

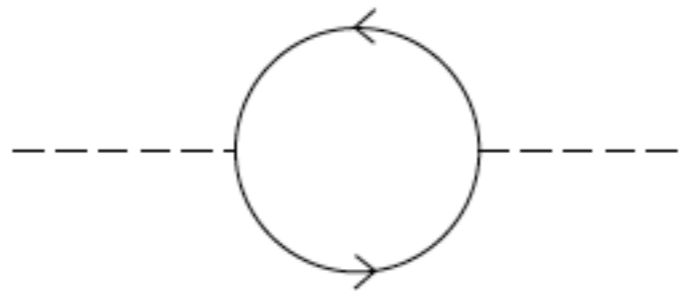
Why construction of ILC must begin immediately

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Perspective from the SUSY BSM group
reporting to Snowmass 2013

- Discovery of Higgs $h(125)$ by LHC completes the Standard Model!
- BUT....

- Problem #1: quantum corrections to $m(h)^2$ are quadratically divergent!



$$\delta m_h^2 = -\frac{3y_t^2}{8\pi^2} \Lambda^2$$

Supersymmetry provides solution by canceling to all orders via superpartners

Story for past 32 years: need weak scale SUSY for realistic model

- Problem #2: strong CP
- why 3 not 4 light pions?
- t'Hooft solution: QCD theta vacuum:

$$\mathcal{L} \ni \bar{\theta} \frac{g_s^2}{32\pi^2} F_A^{\mu\nu} \tilde{F}_{A\mu\nu}$$

- But: nEDM $\Rightarrow \bar{\theta} < 10^{-10}$
- Solution: axion

MSSM: simplest SUSY realization

- gauge coupling unification
- $m(t) \sim 173$ GeV: radiative EWSB
- EWPO (heavy SUSY)
- $m(h) \sim 115-135$ GeV

Conundrum: no sign of SUSY at LHC in spite of theorists telling us for 30 years that $m(\text{sparticle}) \sim m(\text{weak}) \sim m(Z) \sim m(h)$!

What is basis for this statement?

Scalar potential minimization in MSSM:

$$\frac{M_Z^2}{2} = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2$$

EWFT condition:

no large uncorrelated contributions to $m(Z)$!

- $\mu^2 \sim m_Z^2/2$
- $-m_{H_u}^2 \sim m_Z^2/2$
- $\Sigma(\tilde{t}_{1,2}) \sim m_Z^2/2$

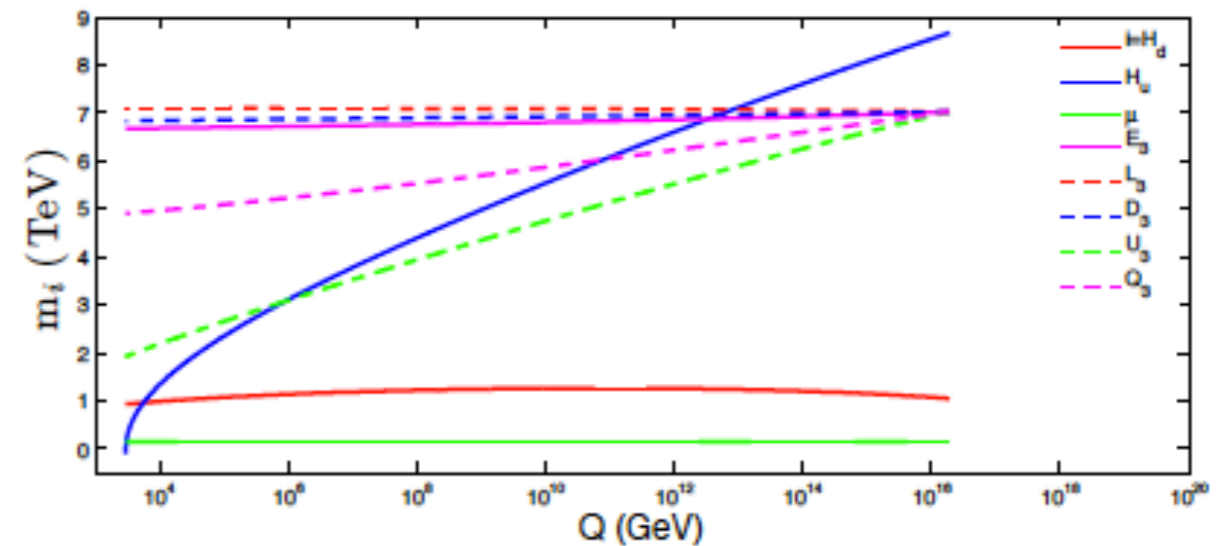
$$\Delta_{EW} \equiv \max_i (C_i) / (M_Z^2/2)$$

