

# Dark Matter: Colliders (and Direct Detection)

M. Perelstein, LEPP/Cornell U.

LCWS-12, Arlington TX

October 25, 2012



Cornell University  
Laboratory for Elementary-Particle Physics

# Dark Matter Puzzle:

- About 25% of the energy in the universe is dark, non-relativistic matter
- Non-particle explanations unlikely
- $\chi$  has to be stable (or at least  $\tau \geq 10$  bln. years)
- $\chi$  cannot have strong interactions (otherwise  $p\chi$  exotic nuclei) or electric charge (dark)
- $\chi$  cannot be a Standard Model neutrino (free streaming)
- Have to invent (at least one) new particle

# WIMP: a Perfect Fit

•  $\chi$ 's interact with the SM matter via **weak forces** (or a new interaction of similar strength/range)

•  $\chi$  is **massive** (1 GeV – 10 TeV range)  $\Rightarrow$   $\chi$ 's are in **thermal equilibrium** with the SM matter as long as  $T > M(\chi)$ :  
$$n_\chi \sigma v > H$$

• When  $T < M(\chi)$ ,  $n_\chi \propto \exp(-M/T)$  (Boltzmann suppression) and  $\chi$ 's **decouple**

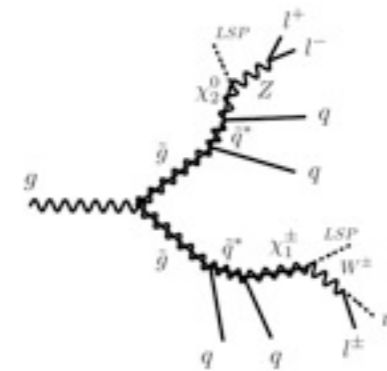
• Energy density of  $\chi$ 's today: 
$$\rho_\chi \approx \frac{T_0^3}{M_{\text{pl}} \sigma} \sim \rho_c$$

# WIMPs at Colliders

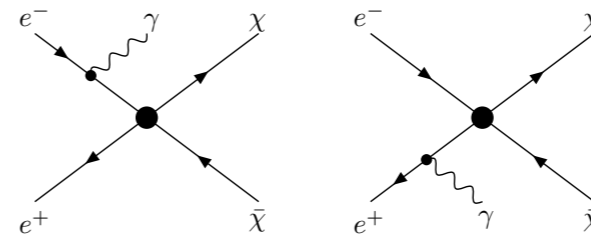
- Much of the reasonable mass range for WIMPs is within reach of the **LHC** and **ILC/CLIC**

- Two basic ways to produce WIMPs at colliders:

- **In decay** of heavier exotic particles: for example



- **Direct production**: for example



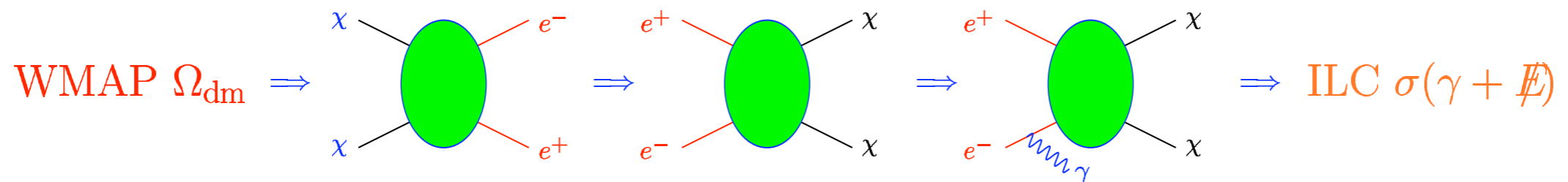
- Production in decay can dominate (e.g. if decaying particle is colored → high rate) but is **more model-dependent** (assumptions beyond WIMP!)

- **Strong LHC bounds** on colored exotic states decaying to MET+SM shrinking parameter space for observing WIMPs in decays...

- **Direct production** is less model-dependent, and is not yet strongly constrained. Will be my focus in this talk.

# Predicting WIMP Signatures: “Model-Independent” Approach

- Many particle physics models contain WIMPs: [SUSY](#), [Extra Dimensions](#), [Little Higgs](#), etc.
- Direct (radiative) WIMP production can be described within a **model-independent formalism** [[Birkedal, Matchev, MP, hep-ph/0403004](#)]



# Assumptions:

- Assume **generic** mass spectrum (no resonances, no coannihilations)

- At the time of  $\chi$  decoupling, the only important reactions are  $\chi\chi \leftrightarrow X_i\bar{X}_j$ , where  $X_i$  is SM

- For non-relativistic WIMPs, can be expanded as:

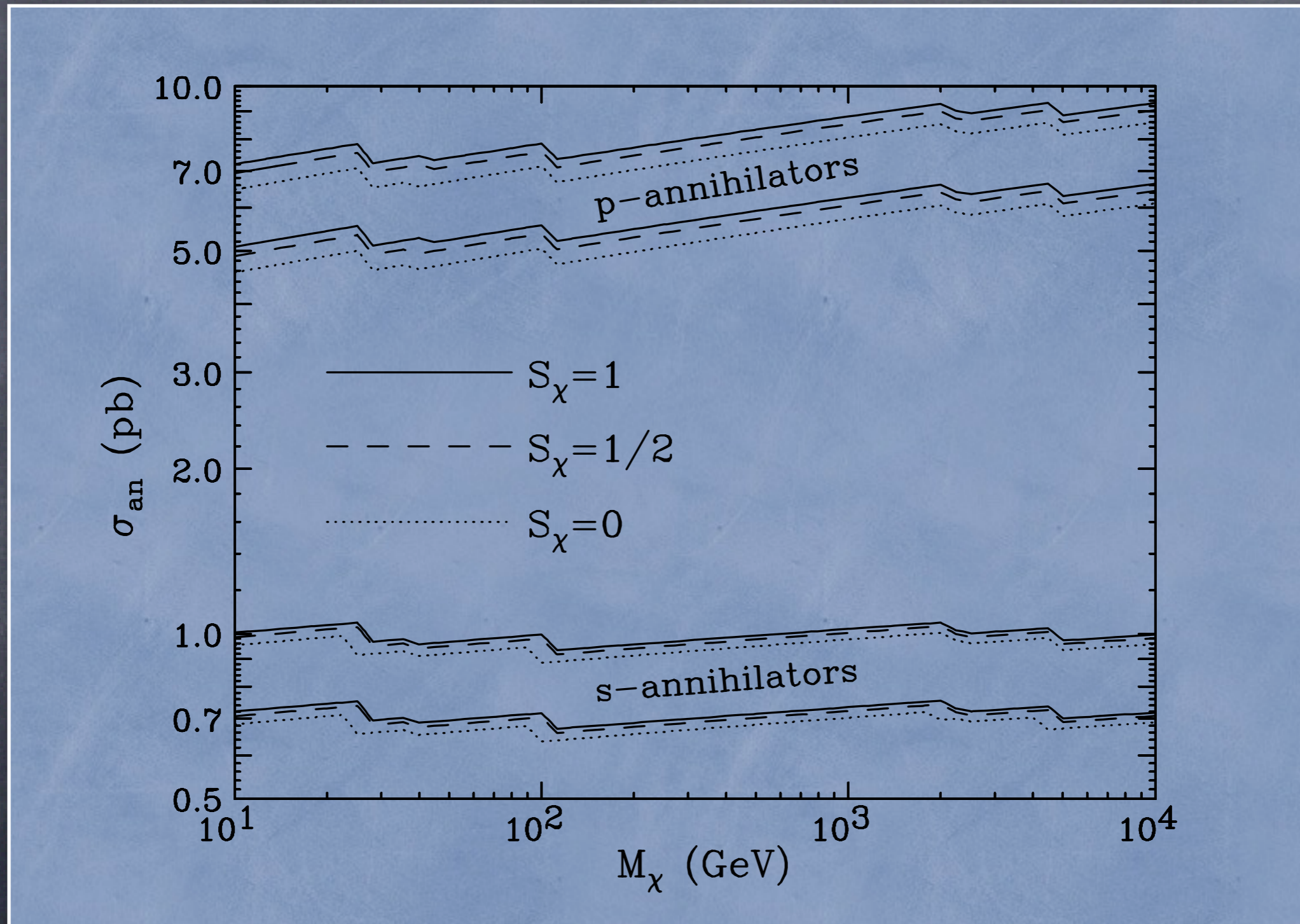
$$\sigma_i v = \sigma_i^{(0)} + \sigma_i^{(1)} v^2 + \dots$$

