



Damping Rings



HeadTail simulations for the impedance in the CLIC-DRs

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CERN



Damping Rings

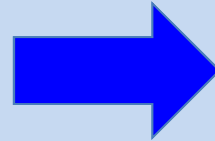


Outline

- DRs parameters
- Head tail simulations for impedance
- Results
- Conclusions
- Studies on the resistive wall from a chamber with coating
- Future steps

Updated list of parameters with the new lattice design at 3 TeV

Parameter	Symbol	Value
Energy	p_0 (GeV)	2.424
Norm. transv. emitt.	$\epsilon_{xn,yn}$ (nm)	381, 4.1
Bunch length	σ_z (mm)	1.53
Momentum spread	σ_δ	1.43×10^{-3}
Bunch spacing	ΔT_b (ns)	0.5
Bunch population	N_b	4.1×10^9
Circumference	C (m)	365.2
Coupling	(%)	0.13
Mom. compact.	α	8×10^{-5}
Number of bunches	n_b	312
Tunes	$Q_{x,y,s}$ (m)	69.82, 33.80
Store time/train	T_{st} (ms)	20
Energy loss	ΔE (MeV/turn)	3.857
Damping times	$\tau_{x,y,z}$ (ms)	1.5, 1.5, 0.74
RF frequency	f_{rf} (GHz)	2
RF voltage	V_{rf} (MV)	4.115
Bend length	L_{bend} (m)	0.545
Bend chamber rad.	R_{bend} (cm)	2
Number of bends	N_{bend} (m)	96
Wiggler length	L_w (m)	2
Wiggler field	B_w (T)	2.5
Number of wigglers	N_w (m)	76
Wiggler radius	r_w (mm)	9



With combined function magnets

Parameter	Symbol	Value
Energy	p_0 (GeV)	2.86
Norm. transv. emitt.	$\epsilon_{xn,yn}$ (nm)	480, 4.7
Bunch length	σ_z (mm)	1.4
Momentum spread	σ_δ	1×10^{-3}
Bunch spacing	ΔT_b (ns)	0.5
Bunch population	N_b	4.1×10^9
Circumference	C (m)	493.05
Coupling	(%)	0.1
Mom. compact.	α	6×10^{-5}
Number of bunches	n_b	312
Tunes	$Q_{x,y,s}$ (m)	58.2, 18.8
Store time/train	T_{st} (ms)	20
Energy loss	ΔE (MeV/turn)	5.9
Damping times	$\tau_{x,y,z}$ (ms)	1.6, 1.6, 0.8
RF frequency	f_{rf} (GHz)	2
RF voltage	V_{rf} (MV)	7.2
Bend length	L_{bend} (m)	0.4
Bend chamber rad.	R_{bend} (cm)	1
Number of bends	N_{bend} (m)	96
Wiggler length	L_w (m)	2
Wiggler field	B_w (T)	2.5
Number of wigglers	N_w (m)	76
Wiggler radius	r_w (mm)	9

⇒ Advantages: DA increased, magnet strength reduced to reasonable, reduced IBS

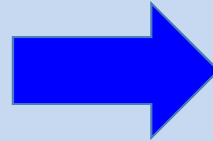
⇒ Relative to collective effects (main changes):

- Higher energy, larger horizontal emittance (good)
- Longer circumference (bad)

**From Y. Papaphilippou,
G. Rumolo, CLIC'09³**

New update of the lattice design at 3 TeV

Parameter	Symbol	Value
Energy	p_0 (GeV)	2.86
Norm. transv. emitt.	$\epsilon_{x,y,n}$ (nm)	480, 4.7
Bunch length	σ_z (mm)	1.4
Momentum spread	σ_δ	1×10^{-3}
Bunch spacing	ΔT_b (ns)	0.5
Bunch population	N_b	4.1×10^9
Circumference	C (m)	493.05
Coupling	(%)	0.1
Mom. compact.	α	6×10^{-5}
Number of bunches	n_b	312
Tunes	$Q_{x,y,s}$ (m)	58.2, 18.8
Store time/train	T_{st} (ms)	20
Energy loss	ΔE (MeV/turn)	5.9
Damping times	$\tau_{x,y,z}$ (ms)	1.6, 1.6, 0.8
RF frequency	f_{rf} (GHz)	2
RF voltage	V_{rf} (MV)	7.2
Bend length	L_{bend} (m)	0.4
Bend chamber rad.	R_{bend} (cm)	1
Number of bends	N_{bend} (m)	96
Wiggler length	L_w (m)	2
Wiggler field	B_w (T)	2.5
Number of wigglers	N_w (m)	76
Wiggler radius	r_w (mm)	9



1 GHz option

Parameter	Symbol	Value
Energy	p_0 (GeV)	2.86
Norm. transv. emitt.	$\epsilon_{x,y,n}$ (nm)	180, 1.5
Bunch length	σ_z (mm)	1.6
Momentum spread	σ_δ	1.3×10^{-3}
Bunch spacing	ΔT_b (ns)	1
Bunch population	N_b	4.1×10^9
Circumference	C (m)	420.56
Coupling	(%)	0.1
Mom. compact.	α	7.6×10^{-5}
Number of bunches per train	n_b	156
Number of trains	n_t	2
Distance between trains	τ_t (ns)	545
Tunes	$Q_{x,y,s}$	55.4, 11.6, 0.00387
Store time/train	T_{st} (ms)	20
Energy loss	ΔE (MeV/turn)	4.2
Damping times	$\tau_{x,y,z}$ (ms)	1.88, 1.91, 0.96
RF frequency	f_{rf} (GHz)	1
RF voltage	V_{rf} (MV)	4.9
Harmonic number	h	1402
Dipole length	L_{dip} (m)	0.43
Dipole chamber rad.	R_{dip} (cm)	1
Number of dipoles	N_{dip} (m)	102
Wiggler length	L_w (m)	2
Wiggler field	B_w (T)	2.5
Number of wigglers	N_w (m)	52
Wiggler gap	r_w (mm)	13
Wiggler width	h_w (mm)	65
Average β_x in wigglers	$\langle \beta_{xw} \rangle$ (m)	4.787
Average β_y in wigglers	$\langle \beta_{yw} \rangle$ (m)	4.185

⇒ Lattice has been redesigned to reduce the space charge effect (ring circumference shortened).
However, higher order cavities will also help in this sense (simulations foreseen)

⇒ The 1 GHz option has been considered because:

- it is better for the RF design (less impedance)
- it could relieve constraints due to e-cloud, ions, coupled-bunch instabilities, ...



Damping Rings



HeadTail simulations for impedance

- Use of HeadTail code
- Simulates single bunch phenomena
- Broadband impedance model
- Tuneshift in horizontal and vertical plane
- Transverse shunt Impedance range: 0-20 M Ω m/m
- 0 and different positive values in chromaticity
- Round and flat geometry



Damping Rings



CLIC-DR parameters used in the simulations

1 GHz
or 2GHz

Tunes $Q_x/Q_y/Q_s$	55.4/11.6/0.00387
Ring circumference (m)	420.56
Number of turns	20000
Energy (GeV)	2.86
N_b	4.1×10^9
Geometry	round/flat
$\langle \beta_x \rangle$ (m) wigglers	4.787
$\langle \beta_y \rangle$ (m) wigglers	4.185

- **Transverse**

In a round chamber the TMCI threshold is given (chromaticity 0, coupling mode 0 and -1 assumed) :

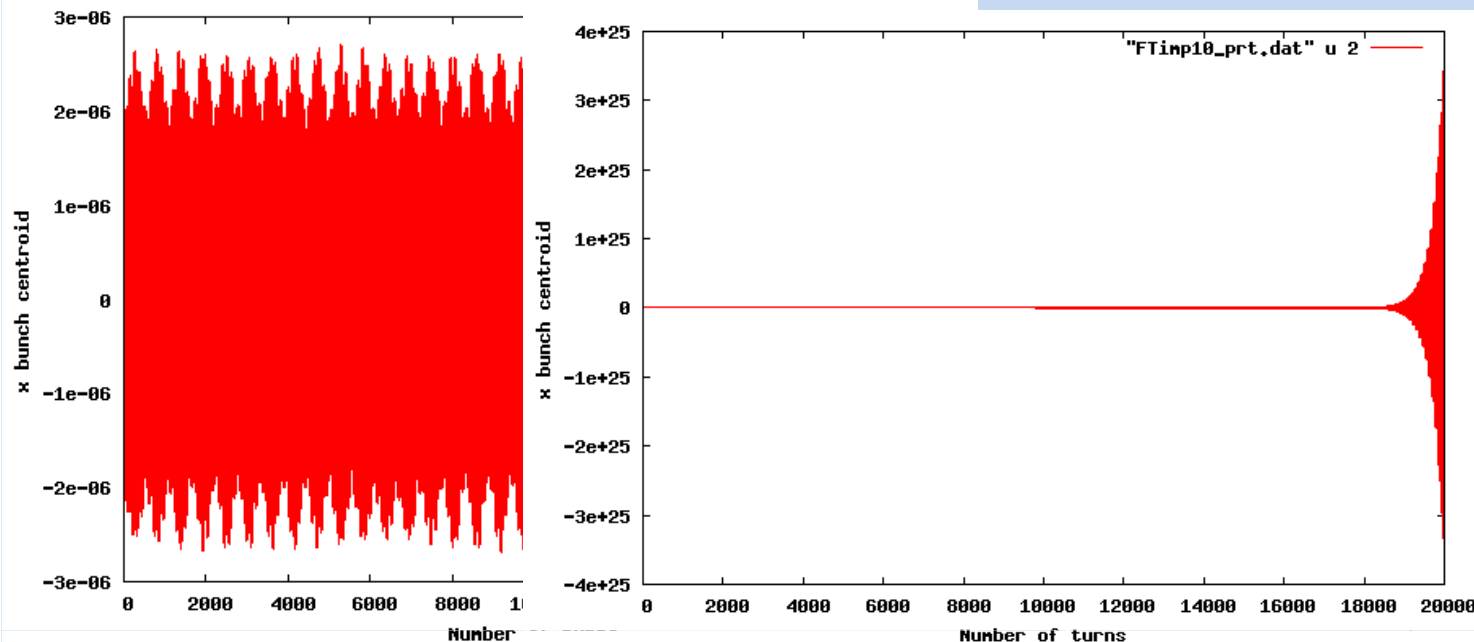
$$\xi < \frac{Q_s}{\omega_r \sigma_t} \quad \text{if } \omega_r \sigma_t \leq 1 \quad \longrightarrow \quad \text{short bunch}$$
$$\xi < \sqrt{2} Q Q_s (\omega_r \sigma_t)^2 \quad \text{if } \omega_r \sigma_t \gg 1 \quad \longrightarrow \quad \text{long bunch}$$

where

$$\xi = \frac{\omega_r / 2\pi <\beta_y> R_T N_b e}{3.75 Q E / e}$$

CLIC-DRs: impedance value of $\approx 12 \text{ MOhm/m}$ if $\omega_r = 2\pi \times 7 \text{ GHz}$

Horizontal and vertical motion in a round chamber



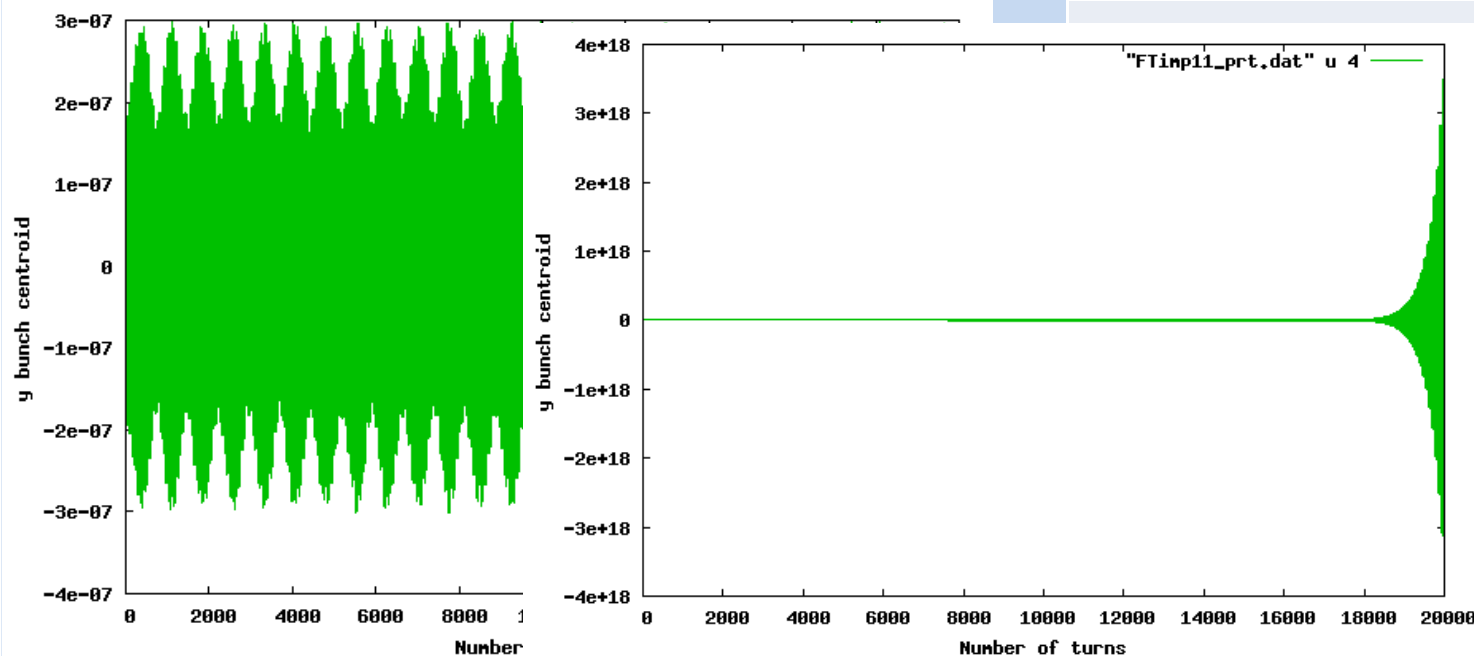
- Centroid evolution in x and y over the number of turns, for different values in impedance

0

0

Round

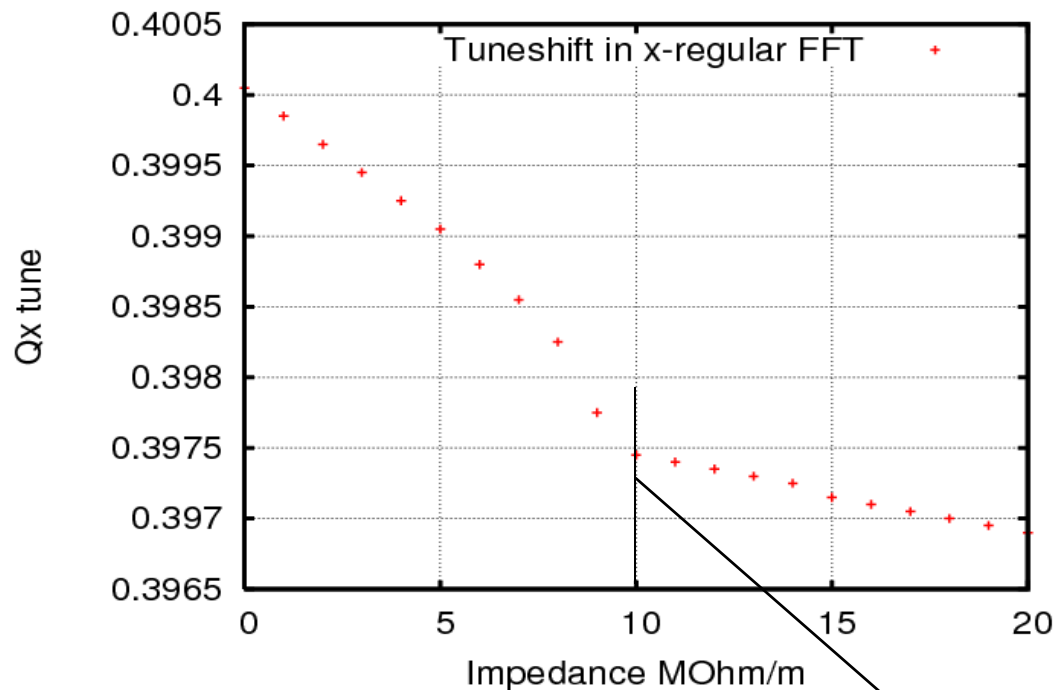
vert. chrom. Q y



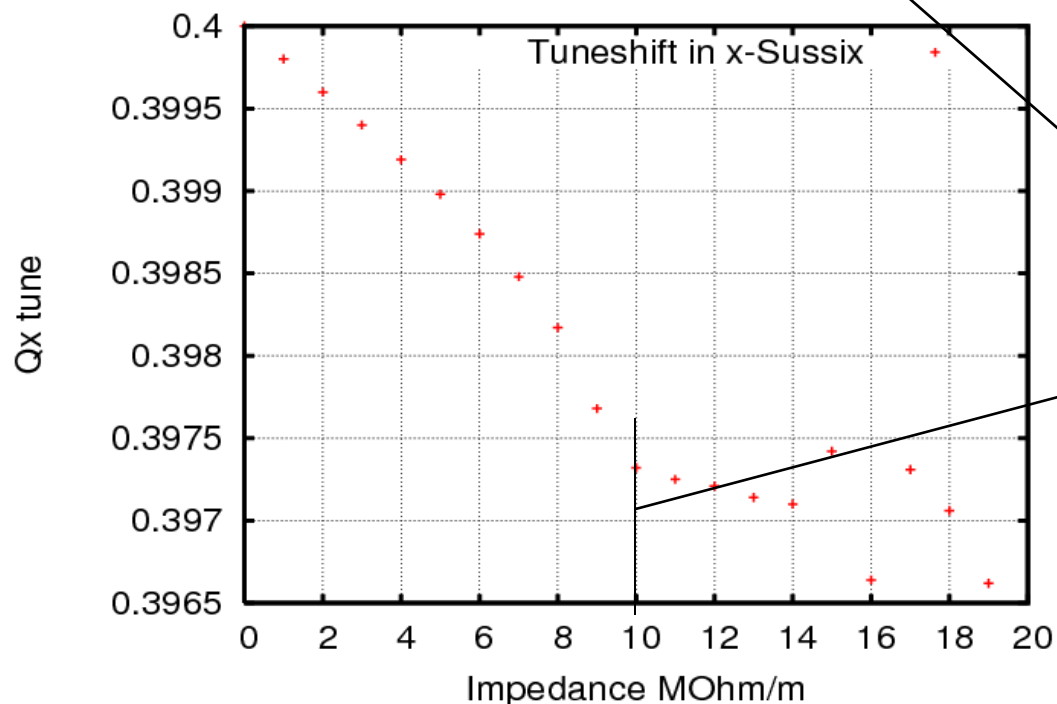
- As the impedance increases an instability occurs

Horizontal plane

- To observe a TMCI, the betatron frequency is measured , i.e. the frequency of the mode 0
- Classical and Sussix algorithms used for Fourier analysis of the coherent horizontal motion



Regular FFT

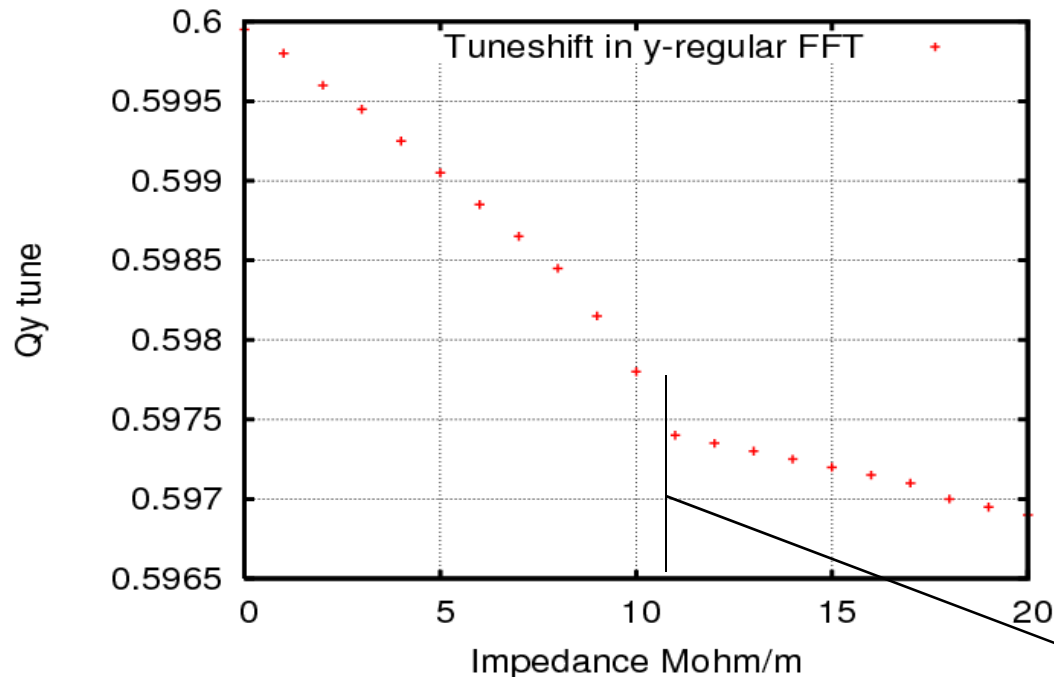


Sussix FFT

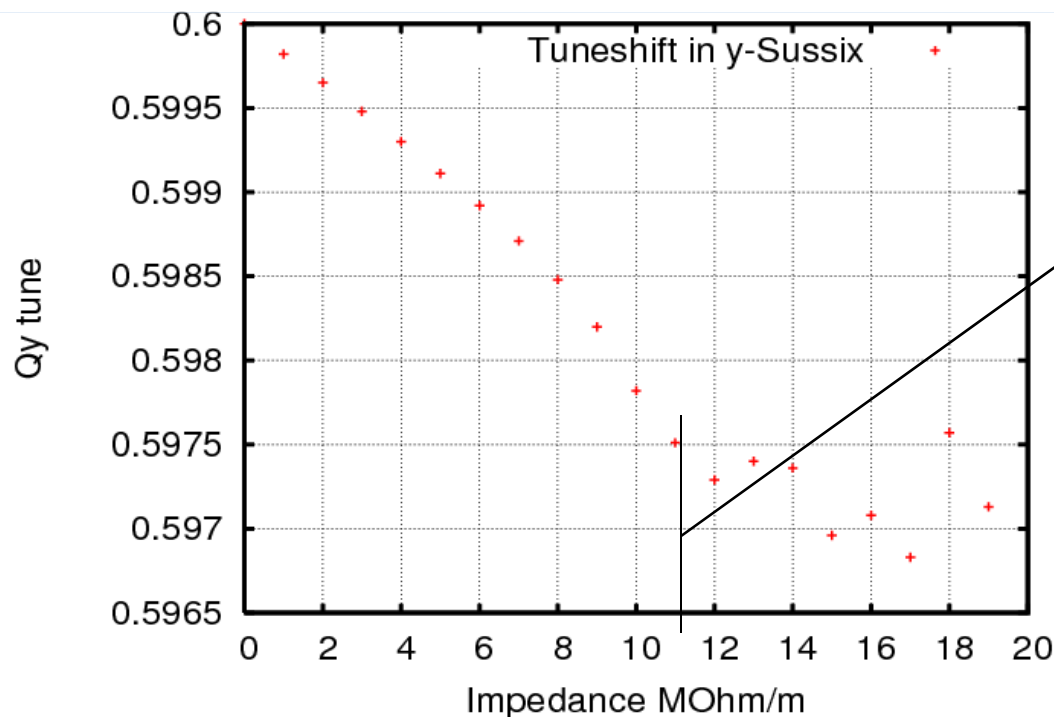
- Instability threshold (TMCI)
10 MOhm/m