- $\mu \sim 100\text{-}200 \text{ GeV}$
(FP region or NUHM models)

$$\mu \sim \lambda m_{3/2}$$

- m_{Hu}^2 doesn't run too much negative

NUHM model



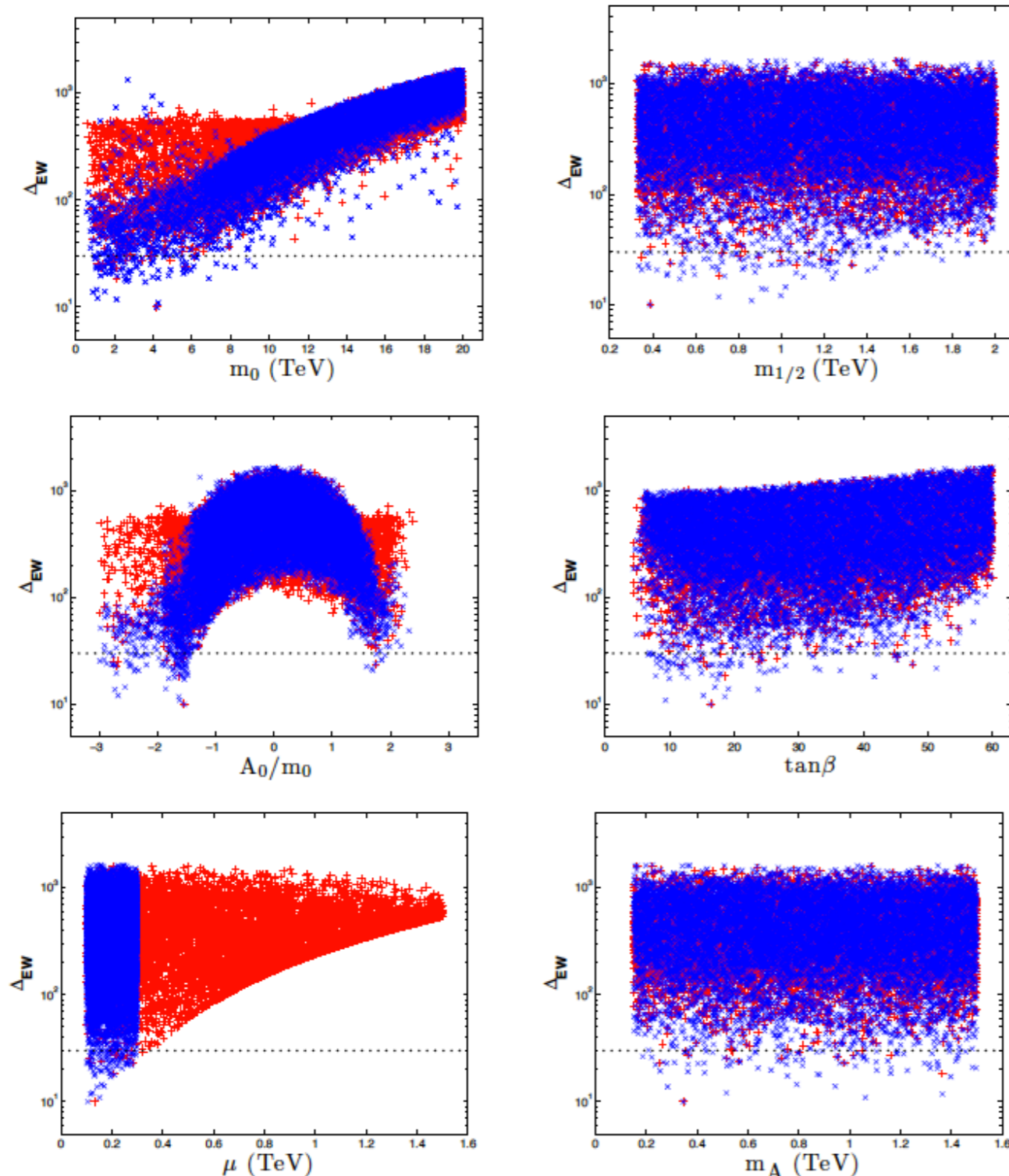
- radiative corrections not too big (no FP)

$$F(m^2) = m^2 (\log(m^2/Q^2) - 1)$$

$$\Sigma_u^u(\tilde{t}_{1,2}) = \frac{3}{16\pi^2} F(m_{\tilde{t}_{1,2}}^2) \times \left[f_t^2 - g_Z^2 \mp \frac{f_t^2 A_t^2 - 8g_Z^2 (\frac{1}{4} - \frac{2}{3}x_W) \Delta_t}{m_{\tilde{t}_2}^2 - m_{\tilde{t}_1}^2} \right]$$

large A_t softens radiative corr'ns while lifting $m(h)$!

Which parameter choices lead to low EWFT and how low can Δ_{EW} be?



$\Delta_{EW} \sim 10$ or 10% *EWFT*

High-scale models with low Δ_{EW} :