- Dominated by either **s-wave** or **p-wave**

- Define

$$\sigma_{\text{an}} = \sum_i \sigma_i^{J_0}$$

# $\Omega_{\text{dm}}$ determines $\sigma_{\text{an}}$



$2\sigma$  constraint using  $\Omega_{\text{dm}}h^2 = 0.112 \pm 0.009$  (WMAP)

# From Cosmology to Colliders

- Cosmology provides a precise, model-independent measurement of  $\sigma_{\text{an}}$
- **Idea:** use this information to predict  $\chi$  production rate at a collider!
- Step 1: **Detailed Balancing** (DB)

$$\frac{\sigma(\chi\chi \rightarrow e^+e^-)}{\sigma(e^+e^- \rightarrow \chi\chi)} = 2 \frac{v_e^2 (2S_e + 1)^2}{v_\chi^2 (2S_\chi + 1)^2}$$

- Define **annihilation fraction:**  $\kappa_e = \sigma_{e^+e^-}^{J_0} / \sigma_{\text{an}}$



# Tagging and Factorization

- Obtain a prediction:

$$\sigma(e^+e^- \rightarrow \chi\chi) = \frac{2^{2(J_0+1)}}{(2S_\chi + 1)^2} \kappa_i \sigma_{\text{an}} \left(1 - \frac{4M_\chi^2}{s}\right)^{1/2+J_0}$$

- This is unobservable (like  $e^+e^- \rightarrow \nu\bar{\nu}$ )

- Consider instead  $e^+e^- \rightarrow \chi\chi + \gamma$

- Step 2: Use soft/collinear factorization:

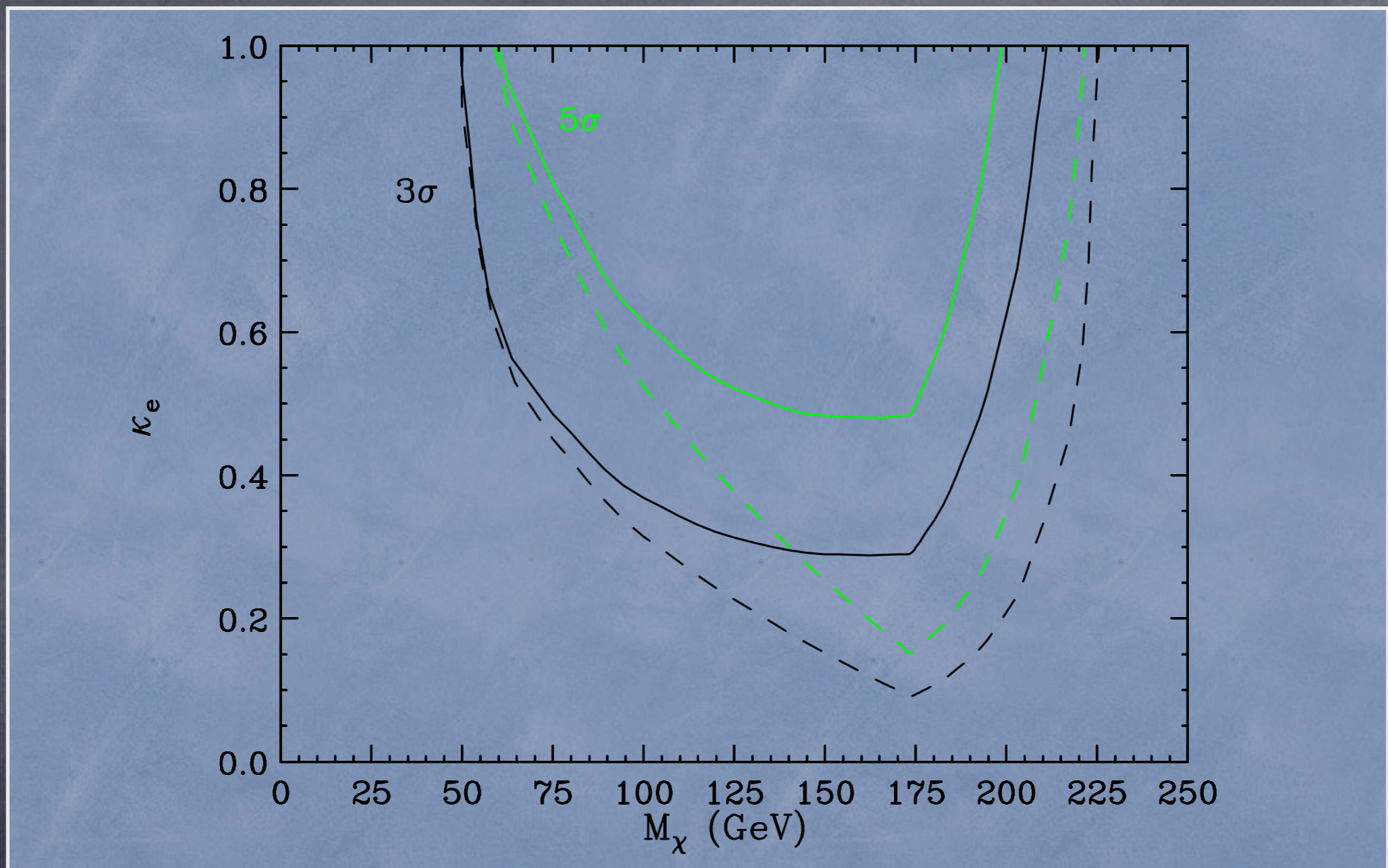
$$\frac{d\sigma(e^+e^- \rightarrow 2\chi + \gamma)}{dx d\cos\theta} \approx \mathcal{F}(x, \cos\theta) \hat{\sigma}(e^+e^- \rightarrow 2\chi)$$

$$\mathcal{F}(x, \cos\theta) = \frac{\alpha}{\pi} \frac{1 + (1-x)^2}{x} \frac{1}{\sin^2\theta}, \quad x = 2E_\gamma/\sqrt{s}$$