Vertical plane

- Classical and Sussix algorithms used for Fourier analysis of the coherent vertical motion



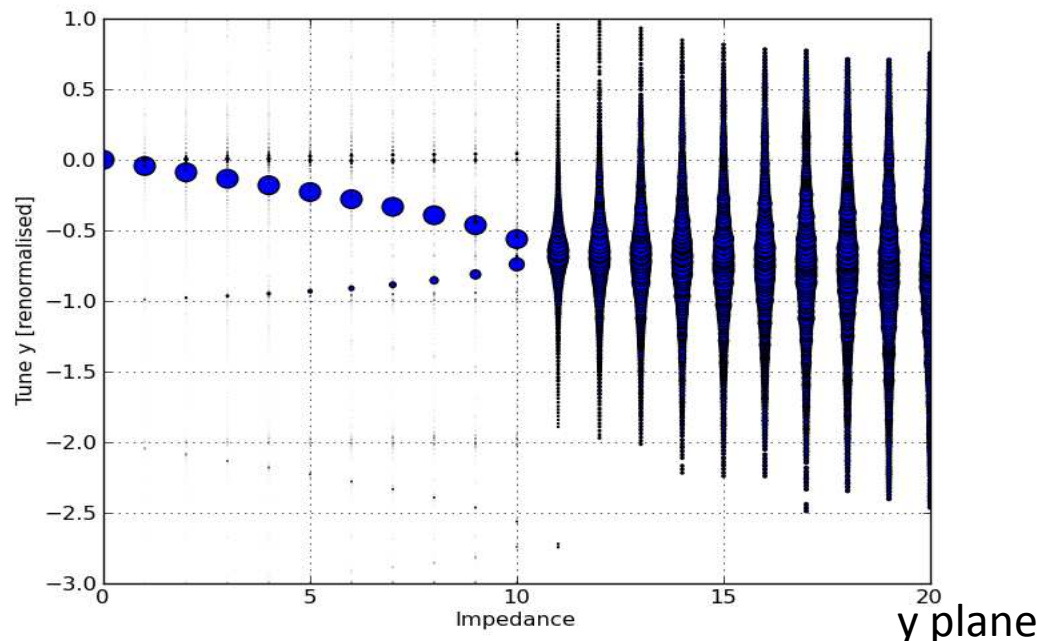
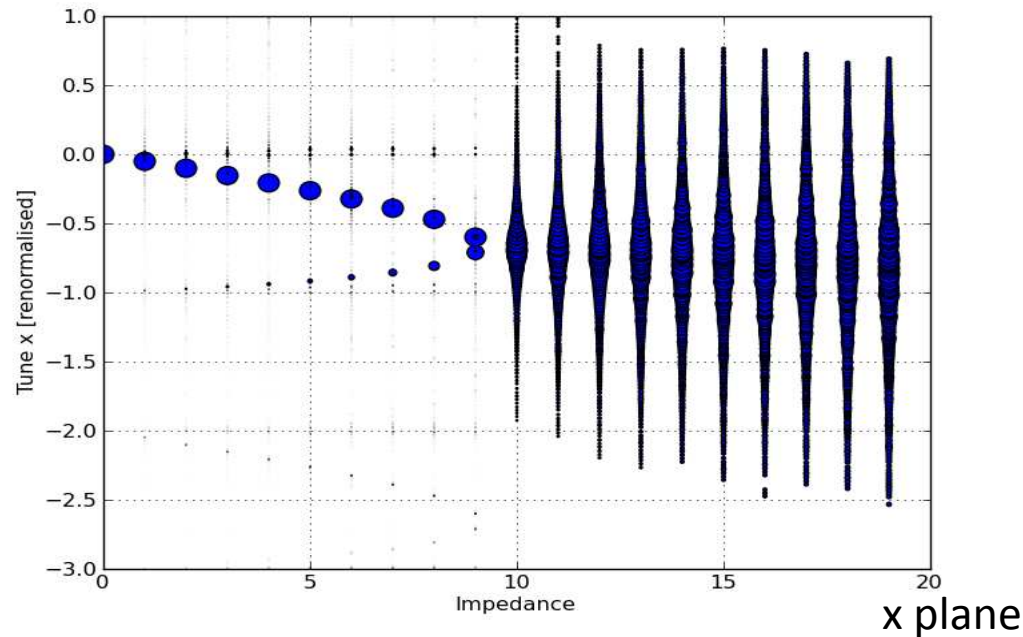
Regular FFT



Sussix FFT

• **Instability threshold (TMCI)**
11 MOhm/m

Mode spectrum of the horizontal and vertical coherent motion as a function of impedance

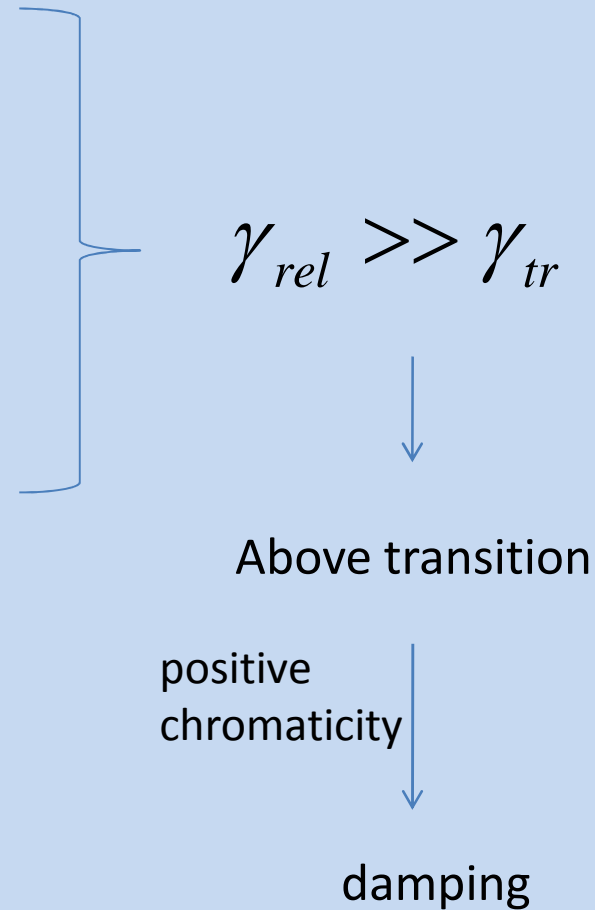


- Plot all the tunes $(Q-Q_x)/Q_s$ and $(Q-Q_y)/Q_s$ with impedance, from the Sussix results
- Mode spectrum represents the natural coherent oscillation modes of the bunch
- The movement of the modes due to impedance can cause them to merge and lead to an instability
- ☐ The mode 0 is observed to couple with mode -1 in both planes
- ☐ Causing a TMCI instability

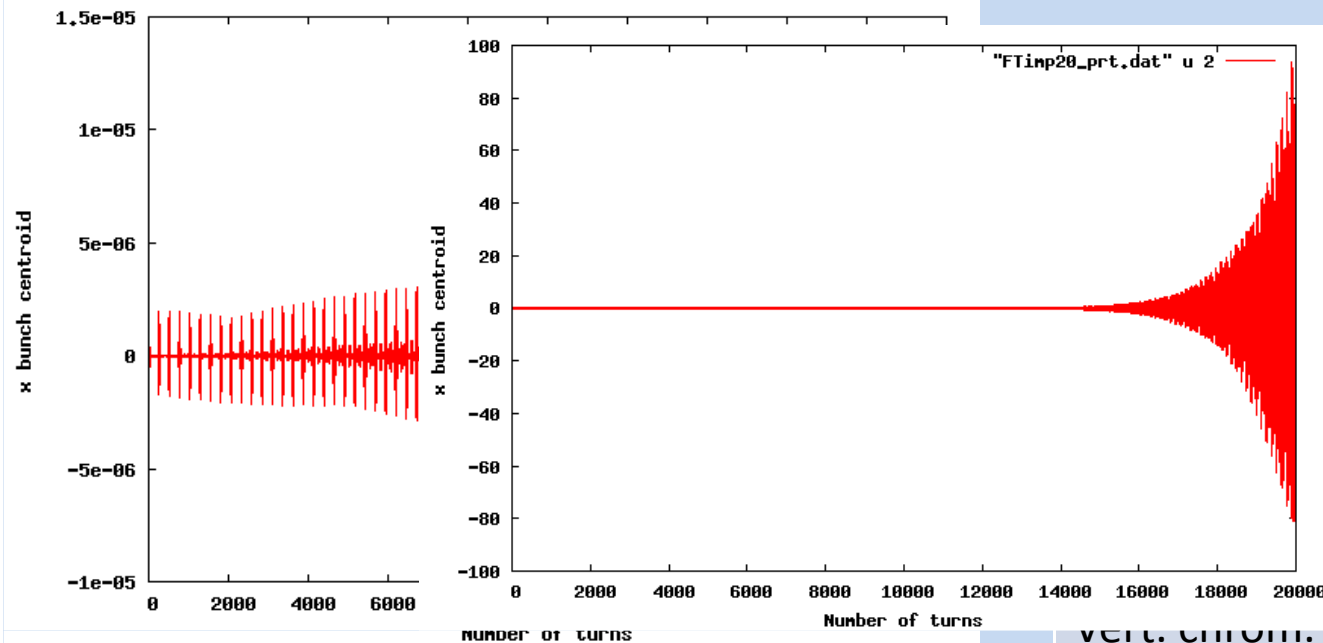
Damping Rings

$$\gamma_{tr} = \frac{1}{\sqrt{a_c}} = \frac{1}{\sqrt{7.6 \times 10^{-5}}} = 115$$

$$\gamma_{rel} = 5597$$

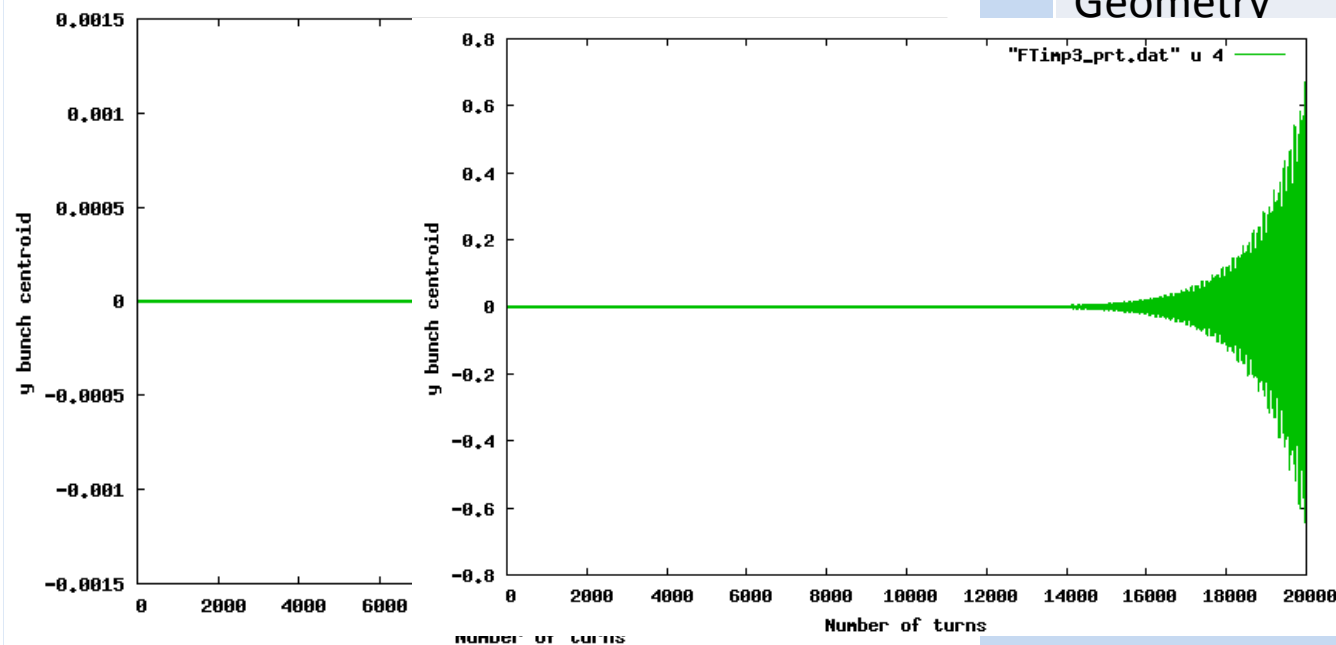


Horizontal and vertical motion in a round chamber

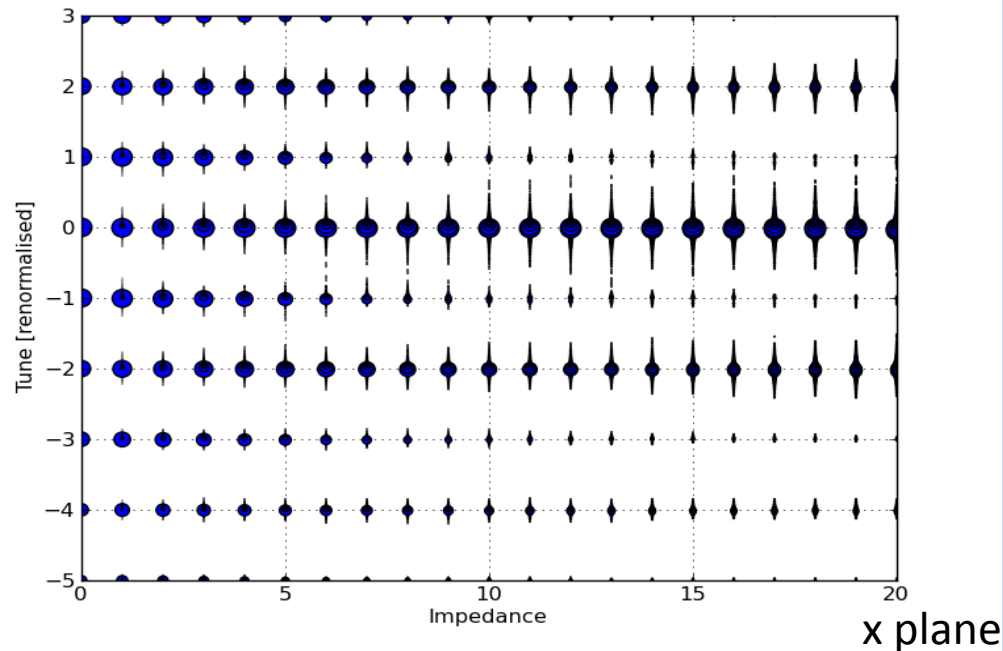


Instability growth

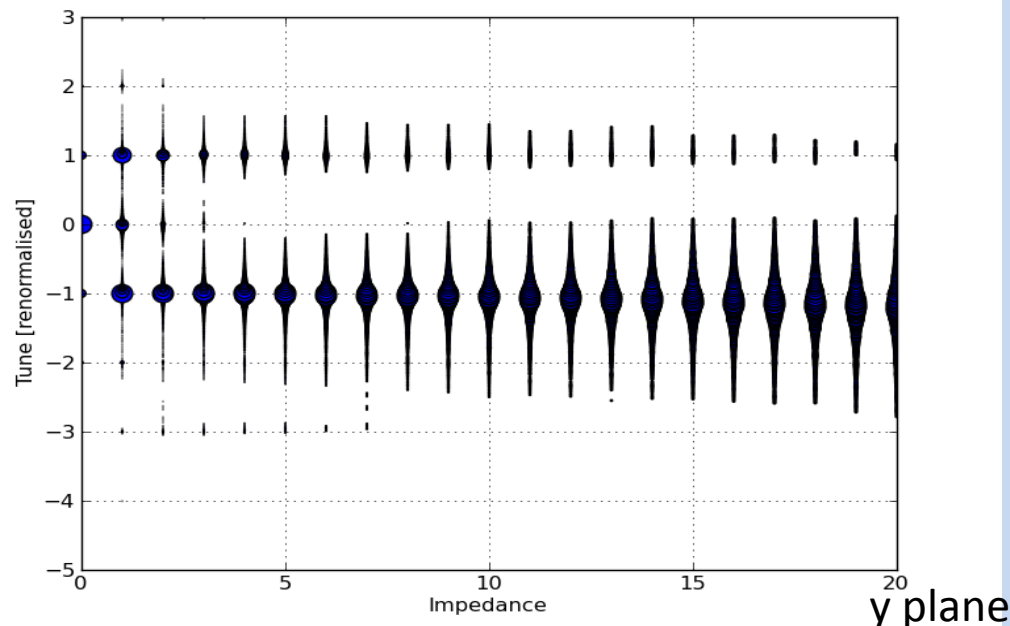
Q' x	9.2
Q' y	1.9
Geometry	Round



Mode spectrum of the horizontal and vertical coherent motion as a function of impedance



- no mode coupling observed, no TMCI

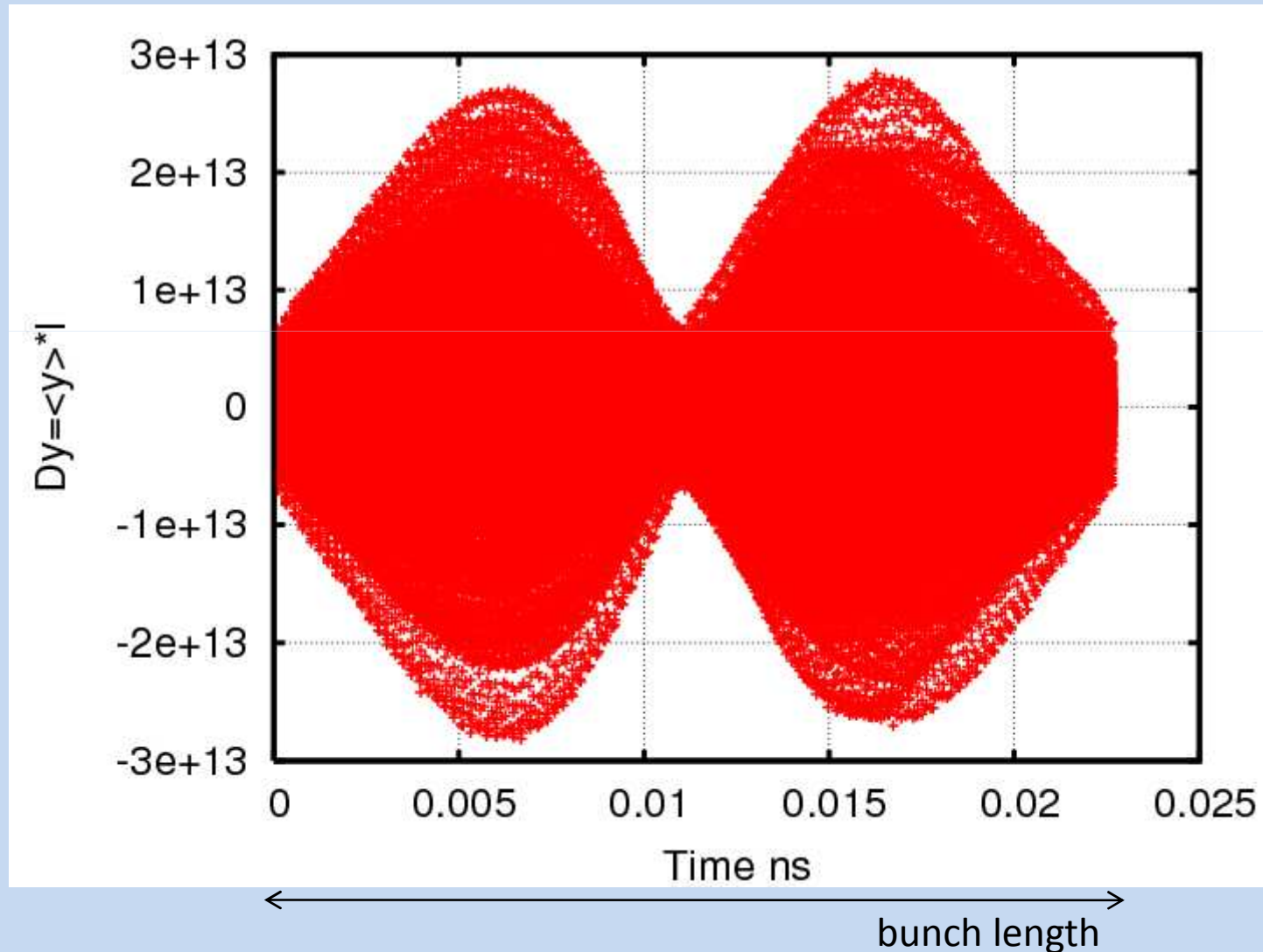


- no mode coupling
- mode 1 is damped
- mode -1 gets unstable

- Presence of chromaticity makes the modes move less, good for coupling
- Another type of instability ¹⁴

