

NSUSY Fits.

Papers discussed:

Lee, Sanz, Trott arXiv:1204.0802

Espinosa, Grojean, Sanz, Trott arXiv:1207.7355

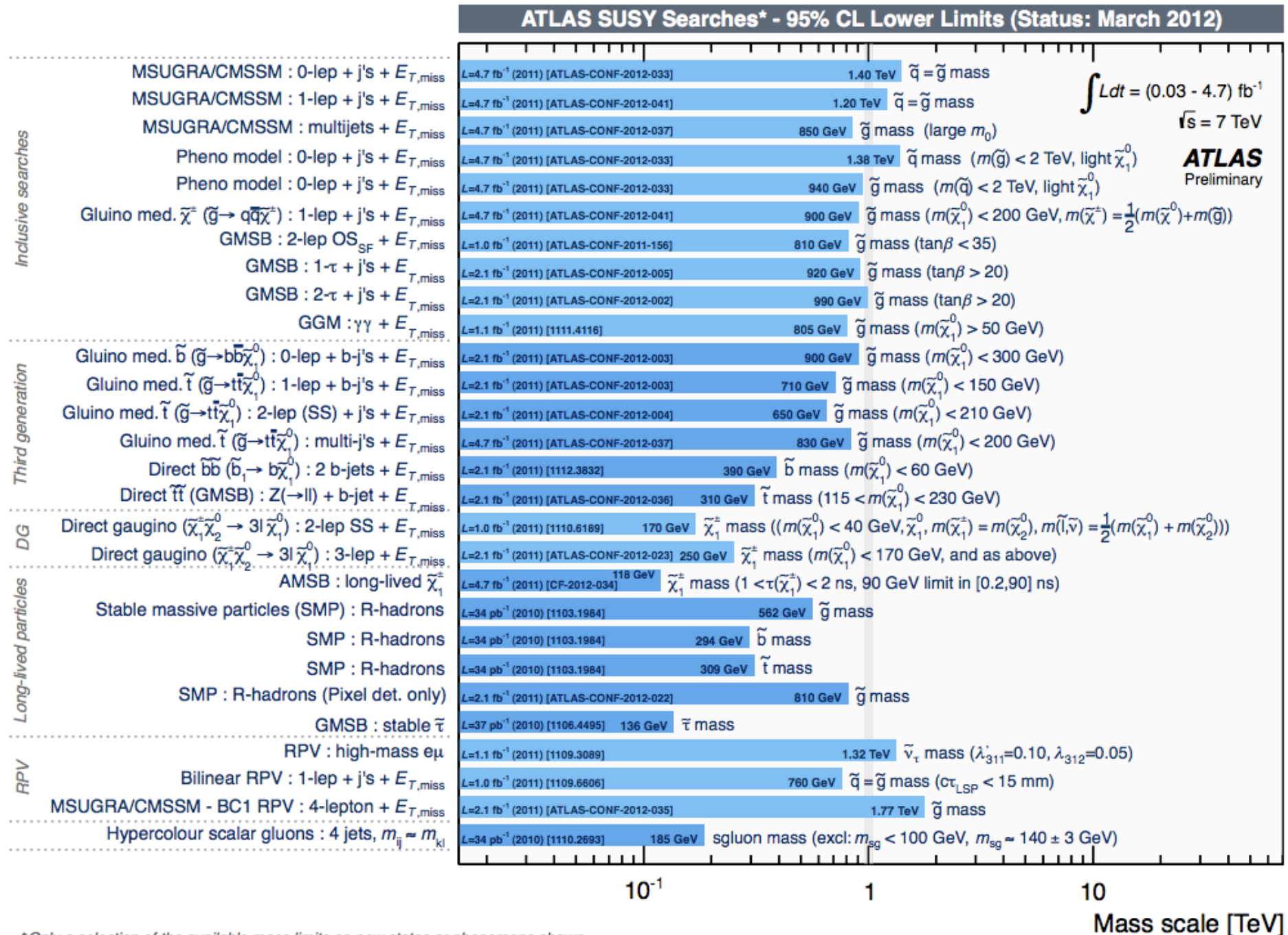
Content (for ref):

fitting sbottom

Higgs data, fits, other constraints

Natural SUSY

In light of recent LHC searches for the lowest lying SUSY fruit being fruitless to date:



Summary of bounds on New physics from ATLAS.

Natural SUSY

Theorists are re-examining what is minimally desired out of a SUSY spectrum.

From an effective field theory point of view what do you want?

One wants SUSY to solve the hierarchy problem by ensuring that the threshold matching corrections of the Higgs mass operator

$$\mu^2 \phi^\dagger \phi$$

effectively cancel up to weak scale matching differences: $(\delta m_B^2 - \delta m_F^2) \phi^\dagger \phi$

So that large NP mass scales can exist and not intrinsically lead to fine tuning in the Higgs mass operator. (See Meade's talk yesterday for exposition)

In solving the Hierarchy problem this way you do not want to re-introduce (significant) fine tuning in the solution.

Natural SUSY

What is the minimal required spectrum for SUSY motivated out of naturalness?



In honour of texas, the SUSY Alamo.

Natural SUSY

Particle theory community by and large



*LHC data
to date*

Light gluino with R parity conservation

Natural SUSY

What is the minimal required spectrum for SUSY motivated out of naturalness?

What do we need to hold out? Lots of options actually.



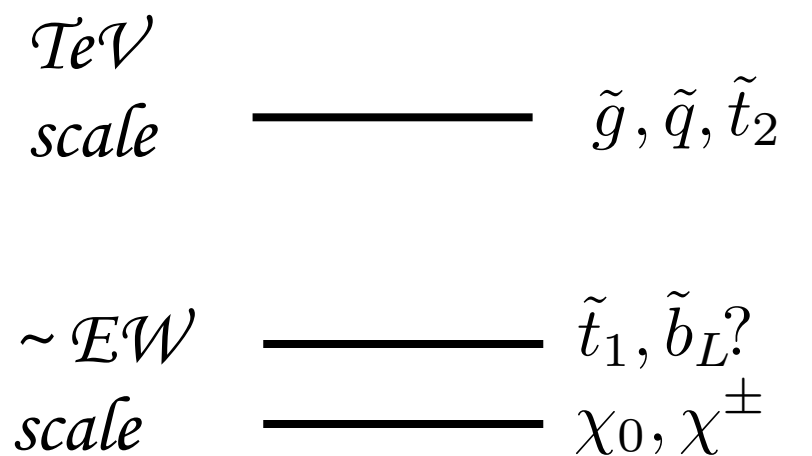
Natural SUSY

What is the minimal required spectrum for SUSY motivated out of naturalness?

What do we need to hold out?



.... we need a spectrum like this perhaps:



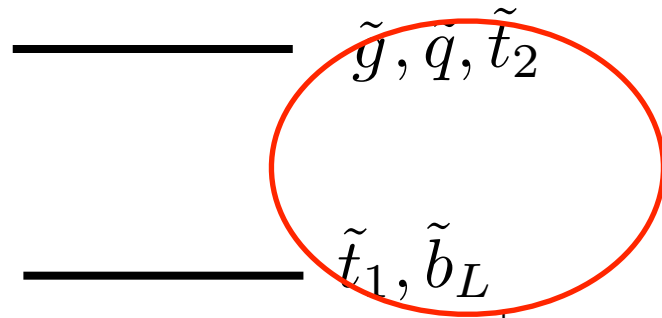
Want this spectrum to not re-introduce fine tuning: higgsinos linked to μ and

$$\delta_{\tilde{t}} m_Z^2 \simeq \frac{3}{8\pi^2 v^2} \left[2m_t^2 (m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2 - 2m_t^2) + \frac{1}{4} (\delta m)^4 \sin^2(2\theta_{\tilde{t}}) \right] \log \left(\frac{2\Lambda^2}{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \right)$$

Natural SUSY

Lets see what the data says for lowish : $\tan \beta \sim 10$

TeV
scale



$\tilde{g}, \tilde{q}, \tilde{t}_2$

\tilde{t}_1, \tilde{b}_L

χ_0, χ^\pm

Custodial sym violation:
precision m_W

What "hitting sbottom" is about.

