

# Little Higgs measurements at the ILC

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Y. Takubo, H. Yamamoto

**Little Higgs** is one of the possibility for TeV new physics.  
 In this talk, we study the phenomenology of **Littlest Higgs model with T-parity (LHT)** at ILC.

**SM**  $g$   $W$   $Z$   $A$ ,  $q$   $l$ ,  $H$

**Partner**  $W_H$   $Z_H$   $A_H$ ,  $q_H$   $l_H$ ,  $\Phi_H$

- SM particles & their partners have same spin.
- $M_{\text{(New particles)}} \sim f \sim 1 \text{ TeV}$ .



We estimate...

- Measurement accuracy of  $M_{WH}$ ,  $M_{ZH}$ ,  $M_{AH}$ ,  $M_{lH}$ .
- Determination of VEV  $f$  & other model parameters
- Determination of couplings with new particles @ ILC.

# Plan

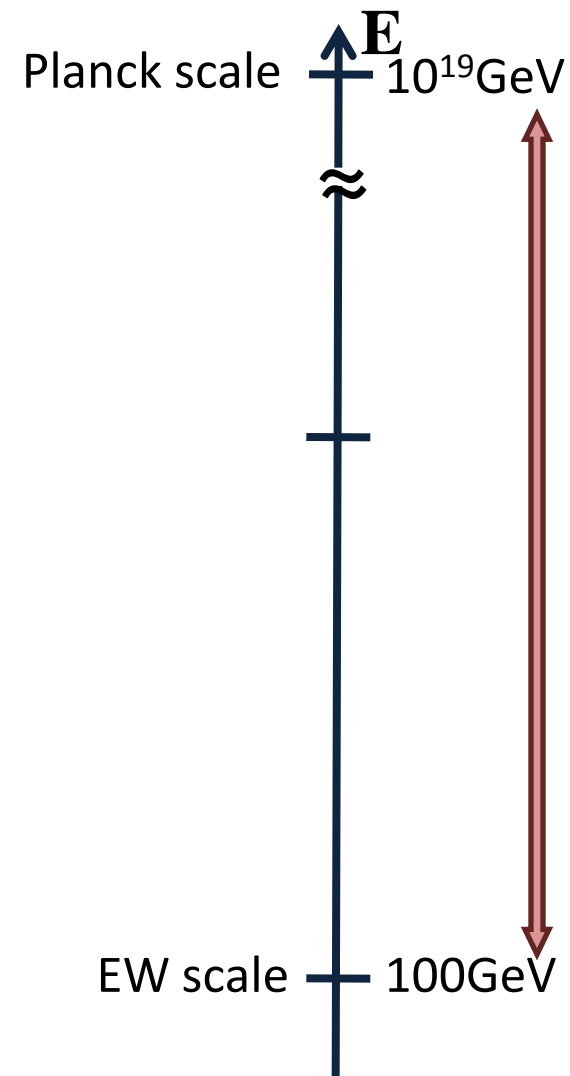
- Introduction
- What is the LHT?
- How to measure the LHT at ILC?
- Simulation results
- Summary and discussion

**SM is the successful model**  
describing physics below  $\sim 100$  GeV.

**But...**

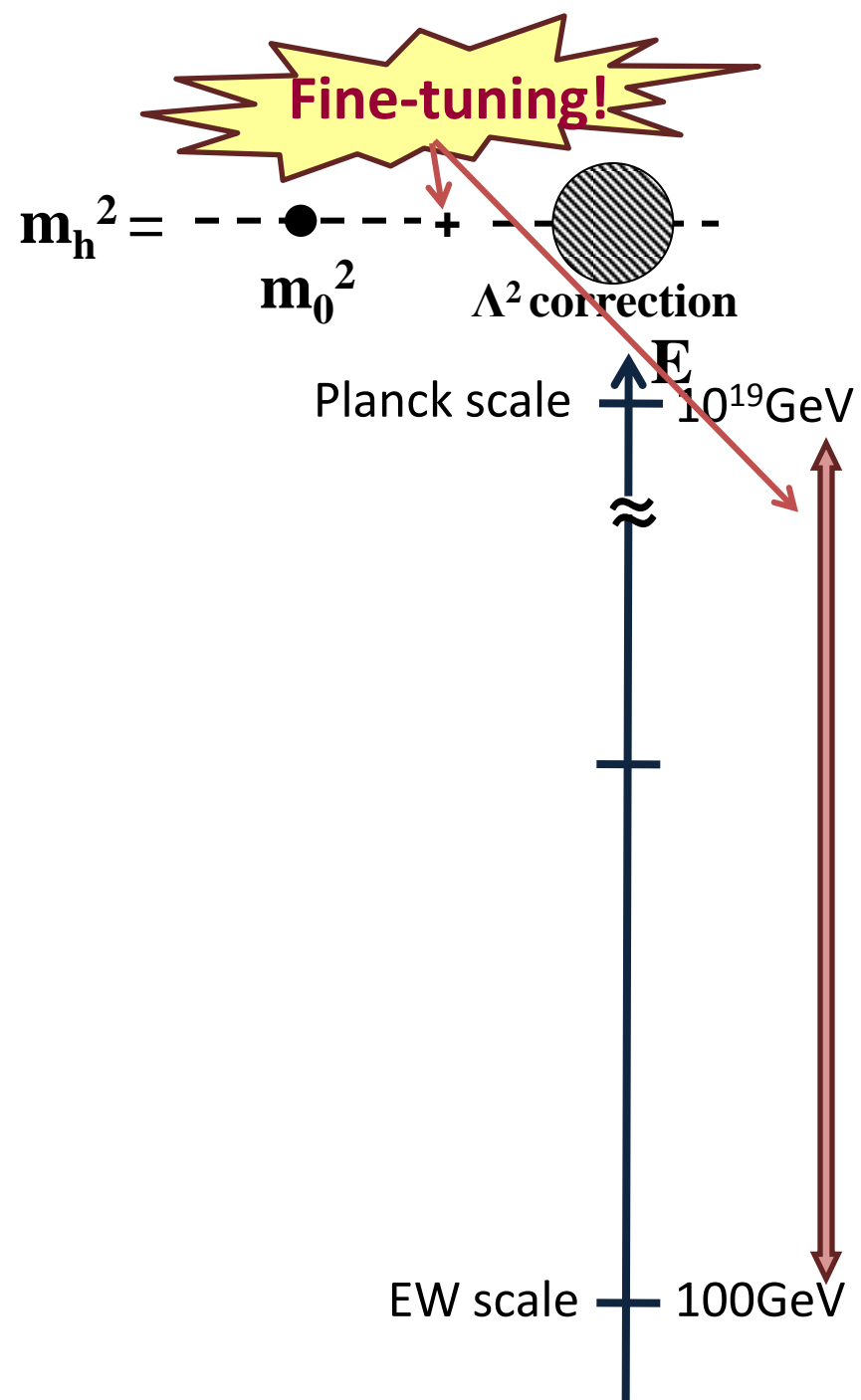
If we assume that SM is valid all the way up to the Planck scale, there is a **hierarchy problem**.

$$\Rightarrow m_h^2 \sim m_0^2 + 3 y_t^2 \Lambda^2 / 4\pi + \dots$$



# • Hierarchy Problem

related to quadratic divergence to the Higgs boson mass term.

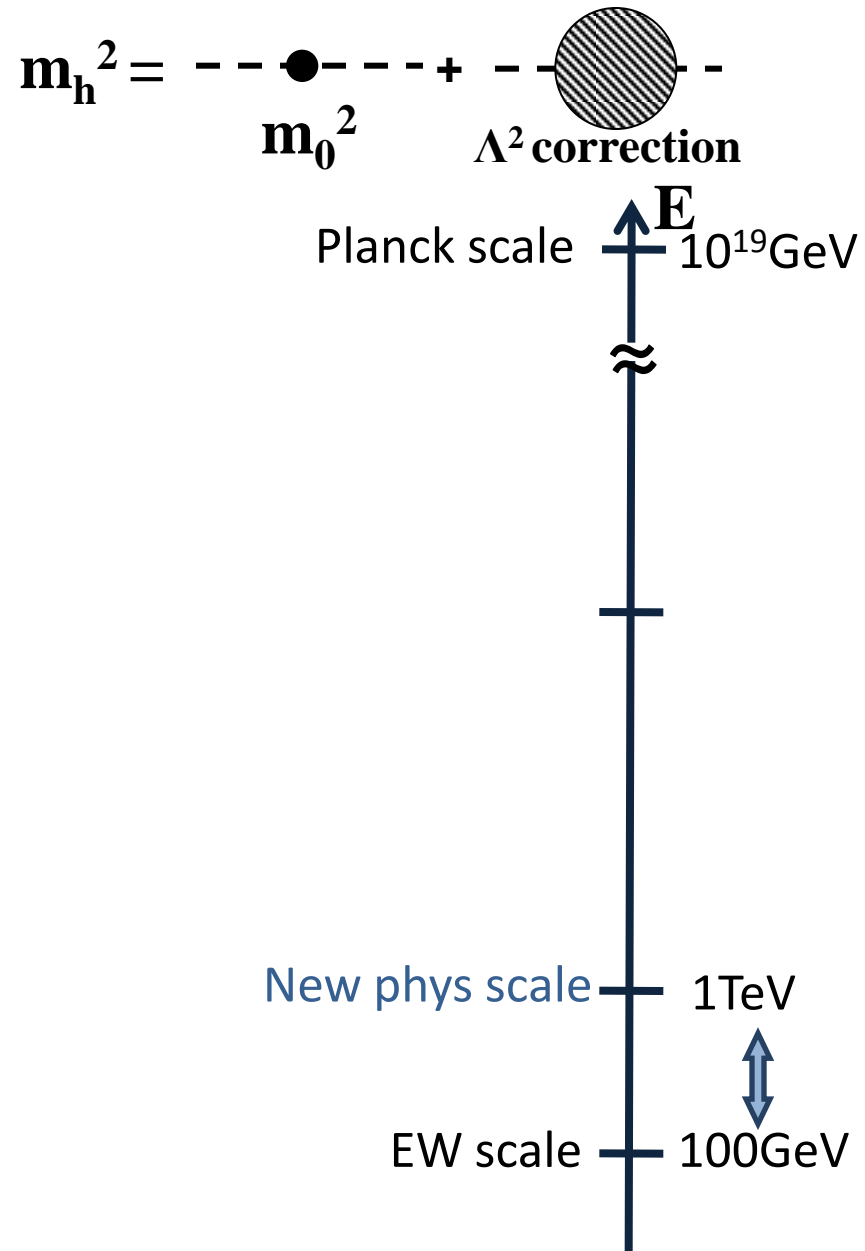


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related to quadratic divergence to the Higgs boson mass term.