Radiatively-driven natural SUSY, or RNS

HB, Barger, Huang, Mickelson, Mustafayev, Tata,
arXiv:1212.2655

Compare RNS to mSUGRA for similar parameters

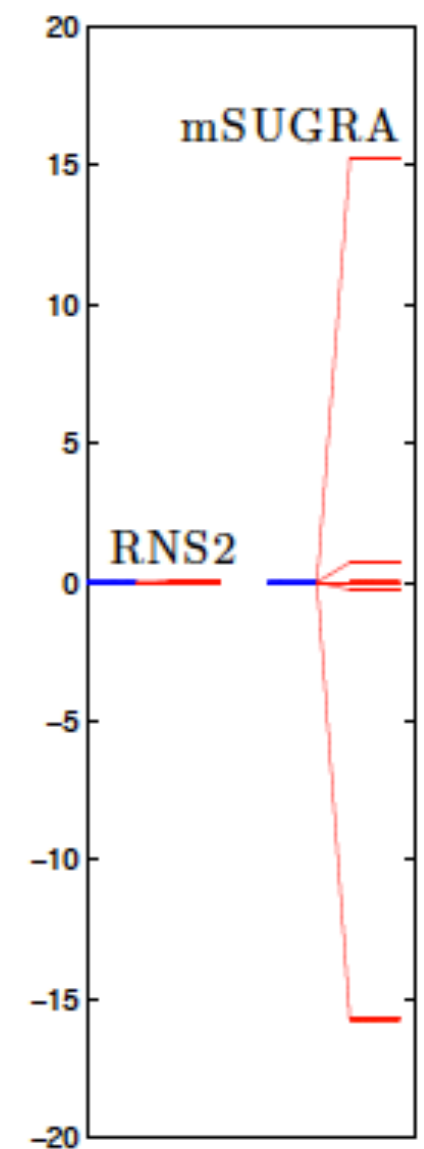
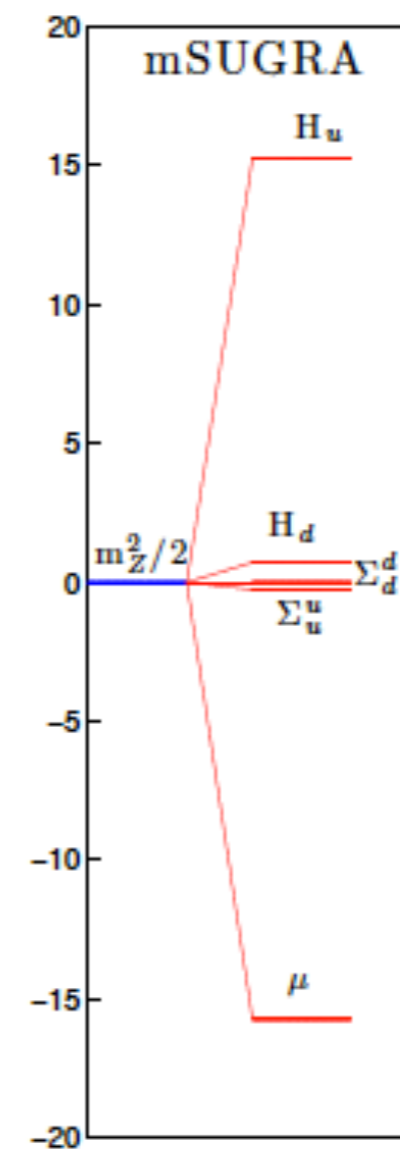
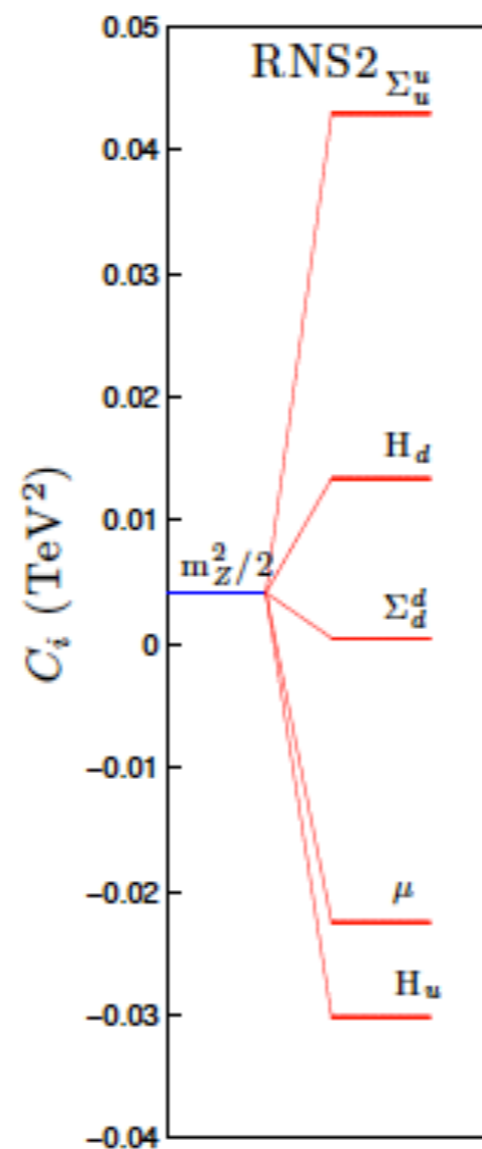
$m_0 = 7025 \text{ GeV}$, $m_{1/2} = 568.3 \text{ GeV}$, $A_0 = -11426.6 \text{ GeV}$, $\tan\beta = 8.55$ with $\mu = 150 \text{ GeV}$ and $m_A = 1000 \text{ GeV}$

RNS

- $C_{\Sigma_u^u} \sim (205 \text{ GeV})^2$
- $C_{H_d} \sim (114 \text{ GeV})^2$
- $C_{\Sigma_d^d} \sim (22 \text{ GeV})^2$
- $C_\mu \sim -(148 \text{ GeV})^2$
- $C_{H_u} \sim -(173 \text{ GeV})^2$
- $m_Z^2/2 \simeq (65 \text{ GeV})^2$

mSUGRA

- $C_{H_u} \simeq (3.87 \text{ TeV})^2$
- $C_\mu \simeq -(3.93 \text{ TeV})^2$



SUSY spectra from radiatively-driven natural SUSY (RNS)

scan NUHM2 space:

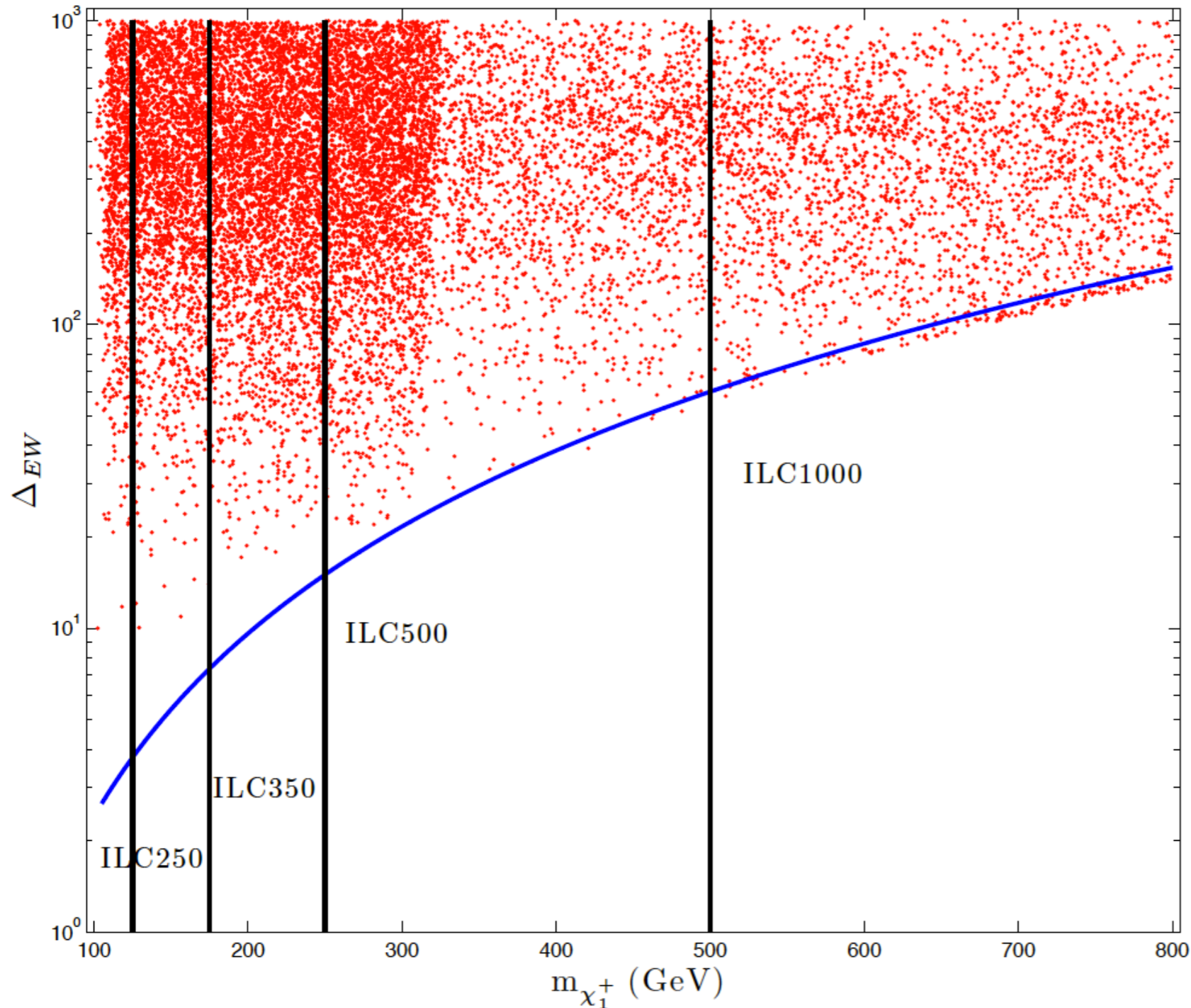
- light higgsino-like \tilde{W}_1 and $\tilde{Z}_{1,2}$ with mass $\sim 100 - 300$ GeV,
- gluinos with mass $m_{\tilde{g}} \sim 1 - 4$ TeV,
- heavier top squarks than generic NS models: $m_{\tilde{t}_1} \sim 1 - 2$ TeV and $m_{\tilde{t}_2} \sim 2 - 5$ TeV,
- first/second generation squarks and sleptons with mass $m_{\tilde{q},\tilde{\ell}} \sim 1 - 8$ TeV. The $m_{\tilde{\ell}}$ range can be pushed up to 20-30 TeV if non-universality of generations with $m_0(1,2) > m_0(3)$ is allowed.