# Experimental Strategy for a Model-Independent WIMP Search at the ILC

- Look for **photon+missing energy** events
- Impose  $p_T^{\min}(\gamma)$  cut to eliminate fakes (mainly Bhabha)
- Impose  $E_\gamma^{\min}$  cut to ensure **non-relativistic WIMPs**
- Compute and subtract the **irreducible** background (mainly  $e^+e^- \rightarrow \nu\bar{\nu}\gamma$ )
- Look for deviations from zero!

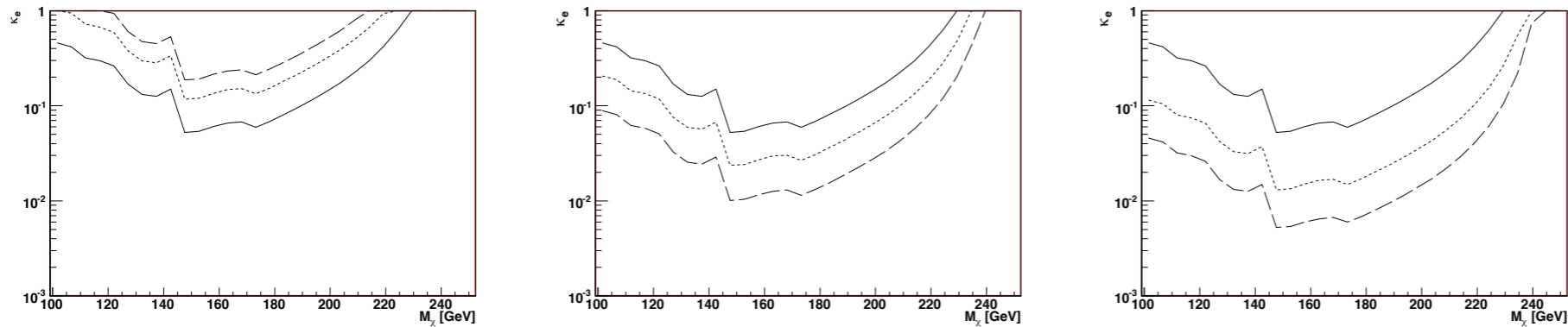
# The Reach of a 500 GeV LC



**Dash** – stat. only ( $\mathcal{L} = 500 \text{ fb}^{-1}$ ), **Solid** – stat. + 0.3% syst.

**Cuts:**  $\sin \theta > 0.1$ ,  $p_T^\gamma > 7.5 \text{ GeV}$ ,  $x_\gamma \in [1 - 8M_\chi^2/s, 1 - 4M_\chi^2/s]$

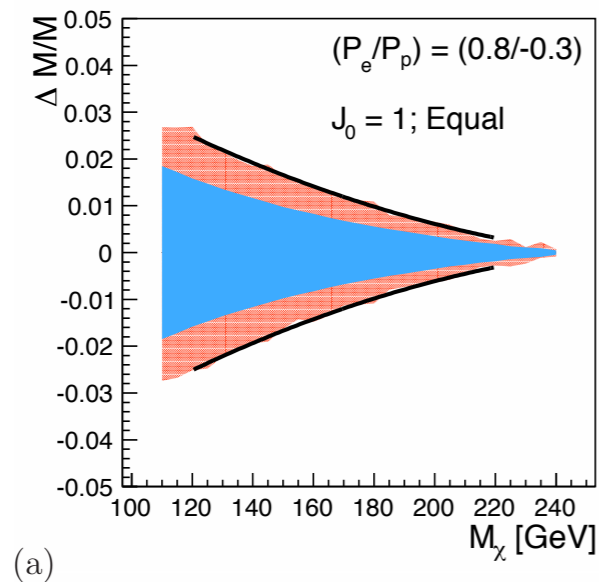
# Detector-Level Studies



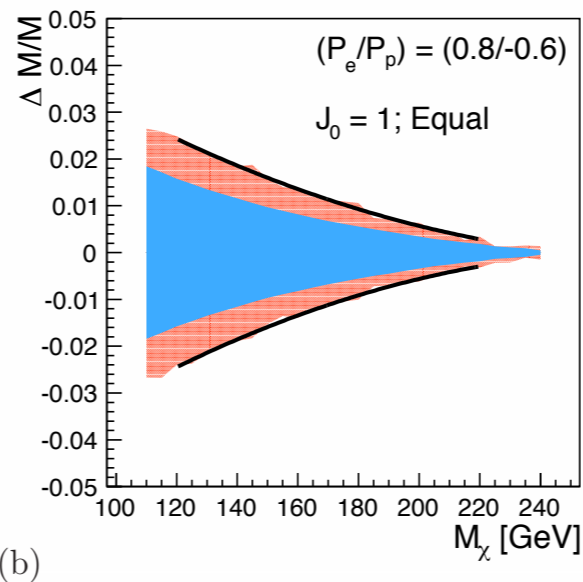
[Bartels, List,  
0709.2629]

Reach down  
to  $\kappa_e \sim 10^{-2}$  !

Figure 4:  $3\sigma$  observation reach of the ILC for a Spin- $\frac{1}{2}$  WIMP in terms of WIMP mass and  $\kappa_e$  for three different assumptions on the chirality of the electron-WIMP coupling, see text. Full line:  $P_{e^-} = P_{e^+} = 0$ , dotted line:  $P_{e^-} = 0.8, P_{e^+} = 0$ , dashed line :  $P_{e^-} = 0.8, P_{e^+} = 0.6$ . Regions above the curves are accessible.



(a)



(b)

[Bartels, Berggren,  
List, 1206.6639]

Percent-level mass  
measurement!