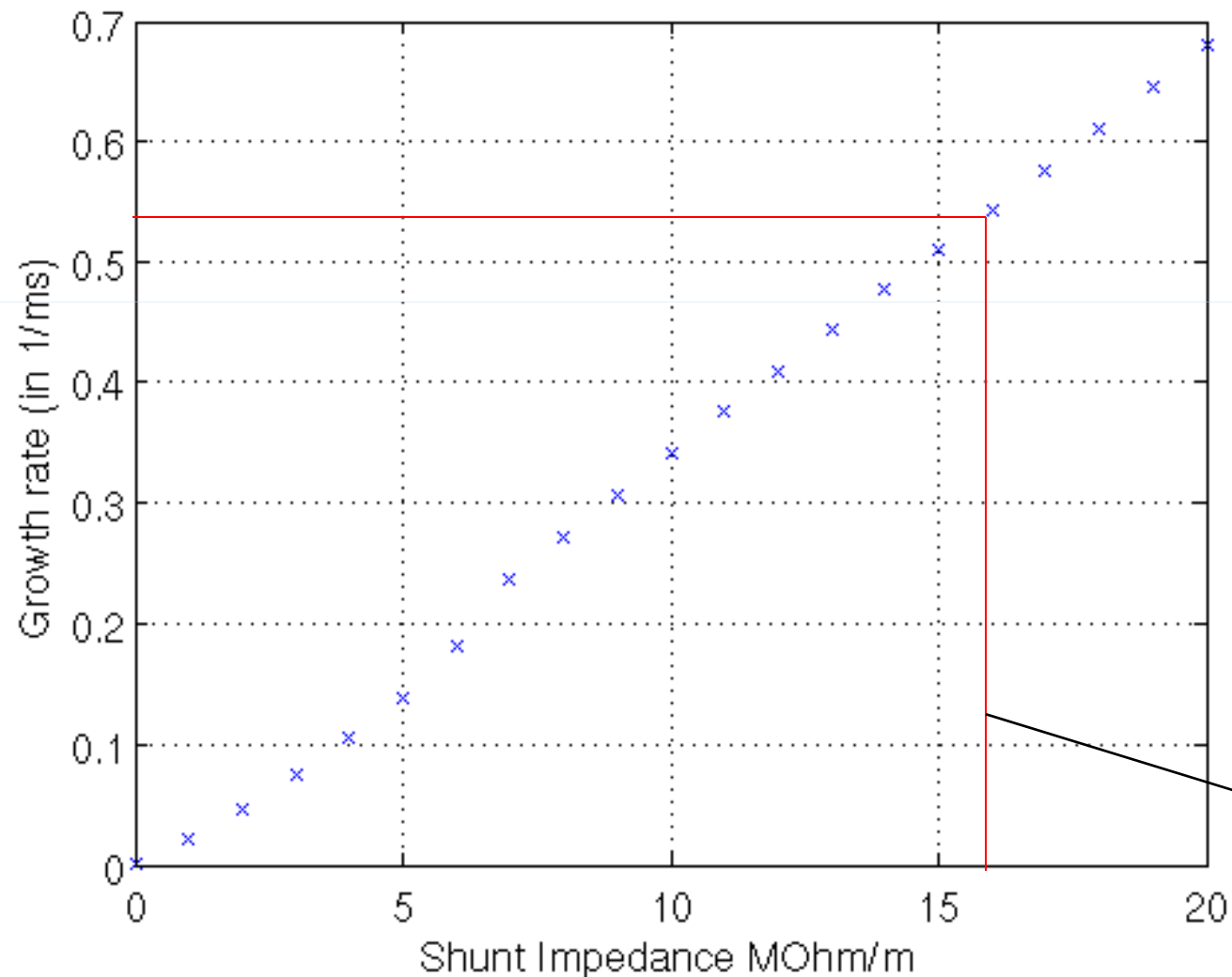
Damping Rings

Looking at hdtl.dat file (information along the bunch)



- y plane
- mode -1

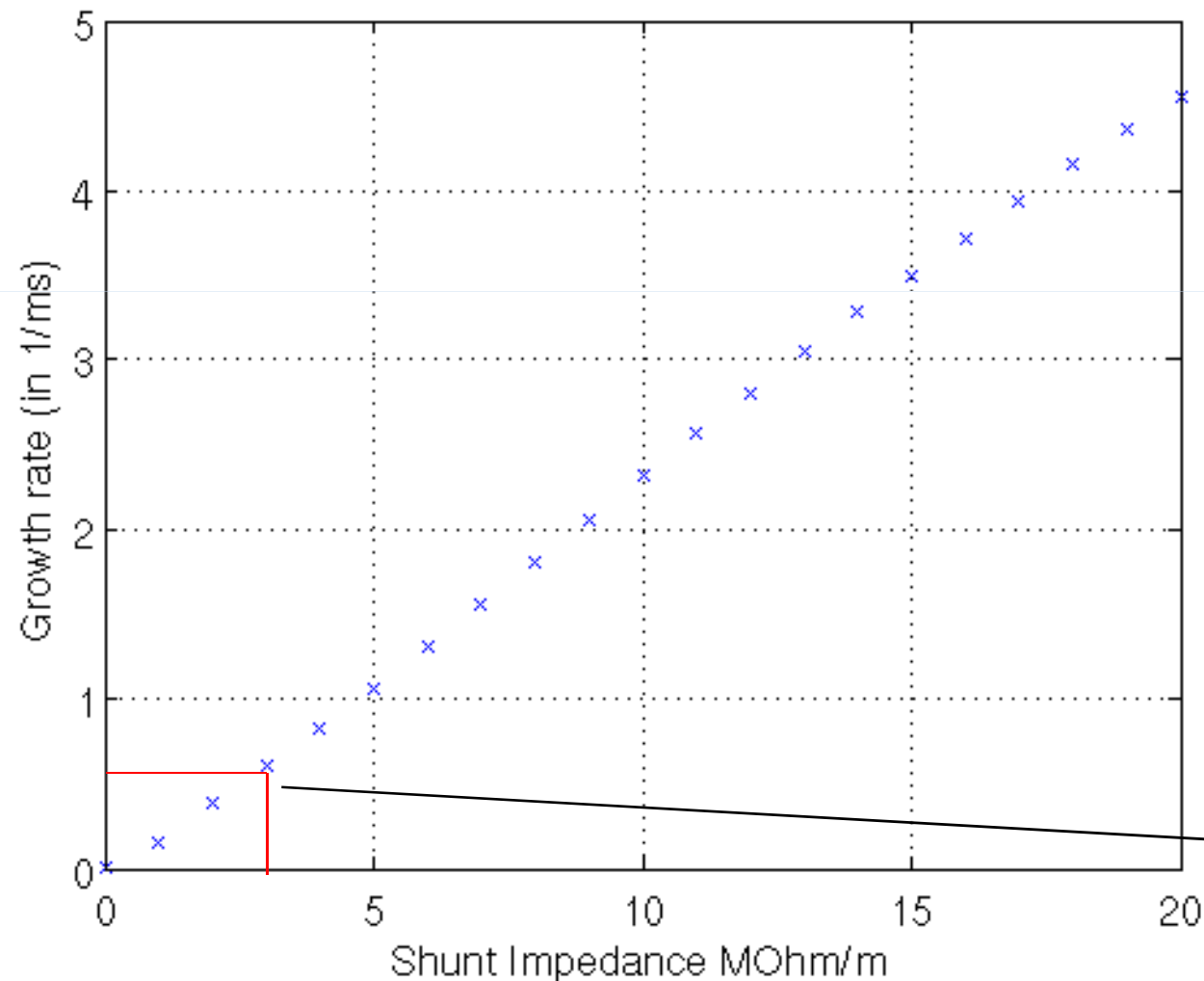
Growth rate – x plane



- Damping time $\tau_x = 1.88 \text{ ms}$
- $> \sim 16 \text{ MOhm/m}$
rise time $< \tau_x$
unstable

Threshold
 $\sim 16 \text{ MOhm/m}$

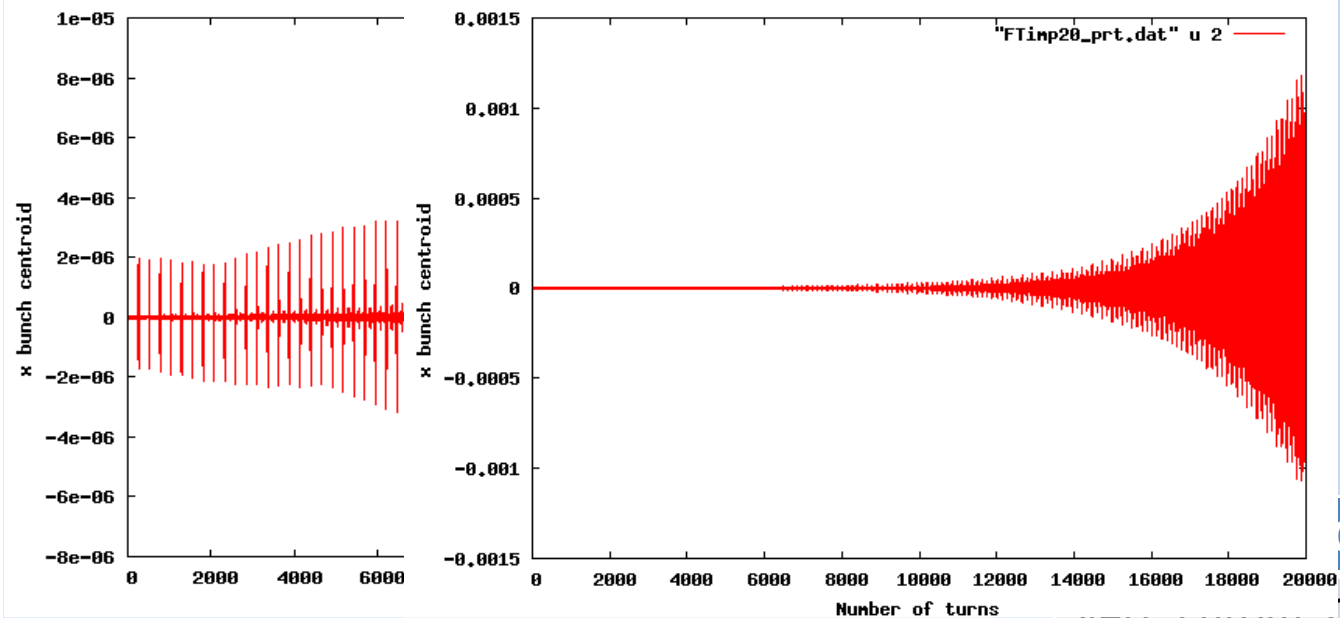
Growth rate – y plane



- Damping time $\tau_y = 1.91 \text{ ms}$
- For impedances above $\sim 3 \text{ MOhm/m}$, rise time $< \tau_y$
- mode -1 is dangerous leading to instability

Threshold
 $\sim 3 \text{ MOhm/m}$

Horizontal and vertical motion in a round chamber



Q'_x

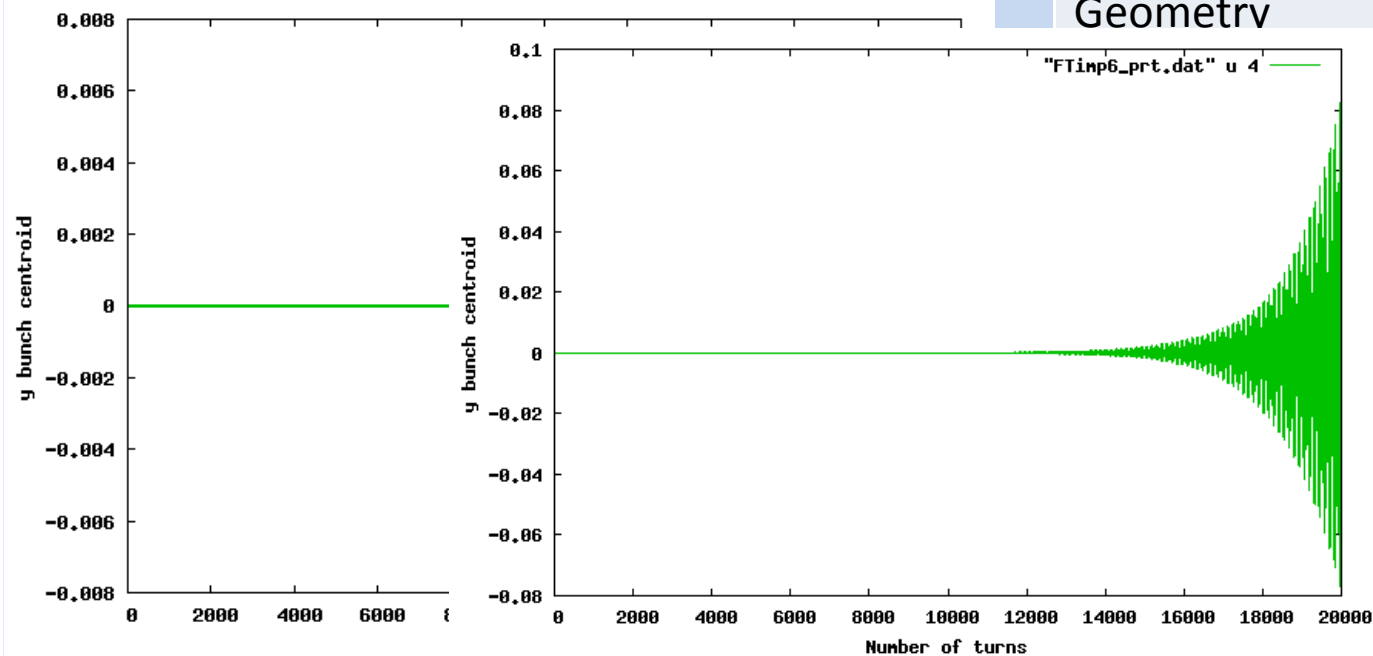
18.4

vert. chrom. Q'_y

3.8

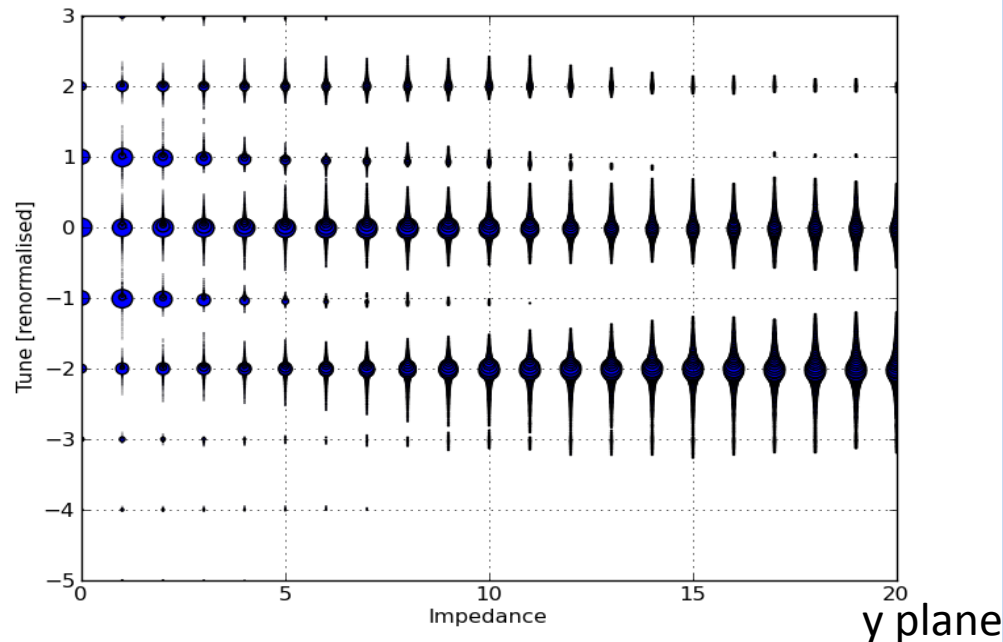
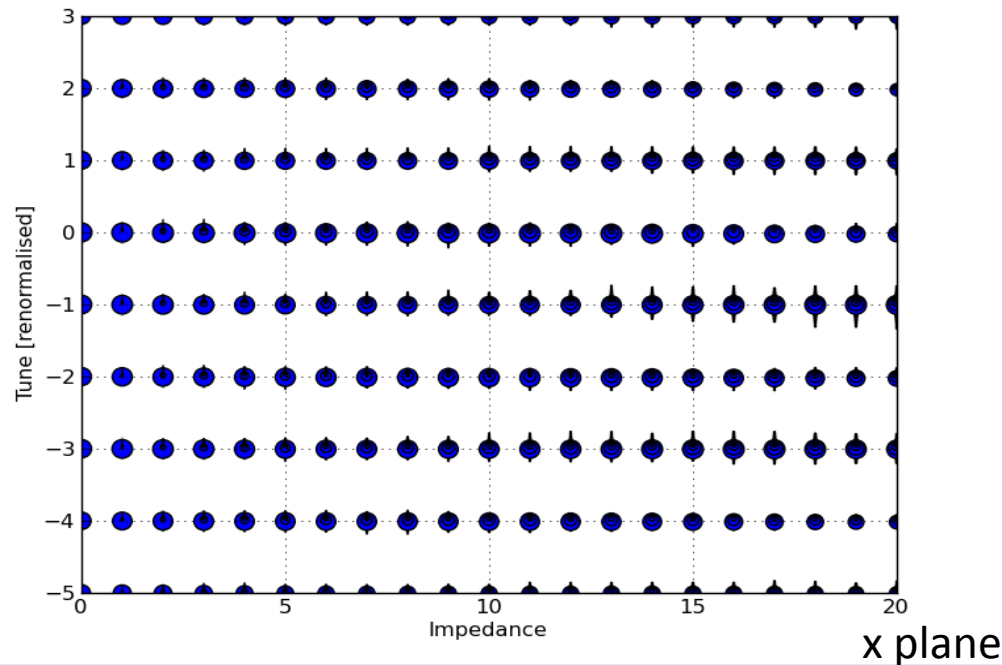
Geometr

Round



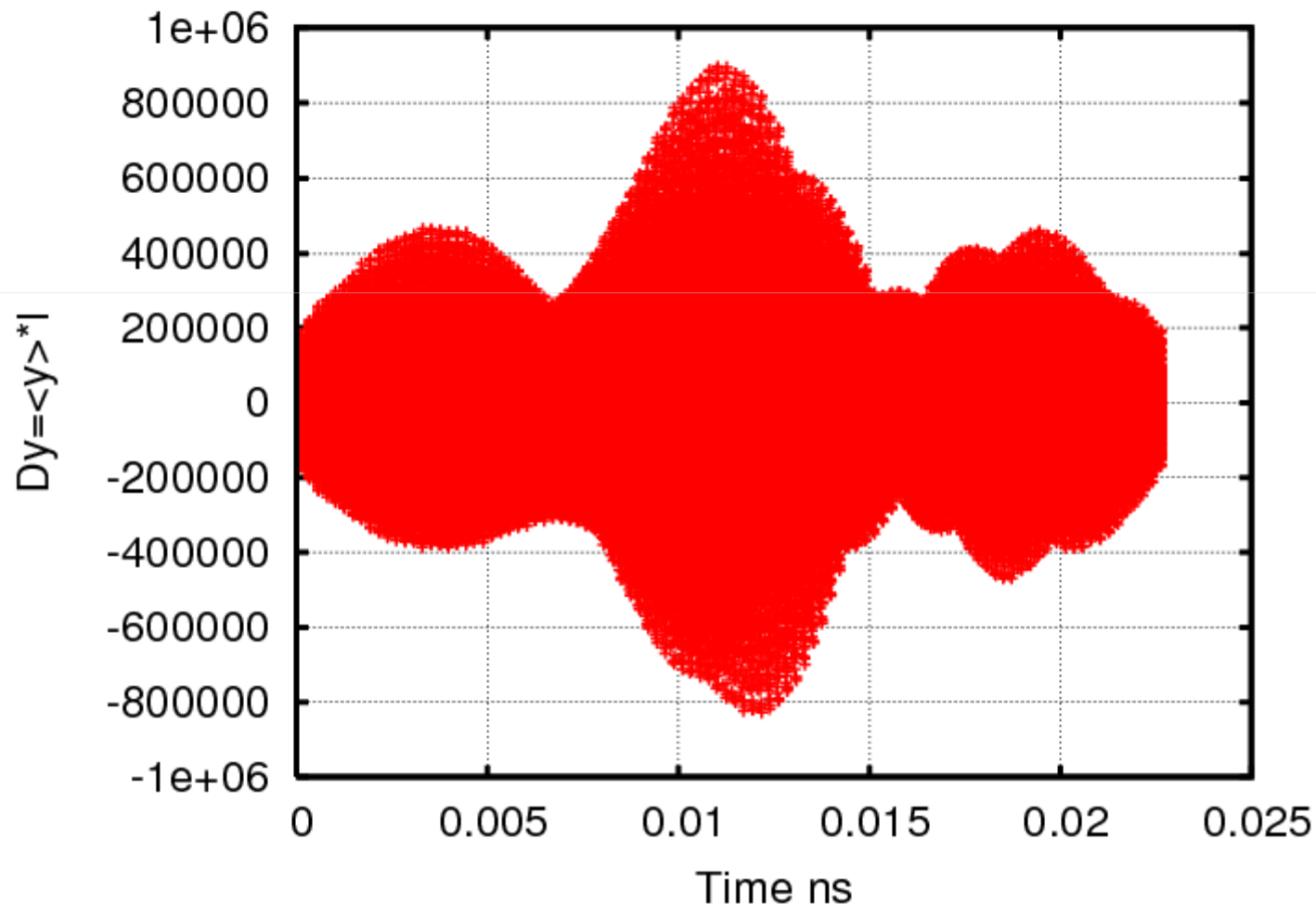
Mode spectrum of the horizontal and vertical coherent motion as a function of impedance

- no mode coupling (TMCI)



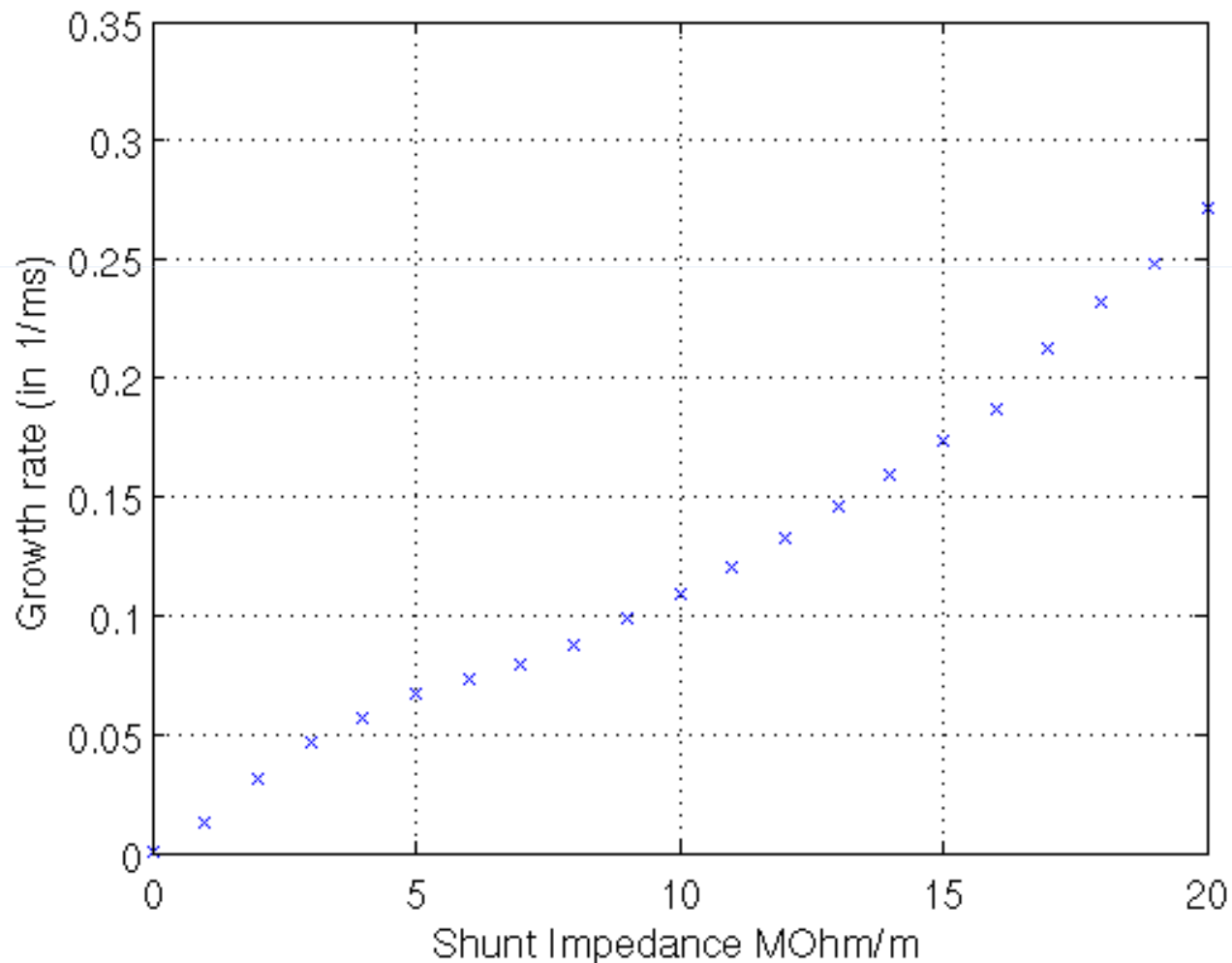
- As the chromaticity is increased , the main unstable mode changes
- mode -2 gets unstable in the y plane