Very hard to see even
at LHC14 when
 $m(\tilde{g}) > 2$ TeV

parameter	RNS1	RNS2	NS2
$m_0(1,2)$	10000	7025.0	19542.2
$m_0(3)$	5000	7025.0	2430.6
$m_{1/2}$	700	568.3	1549.3
A_0	-7300	-11426.6	873.2
$\tan \beta$	10	8.55	22.1
μ	150	150	150
m_A	1000	1000	1652.7
$m_{\tilde{g}}$	1859.0	1562.8	3696.8
$m_{\tilde{u}_L}$	10050.9	7020.9	19736.2
$m_{\tilde{u}_R}$	10141.6	7256.2	19762.6
$m_{\tilde{e}_R}$	9909.9	6755.4	19537.2
$m_{\tilde{t}_1}$	1415.9	1843.4	572.0
$m_{\tilde{t}_2}$	3424.8	4921.4	715.4
$m_{\tilde{b}_1}$	3450.1	4962.6	497.3
$m_{\tilde{b}_2}$	4823.6	6914.9	1723.8
$m_{\tilde{\tau}_1}$	4737.5	6679.4	2084.7
$m_{\tilde{\tau}_2}$	5020.7	7116.9	2189.1
$m_{\tilde{\nu}_\tau}$	5000.1	7128.3	2061.8
$m_{\tilde{W}_2}$	621.3	513.9	1341.2
$m_{\tilde{W}_1}$	154.2	152.7	156.1
$m_{\tilde{Z}_4}$	631.2	525.2	1340.4
$m_{\tilde{Z}_3}$	323.3	268.8	698.8
$m_{\tilde{Z}_2}$	158.5	159.2	156.2
$m_{\tilde{Z}_1}$	140.0	135.4	149.2
m_h	123.7	125.0	121.1
$\Omega_{\tilde{Z}_1}^{std} h^2$	0.009	0.01	0.006
$BF(b \rightarrow s\gamma) \times 10^4$	3.3	3.3	3.6
$BF(B_s \rightarrow \mu^+\mu^-) \times 10^9$	3.8	3.8	4.0
$\sigma^{SI}(\tilde{Z}_1 p)$ (pb)	1.1×10^{-8}	1.7×10^{-8}	1.8×10^{-9}
Δ	9.7	11.5	23.7

Smoking gun signature: 4 light higgsinos at ILC!

$$e^+e^- \rightarrow \tilde{W}_1^+ \tilde{W}_1^-, \tilde{Z}_1 \tilde{Z}_2$$



$$m_{\tilde{W}_1^\pm}, m_{\tilde{Z}_{1,2}}$$

$$\sqrt{s} \sim \sqrt{2\Delta_{EW}m_Z}$$

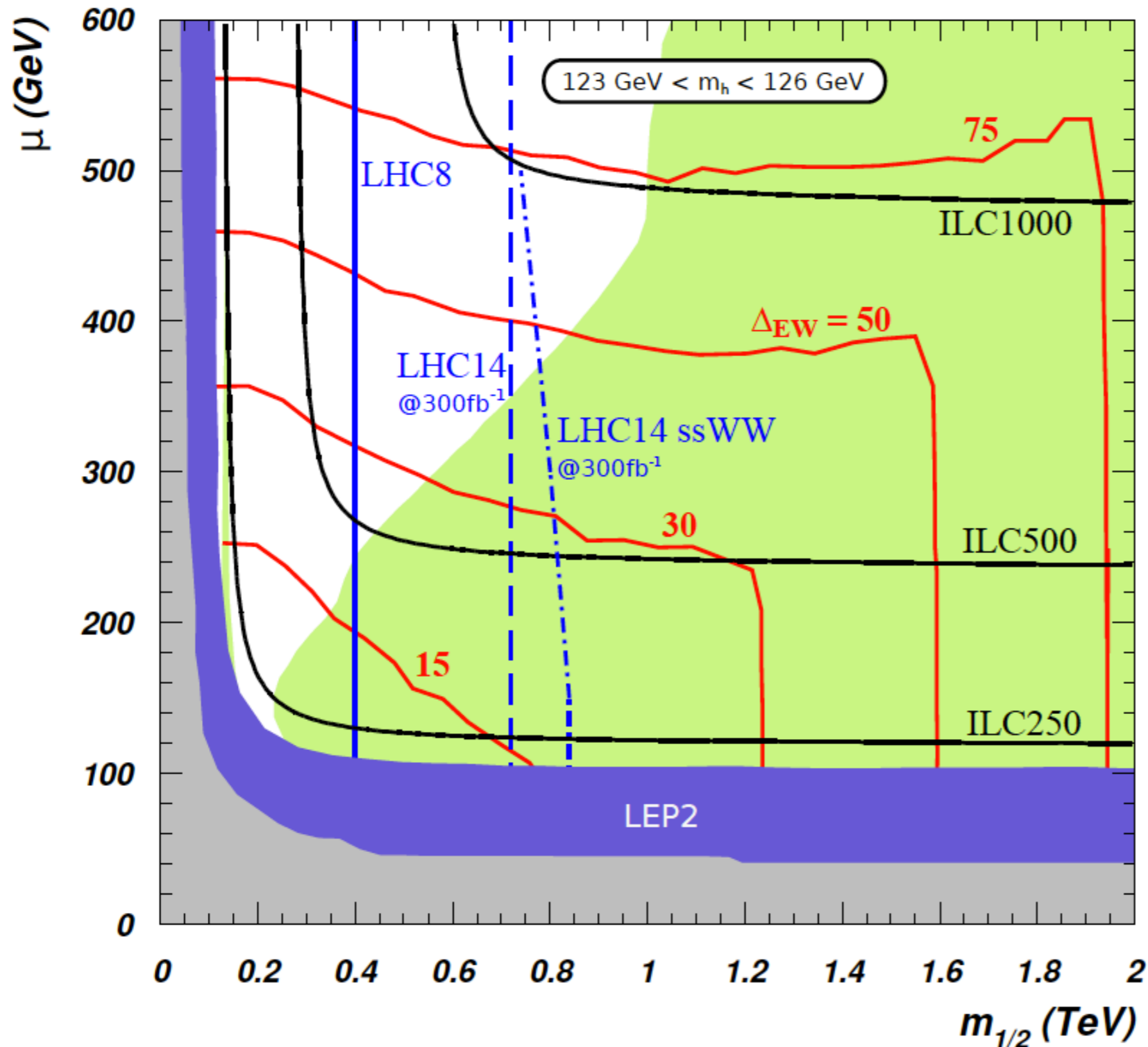
ILC/CLIC have capability to measure SUSY parameters and actually reconstruct

$$\Delta_{EW}$$

measure and check if nature is EWFT'd?