↓ solution

- low energy cut off scenario  
If  $\Lambda \sim 1 \text{ TeV}$ ,  
there are no fine-tuning!



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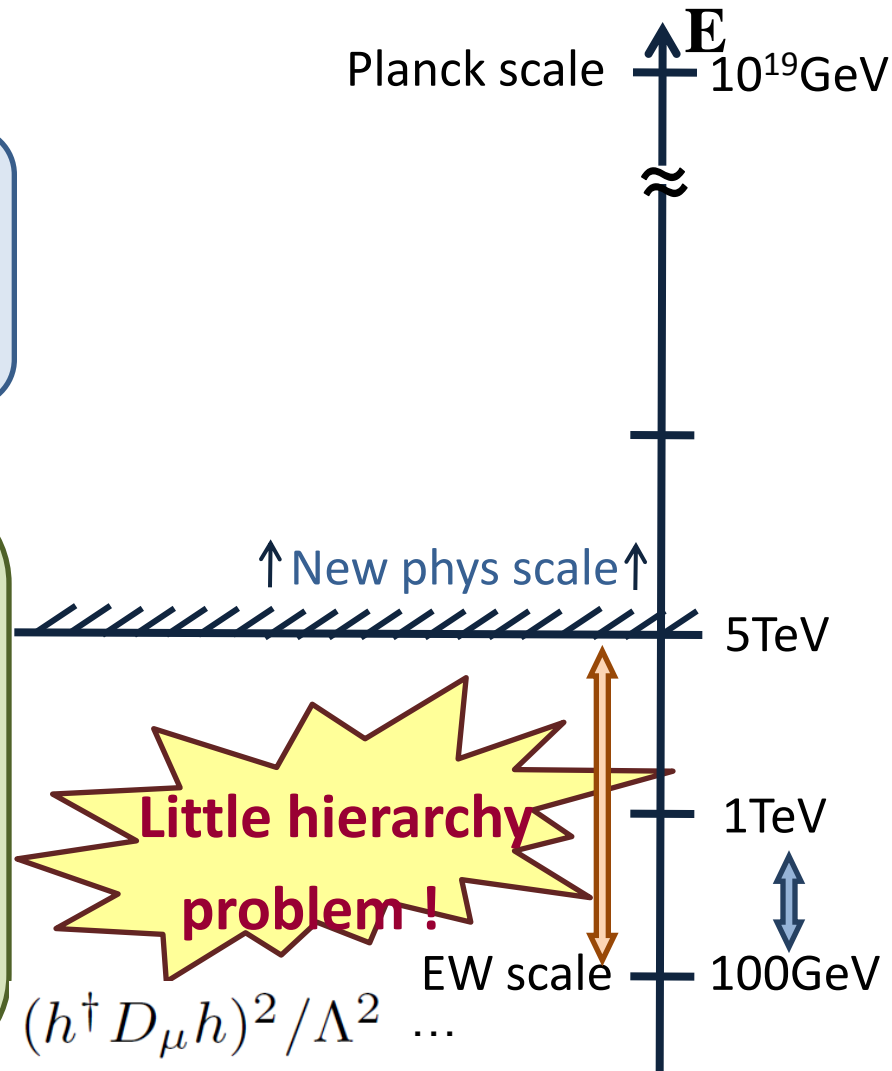
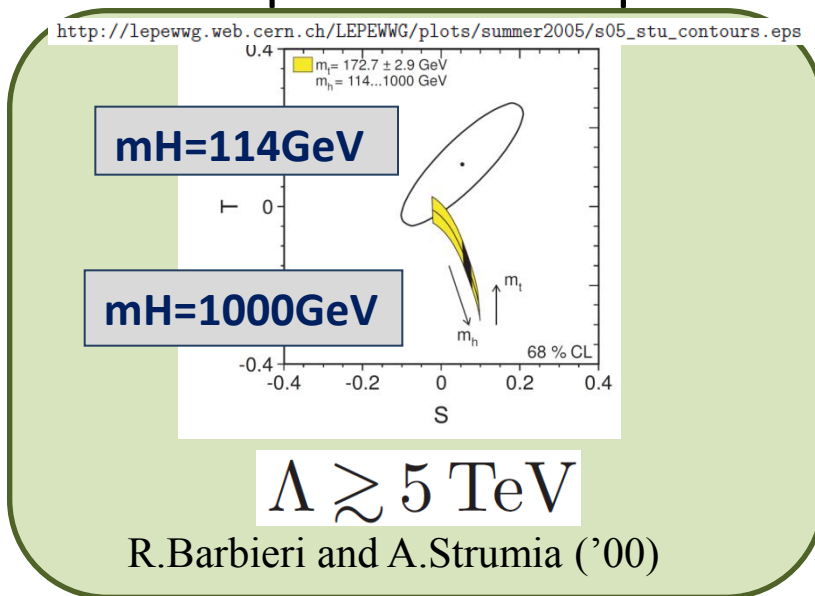
$$m_h^2 = \text{---} \bullet \text{---} + \text{---} \bigcirc \text{---}$$

$m_0^2$                    $\Lambda^2$  correction

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But... LEP experiments require



# Little Hierarchy Problem

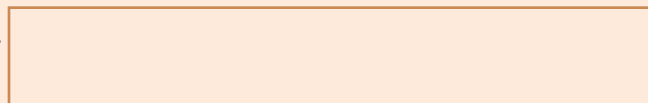
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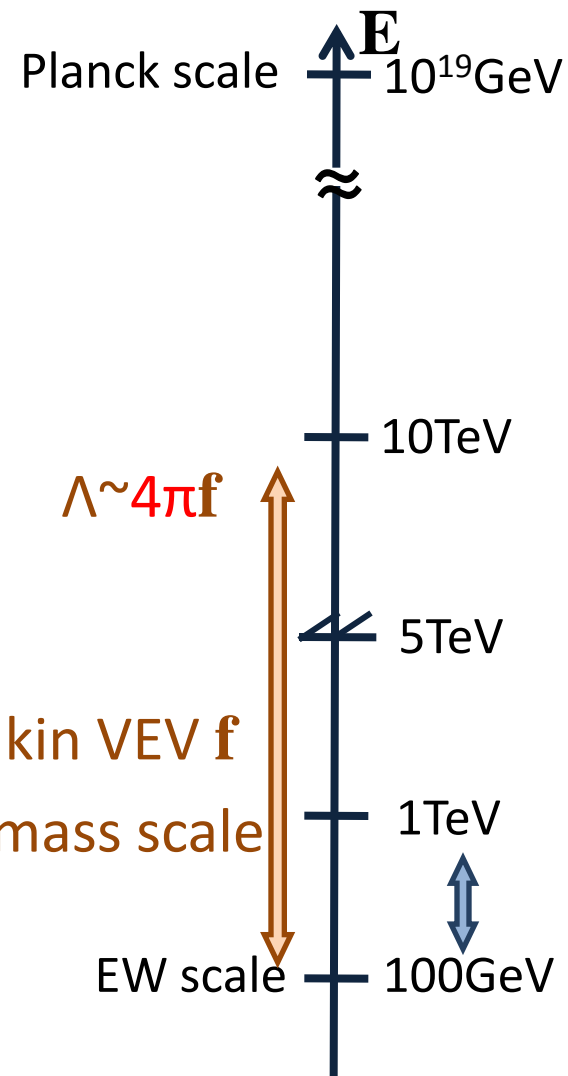
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Even if  $\Lambda \sim 10 \text{ TeV}$ , no fine-tuning!

Because of



Global sym. breakin VEV  $f$   
& new particle mass scale





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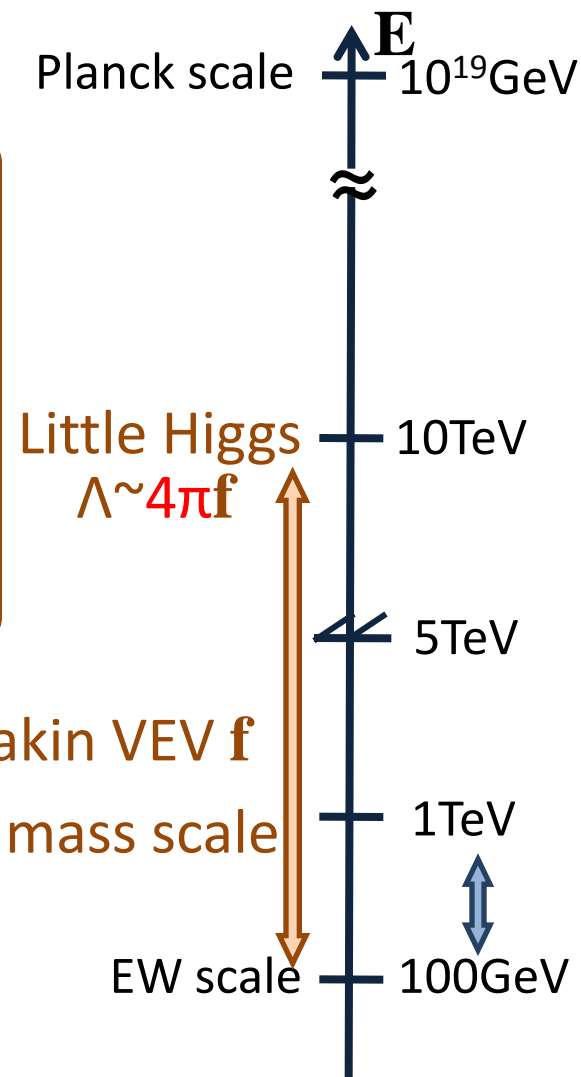
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$m_0^2$                        $\Lambda^2$  correction