Looking at the hdtl.dat file



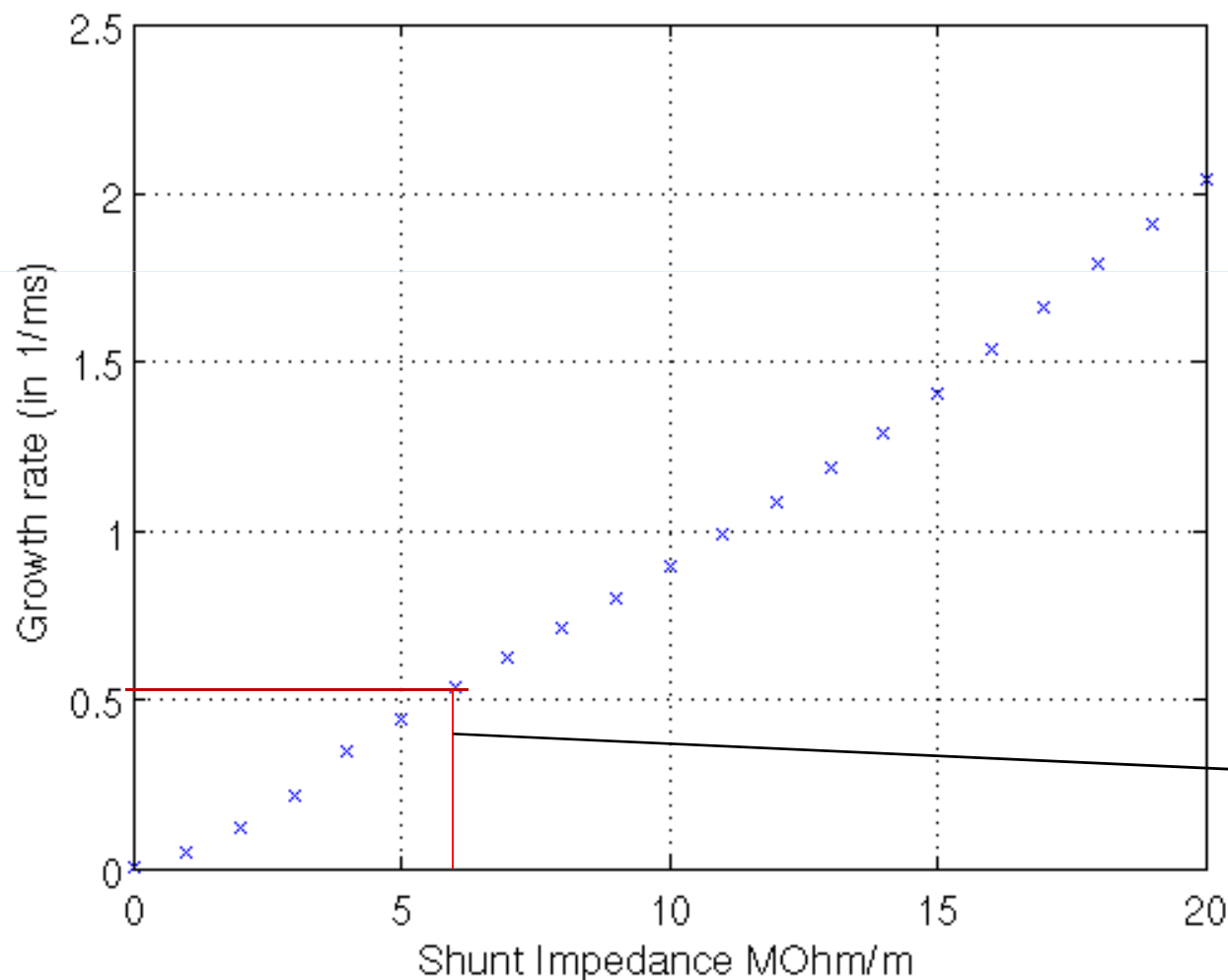
- y plane
- mode -2
- another higher order mode

Growth rate – x plane



- Damping time $\tau_x = 1.88 \text{ ms}$
 - rise time $> \tau_x$
- stable**

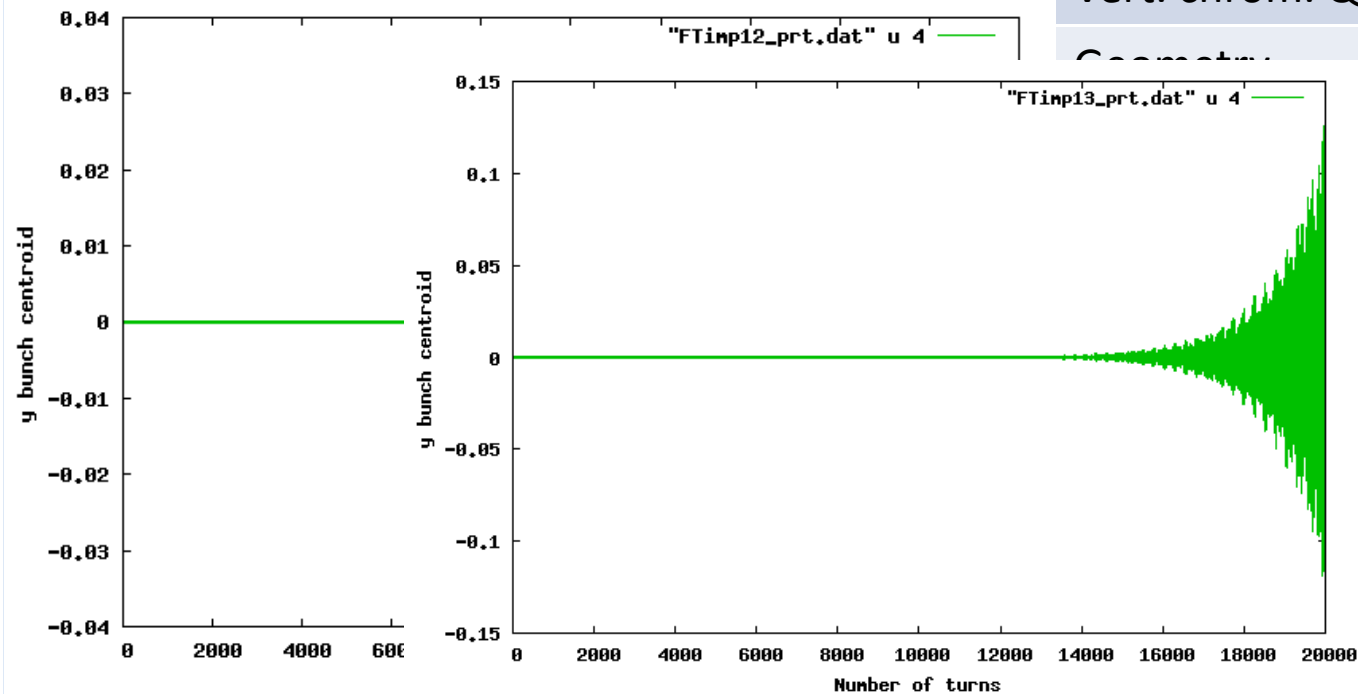
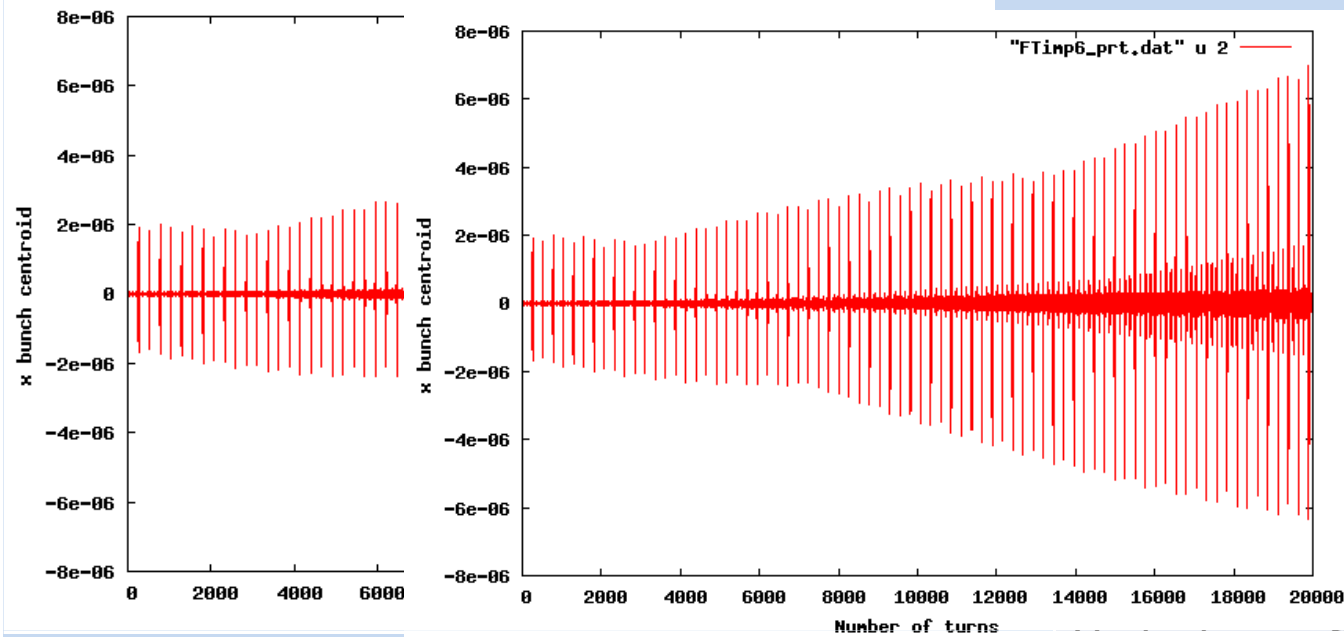
Growth rate – y plane



- Damping time $\tau_y = 1.91 \text{ ms}$
- $> \sim 6 \text{ MOhm/m}$
rise time $< \tau_y$
unstable

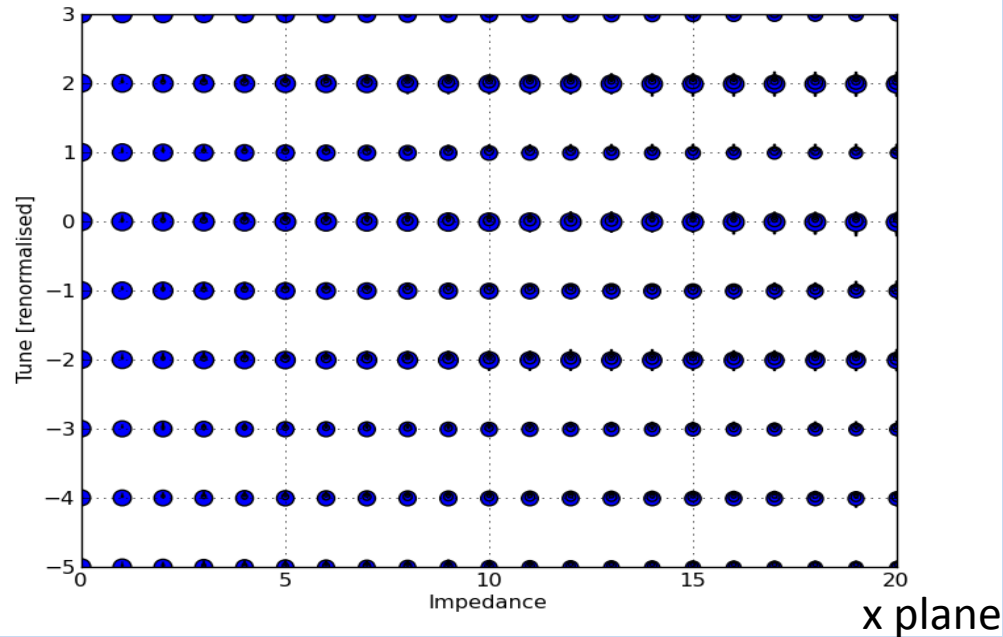
Threshold
~6MOhm/m

Horizontal and vertical motion in a round chamber

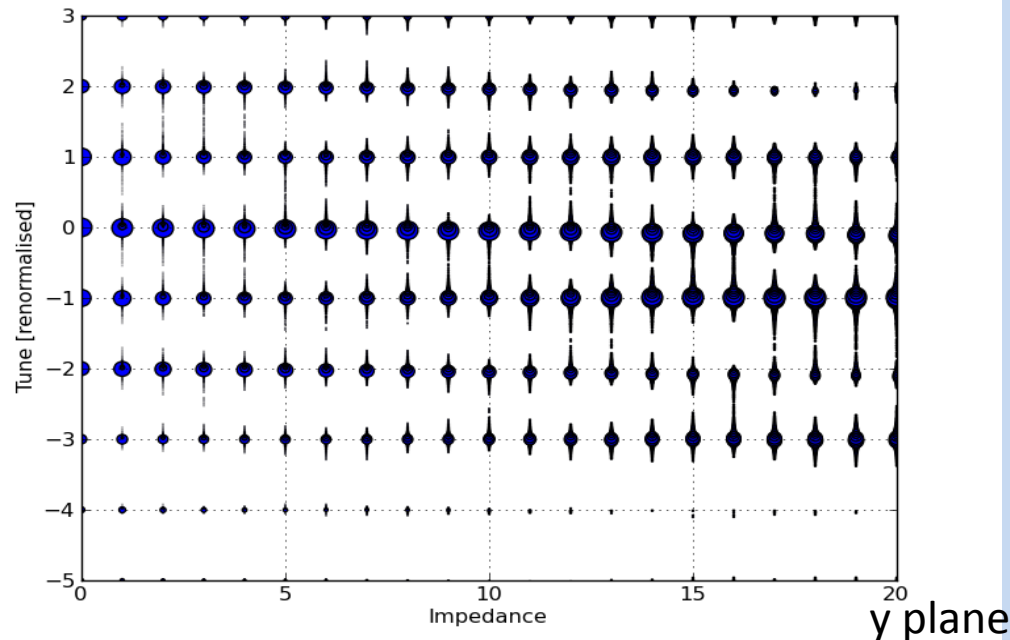


Q'x	27.6
vert. chrom. Q'y	5.7
Geometry	Round

Mode spectrum of the horizontal and vertical motion as a function of impedance

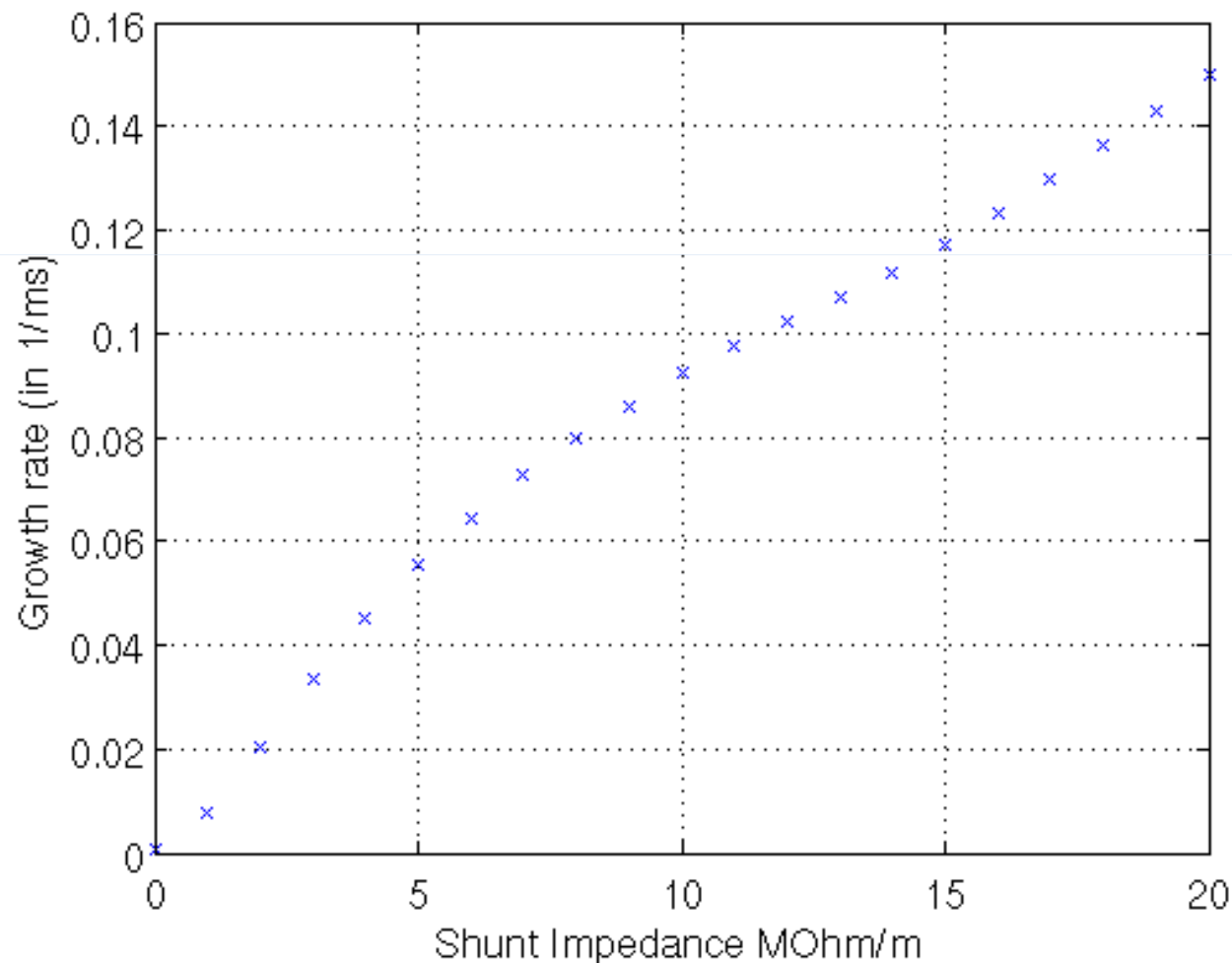


- Gets harder to see the cause of the instability



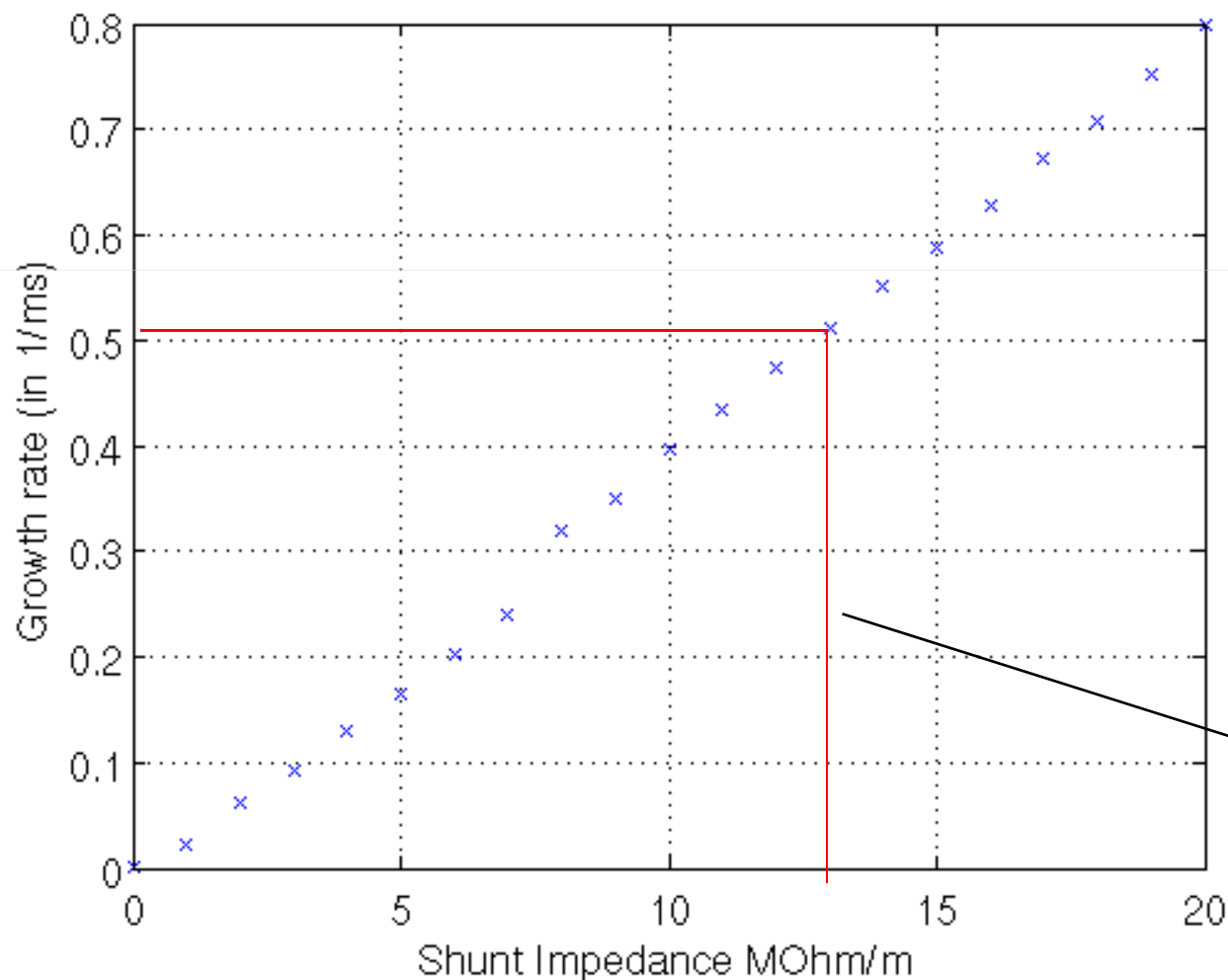
- As the chromaticity is increased, higher order modes are excited