LHC/ILC complementarity

NUHM2: $m_0=5\text{ TeV}$, $\tan\beta=15$, $A_0=-1.6m_0$, $m_A=1\text{ TeV}$, $m_t=173.2\text{ GeV}$



While LHC has some capacity, it will require ILC to draw the story of SUSY electroweak naturalness to a conclusion!

A. Mustafayev plot

- Thus, we see that the only SUSY particles guaranteed $\sim m(\text{weak})$ are higgsinos:
- $m(Z) \sim m(\text{higgs}) \sim m(\text{higgsino})$
- ILC will be higgsino factory in addition to Higgs factory
- $m(\chi^+) / m(\chi^0)$ to ~ 2 GeV;
 $\Delta M(\chi^+, \chi^0)$ to ~ 5 MeV;
 pol. x-sections to 1-2% ; detailed talks: see
 afternoon BSM session: Sert, Rolbiecki
- Either ILC w/ $\sqrt{s} \sim 500\text{-}600$ GeV
 discovers light higgsinos or the idea of
 SUSY weak scale naturalness is dead!

For the lower range of $m(\text{gluino}) \sim < 2 \text{ TeV}$, LHC may discover! But as colored sparticle masses increase, then cross sections lower: precision studies limited by statistics

At the ILC, SUSY is really simplified:
can cover all NSLP-LSP
combinations for arbitrary mass differences
(meta-stable, kinks, prompt,...) like at LEP

Other featured studies:

- * general ewkino scan joint with LHC:

Higgsino, Wino or Bino LSP, scan the other two parameters ->
plot 5sigma discovery range / 2sigma
exclusion range for LHC (Delphes) / ILC (SGV)

- * a complexified SUSY model:

simplified models can give very misleading messages at LHC ->
study a very rich spectrum with many concurring decay chains:
STC benchmark (aka TDR4) both from LHC
(Delphes) and ILC (full sim / SGV)

- * connection to intensity frontier:

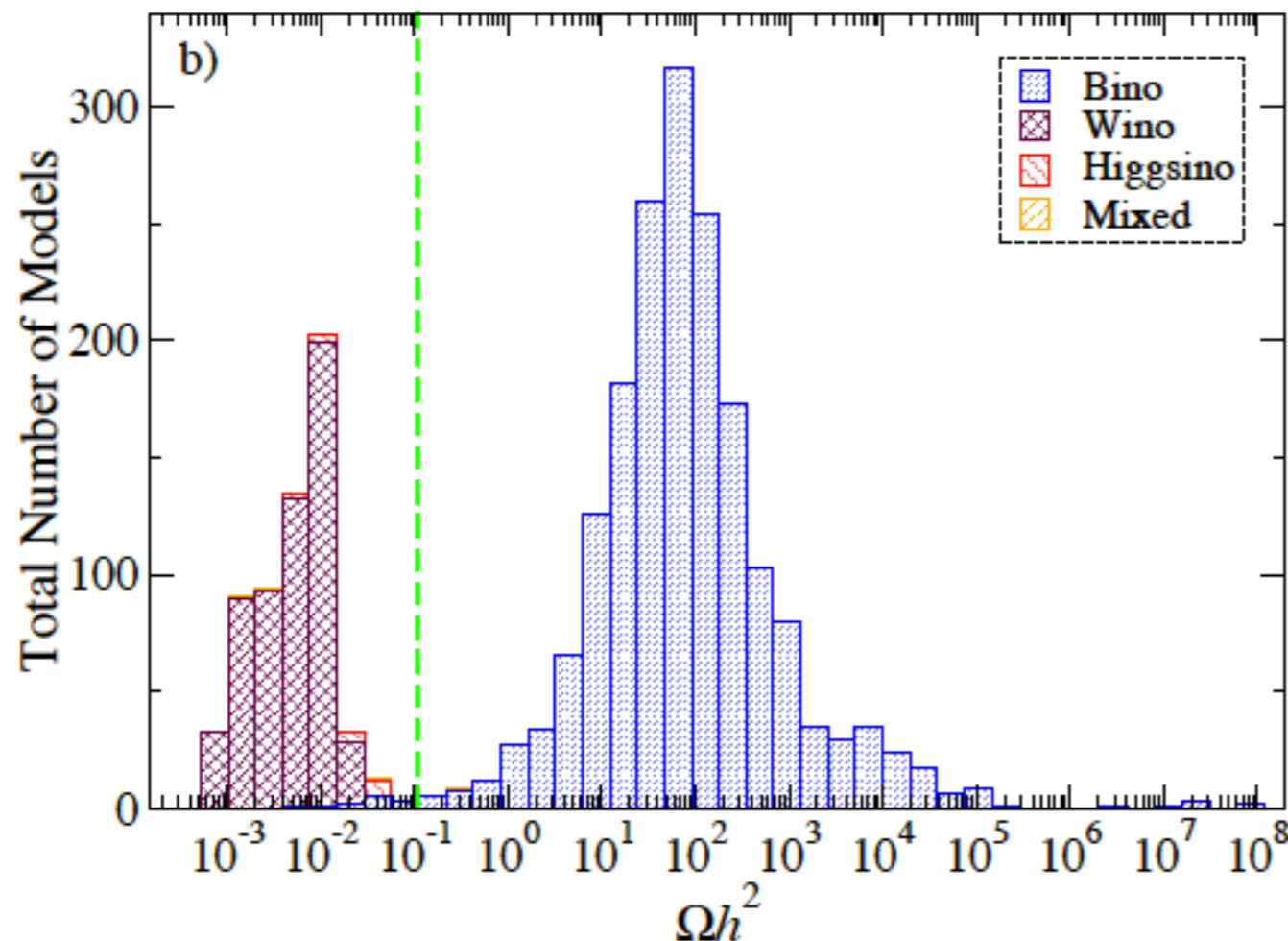
bRPV: can do neutrino physics at colliders:
cf talks by LHC and ILC (Vormwald) later