↓ solution

- Little Higgs model
- Even if  $\Lambda \sim 10 \text{ TeV}$ , no fine-tuning!
- Because of Little Higgs mechanism
- collective symmetry breaking (VEV  $f$ )
- induces cancelation

⇒ Global sym. breakin VEV  $f$   
& new particle mass scale



# • Little Hierarchy Problem

$$m_h^2 = \overset{\bullet}{\text{---}} + \overset{\text{shaded circle}}{\text{---}}$$

$m_0^2$        $\Lambda^2$  correction

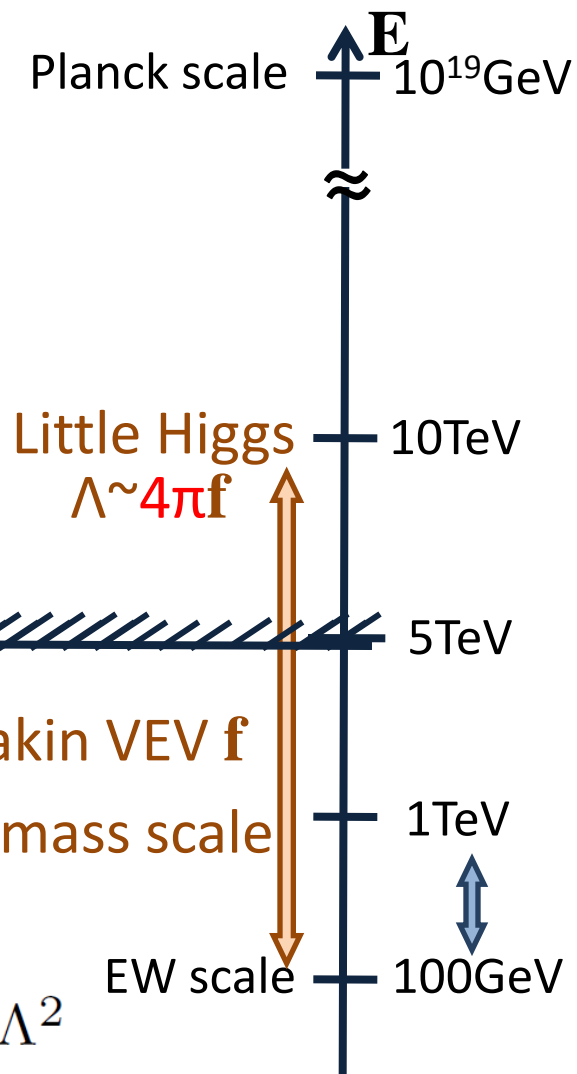
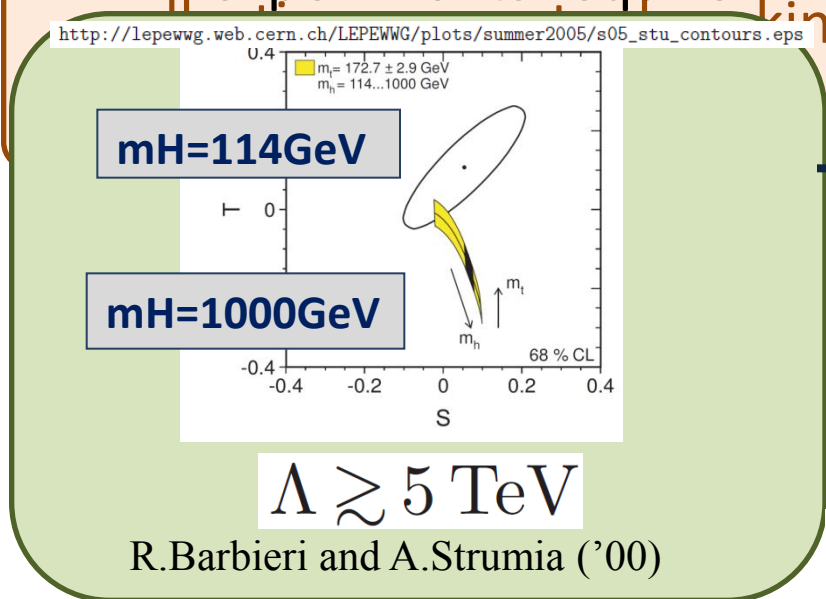
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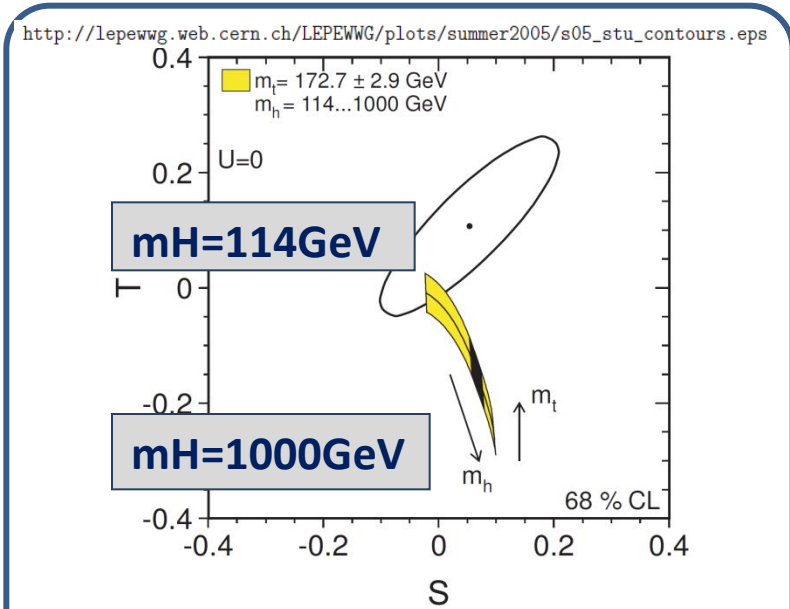
But...