Growth rate – x plane



- Damping time $\tau_x = 1.88 \text{ ms}$
- rise time $> \tau_x$
- **stable**

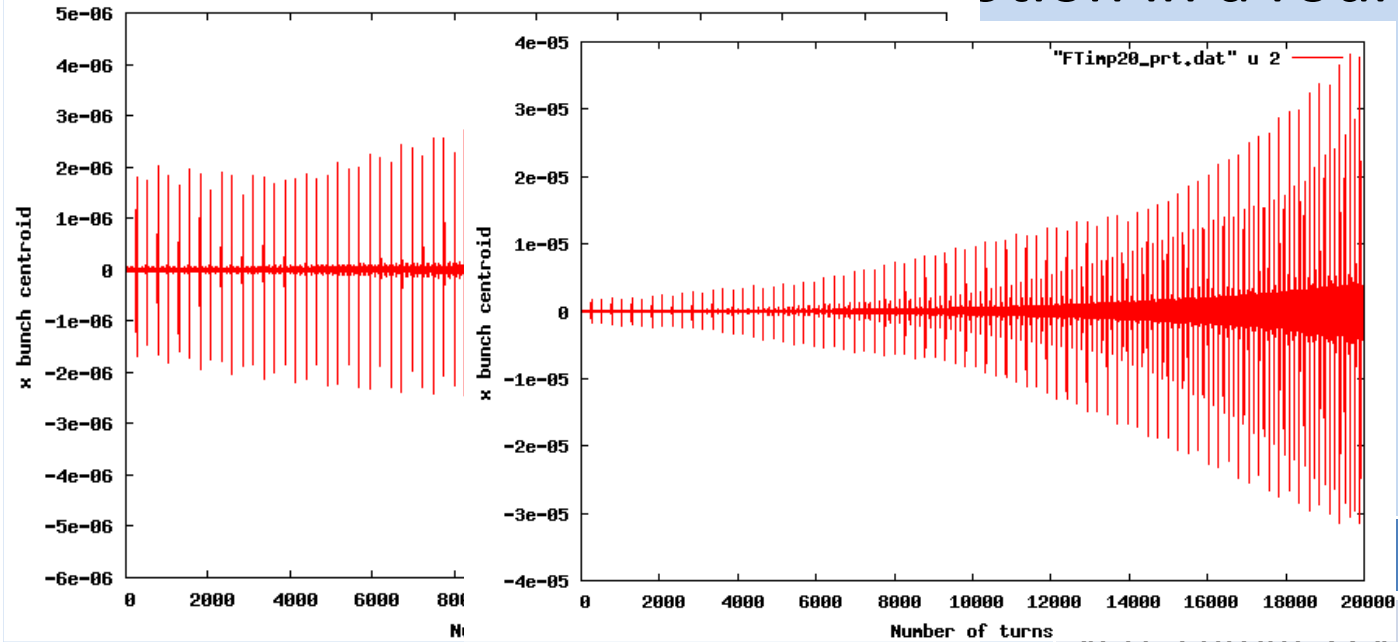
Growth rate – y plane



- Damping time $\tau_y = 1.91 \text{ ms}$
- $> \sim 13 \text{ MOhm/m}$
rise time $< \tau_y$
unstable

Threshold
 $\sim 13 \text{ MOhm/m}$

Horizontal and vertical motion in a round chamber

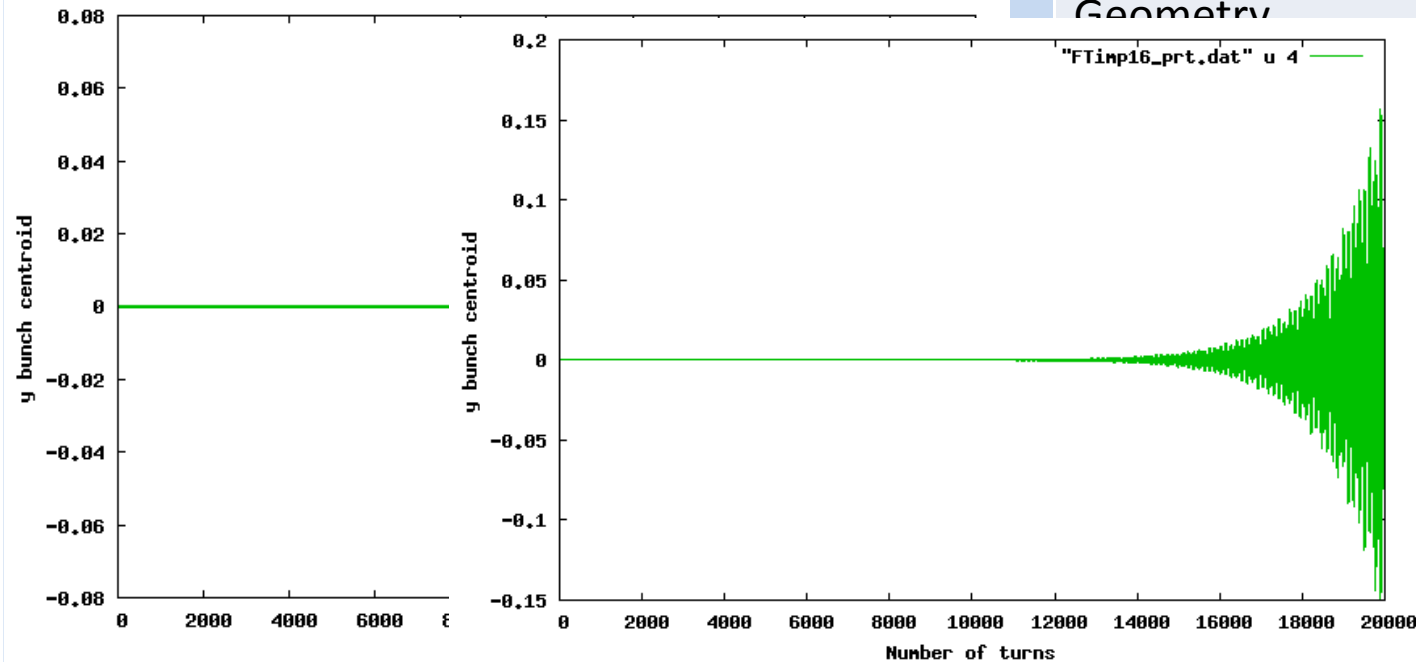


36.8

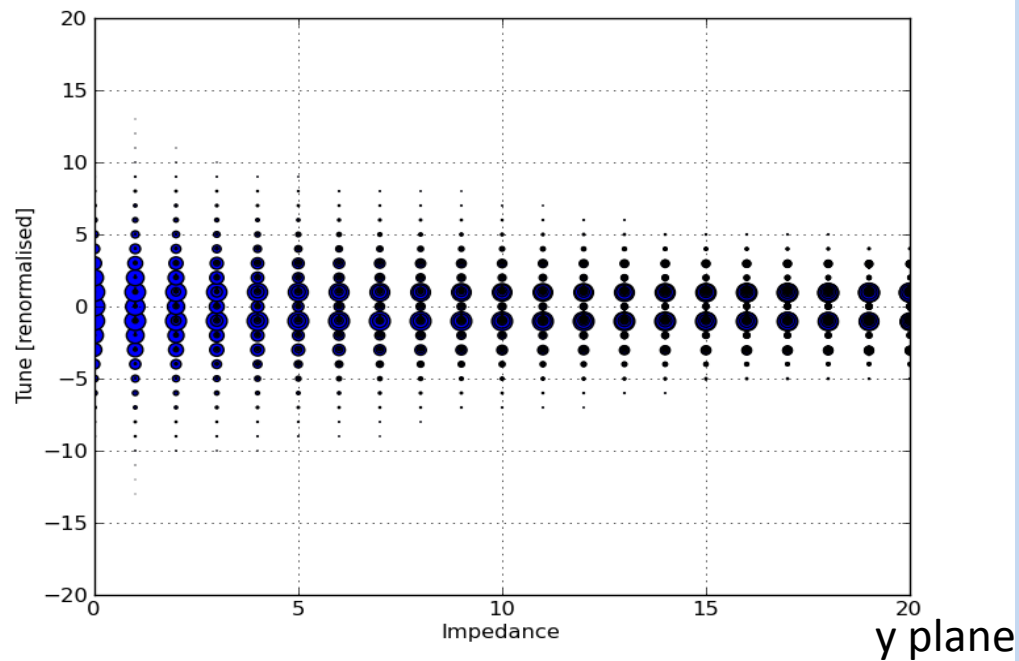
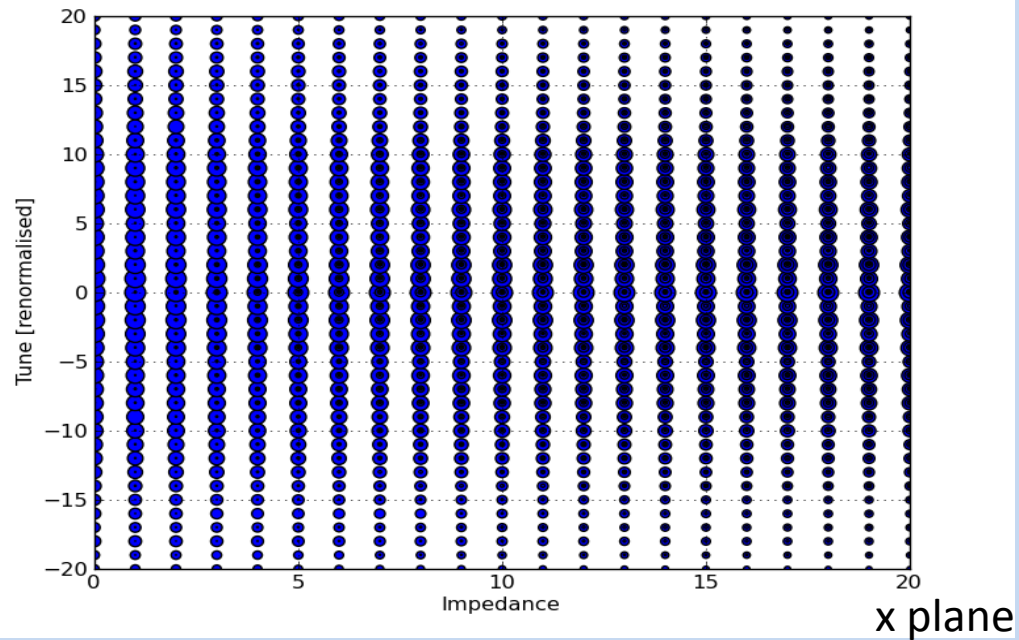
7.6

Geometry

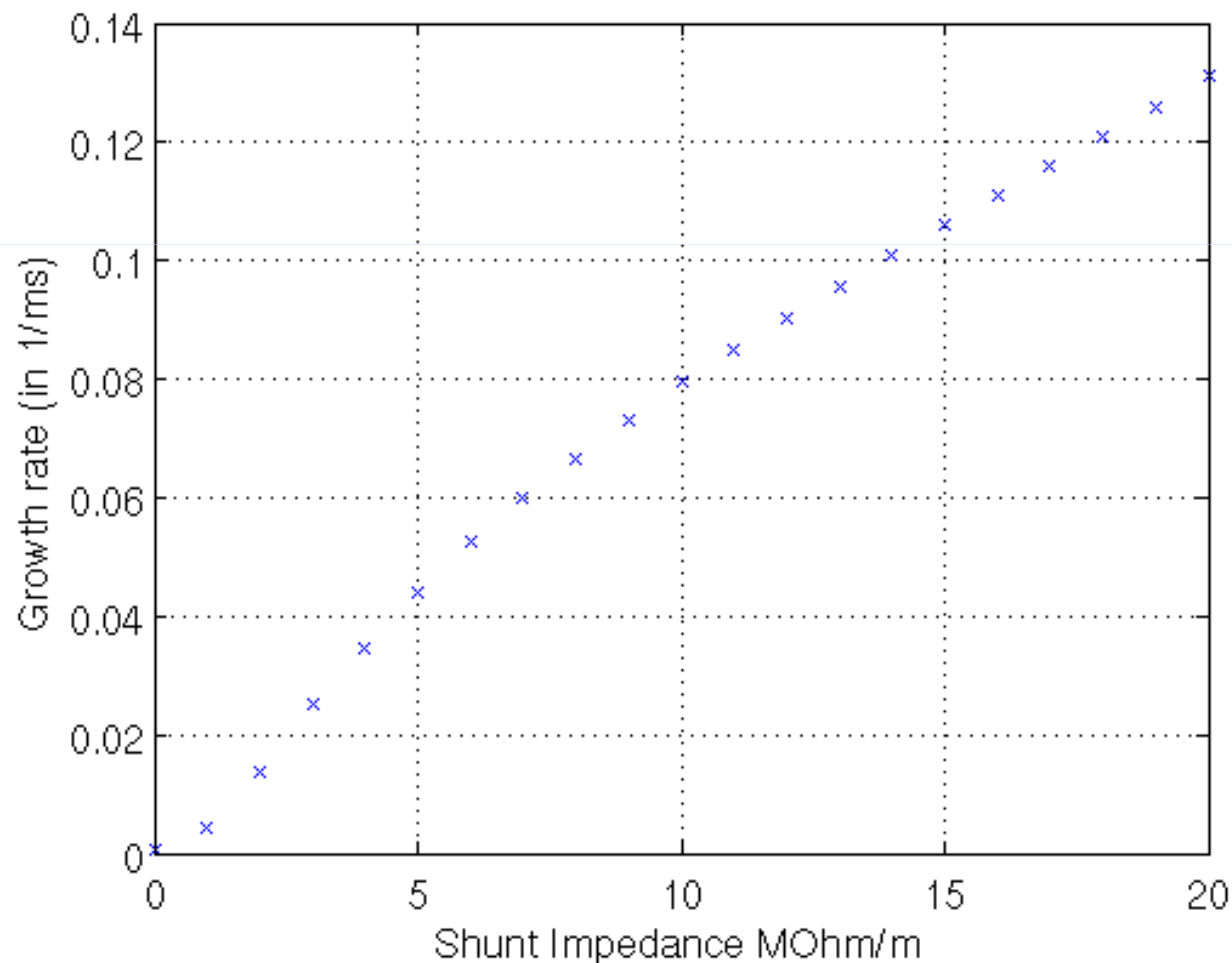
Round



Mode spectrum of the
horizontal and vertical motion
as a function of impedance

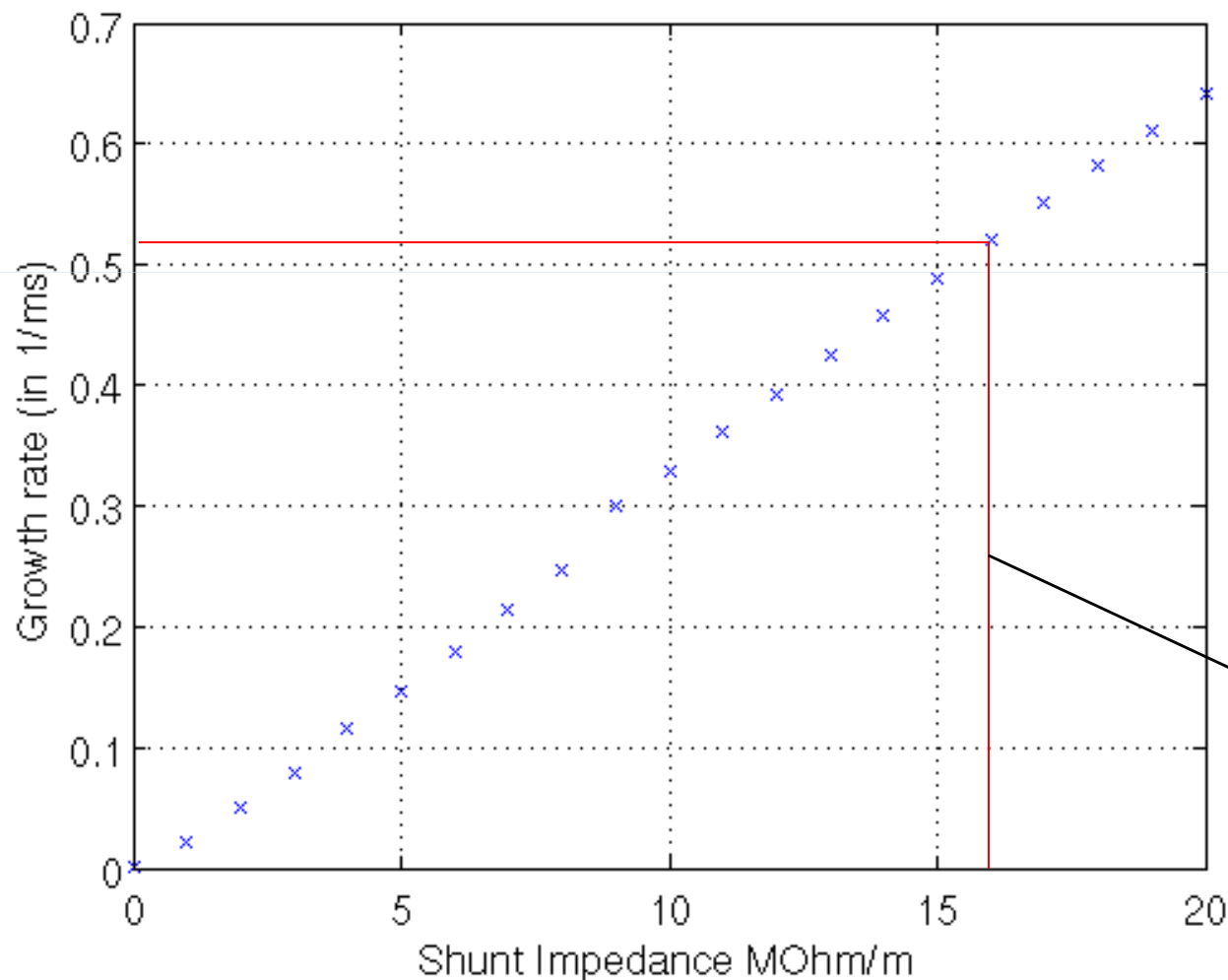


Growth rate – x plane



- Damping time $\tau_x = 1.88 \text{ ms}$
 - rise time $> \tau_x$
- stable**

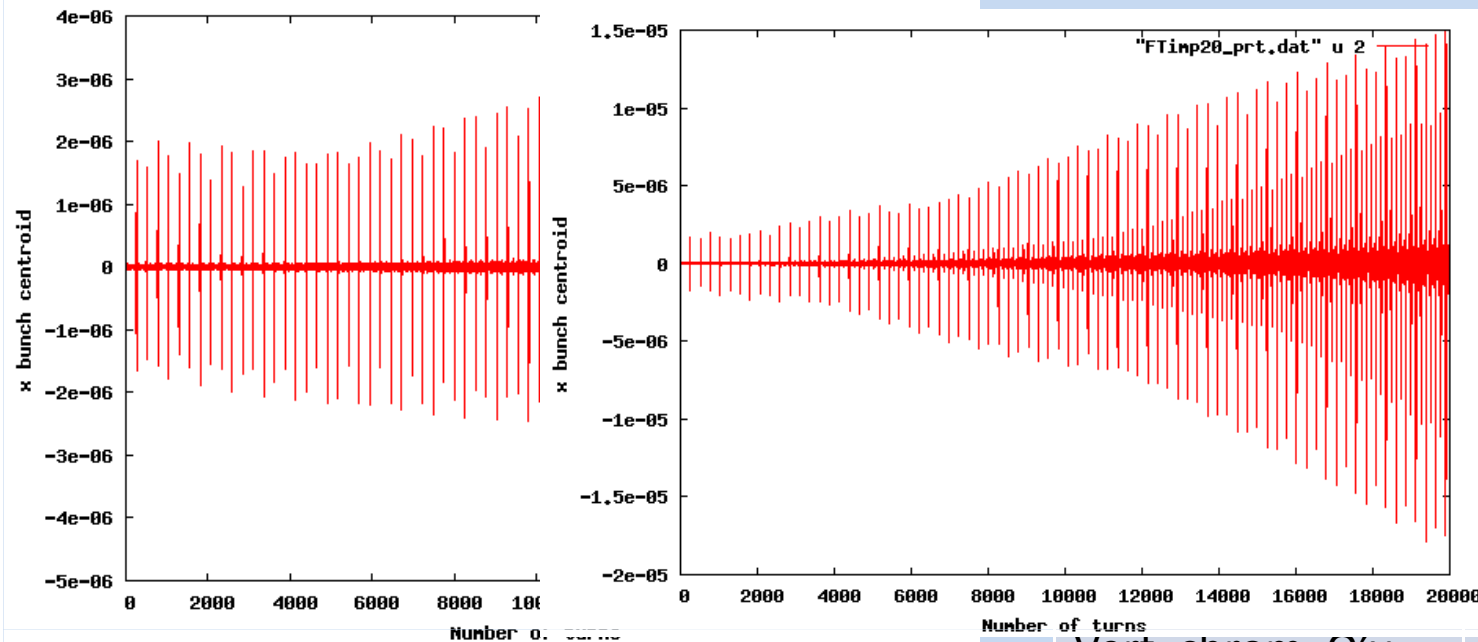
Growth rate – y plane



- Damping time $\tau_y = 1.91 \text{ ms}$
- $> \sim 16 \text{ MOhm/m}$
rise time $< \tau_y$
unstable