Natural SUSY

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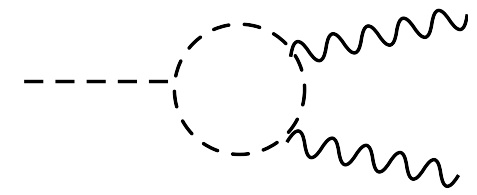
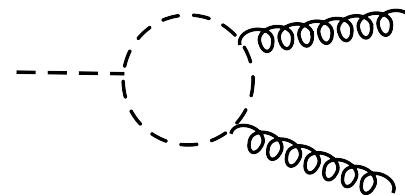
Custodial sym violation:
precision m_W

$\sim EW$
scale

———— \tilde{t}_1, \tilde{b}_L
———— χ_0, χ^\pm

What “hitting sbottom” is about.

Higgs properties



Natural SUSY

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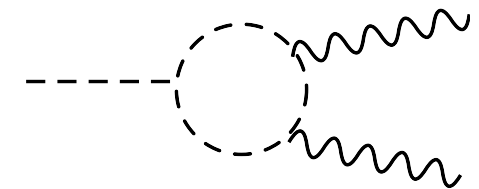
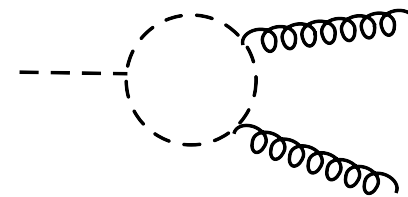
TeV
scale ————— $\tilde{g}, \tilde{q}, \tilde{t}_2$

$\sim EW$
scale ————— \tilde{t}_1, \tilde{b}_L
————— χ_0, χ^\pm

Custodial sym violation:
precision m_W

What "hitting sbottom" is about.

Higgs properties



Precision flavour physics constraints
sensitive to the stop splitting
 $\text{Br}(\bar{B} \rightarrow X_s \gamma)$

Natural Susy

Do a real EFT matching taking into account the QCD latching correction (in the no mixing limit):

Stop contributions to Higgs production:

$$\frac{\sigma(gg \rightarrow h)}{\sigma^{SM}(gg \rightarrow h)} \simeq \frac{\Gamma(h \rightarrow gg)}{\Gamma^{SM}(h \rightarrow gg)} \simeq |1 + r_g|^2 \quad r_g = \frac{C_g(\alpha_s) F_g(m_{\tilde{t}_1}, m_{\tilde{t}_2}, \theta_{\tilde{t}})}{F_g^{SM}(m_t, m_b \dots)}$$

Stop contributions to interesting decays:

$$\frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma^{SM}(h \rightarrow \gamma\gamma)} \simeq |1 + r_\gamma|^2, \quad r_\gamma = \frac{N_c Q_{\tilde{t}}^2 C_\gamma(\alpha_s) F_g(m_{\tilde{t}_1}, m_{\tilde{t}_2}, \theta_{\tilde{t}})}{F_\gamma^{SM}(m_t, W, m_b \dots)}$$

QCD squark
matching
correction

Same loop functions, non abelian nature of QCD irrelevant at leading order in matching.

Natural Susy

Do a real EFT matching taking into account the QCD latching correction (in the no mixing limit):

The Effective Lagrangian:

$$\mathcal{L}_{HD} = -\frac{c_g g_3^2}{2\Lambda^2} H^\dagger H G_{\mu\nu}^A G^{A\mu\nu} - \frac{c_W g_2^2}{2\Lambda^2} H^\dagger H W_{\mu\nu}^a W^{a\mu\nu} - \frac{c_B g_1^2}{2\Lambda^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} - \frac{c_{WB} g_1 g_2}{2\Lambda^2} H^\dagger \tau^a H B_{\mu\nu} W^{a\mu\nu},$$

In terms of operators:

$$\sigma_{gg \rightarrow h} \approx \sigma_{gg \rightarrow h}^{SM} \left| 1 + \frac{2}{F_g^{SM}} \frac{v^2 \tilde{c}_g}{\Lambda^2} \right|^2, \quad \Gamma_{h \rightarrow \gamma\gamma} \approx \Gamma_{h \rightarrow \gamma\gamma}^{SM} \left| 1 + \frac{1}{F_\gamma^{SM}} \frac{v^2 \tilde{c}_\gamma}{\Lambda^2} \right|^2.$$

Matching with no running: $\frac{v^2 \tilde{c}_g}{\Lambda^2} \simeq C_g(\alpha_s) \frac{F_g}{2}, \quad \frac{v^2 \tilde{c}_\gamma}{\Lambda^2} \simeq N_c Q_{\tilde{t}}^2 C_\gamma(\alpha_s) F_g$

$$C_g(\alpha_s) = 1 + \frac{25 \alpha_s}{6\pi}, \quad C_\gamma(\alpha_s) = 1 + \frac{8 \alpha_s}{3\pi}.$$

Natural Susy

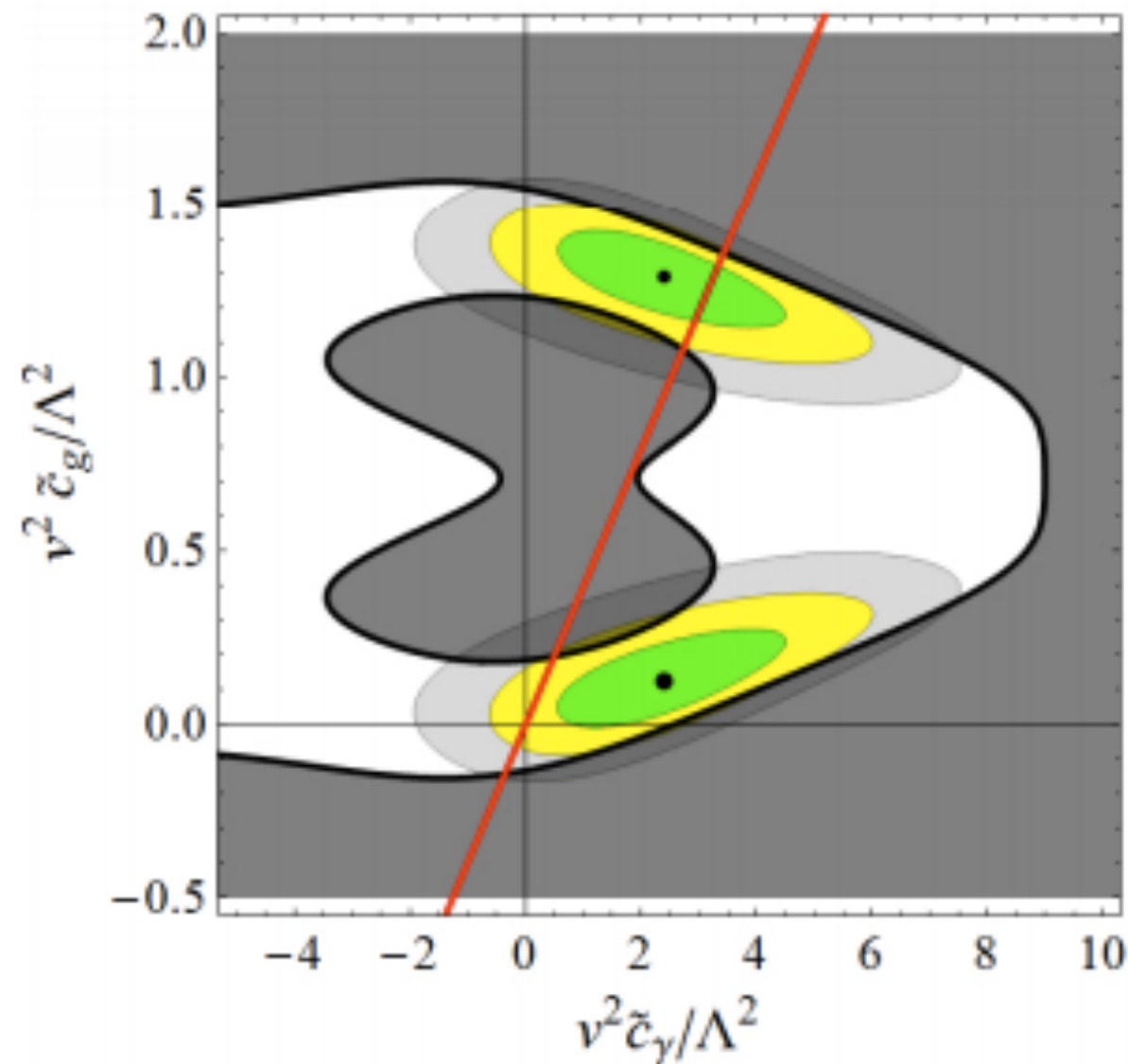
Espinosa, Grojean, Sanz, Trott arXiv:1207.7355