$$(h^\dagger D_\mu h)^2 / \Lambda^2$$

# Little Hierarchy Problem

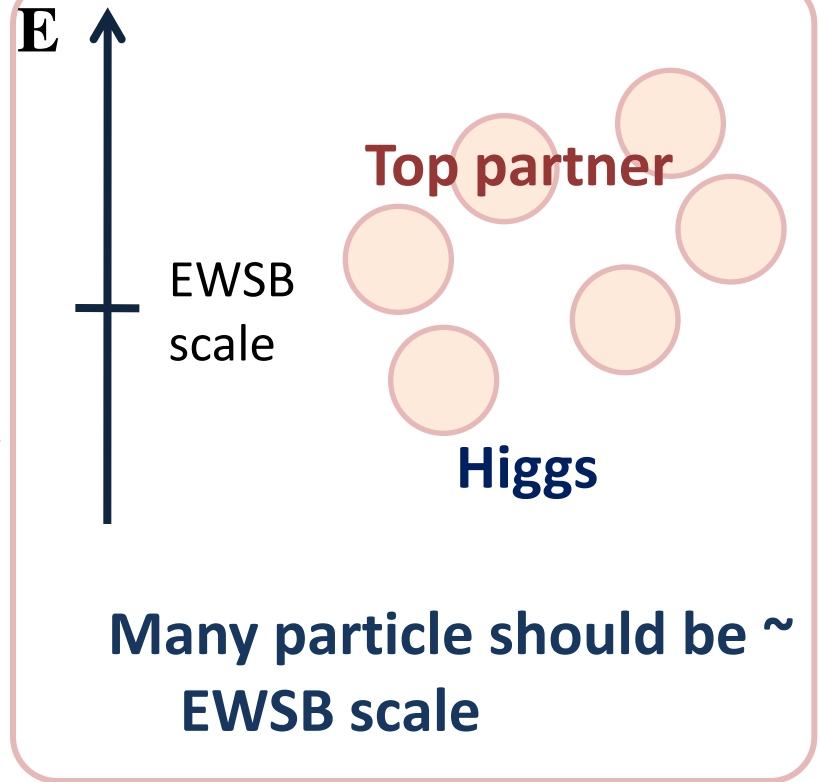
## EWPM



Higgs is light ( $\sim 100 \text{ GeV}$ ) &  
no new physics contribution

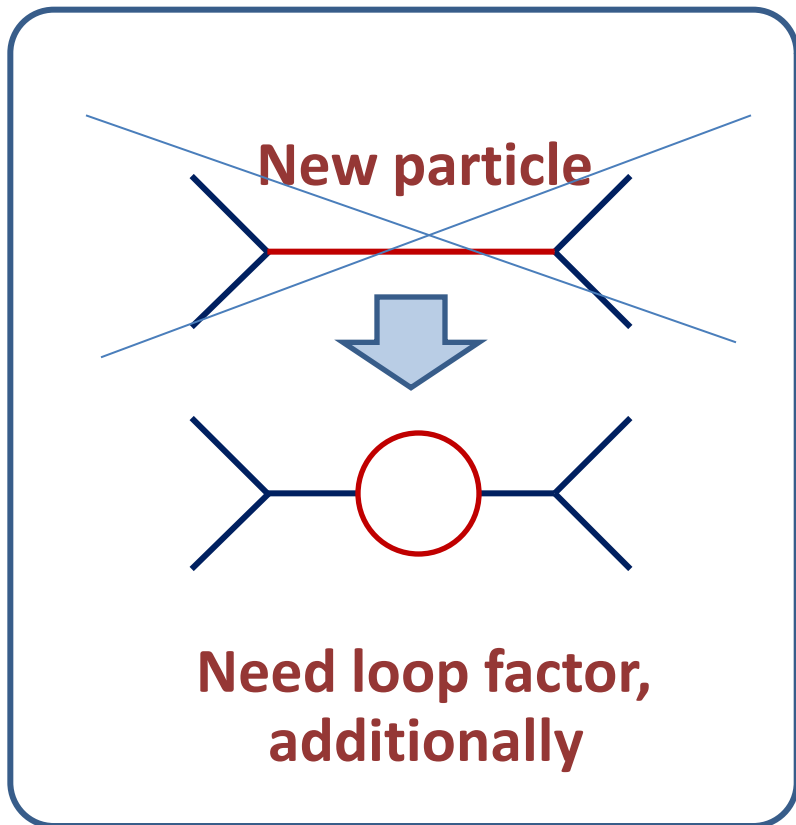


## New physics model

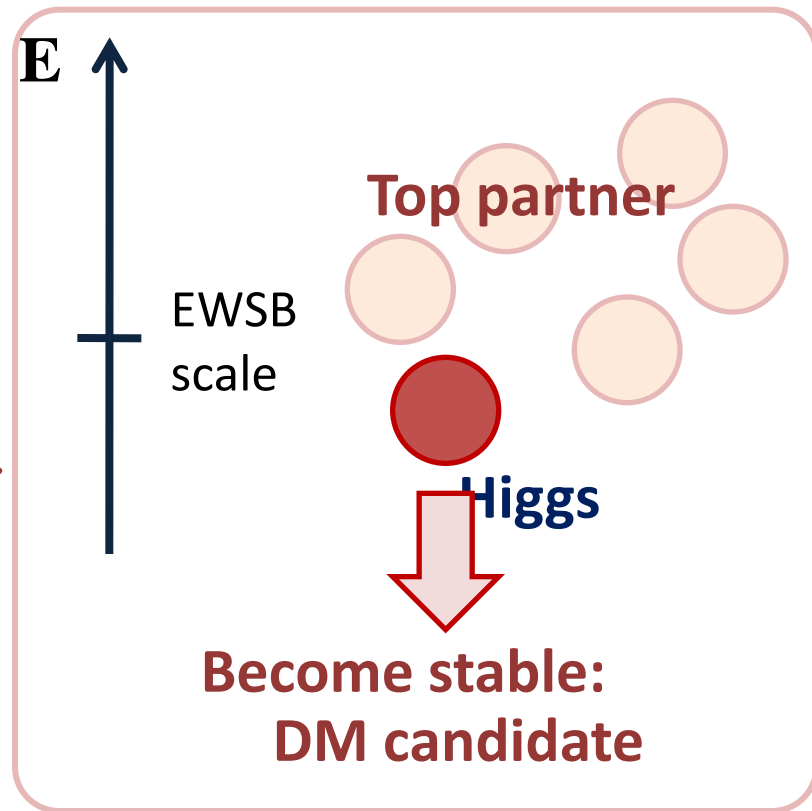


# Little Hierarchy Problem

EWPM



New physics model



**$Z_2$  symmetry!** (which guarantee "no tree level contribution to EWPM" & lightest  $Z_2$ -odd stable particle) Then,  $Z_2$  symmetry would also be a key to solve little hierarchy (& DM).

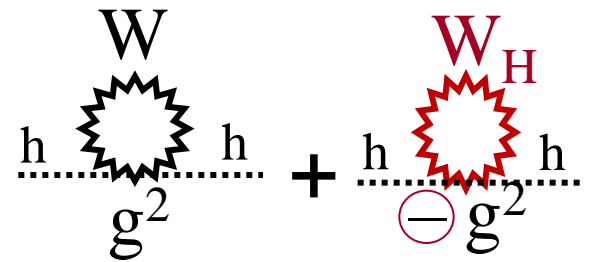
# Little Higgs model with T-parity

## • Little Higgs mechanism

### Collective symmetry breaking (VEV $f$ )

N. Arkani-Hamed, A. G. Cohen, H. Georgi ('01)

- Higgs boson is regarded as **Pseudo NG boson** of a global symmetry at some higher scale.
- Explicit breaking of the global symmetry is specially arranged to **cancel quadratic divergent corrections to  $m_h$**  at 1-loop level by new gauge bosons and fermions.



➔  $\Lambda$  can be  $\sim 10$  TeV without fine-tuning!

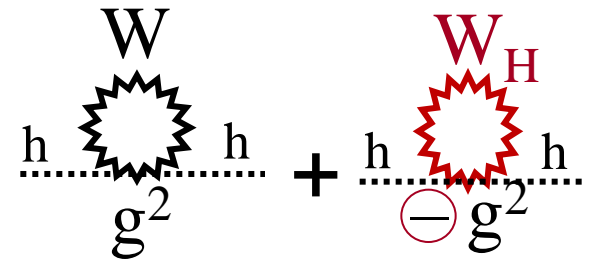
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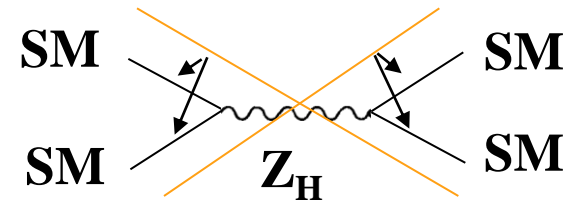
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## • T-parity

H. C. Cheng, I. Low ('03)

In order to avoid constraints from EWPM,  $Z_2$  symmetry called T-parity are introduced.

$$SM \leftrightarrow SM, \quad \text{New} \leftrightarrow -\text{New}$$



**The model contains dark matter candidate! (Heavy photon  $A_H$ )**

What is the LHT?

## In order to Implement the Little Higgs Mechanism...

- Littlest Higgs model with T-parity is based on the non-linear sigma model breaking **SU(5)/SO(5)** symmetry breaking.
- The subgroup **[SU(2) × U(1)]<sup>2</sup>** in SU(5) is gauged, which is broken down to the SM gauge group by vev **f**.

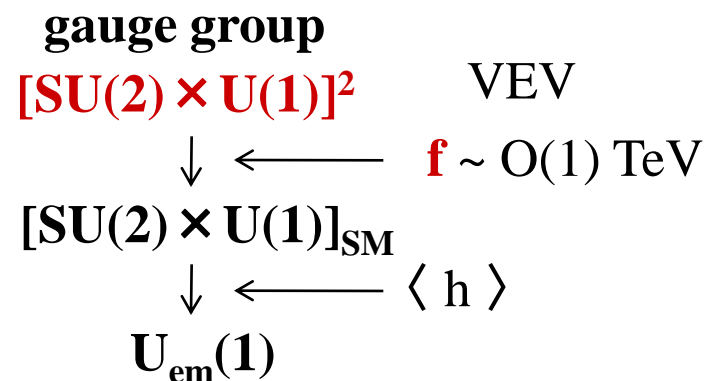
$$\begin{array}{ccc} \text{gauge group} & & \text{VEV} \\ \mathbf{[SU(2) \times U(1)]^2} & & \\ \downarrow \longleftarrow \mathbf{f} \sim \text{O}(1) \text{ TeV} & & \\ \mathbf{[SU(2) \times U(1)]_{\text{SM}}} & & \\ \downarrow \longleftarrow \langle h \rangle & & \\ \mathbf{U_{em}(1)} & & \end{array}$$

[Arkani-Hamed, Cohen, Katz and Nelson (2002)]



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[Arkani-Hamed, Cohen, Katz and Nelson (2002)]

**Kinetic term of non-linear sigma model** — gauge couplings (**SU(2)**)

$$D_\mu \Sigma = \partial_\mu \Sigma - i \sum_j \underline{\underline{g_j}} W_j^a (Q_j^a \Sigma + \Sigma Q_j^{aT}) + \dots, \quad Q_1^a = \begin{pmatrix} \sigma^a/2 & 0 & 0 \\ 0 & \boxed{0 \ 0} \\ 0 & \boxed{0 \ 0} \end{pmatrix} Q_2^a = \begin{pmatrix} \boxed{0 \ 0} & 0 \\ \boxed{0 \ 0} & 0 \\ 0 & 0 & -\sigma^{a*}/2 \end{pmatrix}$$

The global sym. is explicitly broken by gauge & Yukawa coupling **but only collectively**: Each gauged **[SU(2) × U(1)]<sub>i</sub>** commute with different subgroup of **G** that acts on the Higgs field non-linearly.

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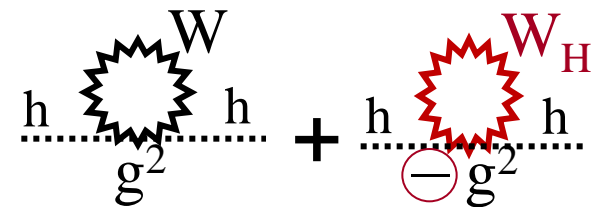
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The global sym. is explicitly broken by gauge & Yukawa coupling **but only collectively**: Each gauged **[SU(2) × U(1)]<sub>i</sub>** commute with different subgroup of **G** that acts on the Higgs field non-linearly. Thus, **if any one of gauge couplings is zero, the Higgs boson is an exact Nambu-Goldstone boson (due to the restoration of the global symmetry).**

Since there are one gauge coupling in 1-loop diagram, the quadratic divergence is vanished in 1-loop level!



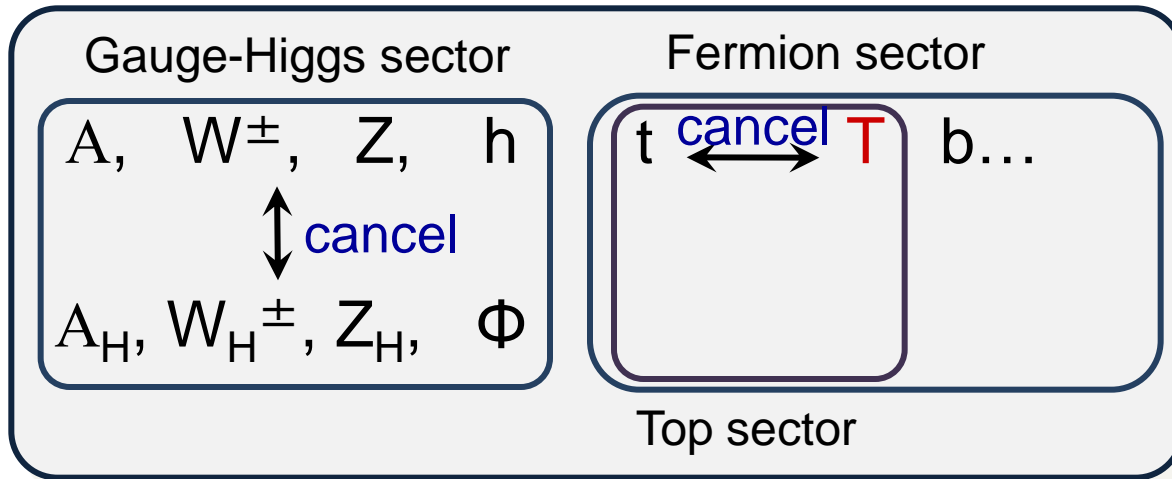
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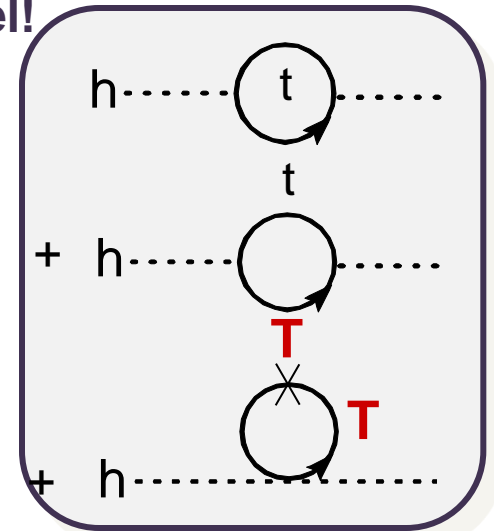
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[Arkani-Hamed, Cohen, Katz and Nelson (2002)]

## Particle contents



## Cancel!



Due to the cancelation of quadratic divergences, partners are introduced. In the fermion sector, only top partner is required to cancel the quadratic divergences because other fermion Yukawa couplings are small.