Threshold
 $\sim 16 \text{ MOhm/m}$

Horizontal and vertical motion in a round chamber

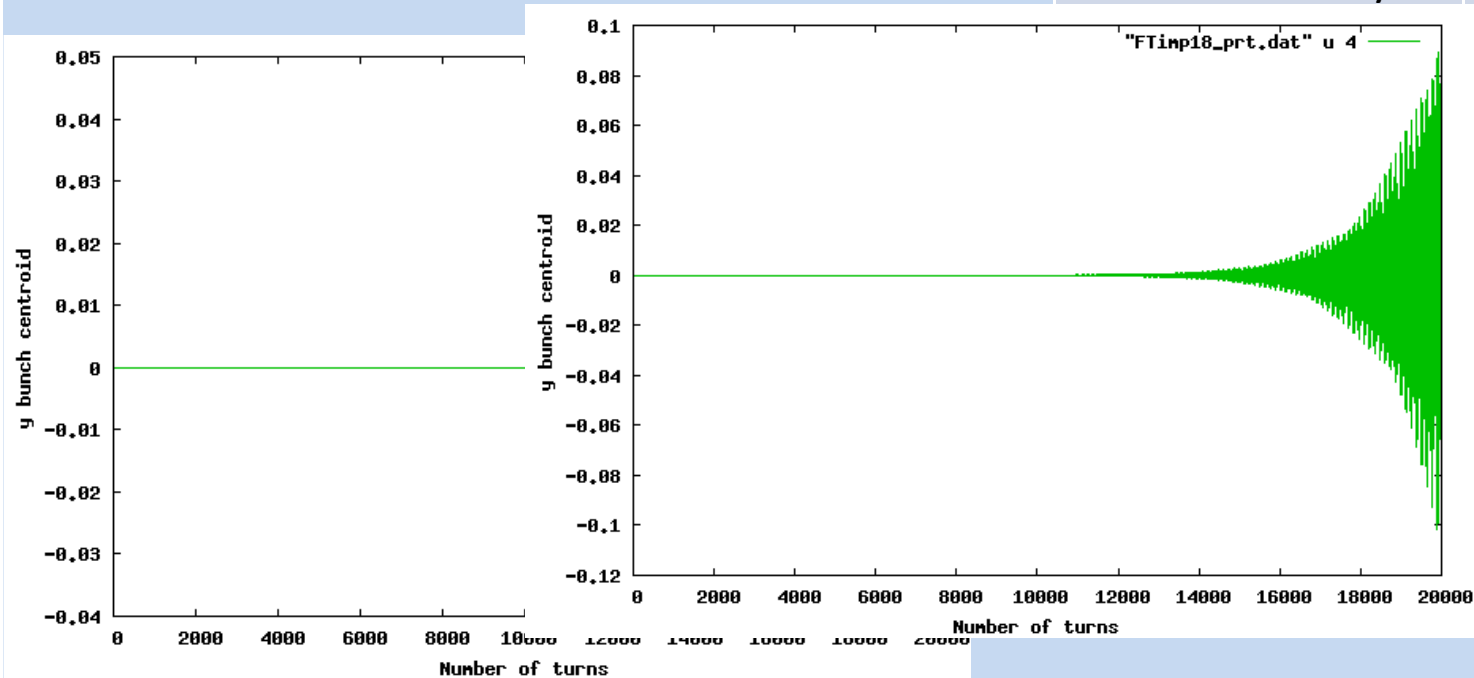


46

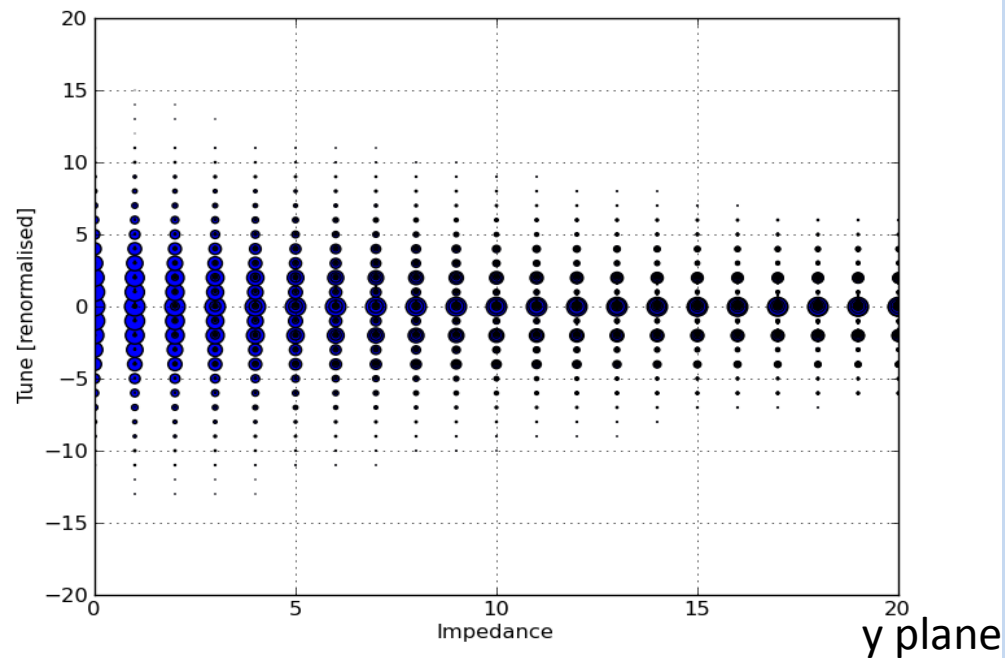
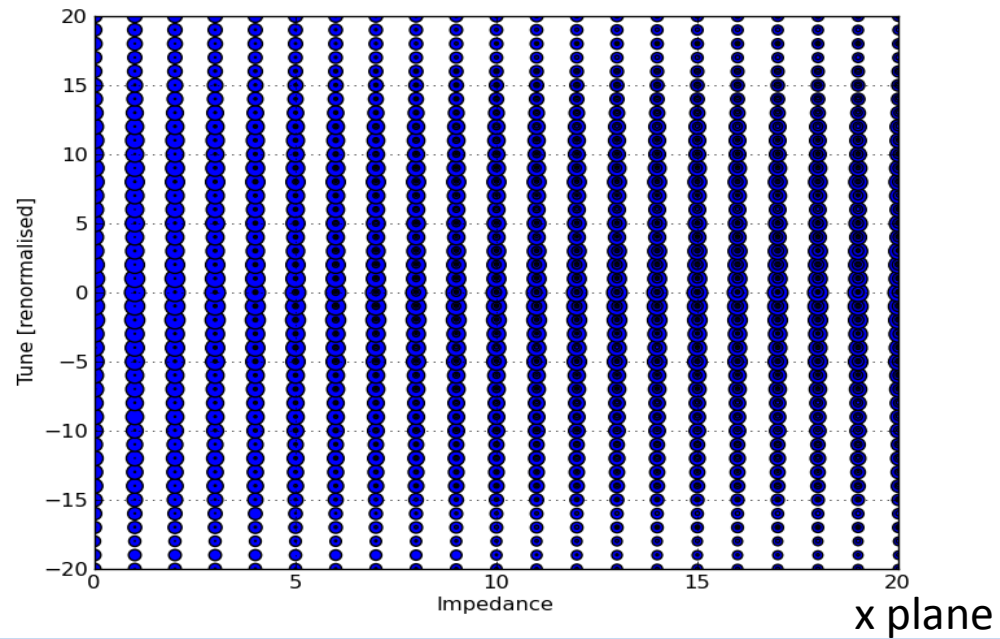
Vert. chrom. Q_y

9.5

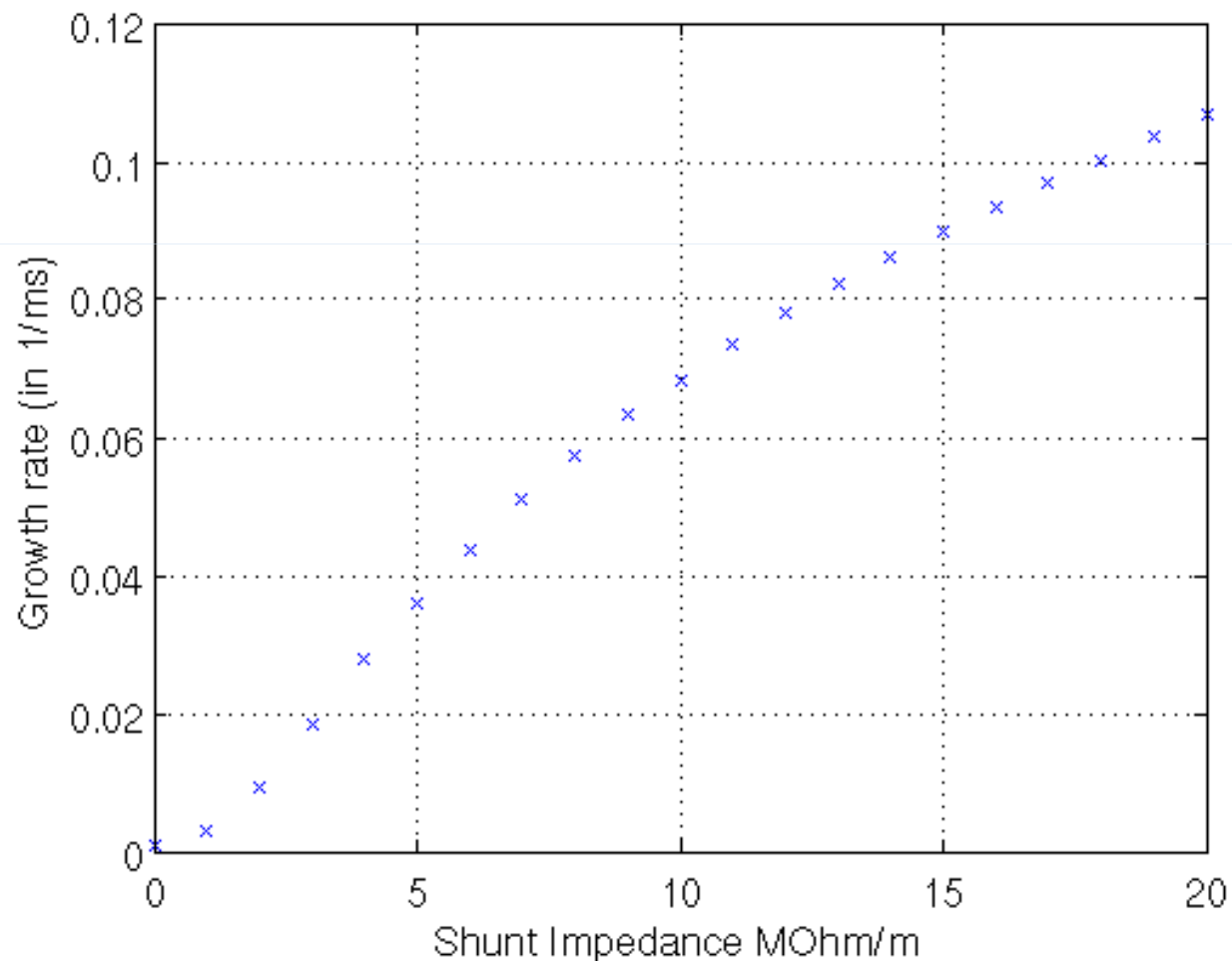
Round



Mode spectrum of the
horizontal and vertical motion
as a function of impedance

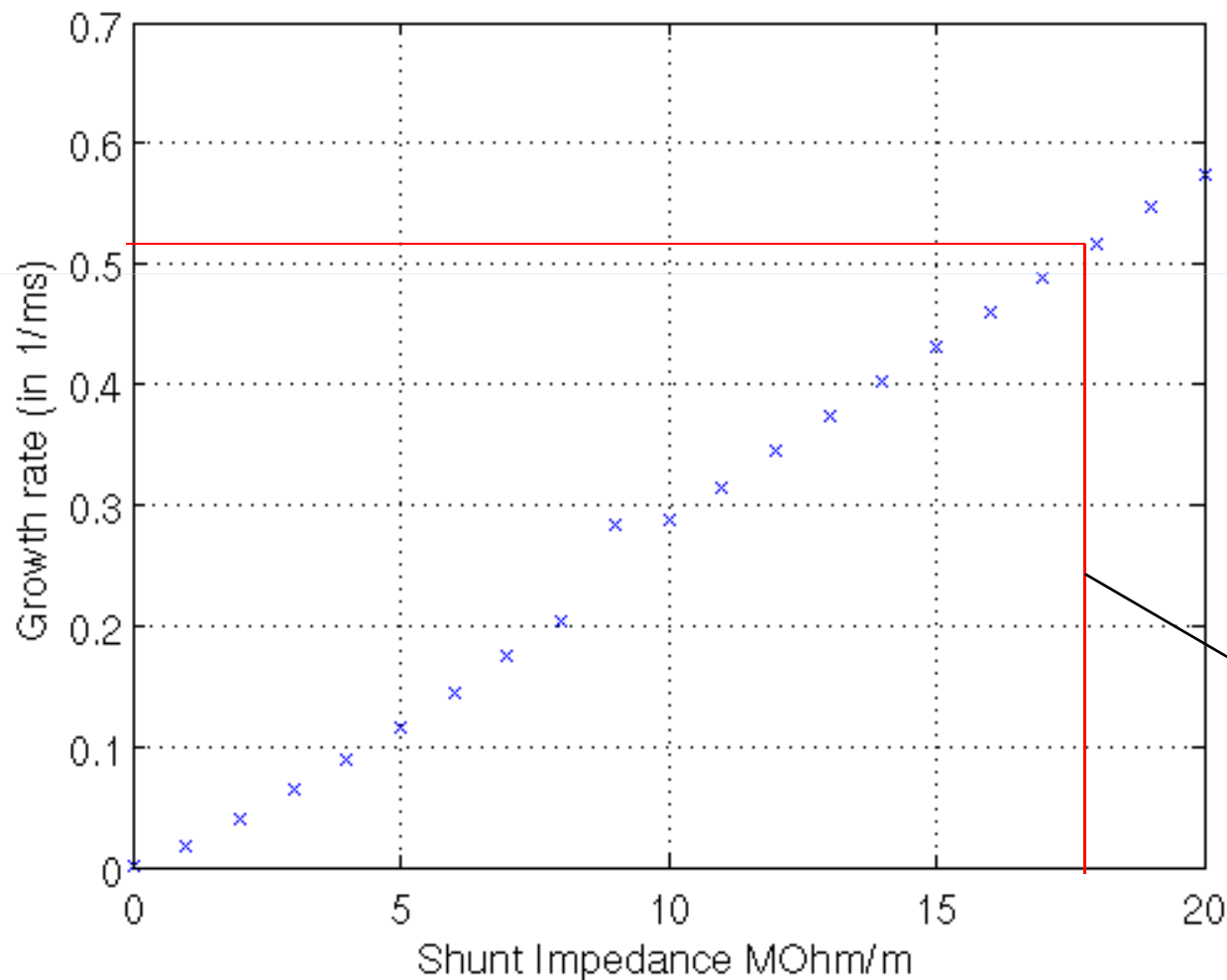


Growth rate – x plane



- Damping time $\tau_x = 1.88 \text{ ms}$
 - rise time $> \tau_x$
- stable**

Growth rate – y plane



- Damping time $\tau_y = 1.91 \text{ ms}$
- $> \sim 18 \text{ MOhm/m}$
rise time $< \tau_y$
unstable

Threshold
 $\sim 18 \text{ MOhm/m}$



Damping Rings



Results for round chamber

	x	y	x	y
Chromaticity Q'_x / Q'_y	Threshold M Ω m/m		Rise time (ms) $\tau_x=1.88, \tau_y=1.91$	
0/0	10	11	0.38	0.49
9.2/1.9	16	3	1.92	2
18.4/3.8	stable	6	stable	1.81
27.6/5.7	stable	13	stable	1.81
36.8/7.6	stable	16	stable	1.81
46/9.5	stable	18	stable	1.81

- For chromaticity 0, the TMCI threshold is at 10 and 11 M Ω m/m for x,y respectively
- For positive chromaticity, there is no TMCI but another instability occurs.
- As the chromaticity is increased, higher order modes get excited, less effect, move to higher instability thresholds



Damping Rings



Results for round chamber

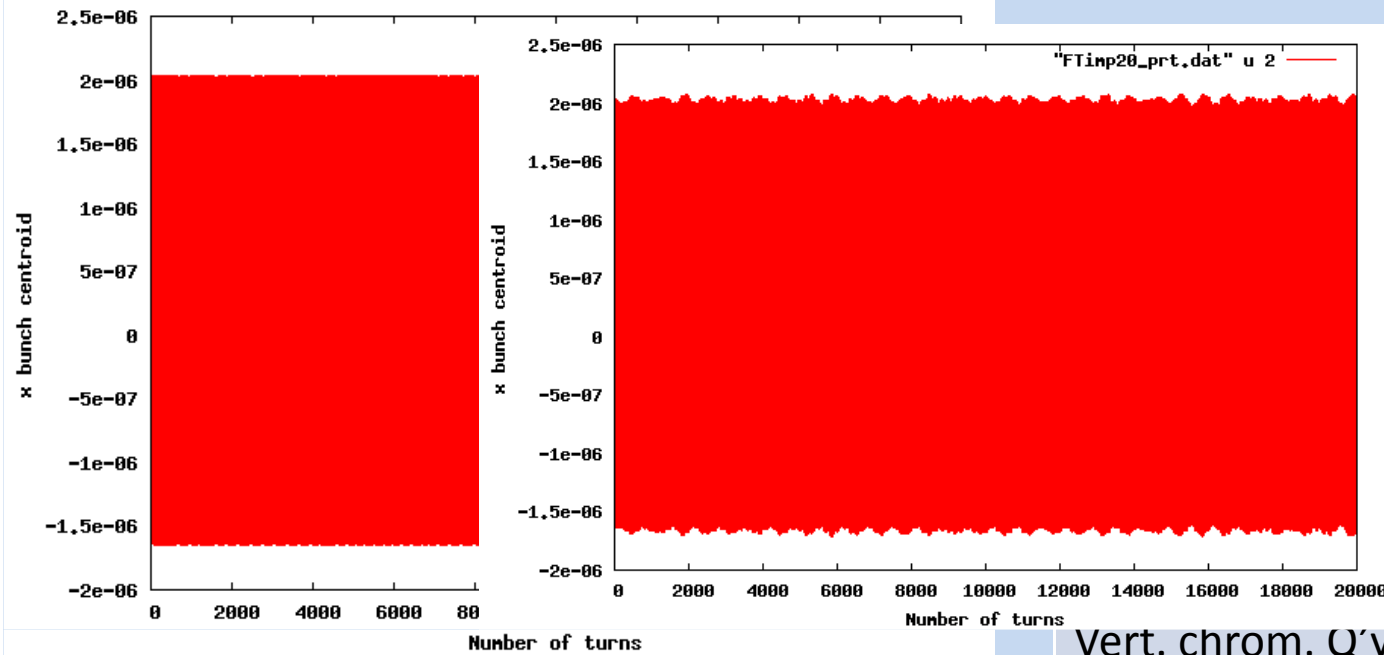
- Impedance cause bunch modes to move and merge, leading to a strong TMCI instability
- Chromaticity make the modes move less, therefore it helps to avoid the coupling (moved to a higher threshold)
- Still some modes can get unstable due to impedance

- As the chromaticity is increased, higher order modes are excited (less effect on the bunch)

Conclusion

- Either we correct the chromaticity and operate below the TMCI threshold or sufficient high positive chromaticity must be given so that mode -1, or -2 or maybe higher order mode is stable for the damping time