# Alternative: Effective Operator Approach

- The formalism I just reviewed makes no reference to a Lagrangian
- Alternative: Model DM-SM couplings with **effective operators** in a Lagrangian [Beltran et. al. 1002.4137; Goodman et. al. 1005.1286; Bai, Fox, Harnik, 1005.3797; Fox et. al. 1109.4398]

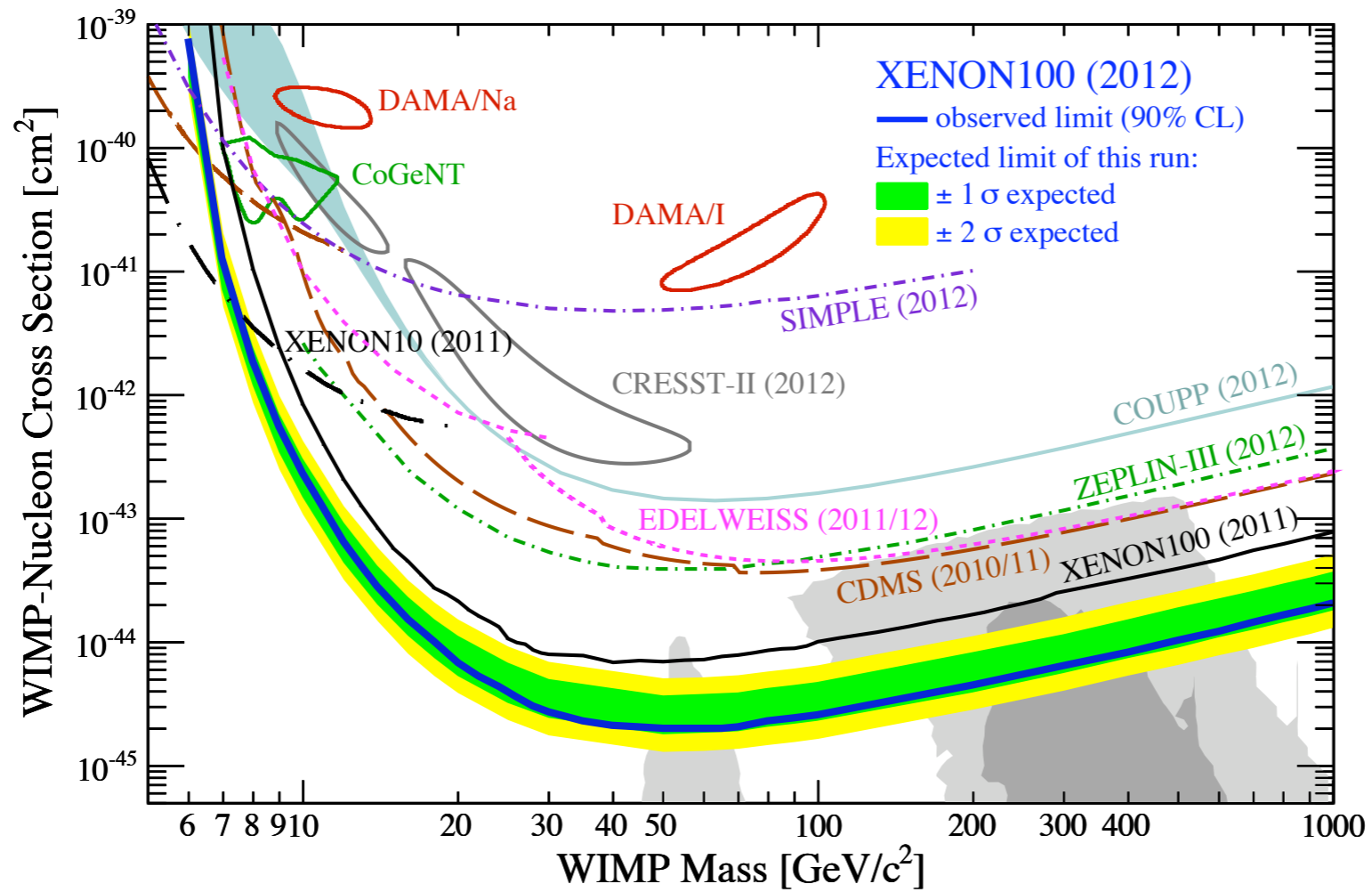
- **Example:** Spin-1/2 Dirac WIMP, some of the possible electron-DM couplings are

$$\begin{aligned} \mathcal{O}_V &= (\bar{\chi}\gamma_\mu\chi)(\bar{\ell}\gamma^\mu\ell), & \mathcal{O}_A &= (\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{\ell}\gamma^\mu\gamma^5\ell), \\ \mathcal{O}_S &= (\bar{\chi}\chi)(\bar{\ell}\ell), & \mathcal{O}_t &= (\bar{\chi}\ell)(\bar{\ell}\chi), \end{aligned}$$

- Parameterizes the effect of heavy particles mediating WIMP-DM interactions (e.g. t-channel selectrons in the MSSM), in a **model-independent** way
- Works if the scale  $\Lambda$  is above the energy scale of the experiment
- Does not require NR WIMPs - **broader** kinematic validity
- Applicable to more processes - e.g.  $q\chi$  elastic scattering (**direct detection!**)

# Direct Detection Status

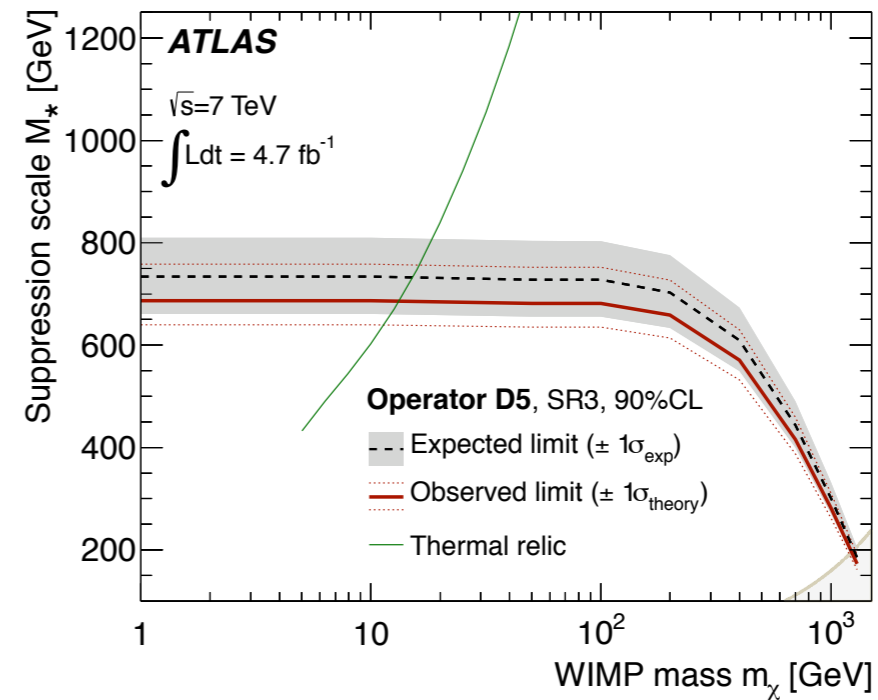
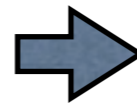
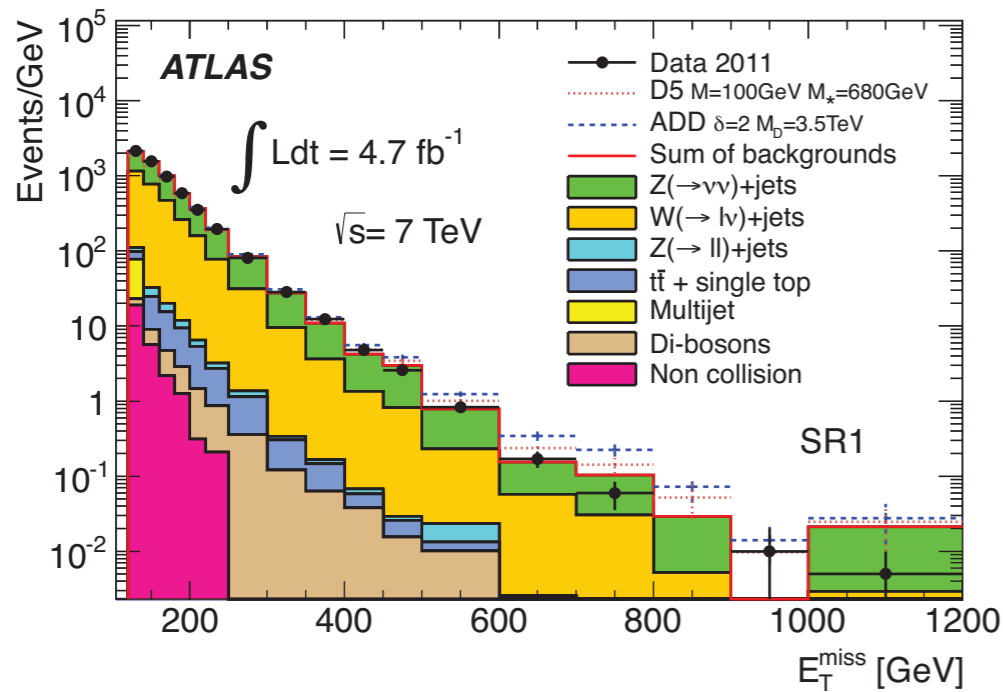
[XENON100, 1207.5988]