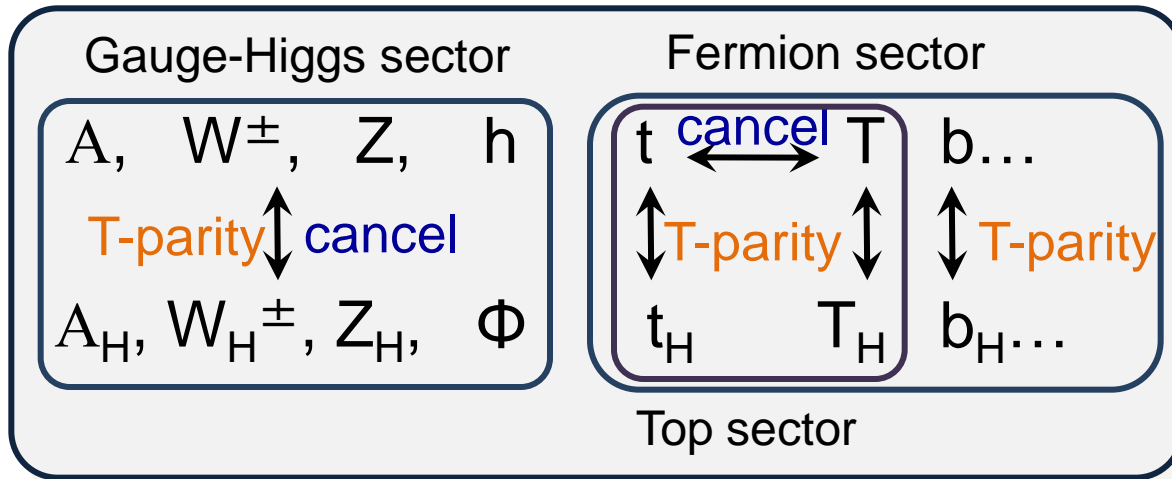
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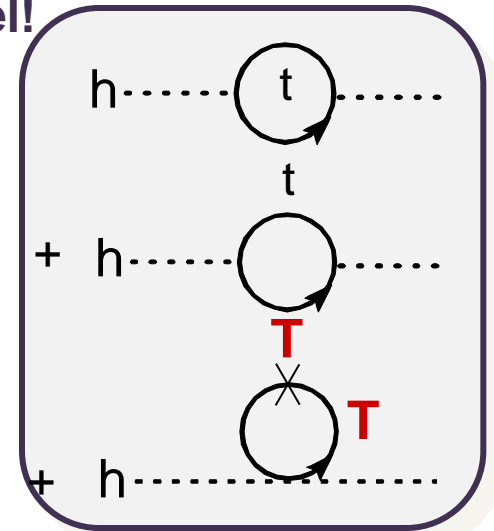
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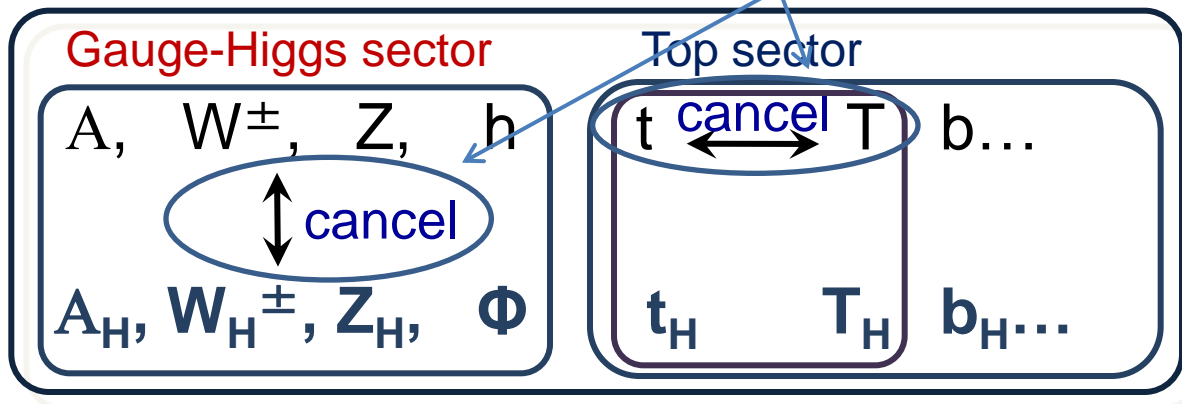


In order to implement the **T-parity**...  $[SU(2) \times U(1)]_1 \leftrightarrow [SU(2) \times U(1)]_2$

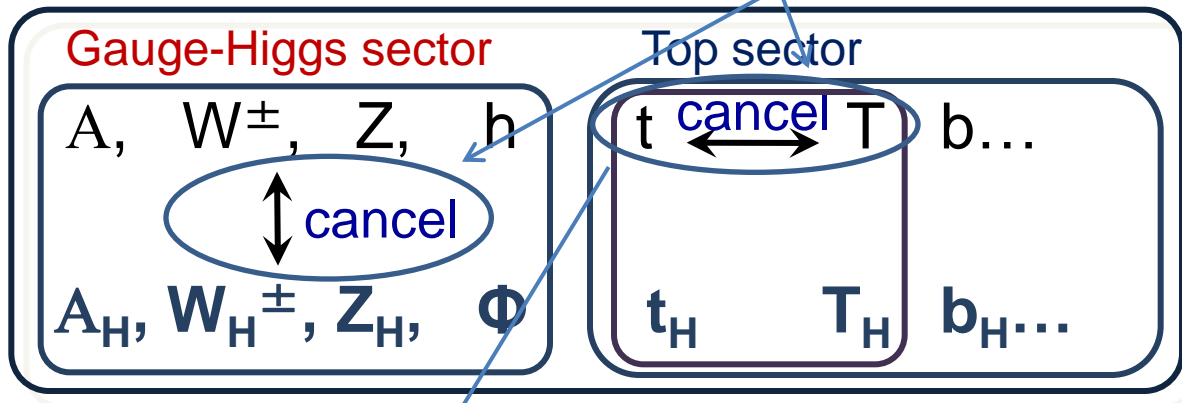
T-odd partners are introduced for each fermions.

How to measure the LHT at ILC?

In order to verify the Little Higgs model,  
we should measure the **Little Higgs mechanism!**

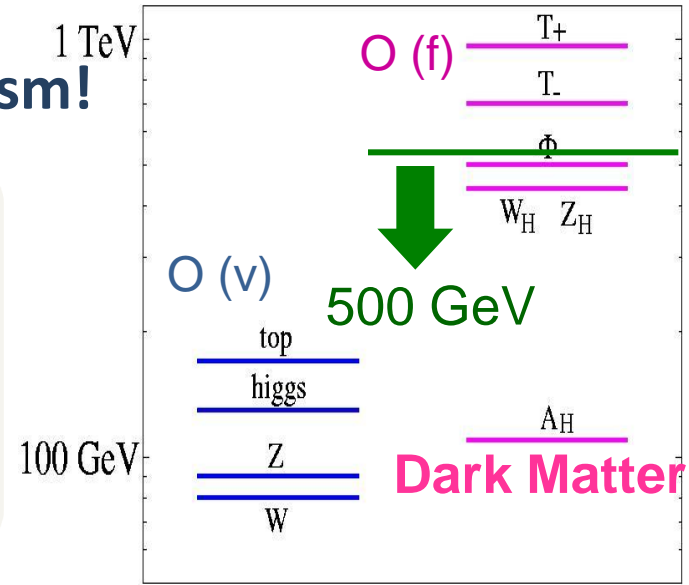


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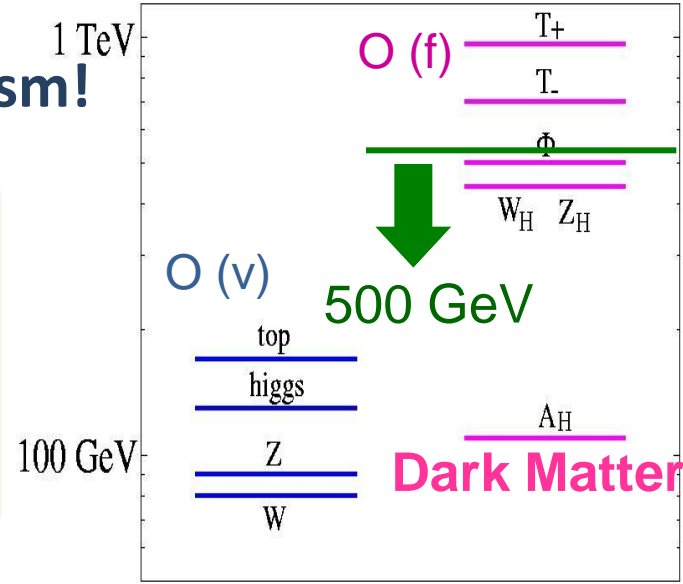
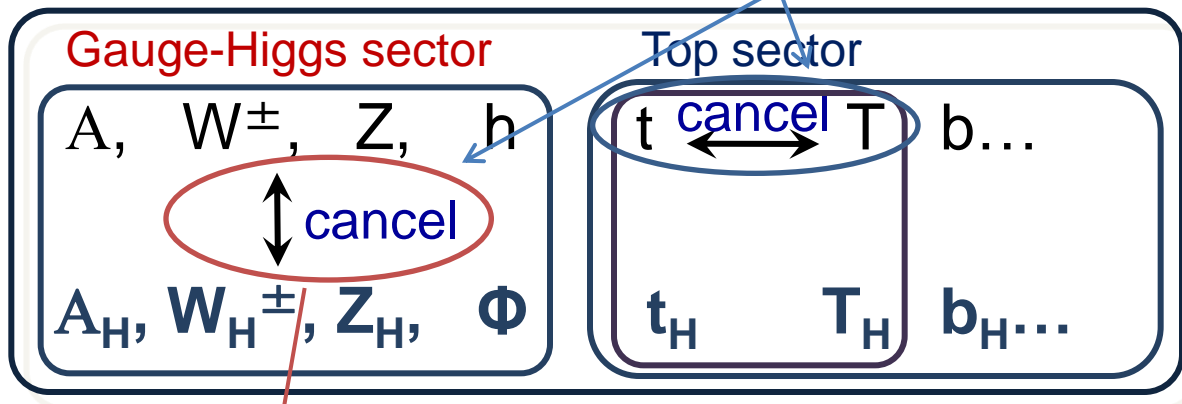
- can be measured at LHC.
- too heavy to produced at ILC.