This is a predictive scenario for the wilson coefficients of the higher d ops:

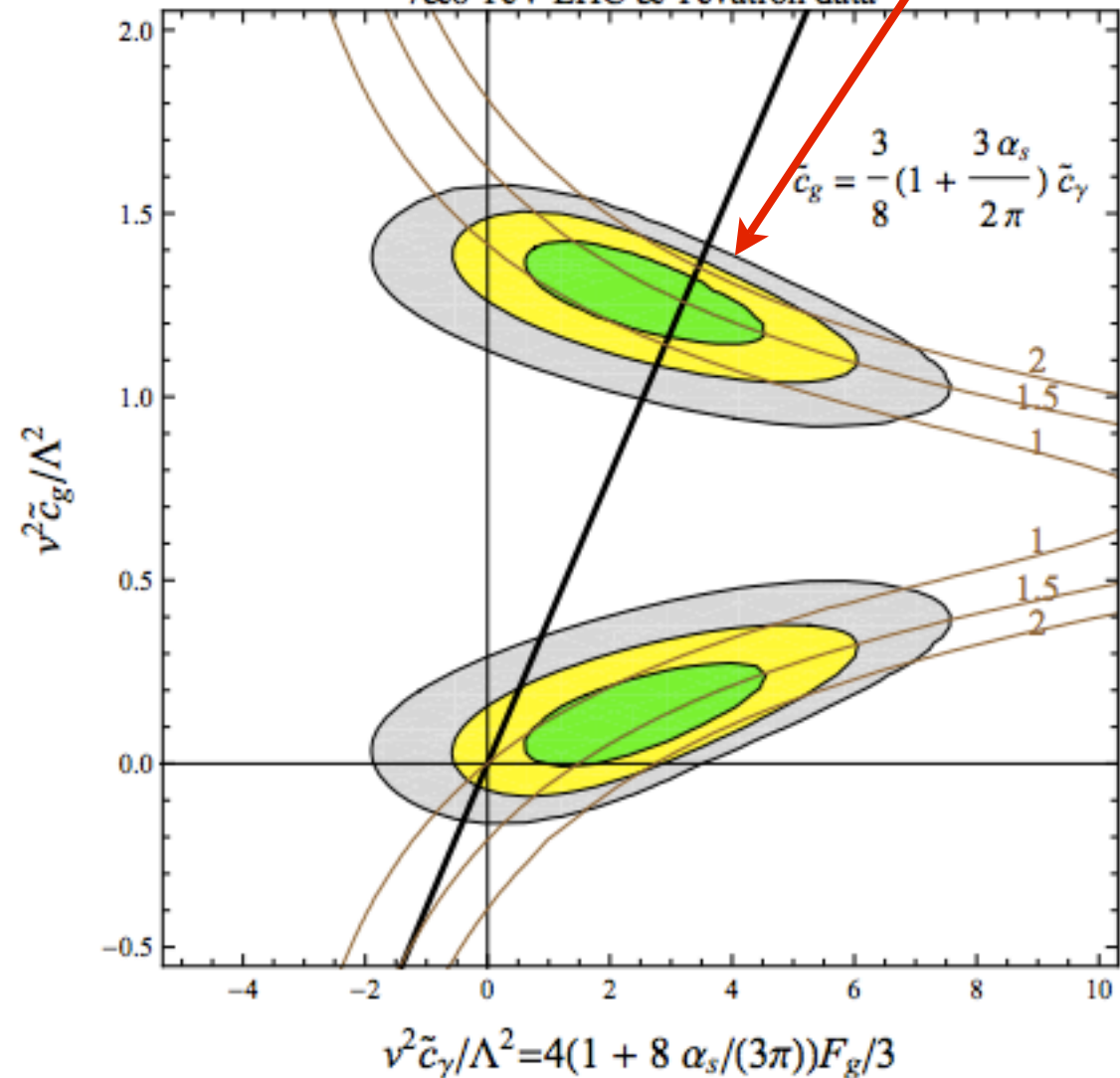
$$\frac{\tilde{c}_g}{\tilde{c}_\gamma} = \frac{1}{2N_c Q_{\tilde{t}}^2} \frac{C_g(\alpha_s)}{C_\gamma(\alpha_s)} = \frac{3}{8} \left(1 + \frac{3\alpha_s}{2\pi} \right)$$