# LHC Limits

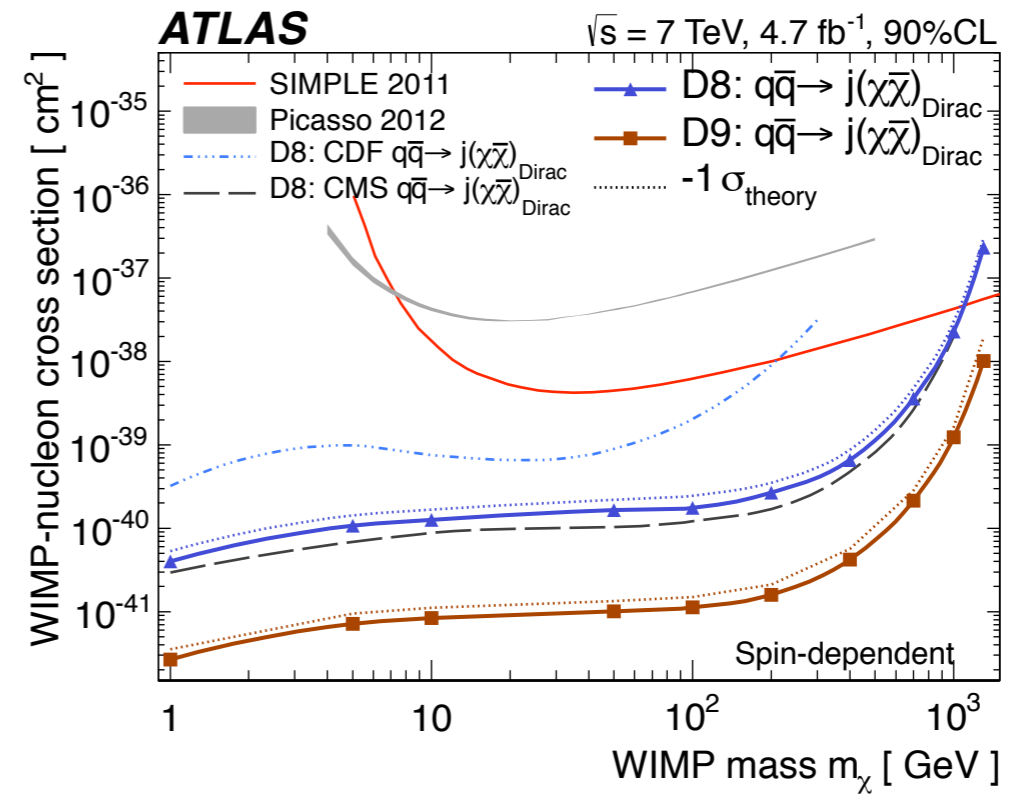
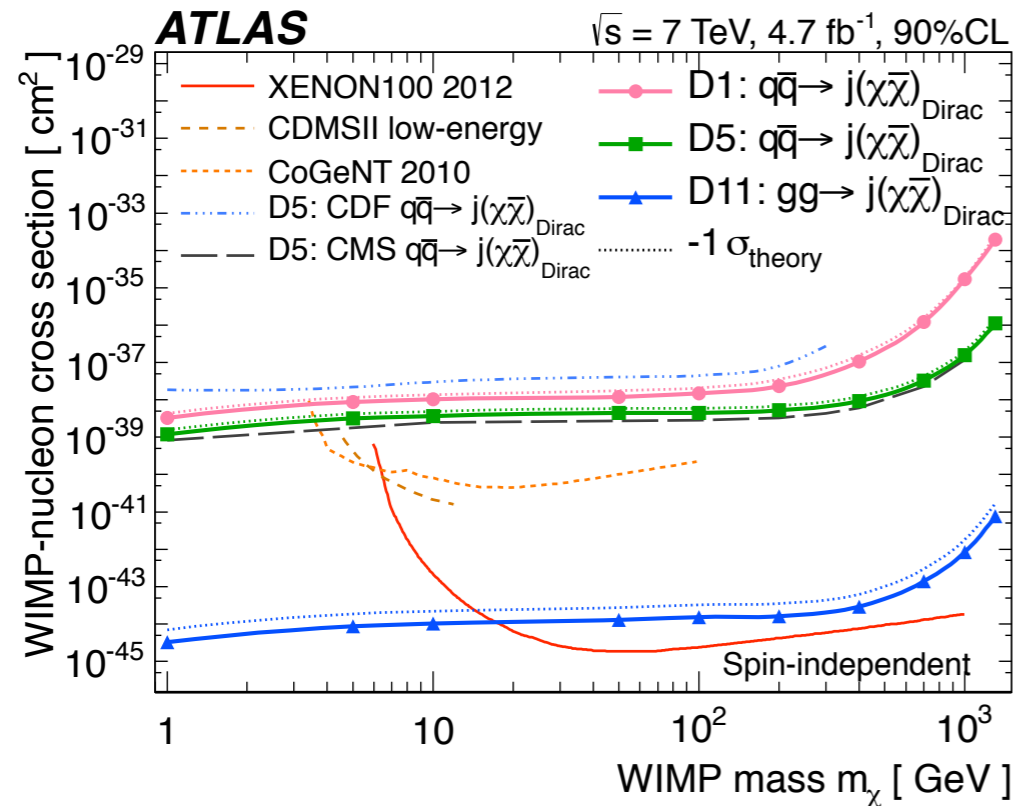
[ATLAS, 1210.4491; see also CMS, 1206.5663]

| Name | Initial state | Type         | Operator  |
|------|---------------|--------------|---|
| D1   | $qq$          | scalar       | $\frac{m_q}{M_*^3} \bar{\chi} \chi \bar{q} q$                                       |
| D5   | $qq$          | vector       | $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$                   |
| D8   | $qq$          | axial-vector | $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$ |
| D9   | $qq$          | tensor       | $\frac{1}{M_*^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$         |
| D11  | $gg$          | scalar       | $\frac{1}{4M_*^3} \bar{\chi} \chi \alpha_s (G_{\mu\nu}^a)^2$                        |