mass spectrum

**S. Matsumoto, T. Moroi, K. Tobe (08) and recent papers .....**

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1. Top sector

- can be measured at LHC.
- too heavy to produced at ILC.

2. Gauge sector

- can be produced at ILC!
- cannot precise measurement at LHC
- masses depend only on vev  $f$ !

S. Matsumoto, T. Moroi, K. Tobe (08) and recent papers .....

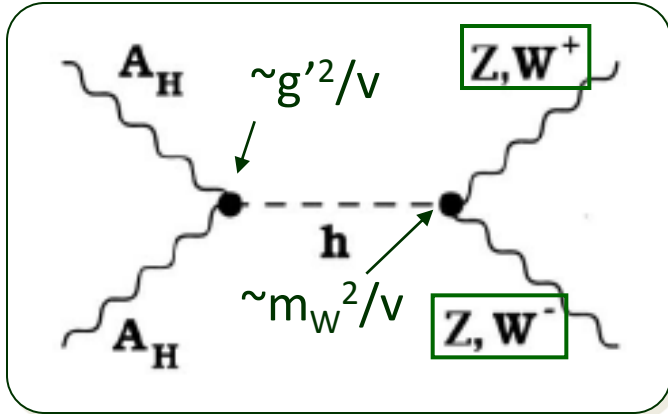
Q. H. Cao, C. R. Chen (07);  
K. Cheung, et. al. (13)

**We should measure the gauge sector in the LHT at ILC.**



Representative point  
of our simulation

# Dark matter annihilation



- The relic density depends only on  $f$  &  $m_h$ !

$$m_{A_H} \simeq \frac{g'}{\sqrt{5}} f$$

M. A, S. Matsumoto, N. Okada, Y. Okada PhysRevD.75.063506

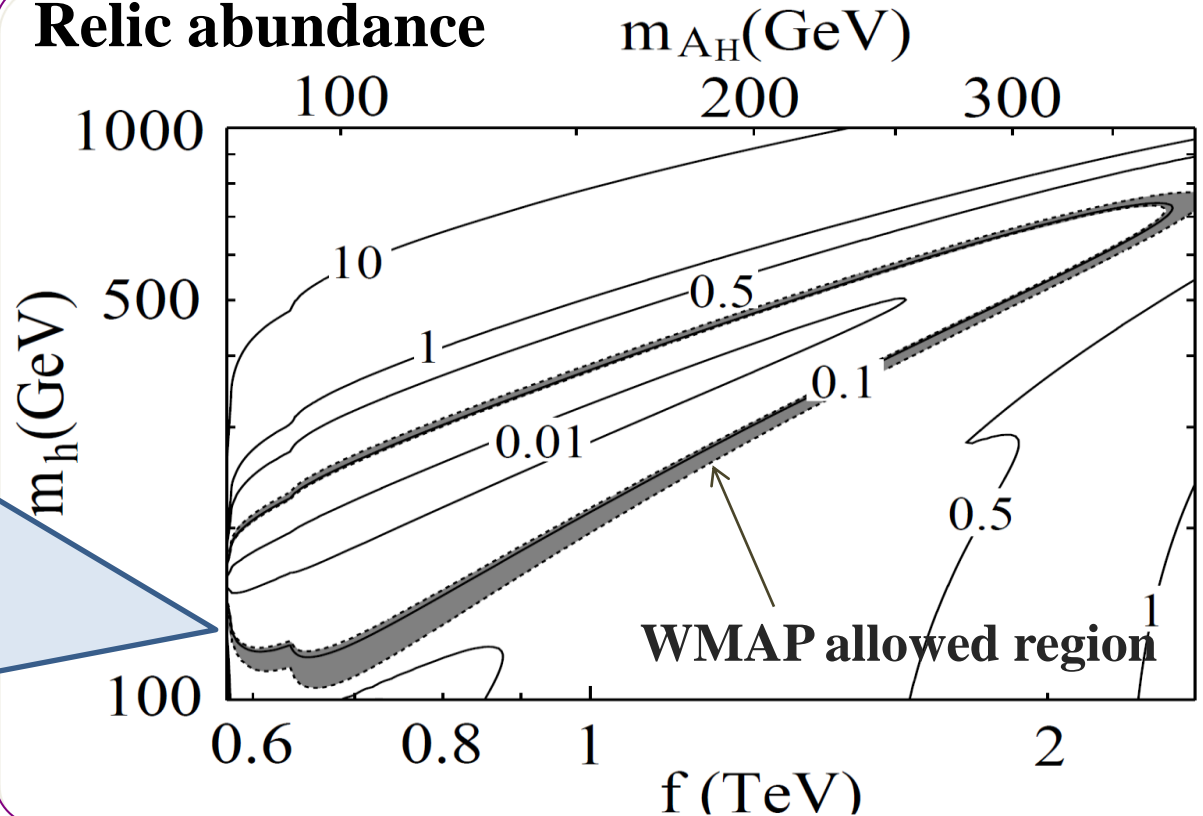
## Dark matter ( $A_H$ )

Shaded area is WMAP allowed region .

Related with s-channel pole.

**At this WMAP allowed region, the model also satisfies other experimental constraints!**

## Relic abundance



# Mode & Representative point

Which particle can produced at the ILC ?  $\rightarrow$  It depends on VEV  $f$ .

Because heavy gauge boson mass depend only on vev  $f$ .

$$e^+e^- \rightarrow A_H Z_H$$

$$e^+e^- \rightarrow W_H W_H$$

If  $f$  is larger than  $\sim 800\text{GeV}$ , the hierarchy problem appears again.

