuracy	$\sigma_{ m acc}$	20	$\mu\mathrm{m}$
BPM resolution	$\sigma_{ m res}$	4	$\mu\mathrm{m}$
Time resolution	$t_{ m res}$	60	ns



## Vacuum system specifications

Collective effects studies for the 100 A drive beam:

#### 1) Fast-ion instability

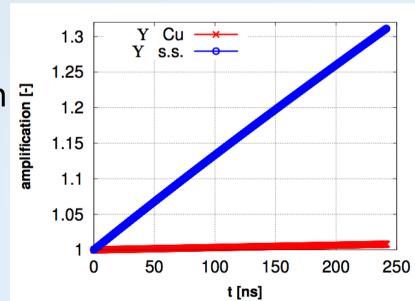
Analytic approximations, neglecting the large energy spread), yields one fast-ion instability rise-time with p = 40 nTorr.

Decelerator: p < 40 nTorr

#### 2) Resisitive wall instability

Analytic calculations yield a significant (unacceptable) amplification of beam offsets for  $\sigma_{res} = \sigma_{s.s}$  while negligible for  $\sigma_{res} = \sigma_{Cu.}$ 

Decelerator:  $\sigma_{res} \sim \sigma_{Cu}$ 



B. Jeanneret, WG4,6, Thursday 08:30

Comparison with requirements for the rest of the drive beam long transfer lines:

Amplification of a coherent beam offset due to the resistive wall wake



# Other specifications

Tolerance	Value	Comment	
PETS offset	100 μm	r <sub>c</sub> < 1 mm fulfilled	
PETS angles	~ 1 mrad	r <sub>c</sub> < 1 mm fulfilled	
Quad angles	~ 1 mrad	r <sub>c</sub> < 1 mm fulfilled	
Quad offset	20 μm	Must be small to be able to transport alignment beam	
BPM accuracy (incl. static misalignment and elec. error)	20 μm	Must be small to be able to perform initial correction	
BPM precision (diff. measurement)	~ 2 μm	Allows efficient suppression envelope growth due to dispersive trajectories	

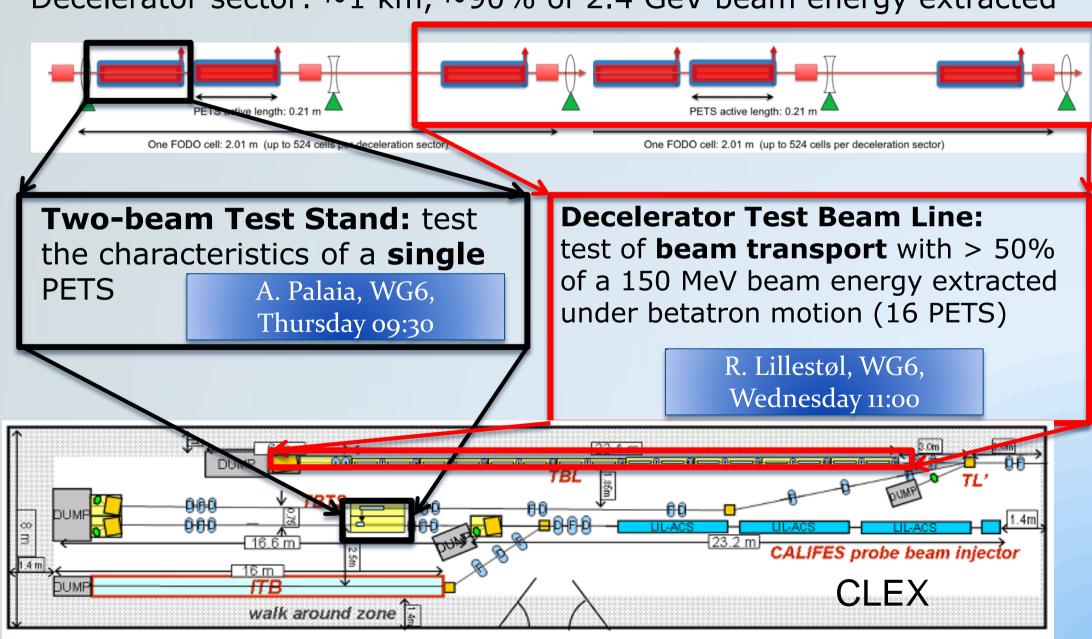
Tolerance	Value	Comment
Quadrupole position jitter	1 μm	r/r <sub>0</sub> < 5 %
Quadrupole field ripple	5· 10 <sup>-4</sup>	r/r <sub>0</sub> < 5 %
Current jitter	< 1%	Stability req. only – RF power constraints might be tighter.
Beta mismatch, dβ/β	~10 %	r/r <sub>0</sub> < 5 %
Injection offset, $y/\sigma_y$	< 0.2	r/r <sub>0</sub> < 5 %