What about DM in RNS?

I heard higgsino-like wimp isn't a good DM candidate?

Lightest neutralino all by itself in general not good DM candidate: too much or too little CDM

Scan over 19 parameters:



HB, Box, Summy
JHEP1010(2010)023

Invoke Peccei-Quinn sol'n to strong CP problem with SUSY

PQMSSM: Axions + SUSY \Rightarrow mixed $a - LSP$ dark matter

- $\hat{a} = \frac{s+ia}{\sqrt{2}} + i\sqrt{2}\bar{\theta}\tilde{a}_L + i\bar{\theta}\theta_L\mathcal{F}_a$ in 4-comp. notation
- Raby, Nilles, Kim; Rajagopal, Wilczek, Turner
- axino is spin- $\frac{1}{2}$ element of axion supermultiplet (R -odd; possible LSP candidate)
- $m_{\tilde{a}}$ model dependent: keV \rightarrow TeV, but $\sim M_{SUSY}$ in gravity mediation
- saxion is spin-0 element: R -even but gets SUSY breaking mass ~ 1 TeV
- axion is usual QCD axion: gets produced via vacuum mis-alignment/coherent oscillations as usual
- additional PQ parameters: $(f_a, m_{\tilde{a}}, m_s, \theta_i, \theta_s,)$ and T_R