# LHC vs. Direct Detection

[ATLAS, 1210.4491; see also CMS, 1206.5663]



Collider Searches are more sensitive in two regimes:

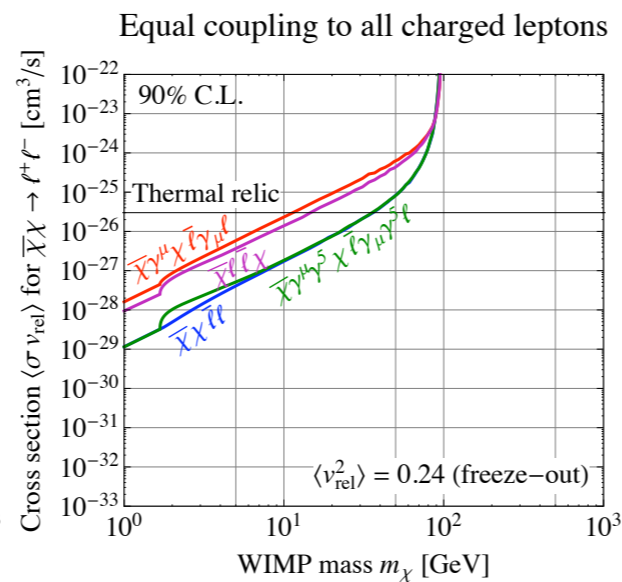
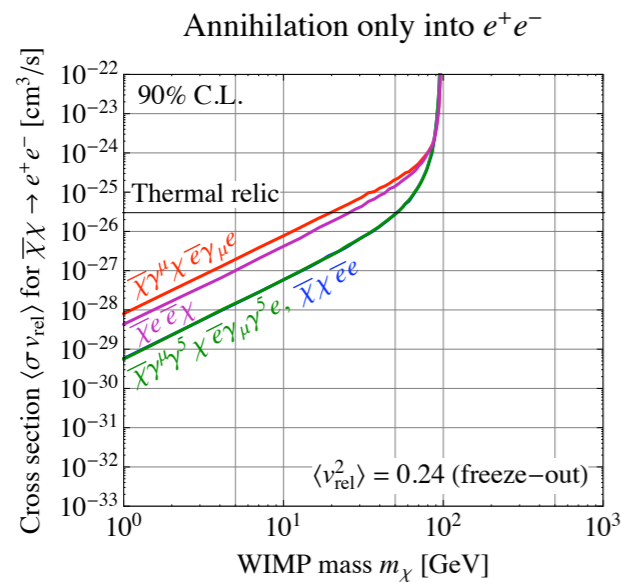
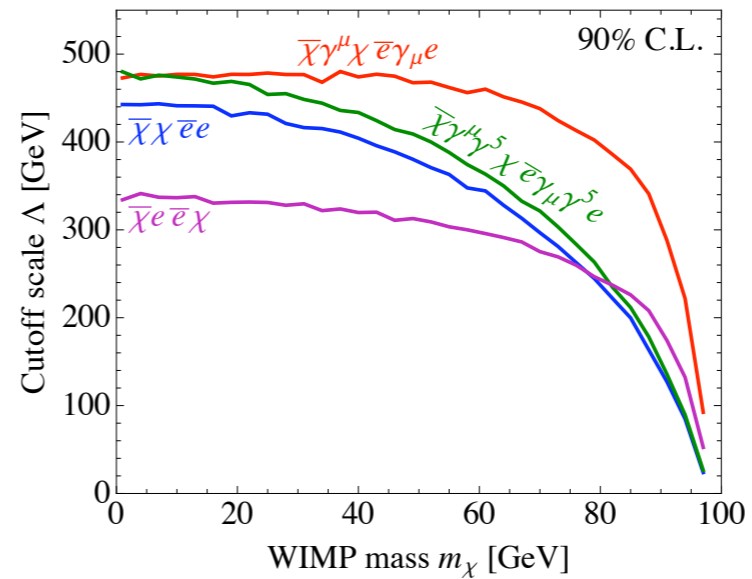
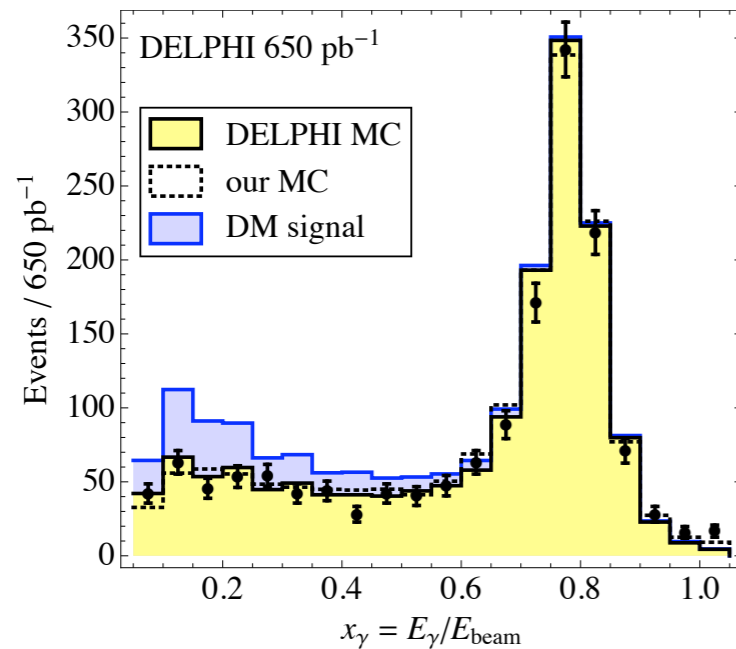
Low WIMP mass ( $< 10 \text{ GeV}$ )

Coupling via Spin-Dep. operators



# LEP-2 Limits

[Fox, Harnik, Kopp, Tsai, I 103.0240]