Horizontal and vertical motion in a flat chamber

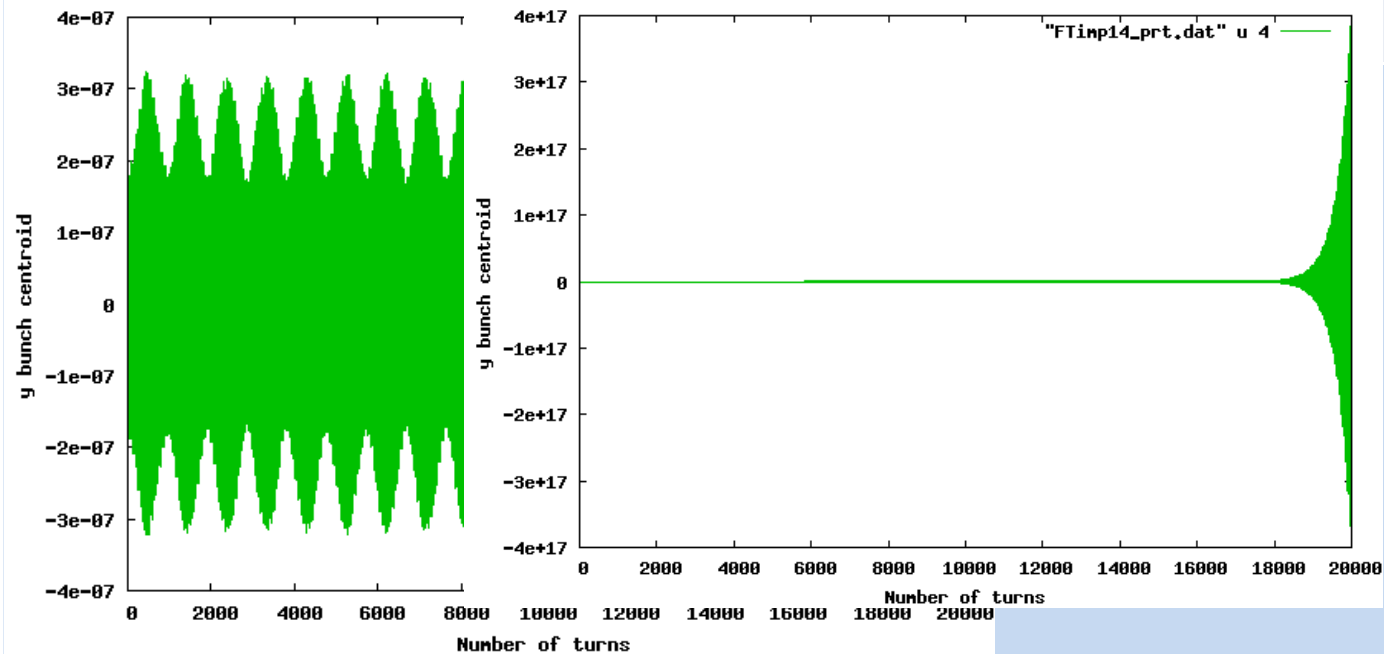


Vert. chrom. Q_y

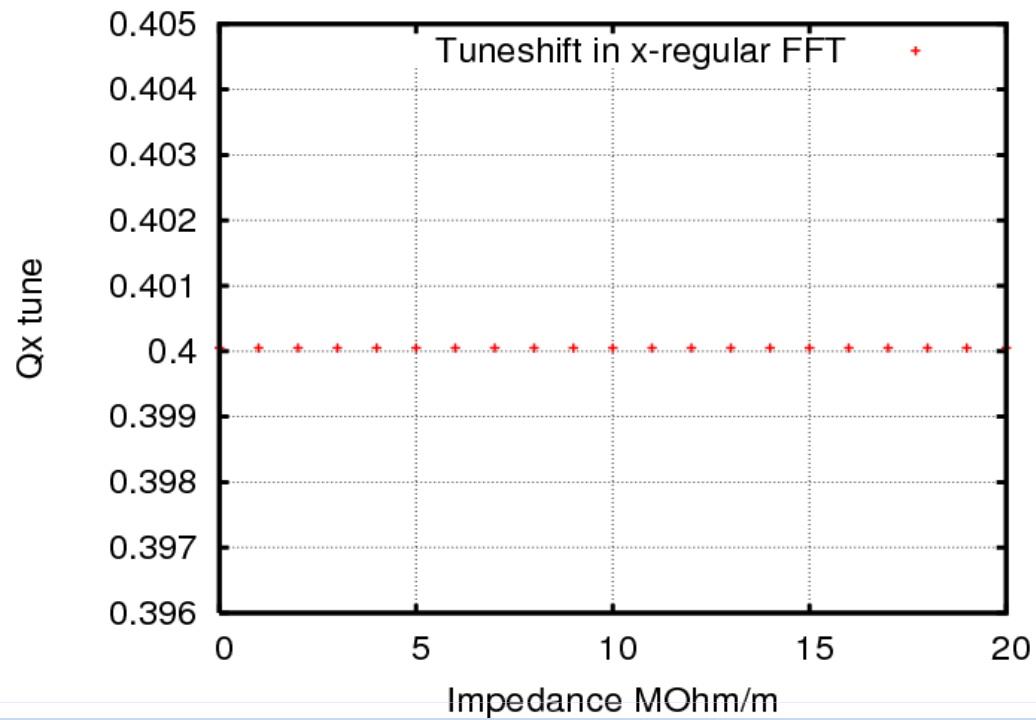
0

0

Flat

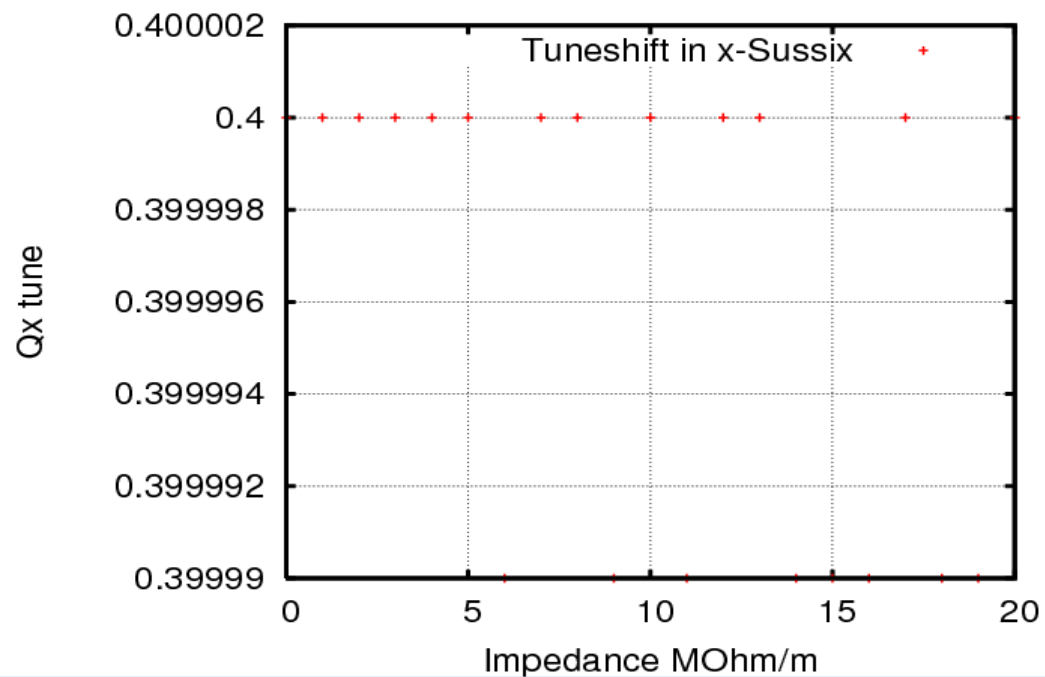


Horizontal plane



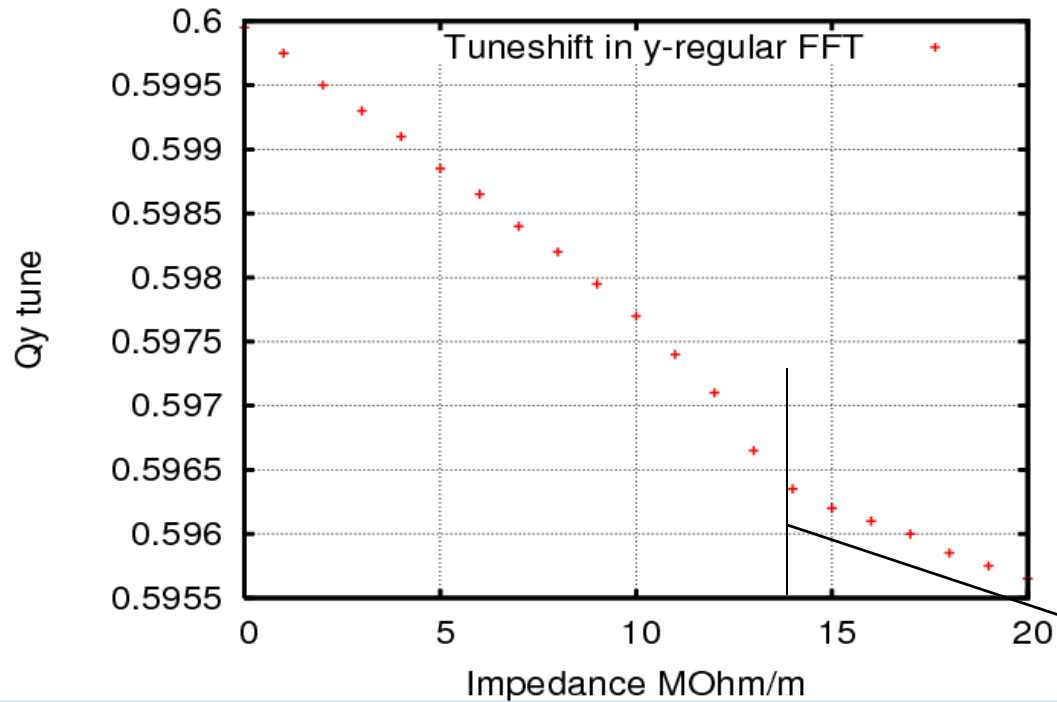
Regular FFT

Stable beam



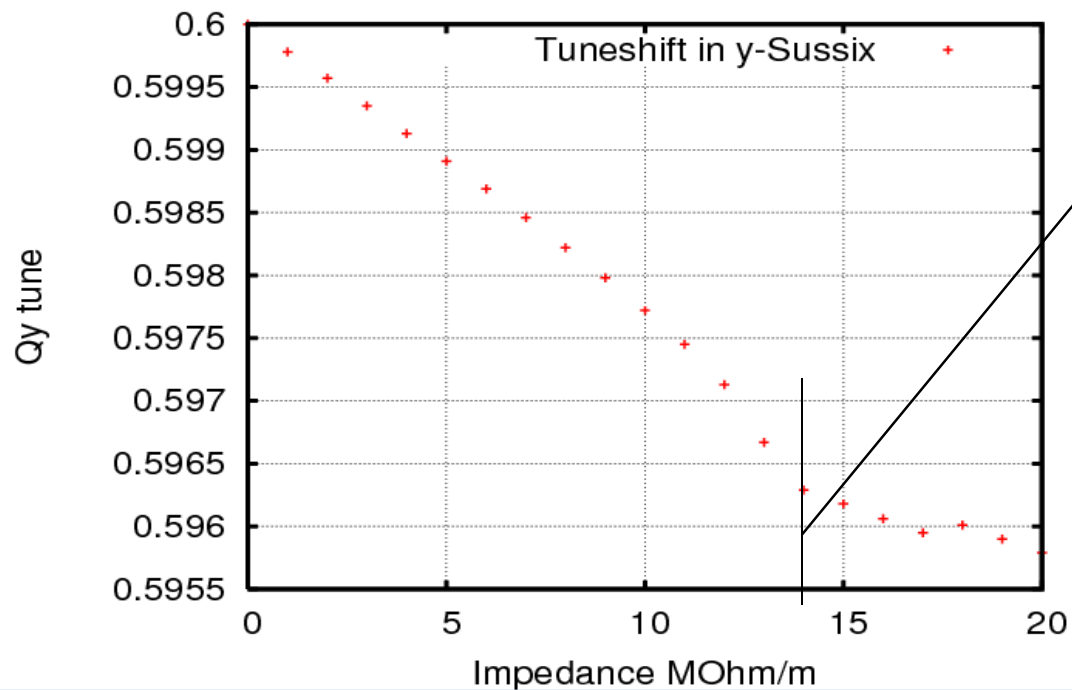
Sussix FFT

Vertical plane



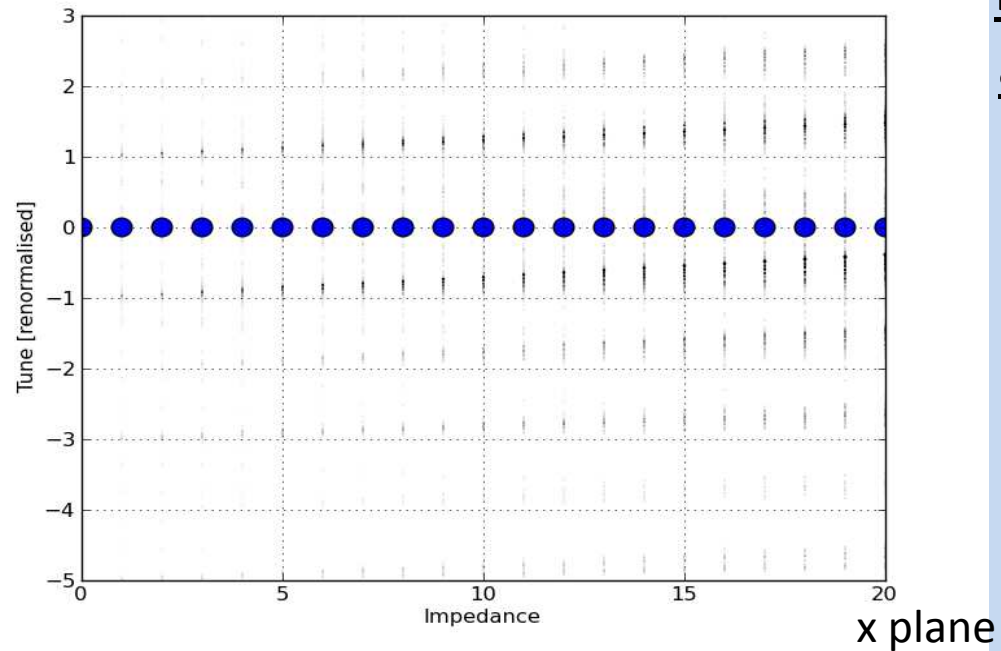
Regular FFT

Instability threshold (TMCI)
14MOhm/m

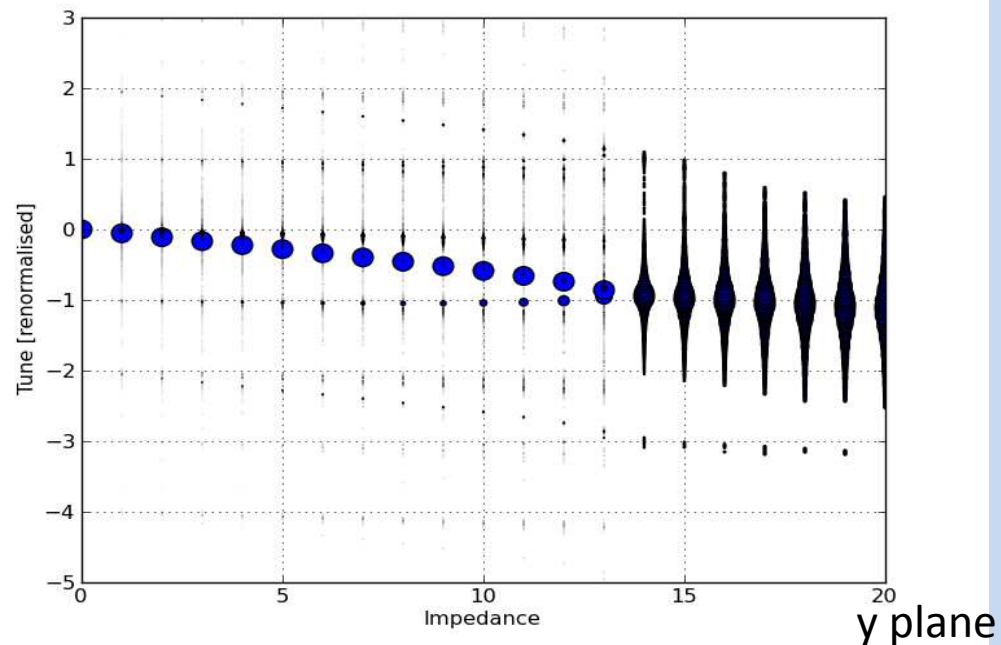


Sussix FFT

Mode spectrum of the horizontal and vertical motion as a function of impedance

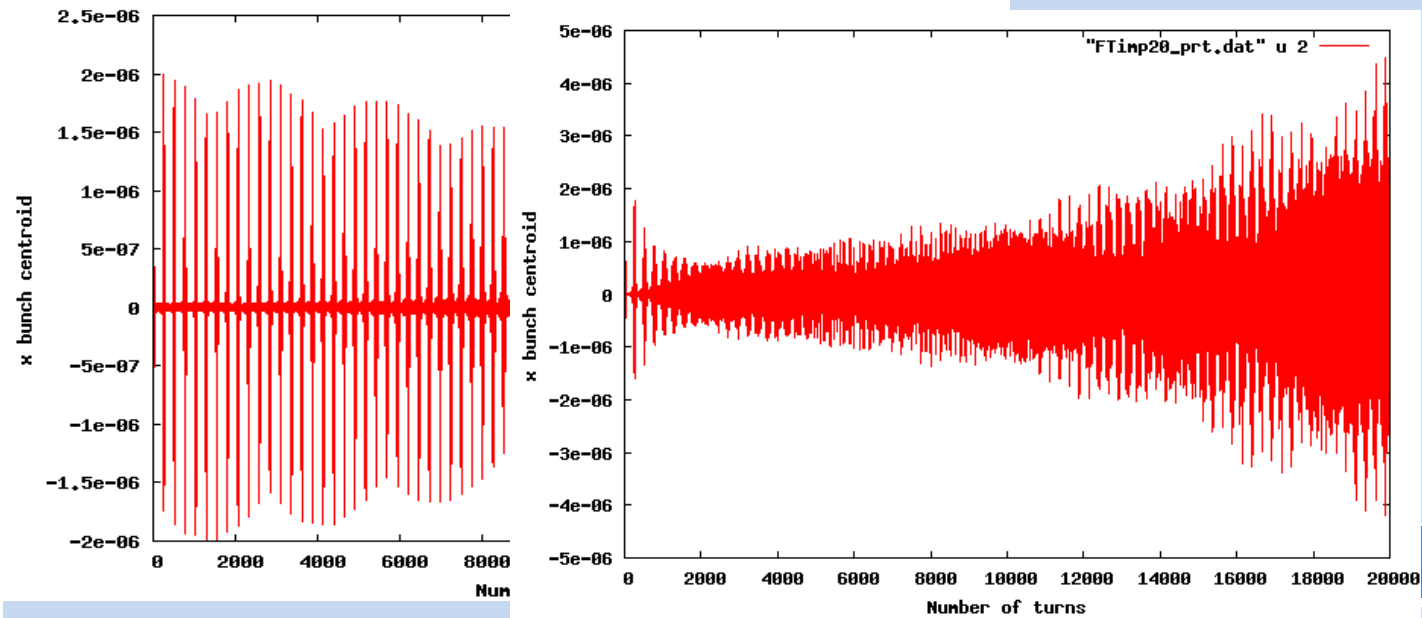


- mode 0 stable
- others shift
- later coupling between 0 and -1



- mode 0 and mode -1 couple
- TMCI instability

Horizontal and vertical motion in a flat chamber

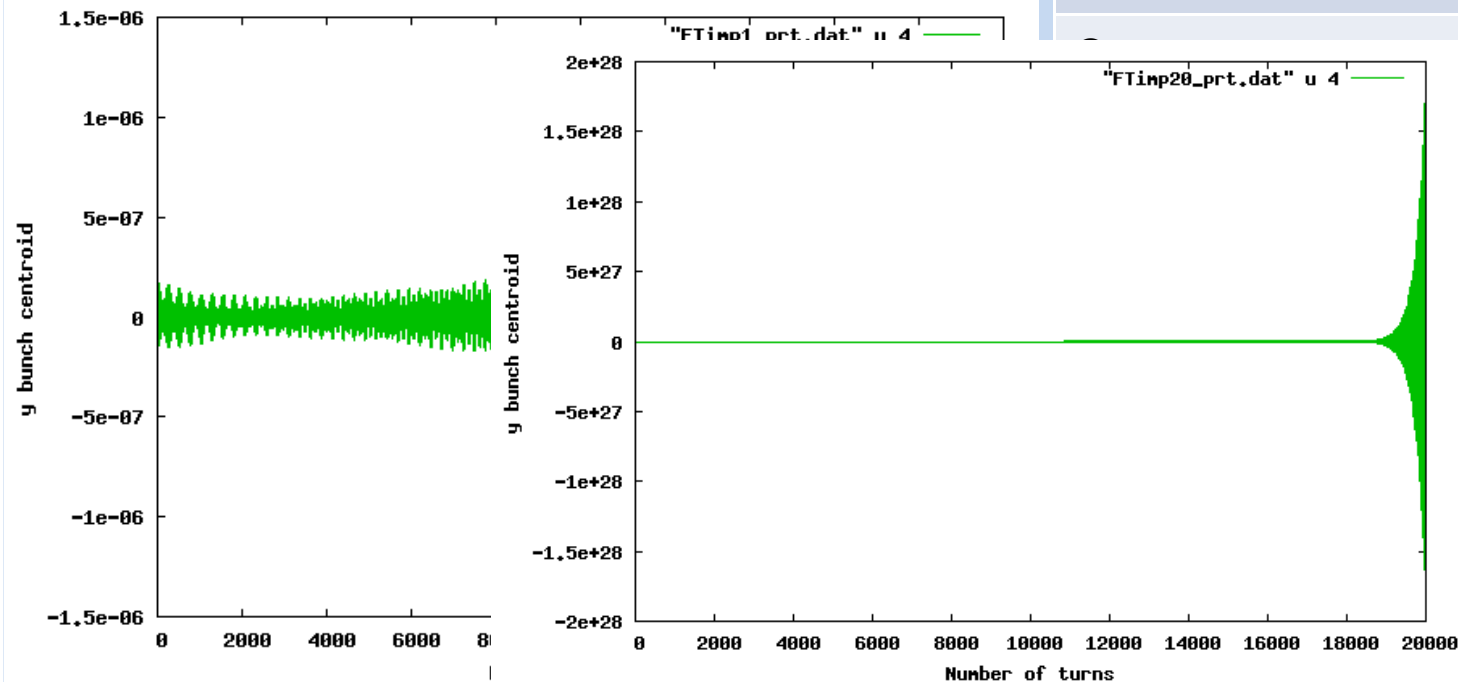


9.2

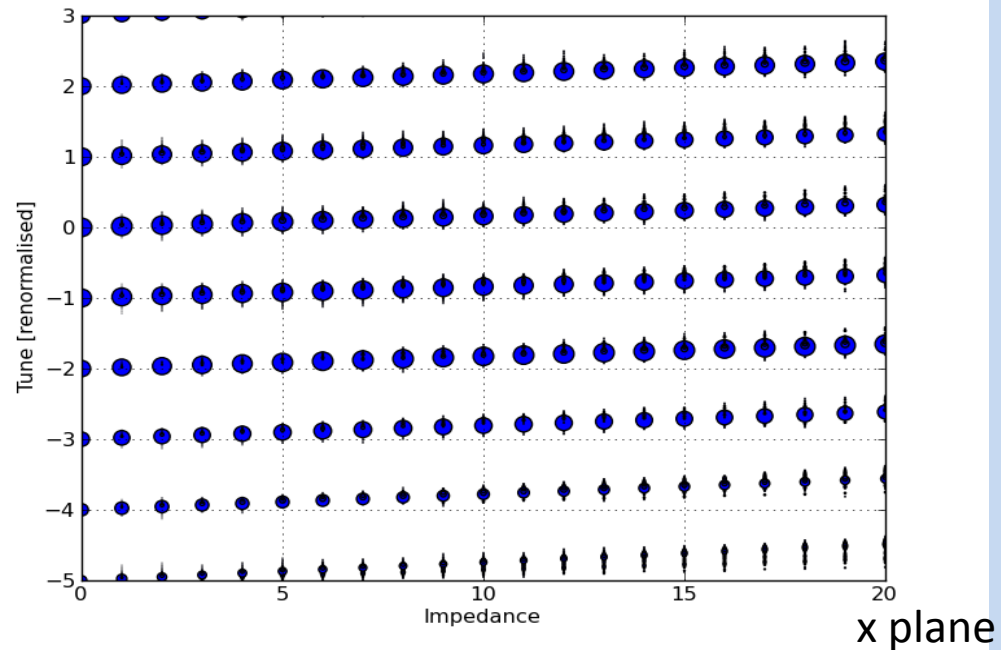
Vert. chrom. Q'_y

1.9

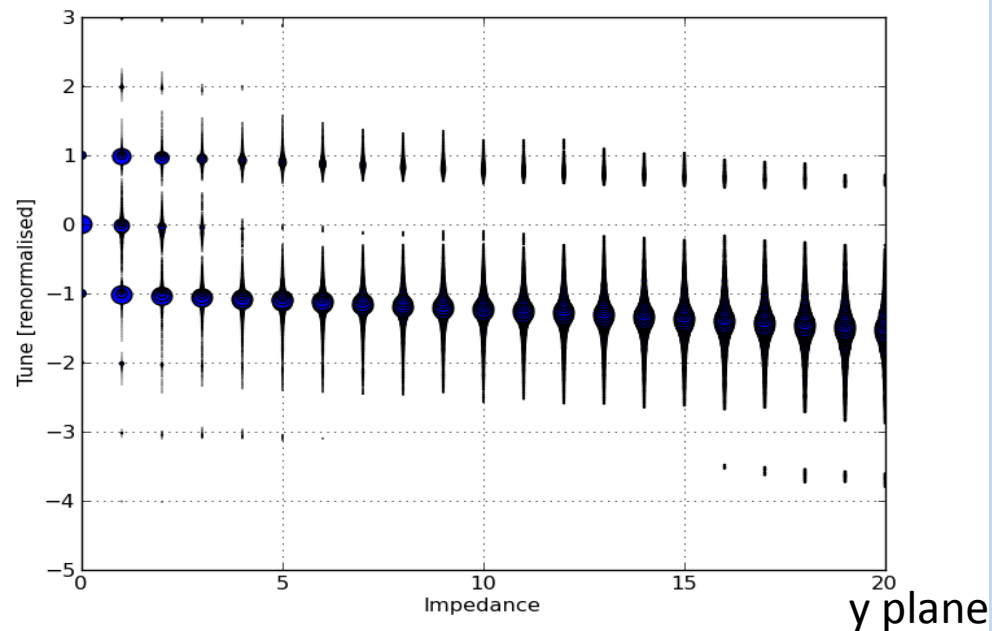
Flat



Mode spectrum of the horizontal and vertical motion as a function of impedance



- no mode coupling
- no TMCI instability
- hard to tell the cause of instability



- mode 1 is damped
- mode -1 gets unstable



Damping Rings



Conclusions for the flat chamber

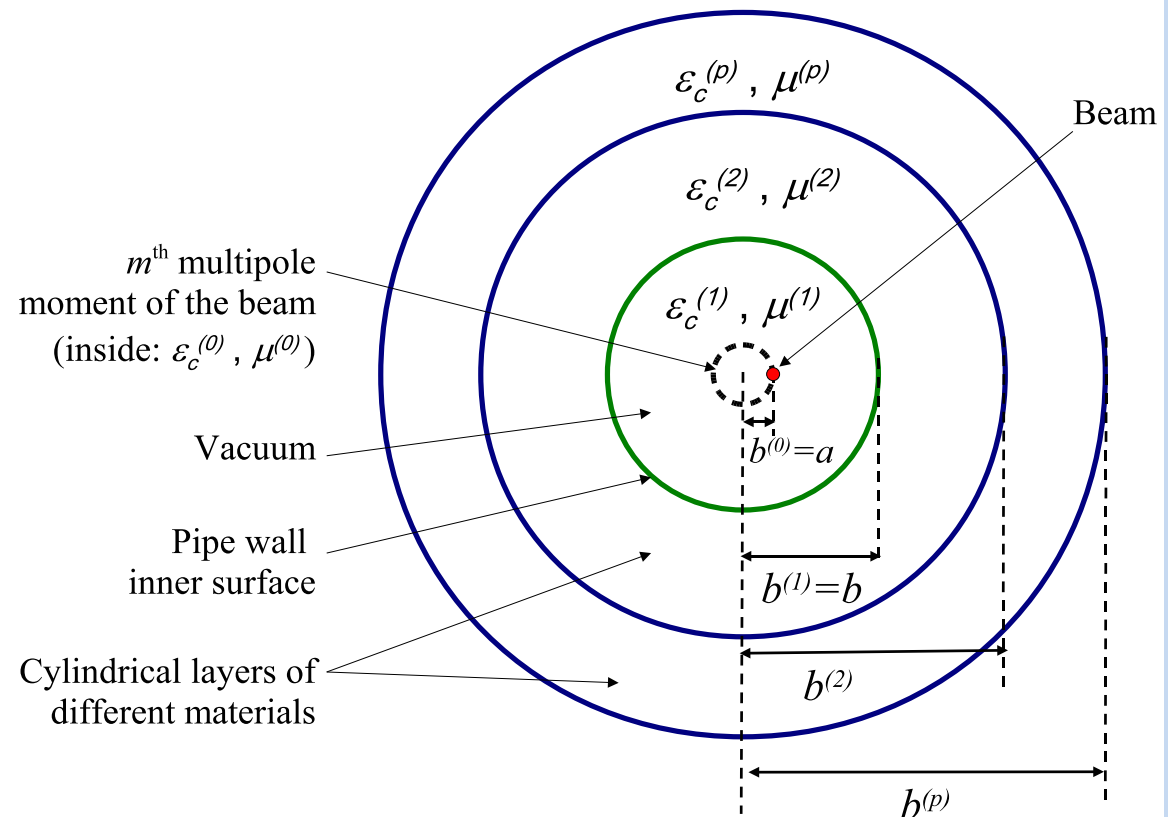
	<i>threshold (MΩm/m)</i>			
	<u>Round</u>		<u>Flat</u>	
Chromaticity	x	y	x	y
0	10	11	stable	14