*monophoton exclusions
limit parameter space
sanz et al. arXiv:1205.1463*

7&8 TeV LHC data & Tevatron



7&8 TeV LHC & Tevatron data



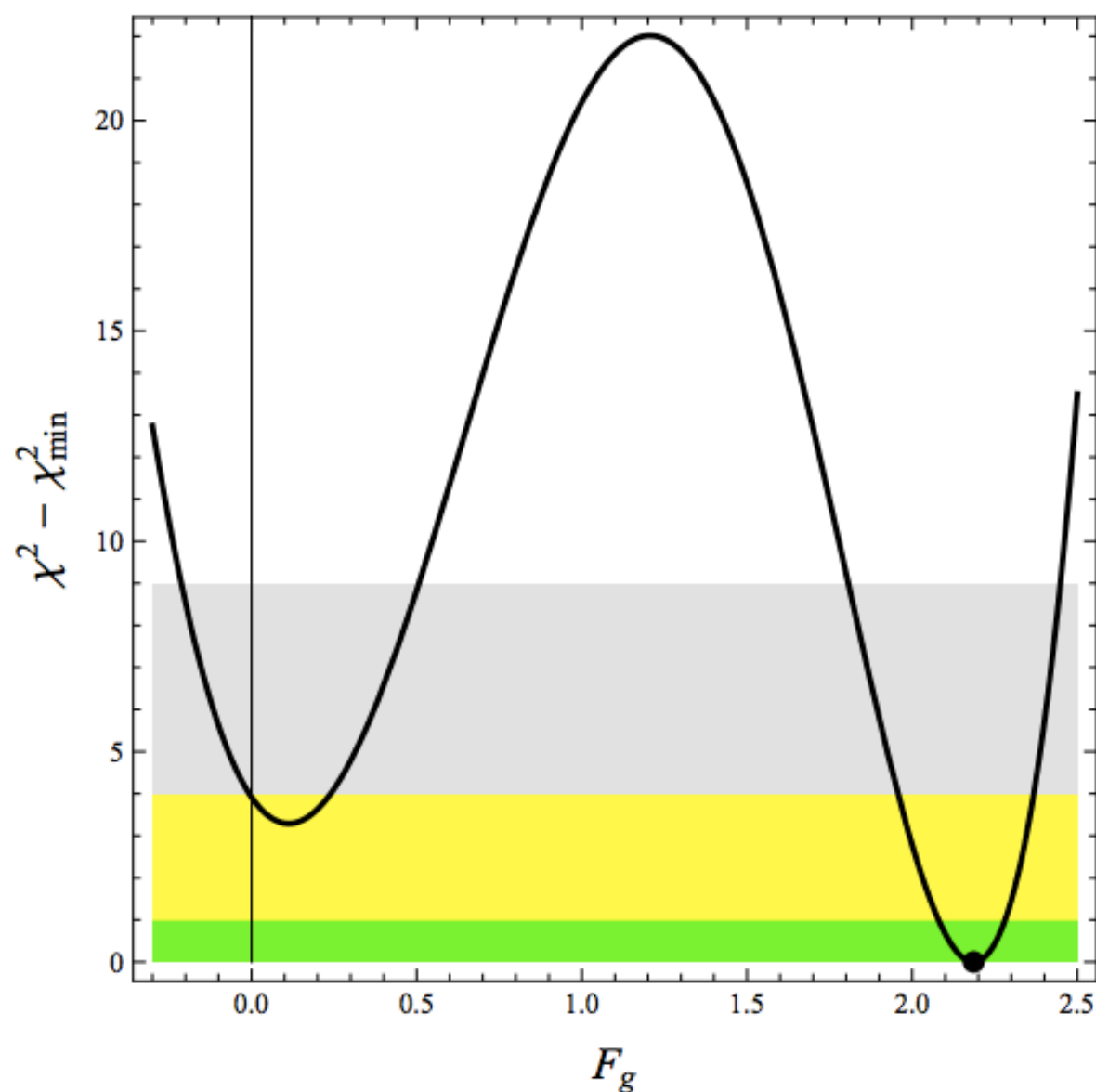
Natural Susy

Espinosa, Grojean, Sanz, Trott arXiv:1207.7355

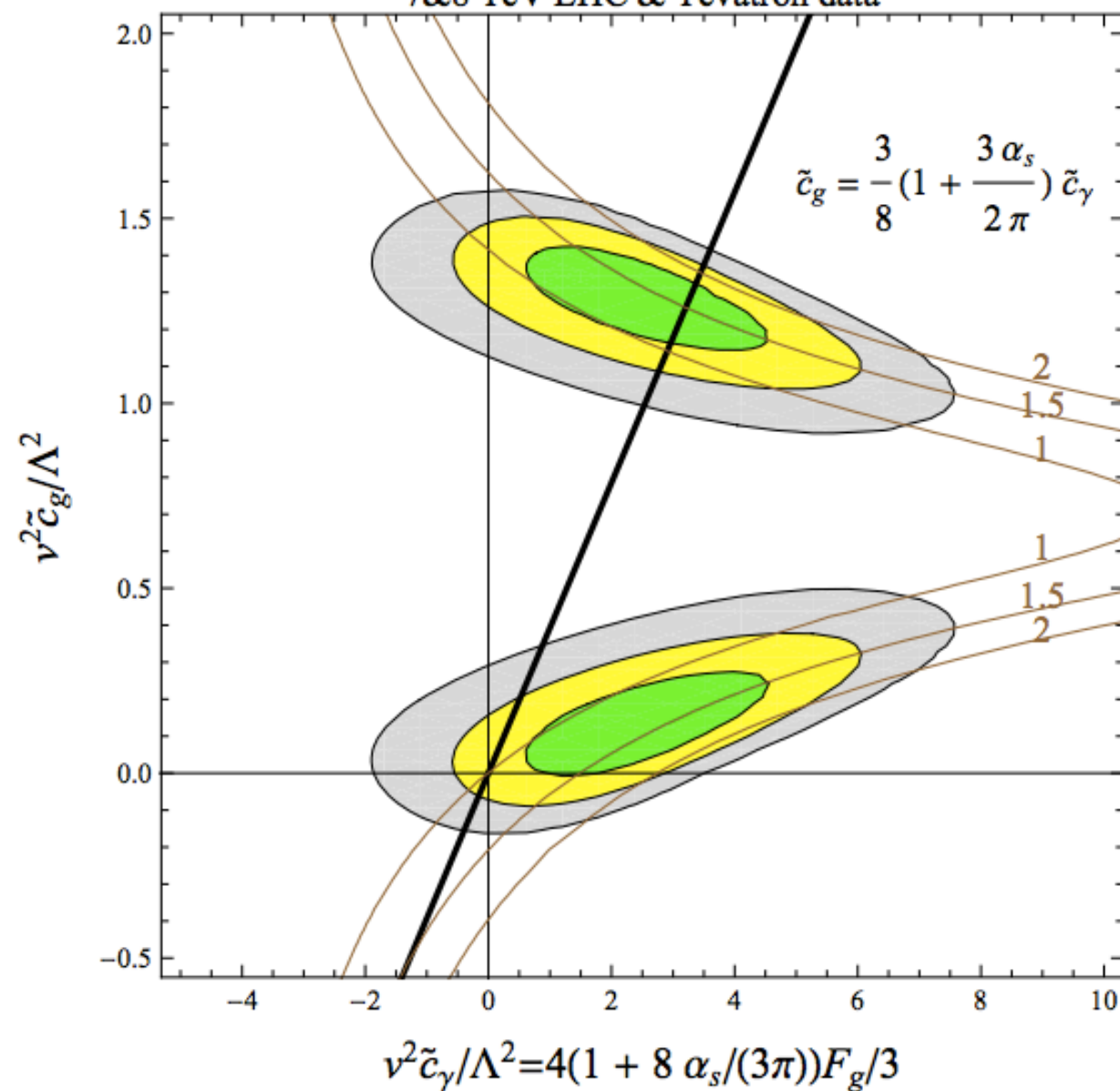
Consider minimal spectrum of stops, left handed sbottom, charginos and neutralinos with large gaugino mass.

This is a predictive scenario for the wilson coefficients of the higher d ops:

Global Higgs Data

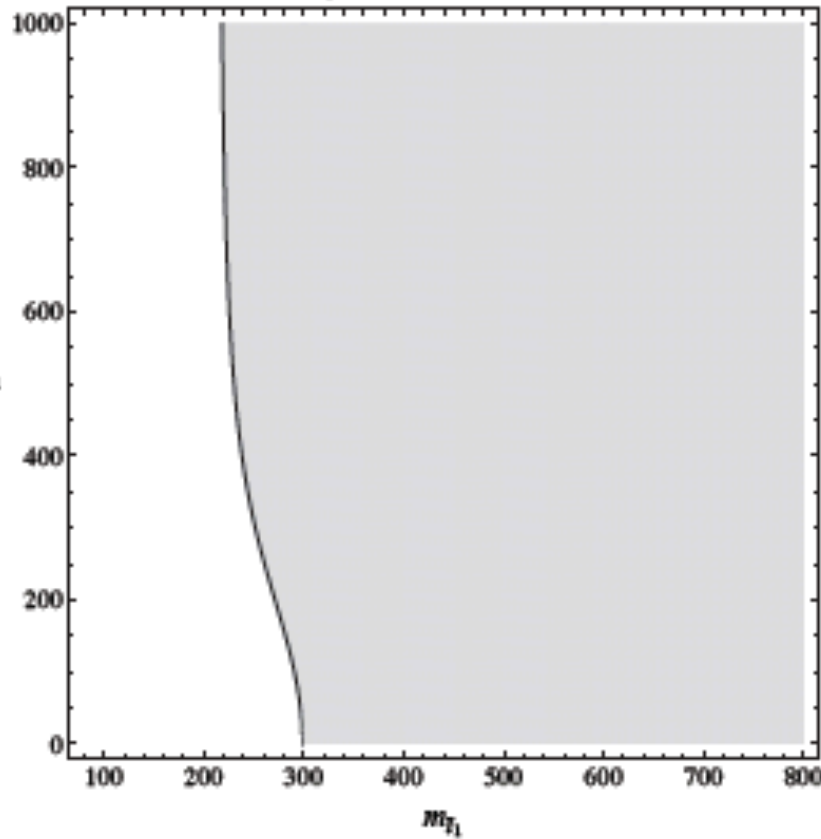


7&8 TeV LHC & Tevatron data

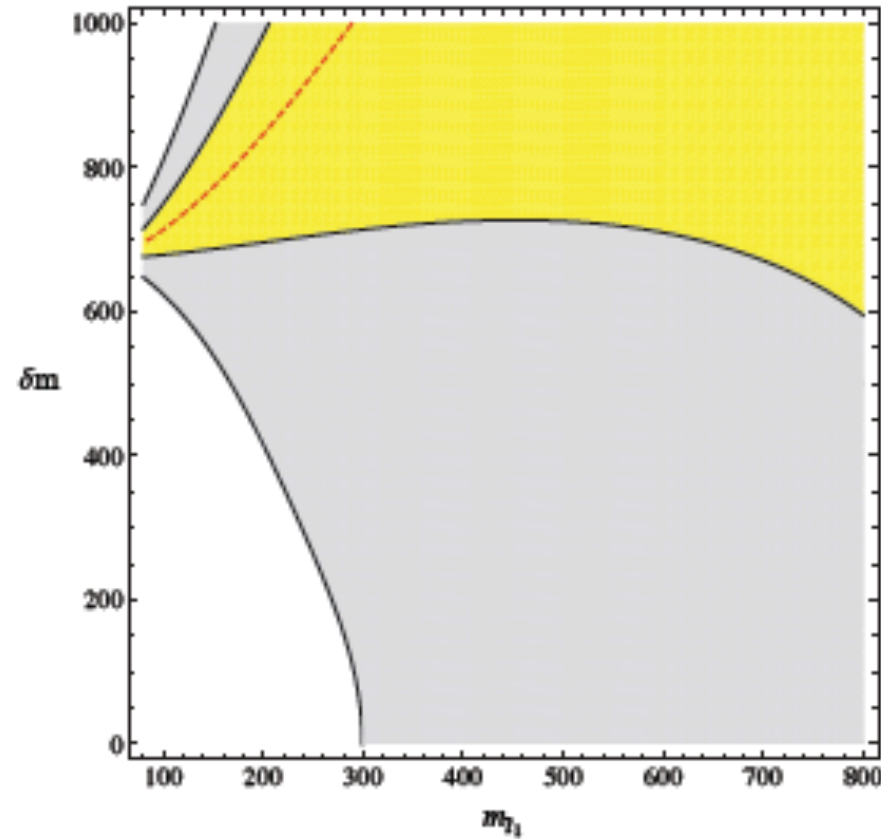


Translate to stop space

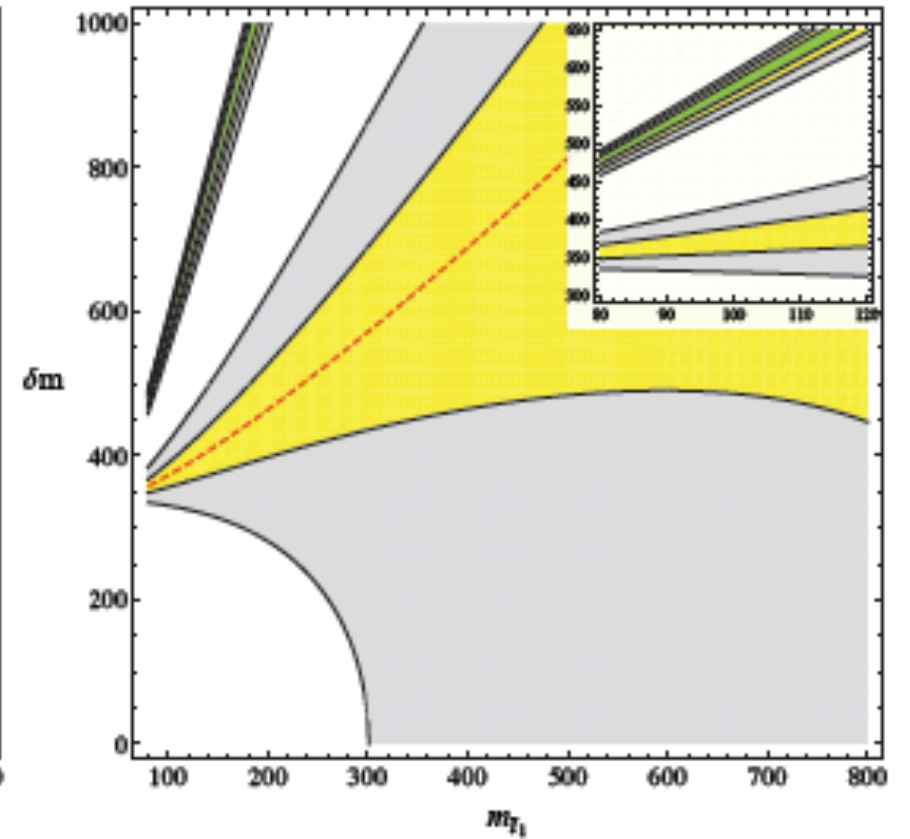
$\theta_t=0$ and $\tan\beta=2-20$



$\theta_t=\frac{\pi}{12}$ and $\tan\beta=2-20$



$\theta_t=\frac{\pi}{4}$ and $\tan\beta=2-20$



- 1σ
- 2σ
- 3σ

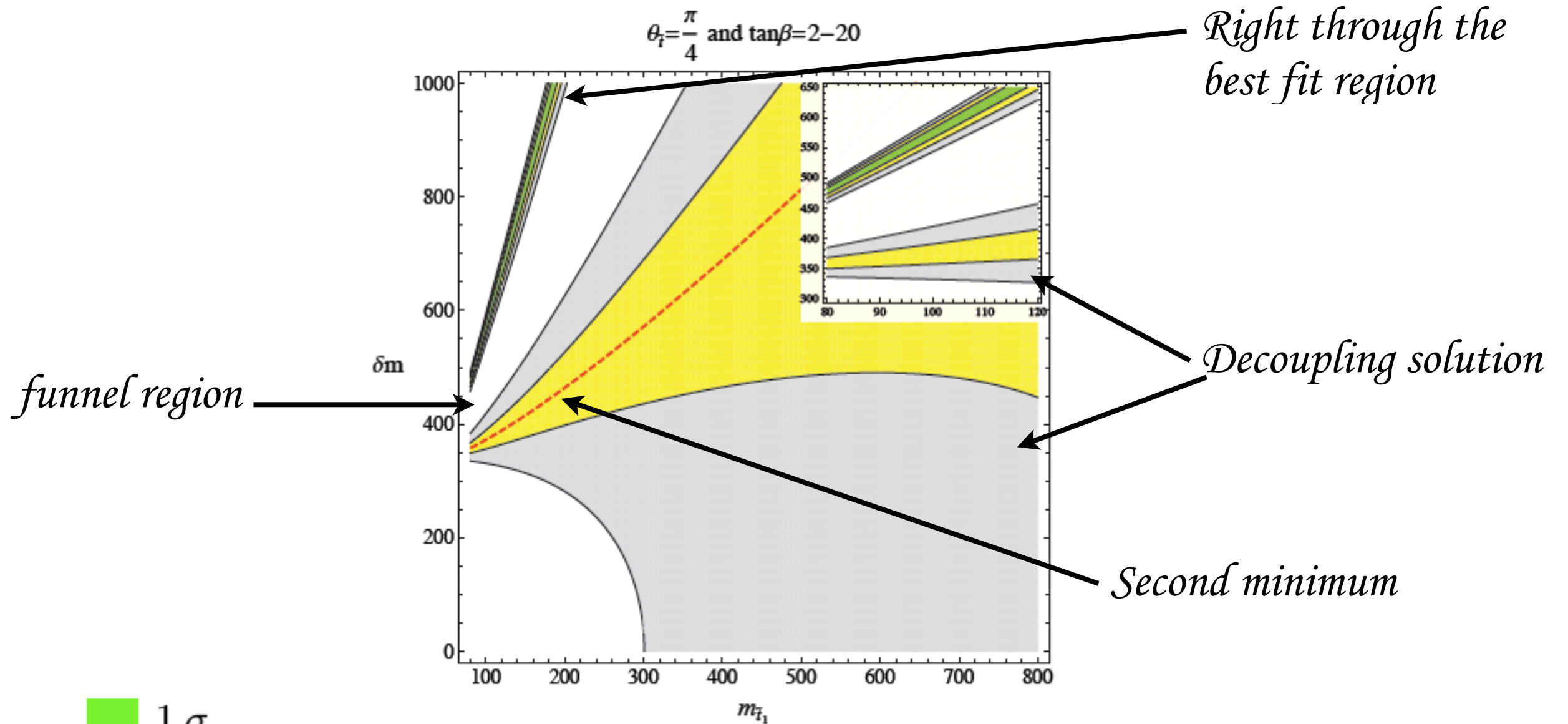
Map to stop space:

$$F_g = -\frac{1}{3} \left[\frac{m_t^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2}{m_{\tilde{t}_2}^2} - \frac{1}{4} \sin^2(2\theta_t) \frac{\delta m^4}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right]$$

Light unmixed stops in bad shape, large mixing preferred in the data.