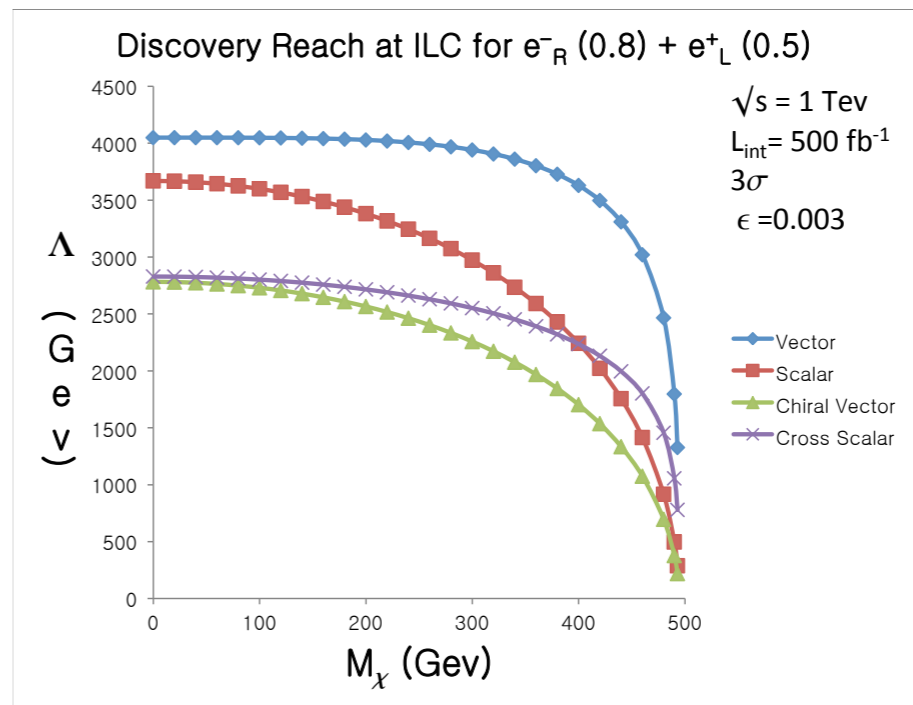
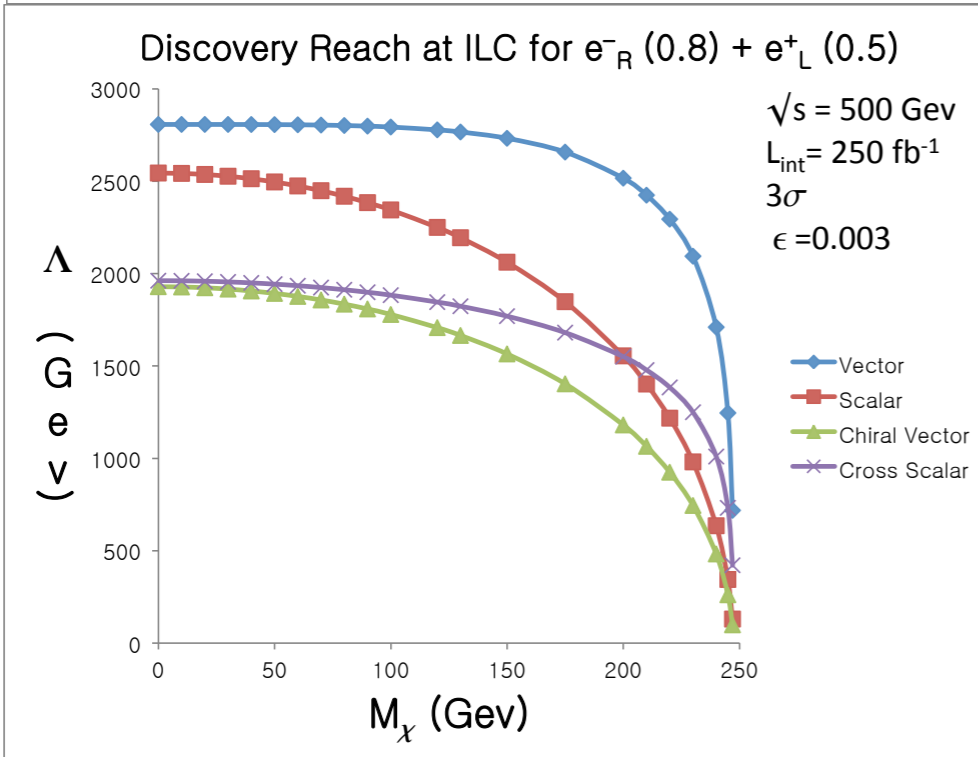
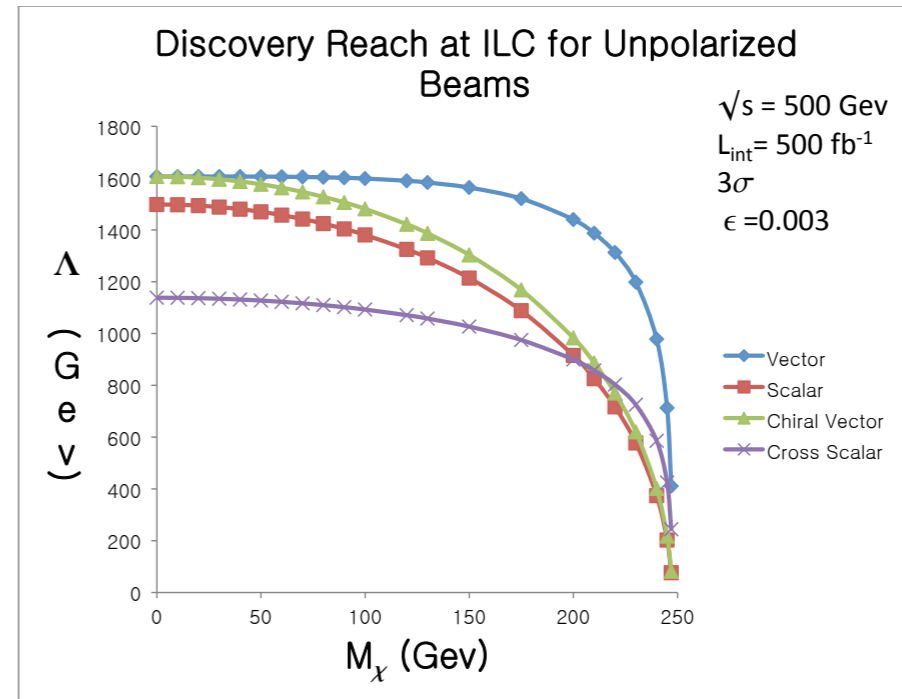
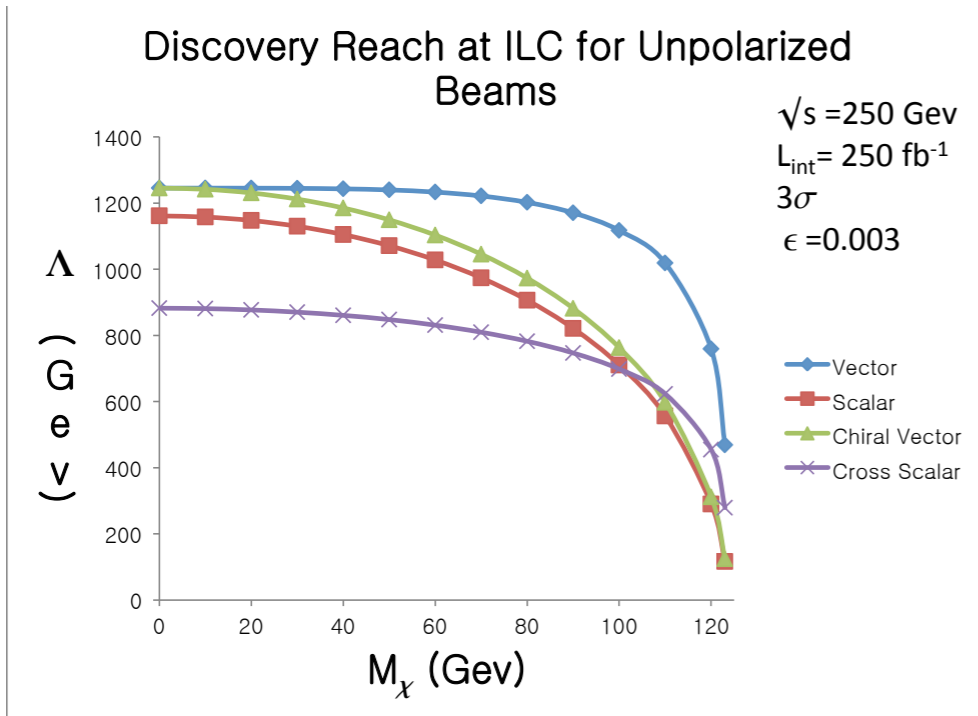
Leptophilic, thermal relic WIMP ruled out below 10-50 GeV dep. on assumptions