M. A, S. Matsumoto, N. Okada, Y. Okada PhysRevD.75.063506

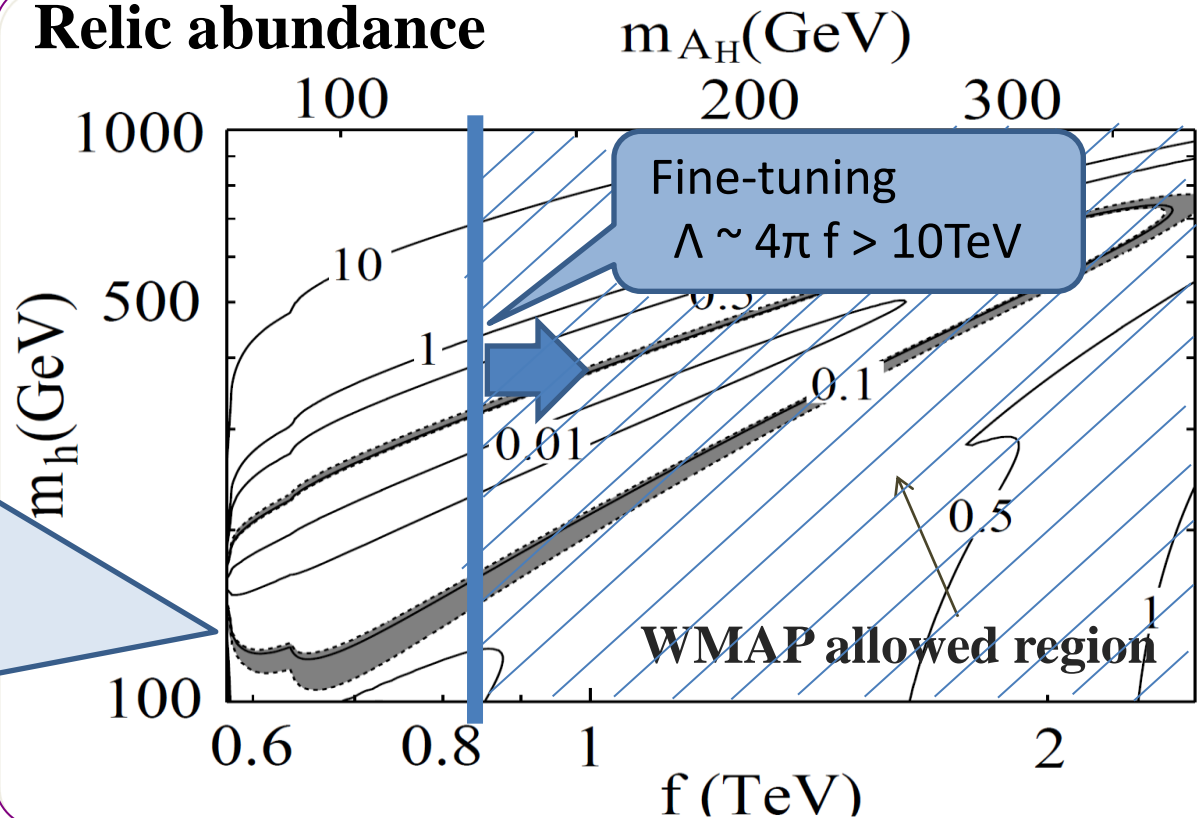
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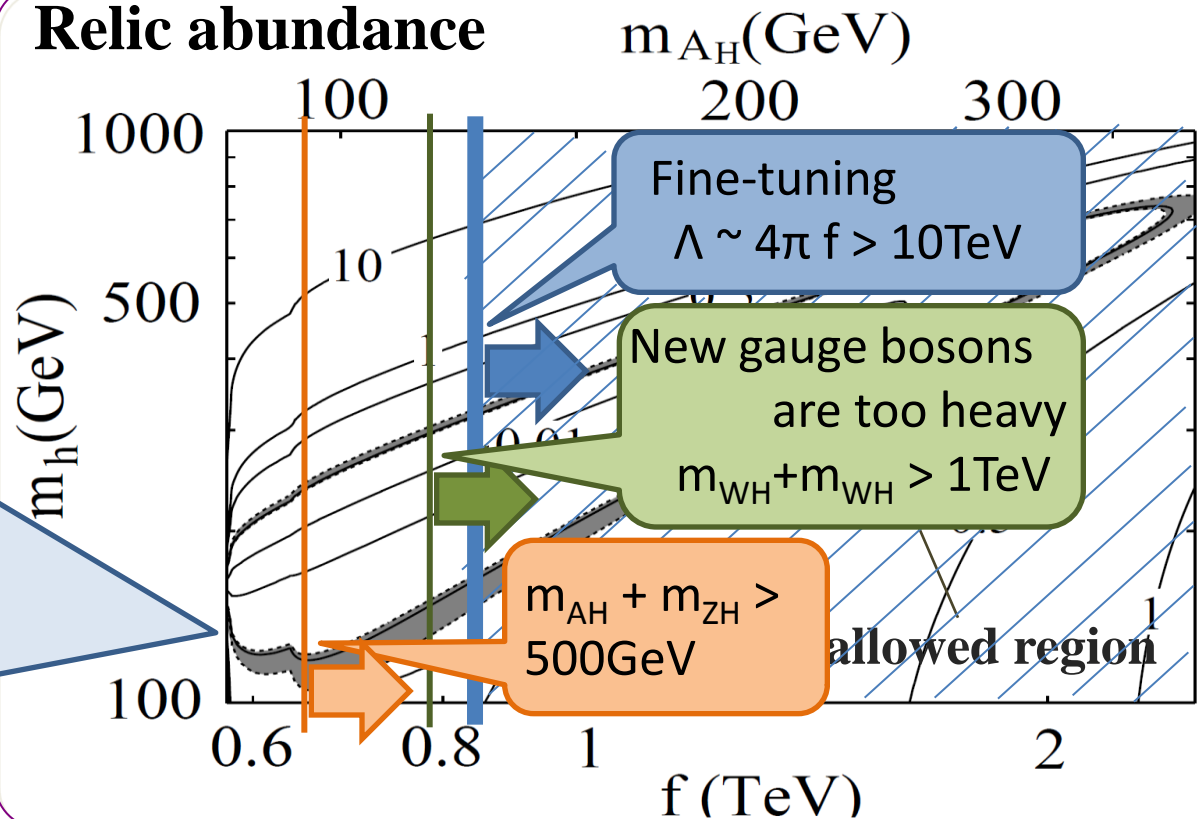
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If Higgs mass is changed as 126 GeV, the result accuracy of this analysis would be better.

M. A, S. Matsumoto, N. Okada, Y. Okada PhysRevD.75.063506

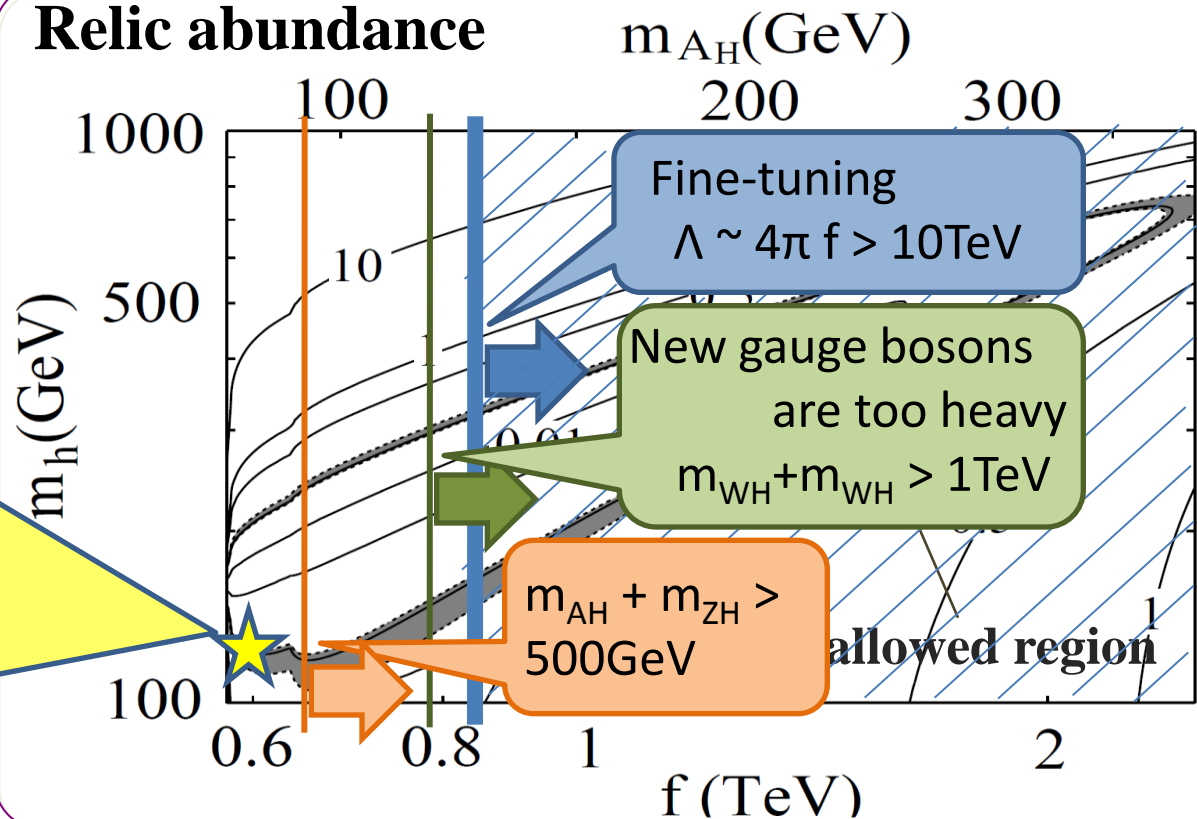
## Mass spectrum

$f$	580
$e_H$	410
$Z_H$	369
$W_H$	368
$A_H$	82
Higgs	134

## Point I (GeV)

Sample points satisfy all experimental & cosmological constraints.

## Relic abundance



# Simulation results

## 1.Event generation

**Physsim** : for **LHT & SM** process

- helicity amplitude calculation, gauge boson polarization effect
- phase space integration and generation of parton 4-momenta

## 2.Hadronization

**PYTHIA** - parton showering and hadronization

**TAUOLA** - for final state tau leptons

## 3.Detector simulation

**JSFQuickSimulator** - create vertex-detector hits

- smear charged-track parameters in central tracker
- simulate calorimeter signals as from individual segments

Simulations has been performed

at  $\sqrt{s} = 1 \text{ TeV}$  with an Integrated luminosity of **500 fb<sup>-1</sup>**

# Heavy gauge boson masses

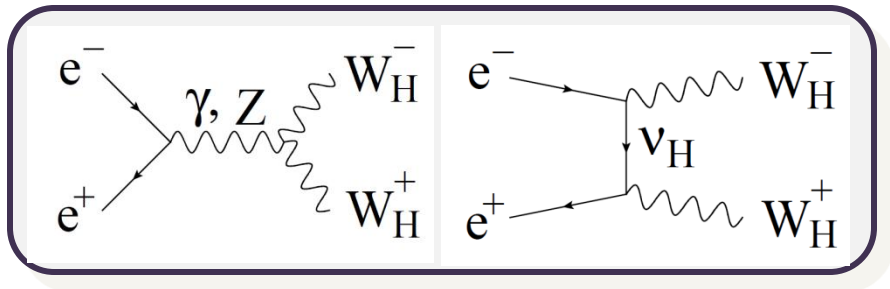
$$e^+e^- \rightarrow W_H W_H$$

$$e^+e^- \rightarrow Z_H Z_H$$



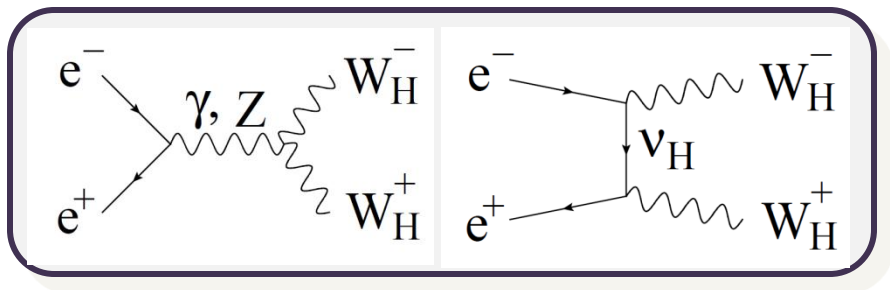
# Event

$$e^+e^- \rightarrow W_H W_H$$



# Event

$$e^+e^- \rightarrow W_H W_H \rightarrow A_H A_H WW$$



$$W \rightarrow q\bar{q} \text{ (BR} = 67.6\%)$$
$$W \rightarrow l\nu \text{ (BR} = 32.4\%)$$