Translate to stop space

$$\theta_{\tilde{t}} = \frac{\pi}{4} \text{ and } \tan\beta = 2-20$$



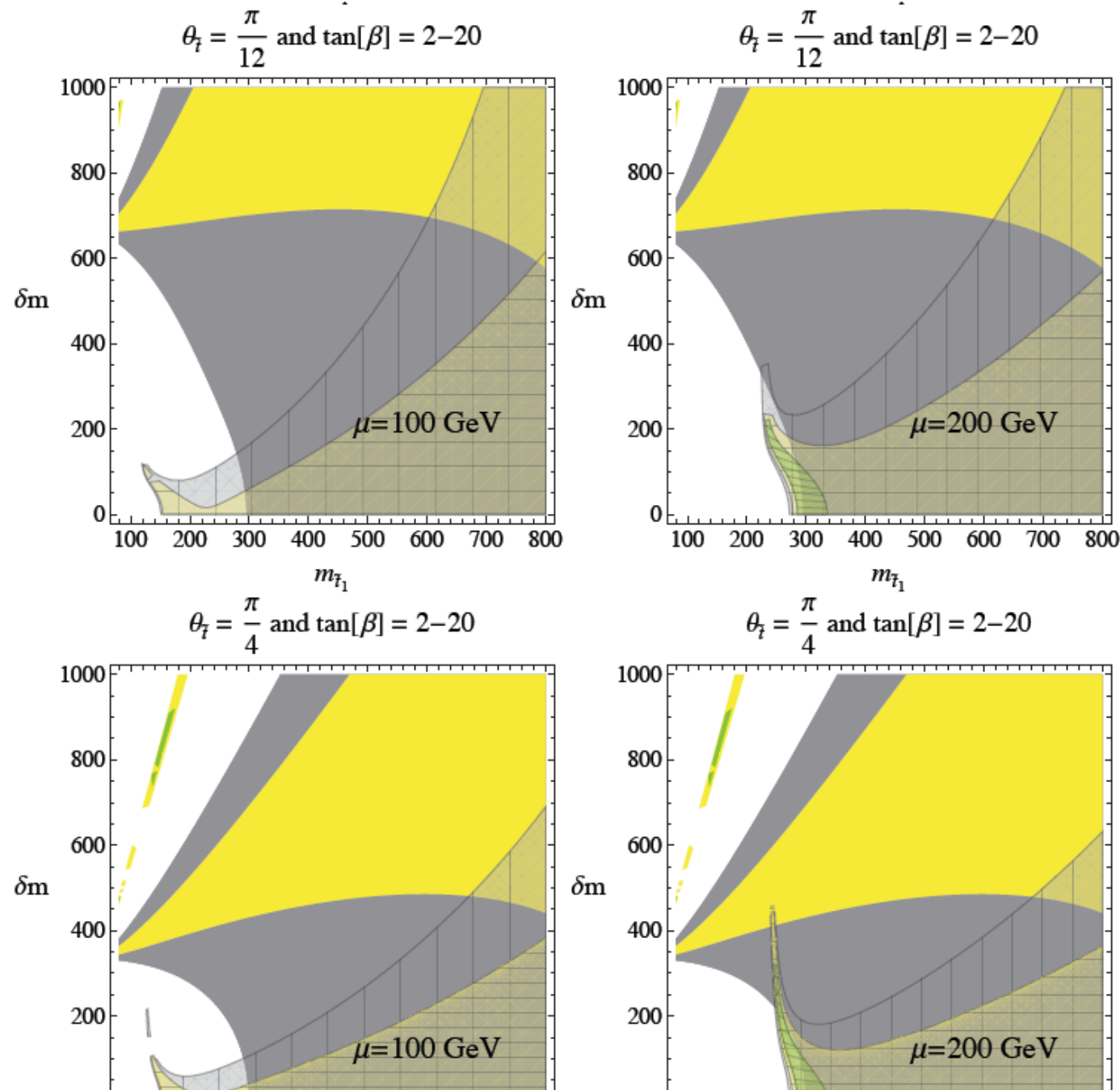
- 1σ
- 2σ
- 3σ

Map to stop space:

$$F_g = -\frac{1}{3} \left[\frac{m_t^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2}{m_{\tilde{t}_2}^2} - \frac{1}{4} \sin^2(2\theta_t) \frac{\delta m^4}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right]$$

Light unmixed stops in bad shape, large mixing preferred in the data.

Consistency of B to s gamma



*Not so good news...
 μ dependance as the
 light chargino
 contributes here in
 NSUSY*

*$\text{Br}(\bar{B} \rightarrow X_s \gamma)$ basically wants
 degenerate stops, this will become
 stronger once the latest BABAR result
 is averaged in.*

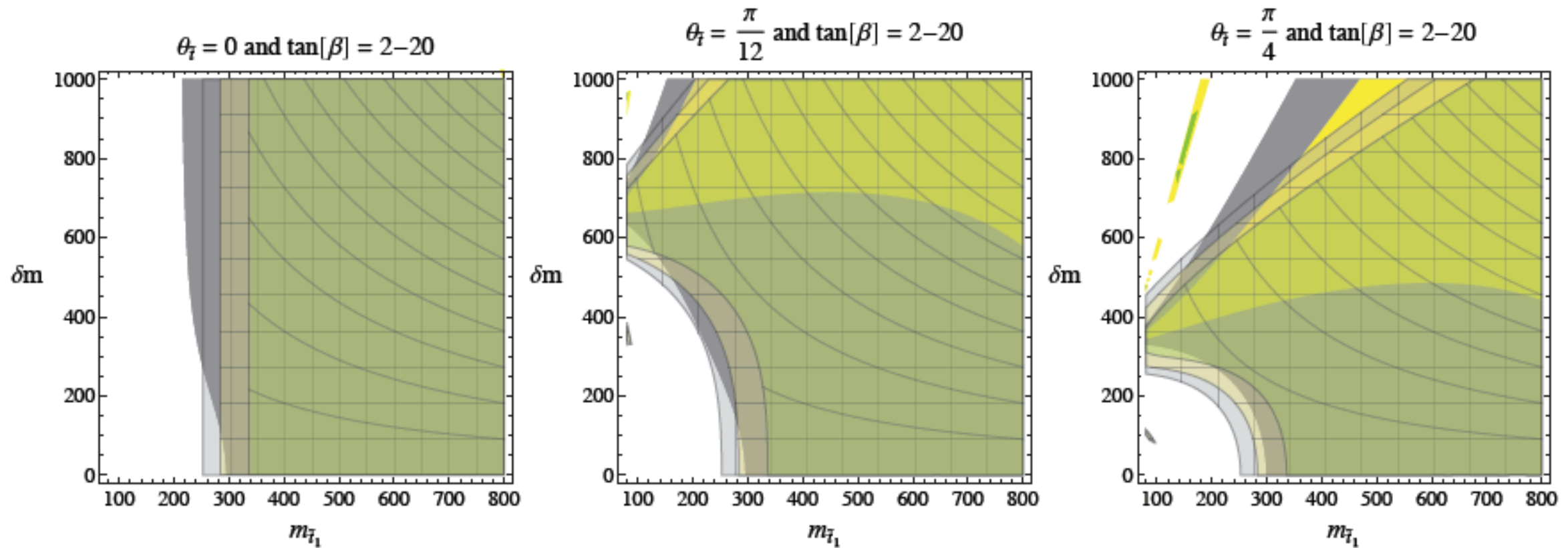
$$\text{BR}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}} = [(3.15 \pm 0.23) - 8.0 \Delta C_7(\mu_0) - 1.9 \Delta C_8(\mu_0)] \times 10^{-4}$$

Theory prediction: B. Grzadkowski and M. Misiak, Phys. Rev. D **78**, 077501 (2008) [hep-ph/0802.1413].

$$\text{BR}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}} = [3.55 \pm 0.24 \pm 0.09] \times 10^{-4}$$

Precision EW

To be fair there might also be good news... what an astonishing coincidence!