Coupled Boltzmann calculation of mixed axion-neutralino abundance

Bae, HB, Lessa, arXiv:1301.7428

Case for dominant $s \rightarrow aa$ decay:
contributes to dark radiation

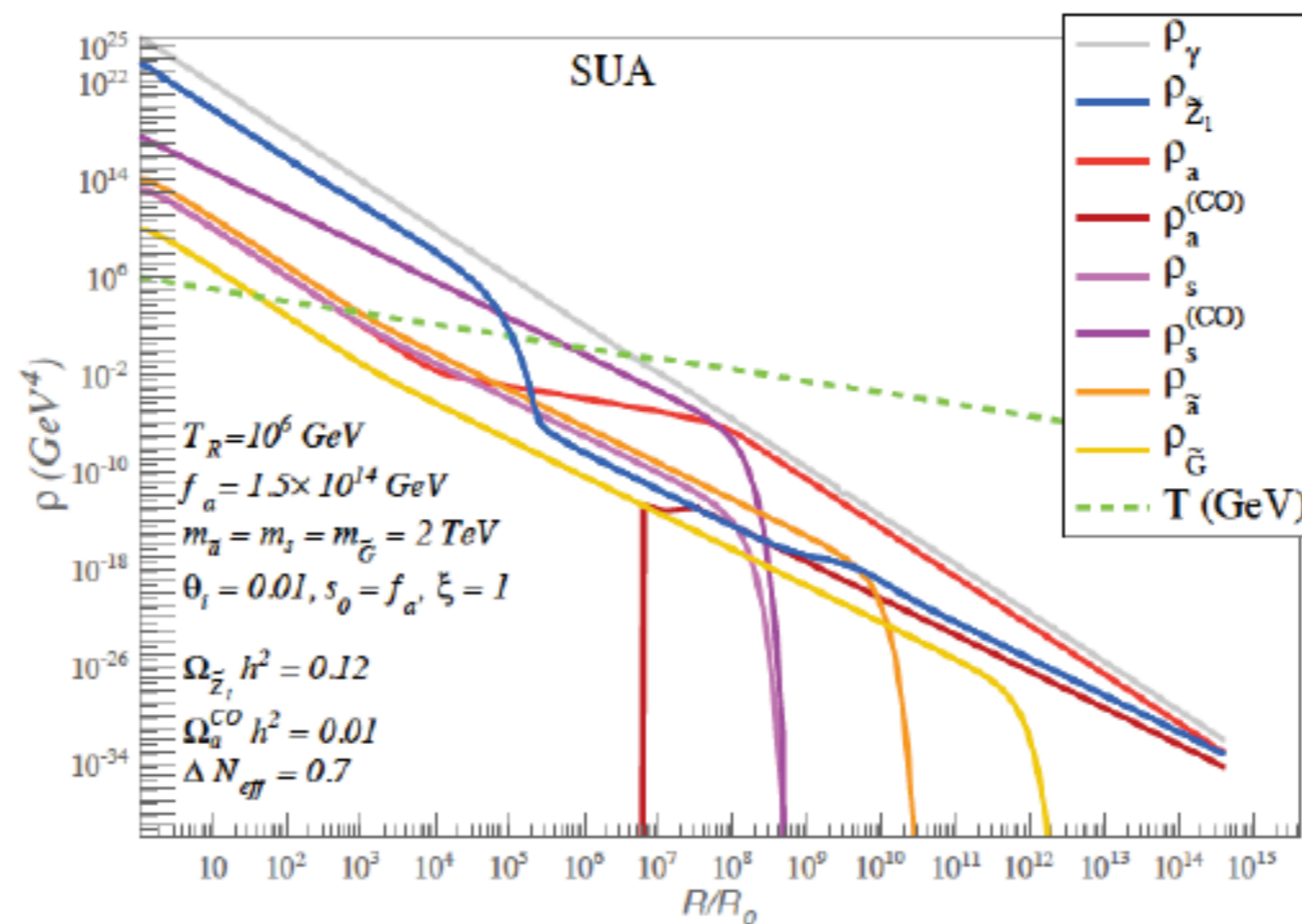


Figure 2: Evolution of various energy densities versus scale parameter R/R_0 for the SUA benchmark.

Conclusions

- Δ_{EW} is more robust measure of Little Hierarchy problem
- Why are $m(Z), m(h) \sim 100$ GeV while sparticle masses are $\gg 1$ TeV?
- $\mu \sim m(Z)$: light higgsinos (ILC!)
- $m(H_u)$ driven somewhat, not grossly, negative
- large mixing in stop sector

Under these conditions, the Little Hierarchy remains
but the “Problem” seems
to melt away and the old paradigm of
SUSY GUTs remains strong:

but with huge implications for collider/dark matter searches!

- The low lying sparticles (higgsinos) have severely compressed spectra: hard to see at LHC (but new signatures e.g. SS dibosons)
- No large cancellations in $m(Z)$, $m(h) \Rightarrow$ ILC is the right machine to build!
- Dark matter production more intricate than usual story: here, we suggest mixed axion-higgsino (co-dark-matter) particles: possibly detect both WIMPS and axions?

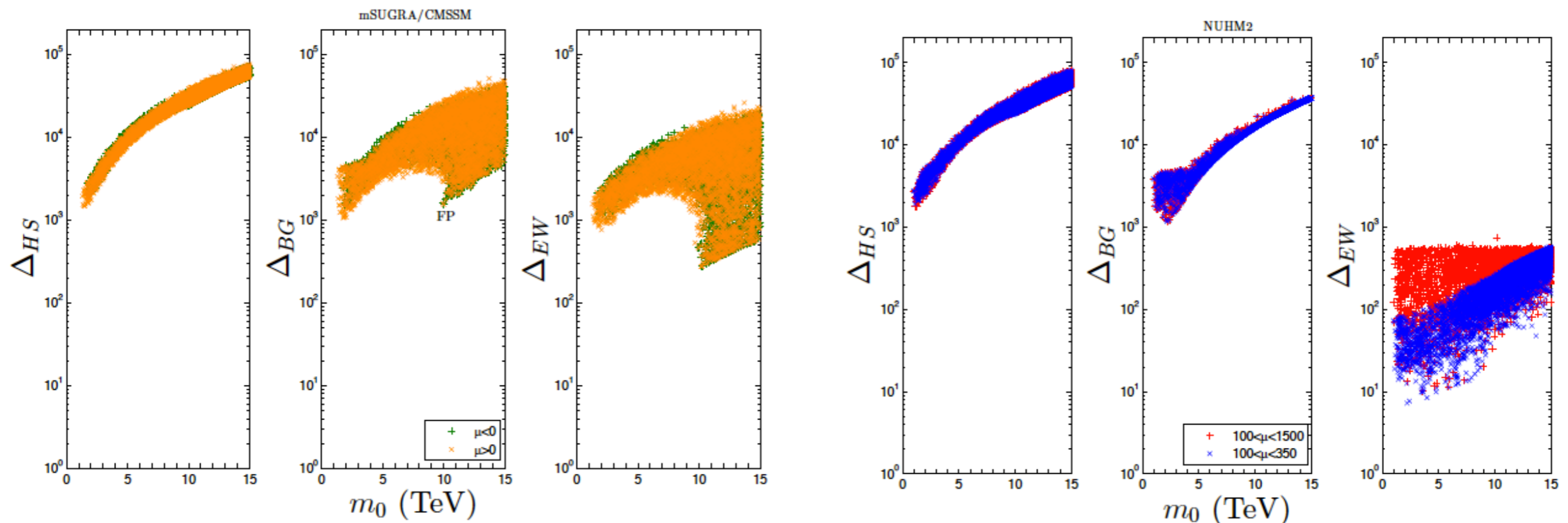
Backup:

$$\frac{m_Z^2}{2} = \frac{(m_{H_d}^2(\Lambda) + \delta m_{H_d}^2 + \Sigma_d^d) - (m_{H_u}^2(\Lambda) + \delta m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - (\mu^2(\Lambda) + \delta \mu^2)$$

$$\Delta_{HS} \equiv \max_i |B_i| / (m_Z^2/2)$$

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 - m_{H_u}^2 \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \mu^2$$

$$\Delta_{BG} \equiv \max_i |c_i| = \max_i \left| \frac{a_i}{m_Z^2} \frac{\partial m_Z^2}{\partial a_i} \right|$$



backup