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### Decelerator test-facilities

Decelerator sector: ~1 km, ~90% of 2.4 GeV beam energy extracted



42.5 m

- Requirements
- Beam physics
- Component specifications
- Test facilities
- Status and conclusions

# clc

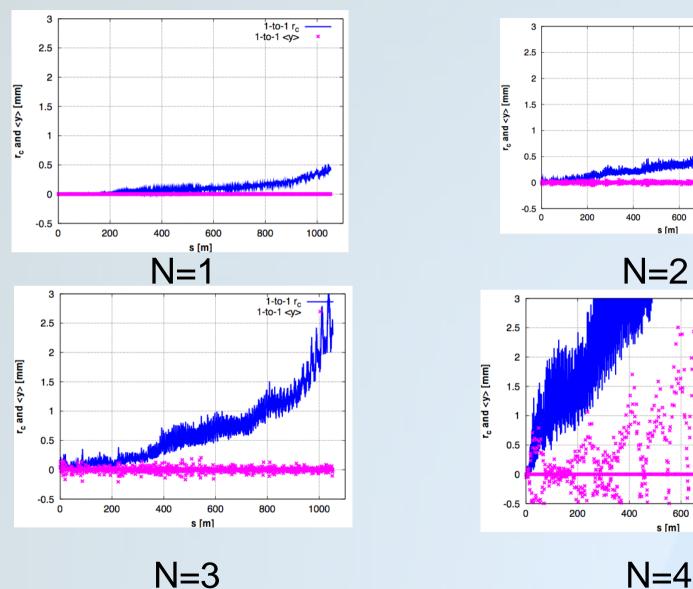
# CLIC decelerator status

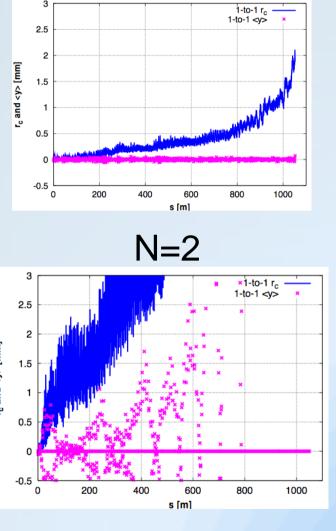
- Simulation framework in place allowing for detailed specification
- A CLIC baseline has been reached where simulations show satisfying beam transport for the baseline parameters
- Component specifications are sometimes tight, but within the feasibility limits
- Items outstanding :
  - Experimental tests of heavily decelerated beam
  - More detailed studies of collective effects for the decelerating beam
  - More detailed machine protection and beam loss studies
  - Benchmarking with other simulation codes would be comforting
  - Further cost optimization (clever component design, further optimization); might be seen in context of a larger iteration of drive beam parameters (TDR?)

- Requirements
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## Effect on reducing number of BPMs

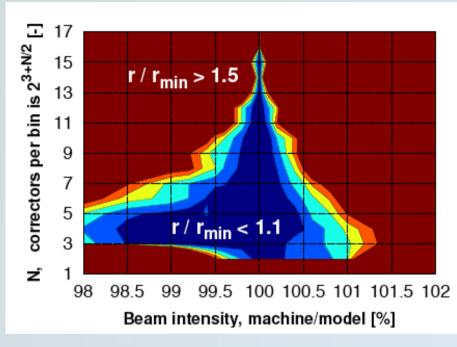


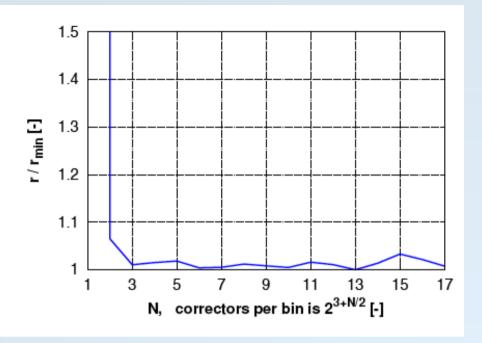


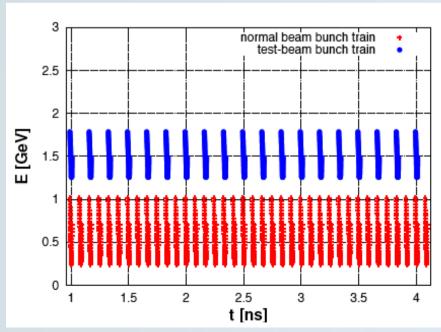
(perfect BPMs and single machine simulated, for illustration purposes)