Recent precise measurements of M_W from the Tevatron: $(m_W)_{exp} = 80.385 \pm 0.015 \text{ GeV}$,
 $(m_W)_{SM} = 80.368 \pm 0.006 \text{ GeV}$

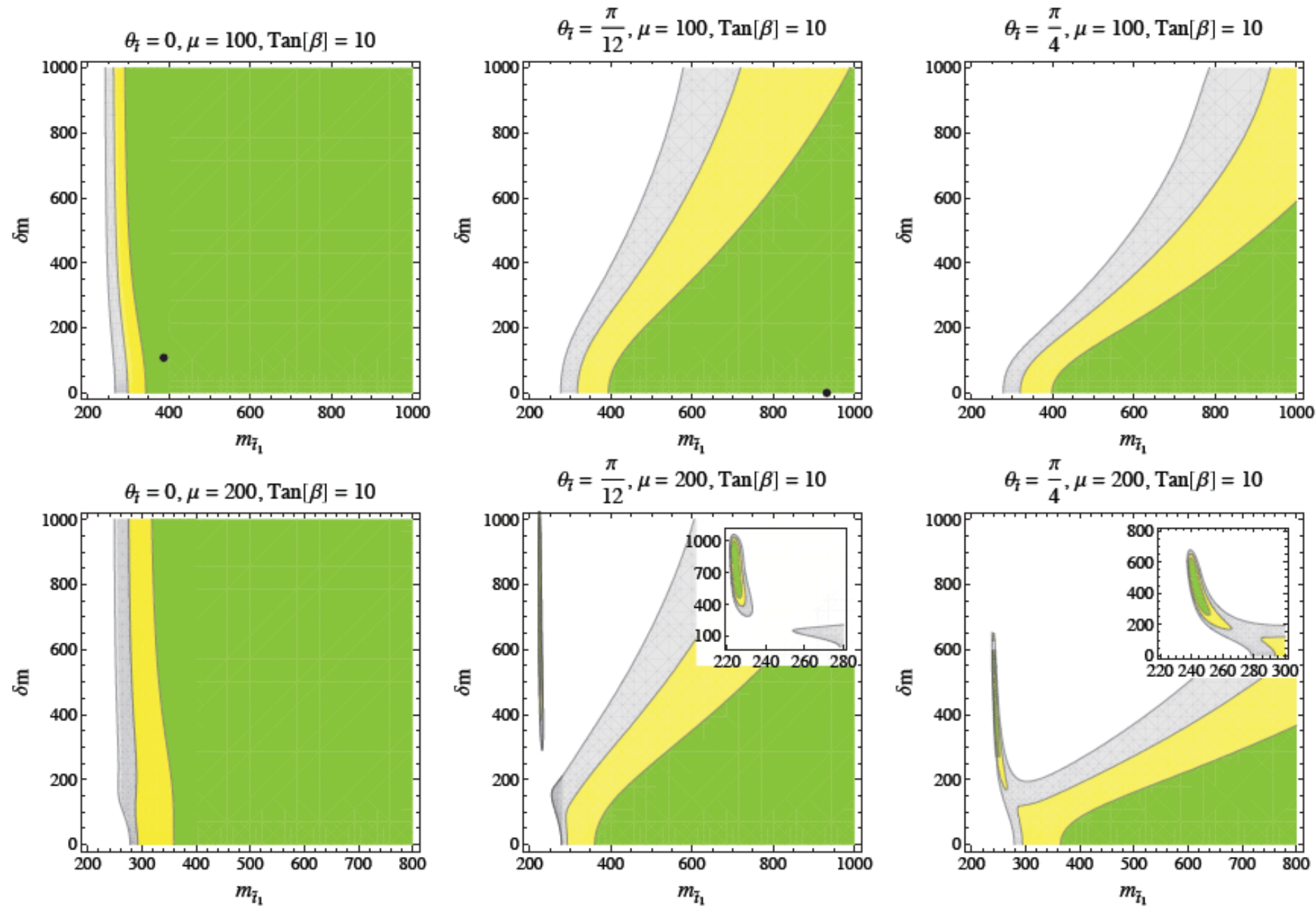
SUSY contribution:

$$(\Delta m_W)^{SUSY} \simeq \frac{m_W c_W^2}{2(c_W^2 - s_W^2)} \Delta \rho^{SUSY}$$

$$\Delta \rho_0^{SUSY} \simeq \frac{3 G_F}{8 \sqrt{2} \pi^2} \left\{ \sum_{i=1,2} |\langle \tilde{t}_L | \tilde{t}_i \rangle|^2 F_0[m_{\tilde{t}_i}^2, m_{\tilde{b}_L}^2] - |\langle \tilde{t}_L | \tilde{t}_1 \rangle|^2 |\langle \tilde{t}_L | \tilde{t}_2 \rangle|^2 F_0[m_{\tilde{t}_1}^2, m_{\tilde{t}_2}^2] \right\}$$

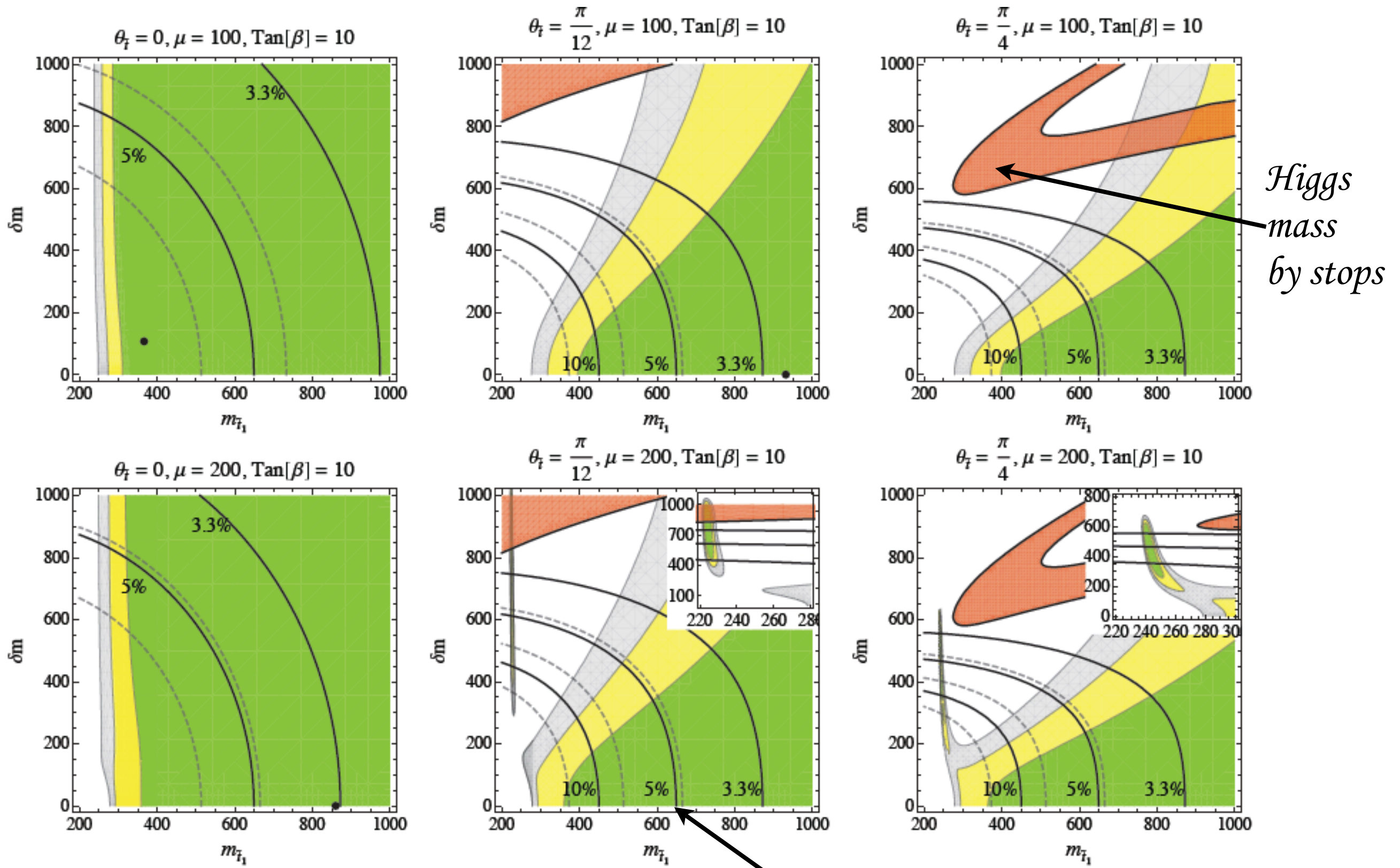
$$\Delta \rho^{SUSY} \lesssim (3.0 \pm 2.8) \times 10^{-4}$$