- Calculate the growth rate for the cases with chromaticity and compare with the round chamber

Resistive wall in the CLIC-DR regime

- Layers of **coating materials** can significantly increase the resistive wall impedance at **high frequency**
 - Coating especially needed in the low gap wigglers
 - Low conductivity, thin layer coatings (NEG, a-C)
 - Rough surfaces (not taken into account so far)

Pipe cross- section:



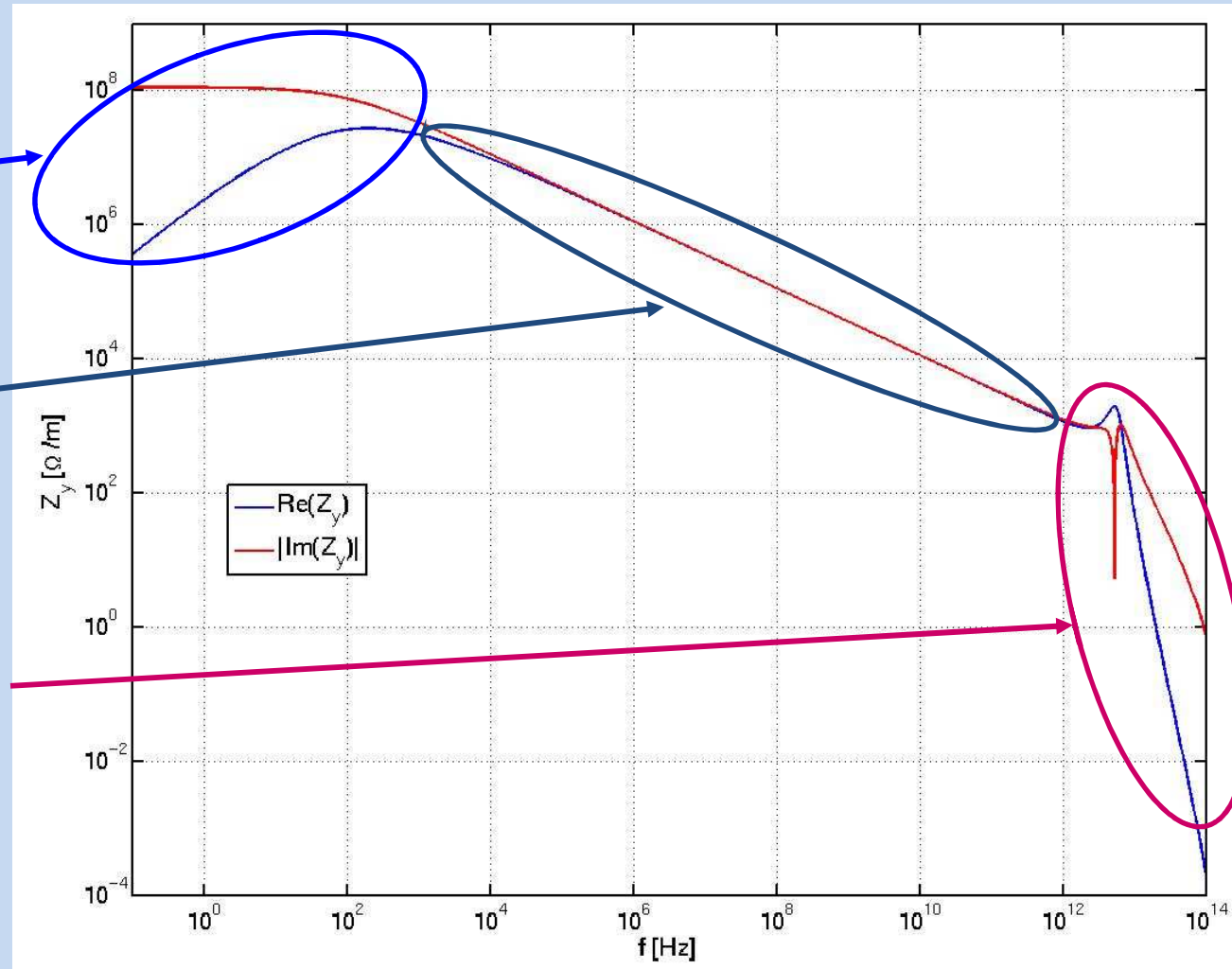
General Resistive Wall Impedance: Different Regimes

- Vertical impedance in the wigglers (3 TeV option, pipe made of copper without coating)

Low frequency or
“inductive-bypass”
regime

“Classic thick-wall”
regime

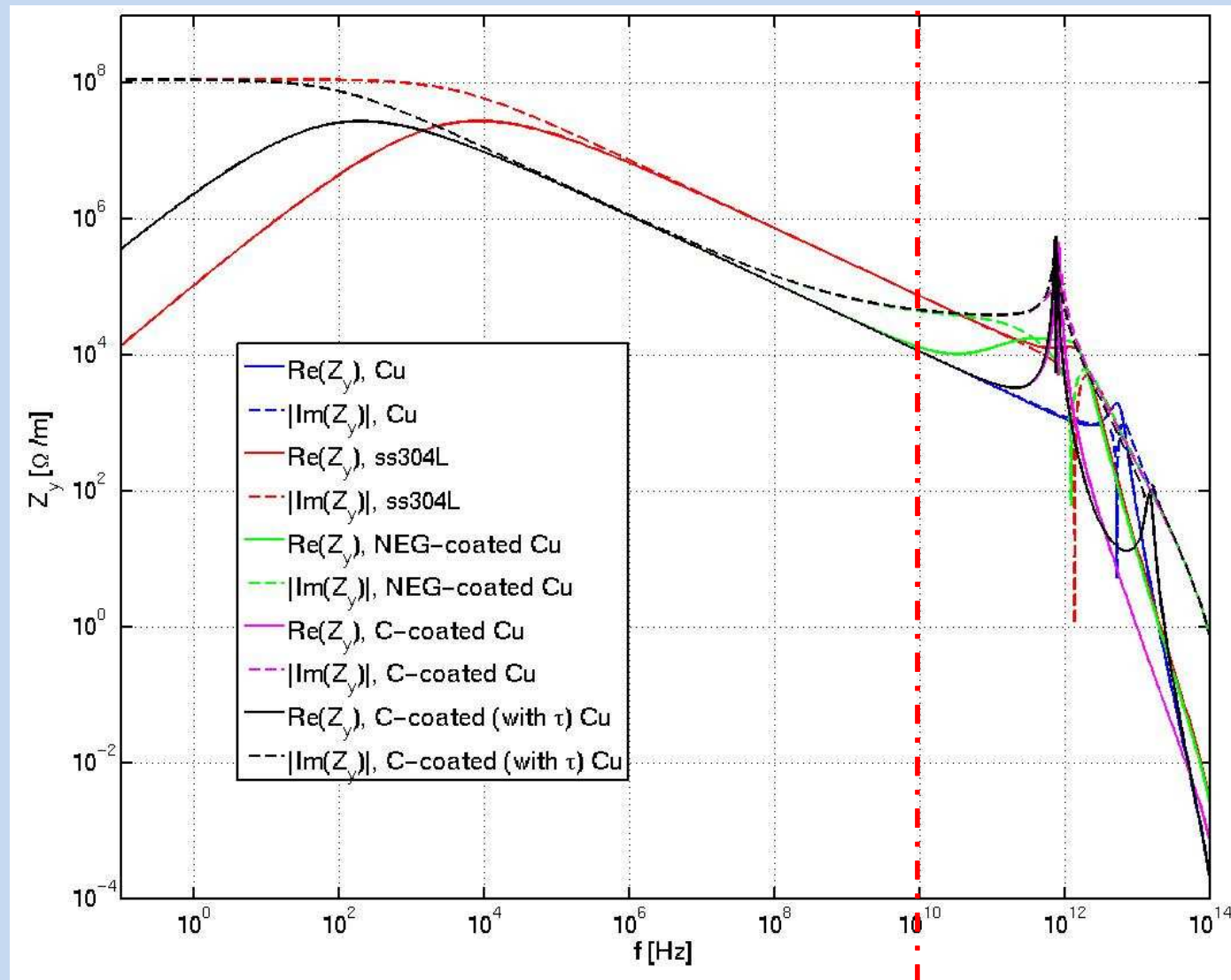
High frequency regime



Note: all the impedances and wakes presented have been multiplied by the beta functions of the elements over the mean beta, and the Yokoya factors for the wigglers

Resistive Wall Impedance: Various options for the pipe

- Vertical impedance in the wigglers (3 TeV option) for different materials



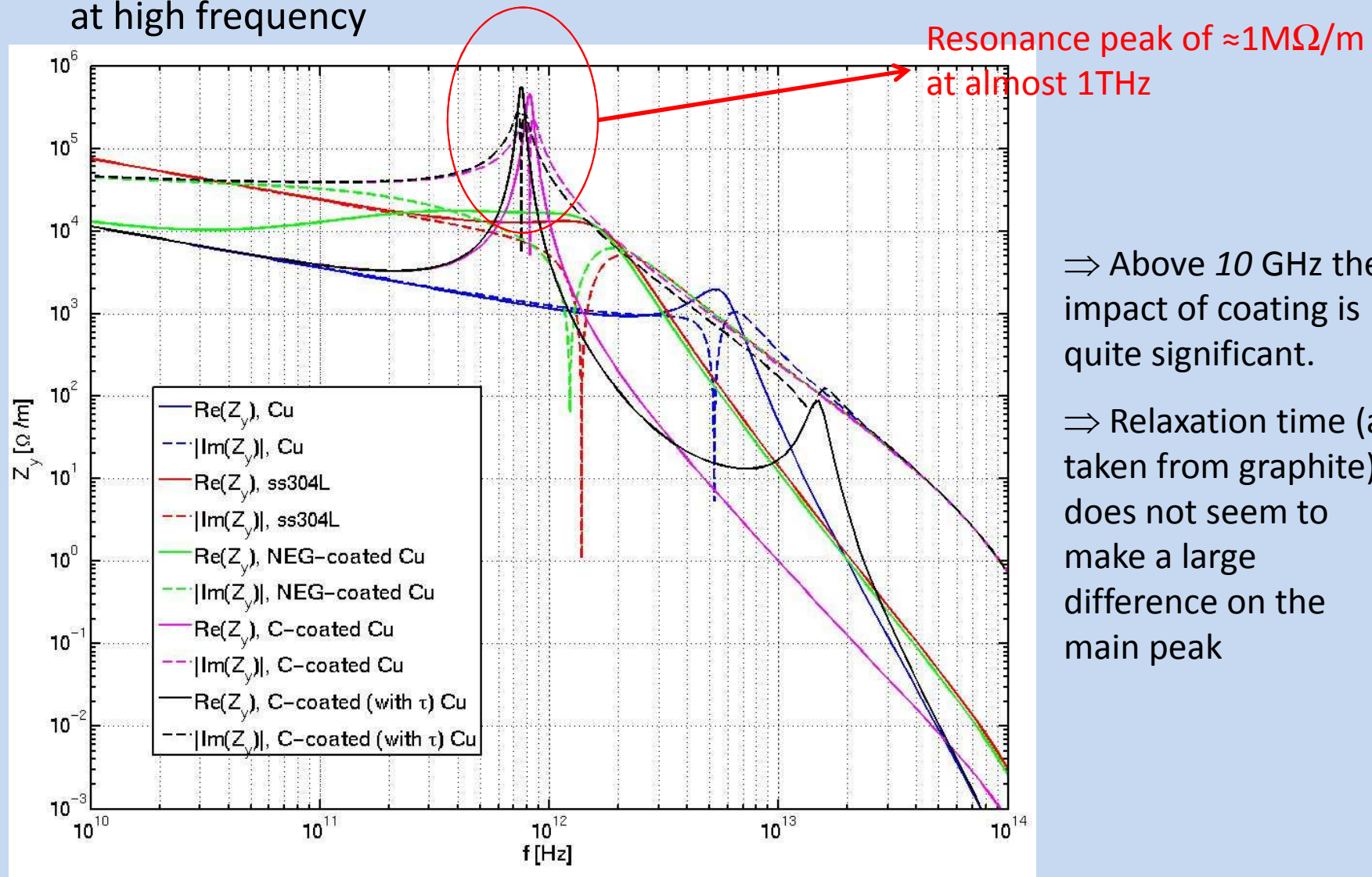
⇒ Coating is “transparent” up to ~ 10 GHz

⇒ But at higher frequencies some narrow peaks appear!!

⇒ So we zoom for frequencies above 10 GHz →

Resistive Wall Impedance: Various options for the pipe

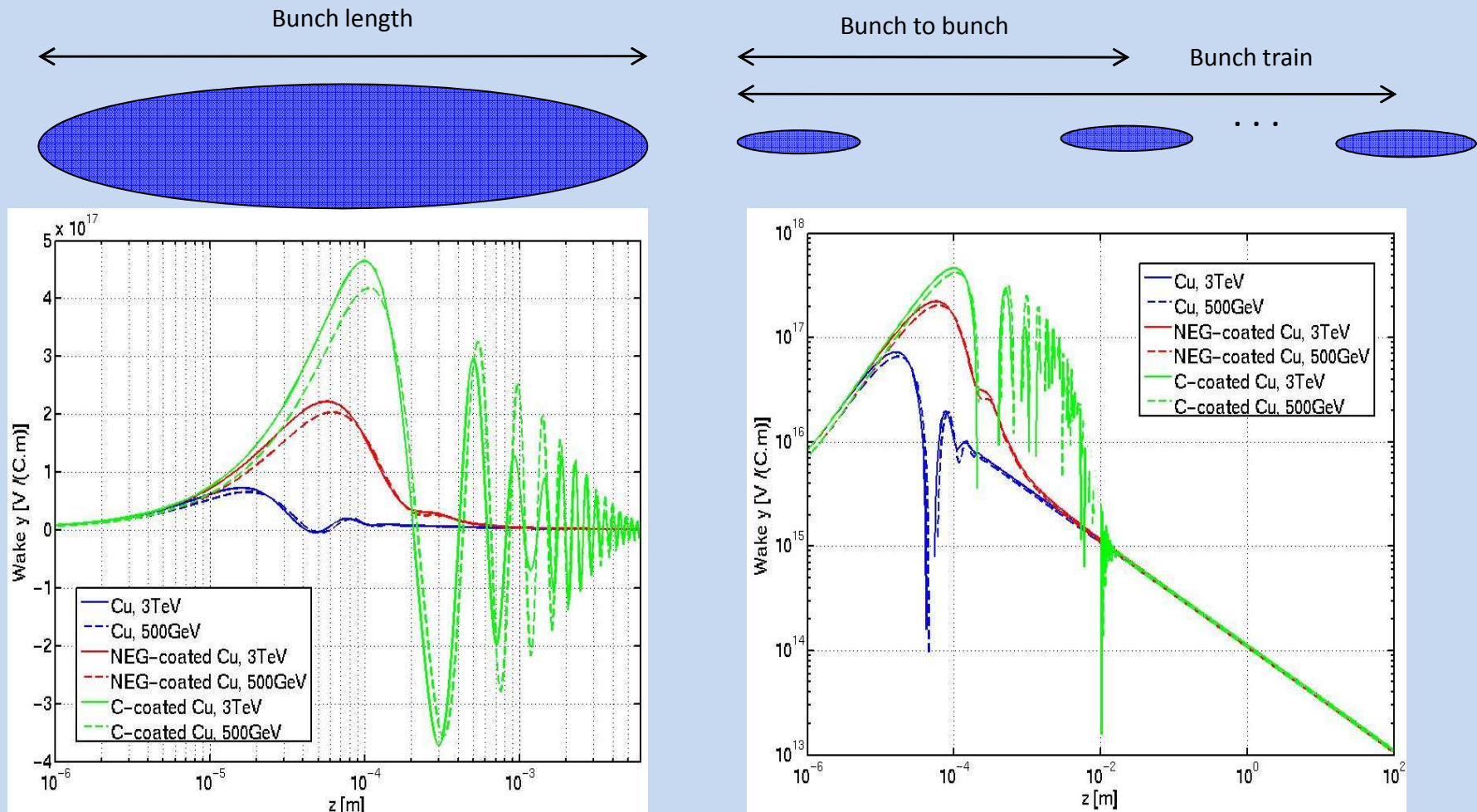
- Vertical impedance in the wigglers (3 TeV option) for different materials: zoom at high frequency



\Rightarrow Above 10 GHz the impact of coating is quite significant.

\Rightarrow Relaxation time (as taken from graphite) does not seem to make a large difference on the main peak

In terms of wake field, we find



- The presence of coatings strongly enhances the wake field on the scale of a bunch length (and even bunch-to-bunch)
- The single bunch instability threshold should be evaluated, as well as the impact on the coupled bunch instability
- This will lower the transverse impedance budget for the DRs



Damping Rings



Next steps...

- Use the average beta functions for the DRs ($\langle\beta_x\rangle=4.568$, $\langle\beta_y\rangle=7.568$)
- Growth rate for the flat chamber
- Theoretical calculation of the tune shift and the growth rate of the headtail modes
- Check for negative chromaticity (mode 0 would be the only one getting unstable but also it's the easiest one to be corrected-damped with a use of a feedback system/ check on the growth time of the instability)
- Include the effect of octupoles in the simulation (detuning with amplitude), but check for the emittance growth
- Include damping in the HeadTail code
- Include the wake field from the resistive wall with coating and other impedance sources in the HeadTail simulation

Damping Rings

Parameter	Symbol	Value
Energy	p_0 (GeV)	2.86
Norm. transv. emitt.	$\epsilon_{xn,yn}$ (nm)	480, 4.5
Bunch length	σ_z (mm)	1.6
Momentum spread	σ_δ	1.3×10^{-3}
Bunch spacing	ΔT_b (ns)	1
Bunch population	N_b	4.1×10^9
Circumference	C (m)	120.56
Coupling	(%)	0.1
Mom. compact.	α	7.6×10^{-5}
Number of bunches per train	n_b	156
Number of trains	n_t	2
Distance between trains	τ_t (ns)	545
Tunes	$Q_{x,y,s}$	55.4, 11.6, 0.00387
Store time/train	T_{st} (ms)	20
Energy loss	ΔE (MeV/turn)	4.2
Damping times	$\tau_{x,y,z}$ (ms)	1.88, 1.91, 0.96
RF frequency	f_{rf} (GHz)	1
RF voltage	V_{rf} (MV)	4.9
Harmonic number	h	1402
Dipole length	L_{dip} (m)	0.43
Dipole chamber rad.	R_{dip} (cm)	1
Number of dipoles	N_{dip} (m)	102
Wiggler length	L_w (m)	2
Wiggler field	B_w (T)	2.5
Number of wigglers	N_w (m)	52
Wiggler gap	r_w (mm)	13
Wiggler width	h_w (mm)	65
Average β_x in wigglers	$\langle \beta_{xw} \rangle$ (m)	4.787
Average β_y in wigglers	$\langle \beta_{yw} \rangle$ (m)	4.185



Damping Rings



Theoretical tune shift

$$\Omega^{(l)} - \omega_\beta - l\omega_s \approx -\frac{1}{4\pi} \frac{\Gamma(l + \frac{1}{2})}{2^l l!} \frac{Nr_0 c^2}{\gamma T_0 \omega_\beta \sigma} i(Z_1^\perp)_{\text{eff}}.$$

$$(Z_1^\perp)_{\text{eff}} = \frac{\sum_{p=-\infty}^{\infty} Z_1^\perp(\omega') h_l(\omega' - \omega_\xi)}{\sum_{p=-\infty}^{\infty} h_l(\omega' - \omega_\xi)}.$$