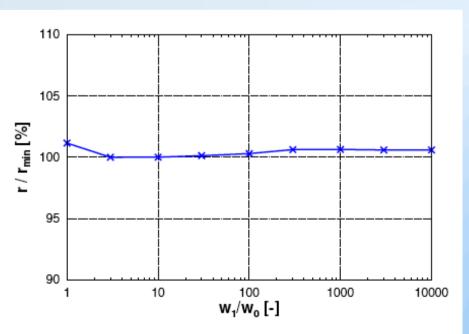


# Dispersion-free steering details



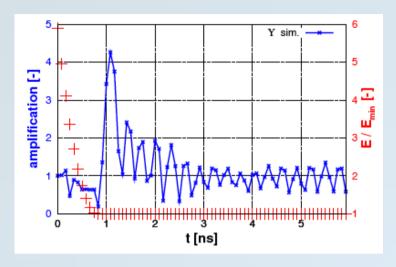


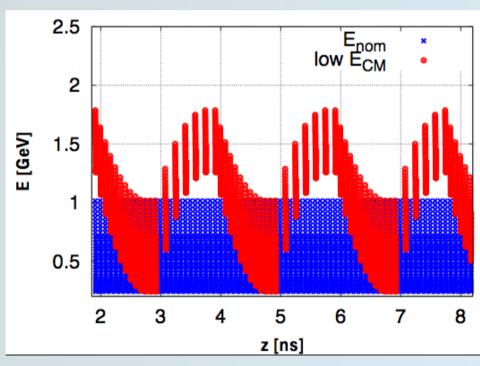


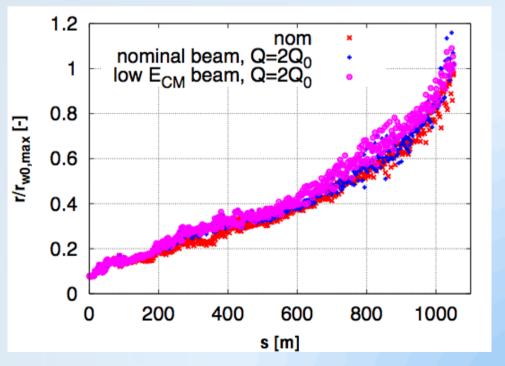




## Delayed switching: low energy running







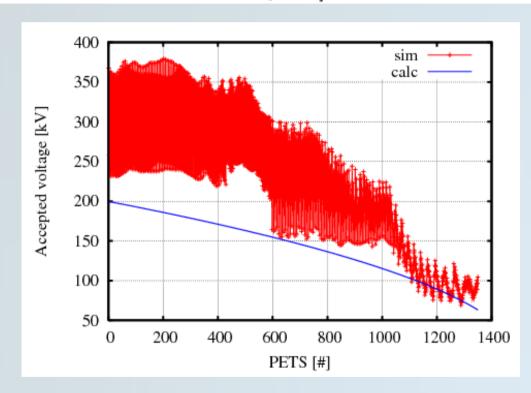


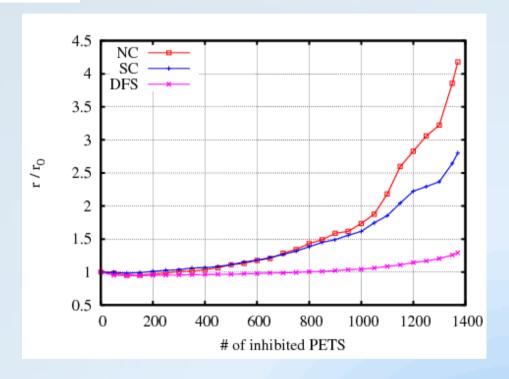
## Power phase-lag in the decelerator



# PETS on/off and kicks

$$U = \Delta y' \times E = \frac{r}{A\hat{\beta}} / \sqrt{\frac{E_i}{E_f}} \times E_i = \frac{r}{A\hat{\beta}} \sqrt{E_i E_f}$$





A number of random PETS inhibited (averaged over 100 seeds)