Put it all together: Indirect stop fit



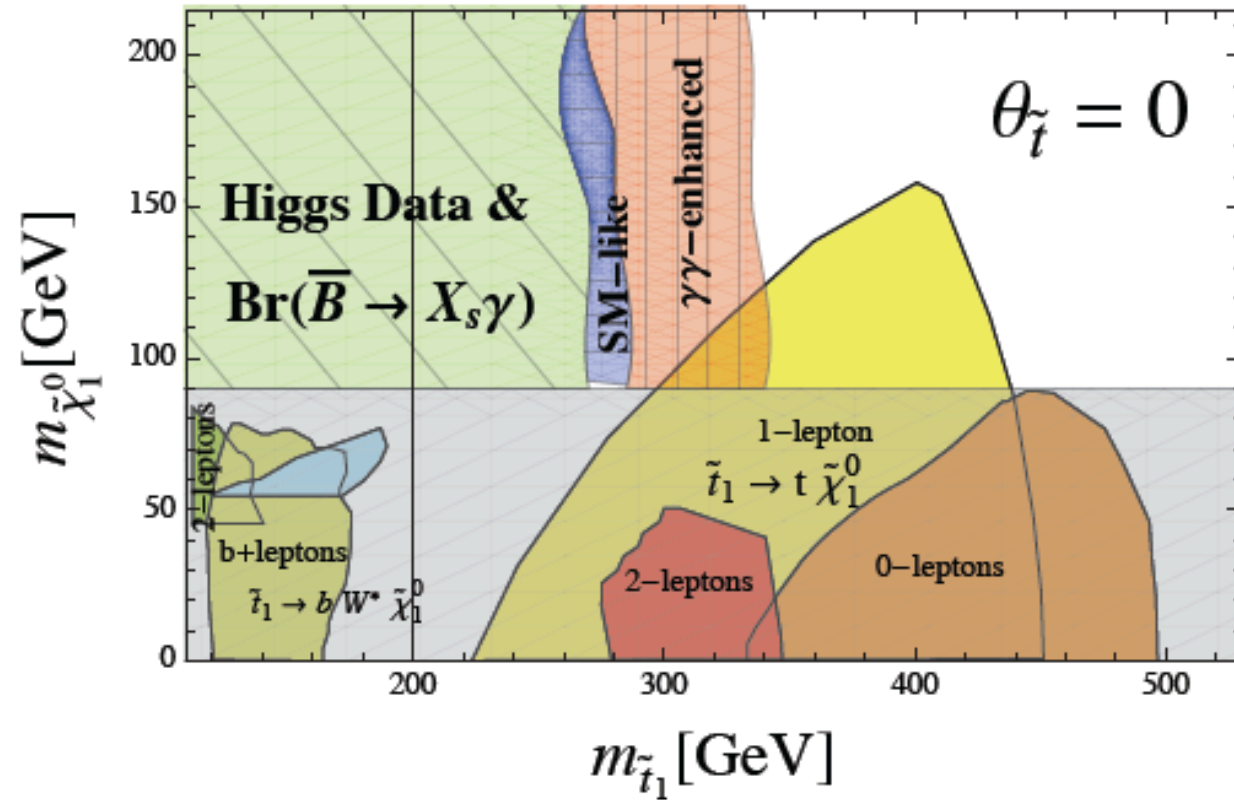
Putting everything together a good fit to the data for ~ 400 geV stop state comparable in quality of fit to the SM, no dramatic improvement. Higgs mass hard to accommodate

Put it all together: Indirect stop fit

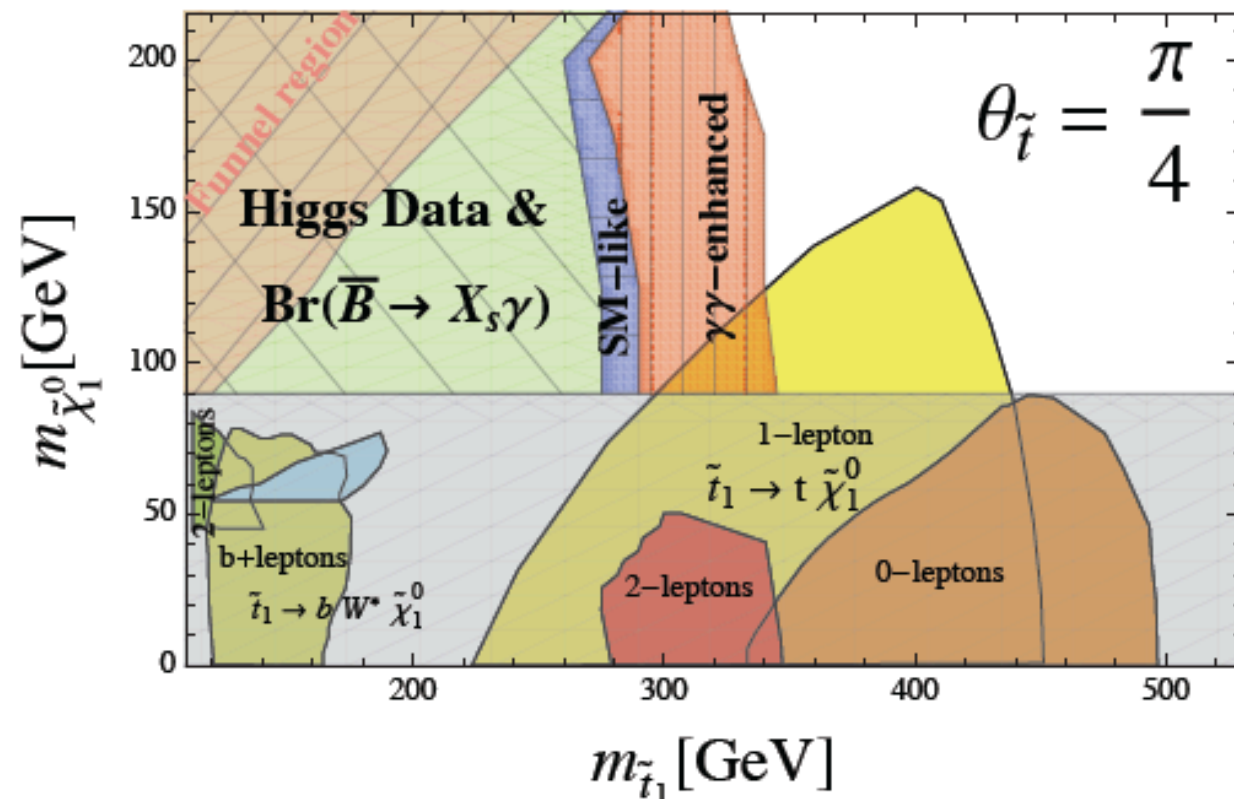


Indirect Exclusion currently and prospects for stops

Espinosa, Grojean, Sanz, Trott arXiv:1207.7355



Combined exclusion with demand that the higgs properties 95 % CL exclude and $\text{Br}(\bar{B} \rightarrow X_s \gamma)$ be within 2 sigma of its experimental value



Indirect probes for stops are powerful tools. Dedicated experimental study warranted and underway now in ATLAS

Conclusions



Need more data.