# Expected ILC Limits

[Yoonseok Chae, MP, to appear]

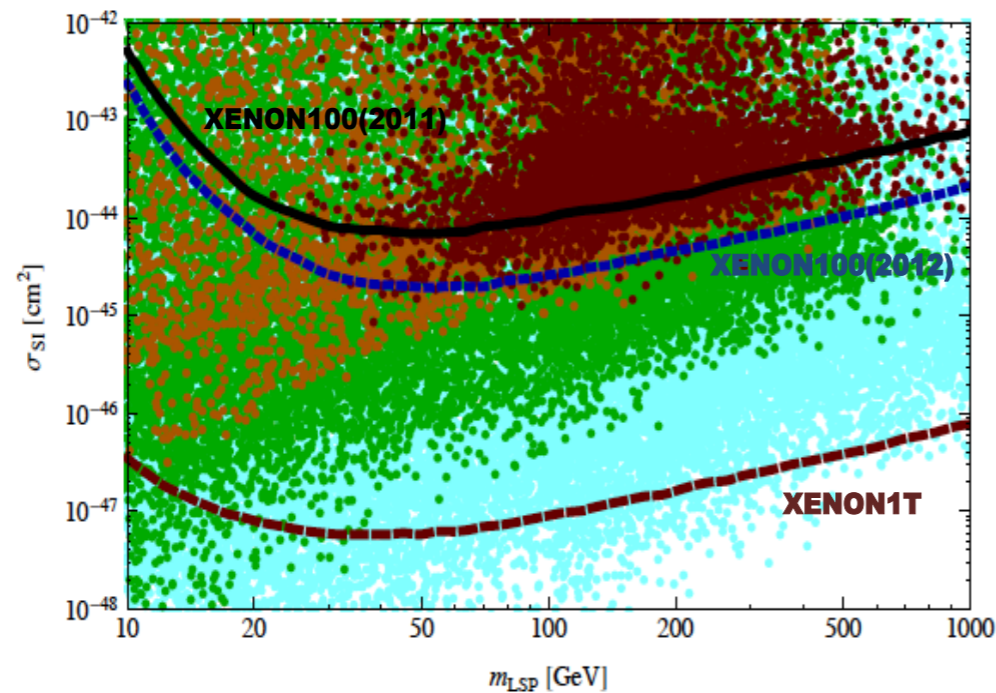
P  
R  
E  
L  
I  
M  
I  
N  
A  
R  
Y  
!

P  
R  
E  
L  
I  
M  
I  
N  
A  
R  
Y  
!

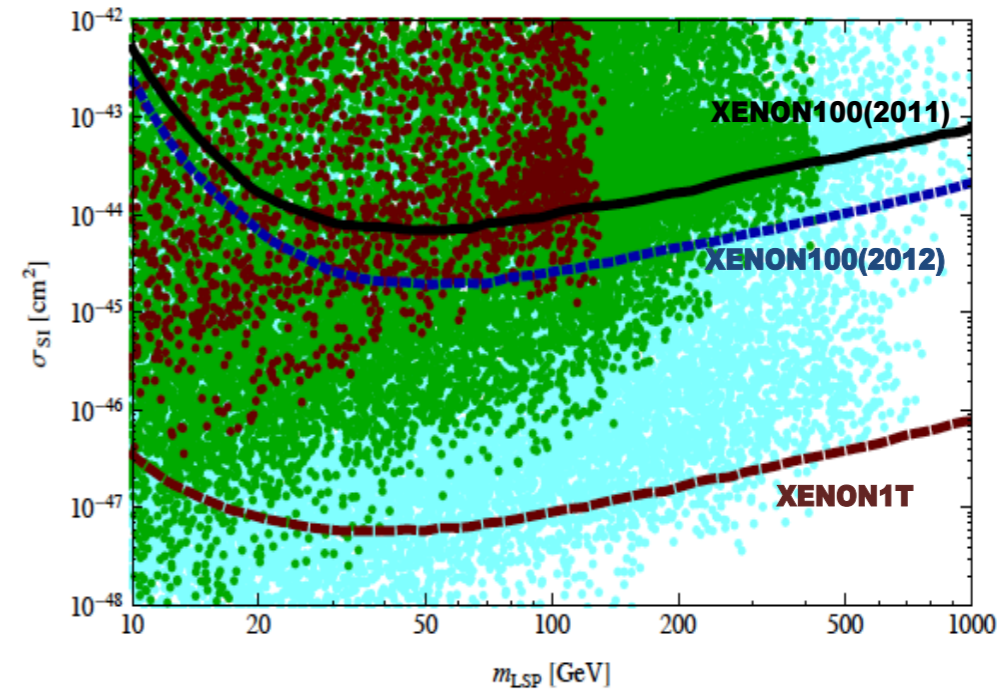


# Direct Detection/Tuning in (N)MSSM

[Shakya, MP, I 107.5048; I 208.0833]



“Purity”: fraction of the subdominant (gaugino or Higgsino) component in the LSP



Fine-tuning in EWSB (tree-level)

Tension is **already developing** in (N)MSSM from null result of direct detection searches!