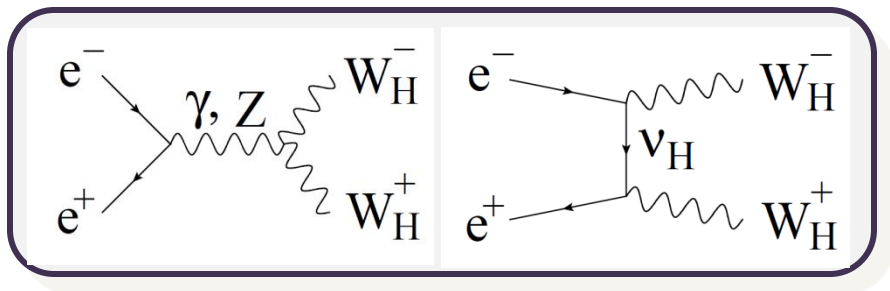


$A_H A_H$  qq $\bar{q}\bar{q}$  as the signal event

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$$e^+e^- \rightarrow W_H W_H \rightarrow A_H A_H WW$$

- Large missing energy
- 4 jet in final state



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$A_H A_H$  qqqq as the signal event

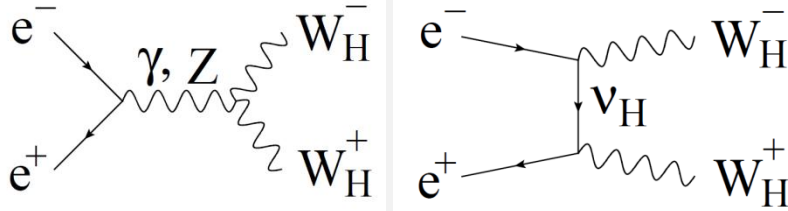
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## Signal

## Background

- Large missing energy
- 4 jet in final state

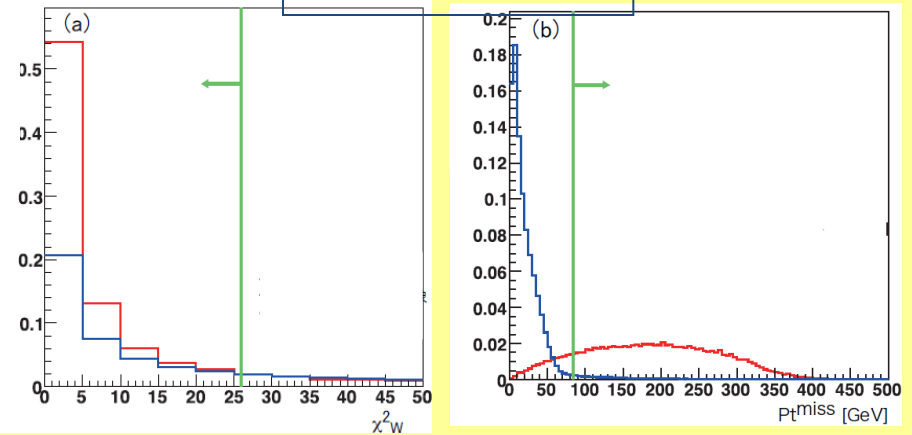


## Selection Cut

- $E_W < 500 \text{ GeV}$  : energy of W
- $\chi_W^2 < 26$  :  $\chi^2$  for  $W^\pm$  reconstruction from jets
- $P_{T, \text{miss}} > 84 \text{ GeV}/c$  : Missing transverse momentum

Process	cross section (fb)
$W_H W_H \rightarrow A_H A_H qqqq$	120.0
$WW \rightarrow qqqq$	1307
$eeWW \rightarrow eeqqqq$	490
$e\nu_e W Z \rightarrow e\nu_e qqqq$	24.5
$\nu\bar{\nu} WW \rightarrow \nu\bar{\nu} qqqq$	18.8
$ZWW \rightarrow \nu\bar{\nu} qqqq$	7.23
$Z_H Z_H \rightarrow A_H A_H qqqq$	5.61

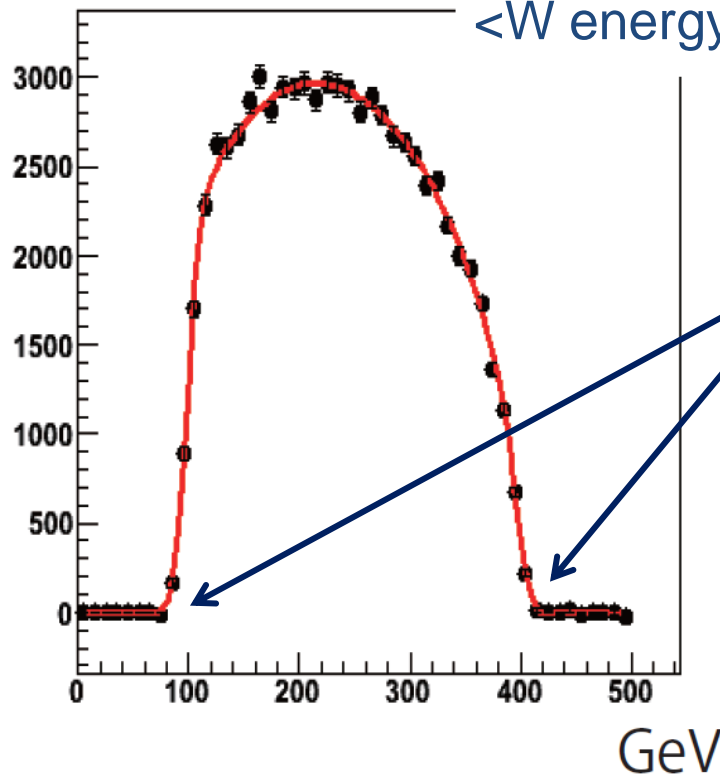
## Just sample



$$\chi_W^2 < 26$$

$$P_t^{\text{miss}} > 84 (\text{GeV})$$

<W energy distribution>



The Energy distribution of W from  $W_H W_H$  after subtraction of the backgrounds, whose number of events was estimated by independent BG samples. The distribution is fitted by a polynomial function.

Results

$$m_{AH} = 81.9 \pm 0.6 \text{ GeV},$$

$$m_{WH} = 368.2 \pm 0.6 \text{ GeV}$$

or

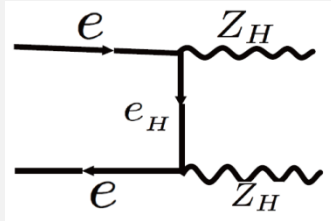
$$m_{AH} = 81.0 \pm 0.6 \text{ GeV},$$

$$m_{WH} = 218.0 \pm 0.7 \text{ GeV}$$

Preliminary

# Event

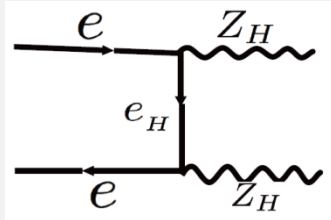
$$e^+e^- \rightarrow Z_H Z_H$$



## Event

$$e^+e^- \rightarrow Z_H Z_H \rightarrow A_H A_H hh$$

- Large missing energy
- 4 jet in final state



Higgs boson mainly decays into  $bb$   
(or  $WW$ )



$A_H A_H$   $qqqq$  as the signal event

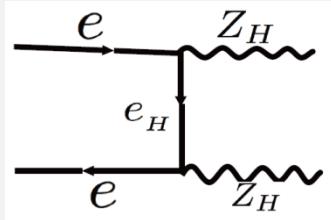
## Event

$$e^+e^- \rightarrow Z_H Z_H \rightarrow A_H A_H hh$$

## Signal

## Background

- Large missing energy
- 4 jet in final state



## Selection Cut

No isolated leptons  
# of b-jet > 1  
Acoplanarity > 20°

$E_h < 500 \text{ GeV}$  : energy of higgs

$\chi_h^2 < 60$  :  $\chi^2$  for higgs reconstruction from jets

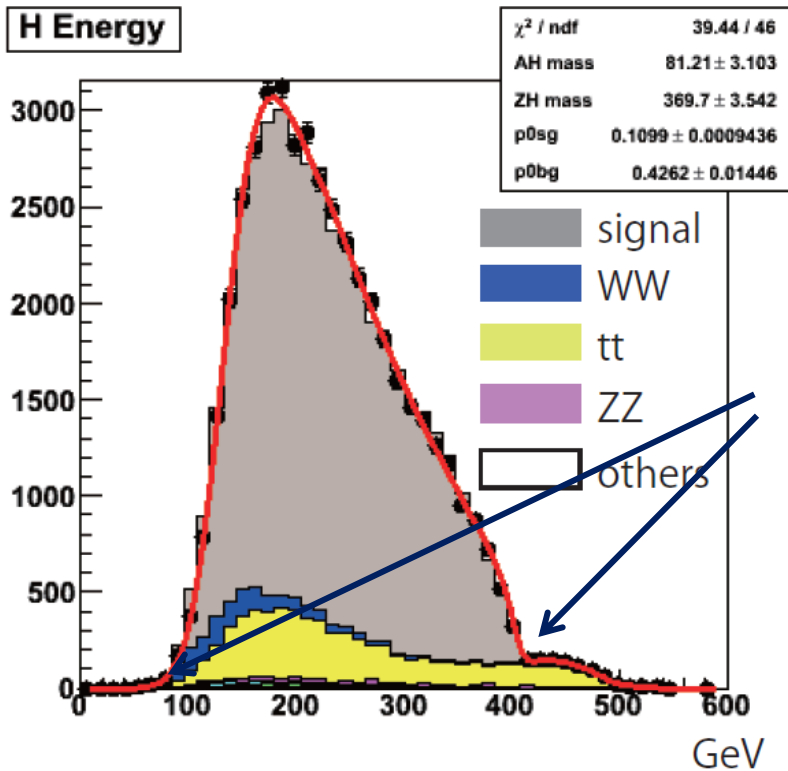
$P_{\text{T}}^{\text{miss}} > 84 \text{ GeV}/c$  : Missing transverse momentum

Process	cross section (fb)
$Z_H Z_H \rightarrow A_H A_H H H$	98.0
$\rightarrow A_H A_H q q q q$	
$W W Z$	63.86
$\nu \nu W W$	14.67
$W W$	3069
$t t$	192.9
$Z Z$	202.2
$Z H$	17.98
$W_H W_H$	108.6



## <Higgs energy distribution>

The Energy distribution of reconstructed higgs boson. The distribution is fitted by a polynomial function.



## Results

$$m_{\text{AH}} = 82.7 \pm 3.5 \text{ GeV,}$$

$$m_{\text{ZH}} = 367.1 \pm 4.7 \text{ GeV}$$

or

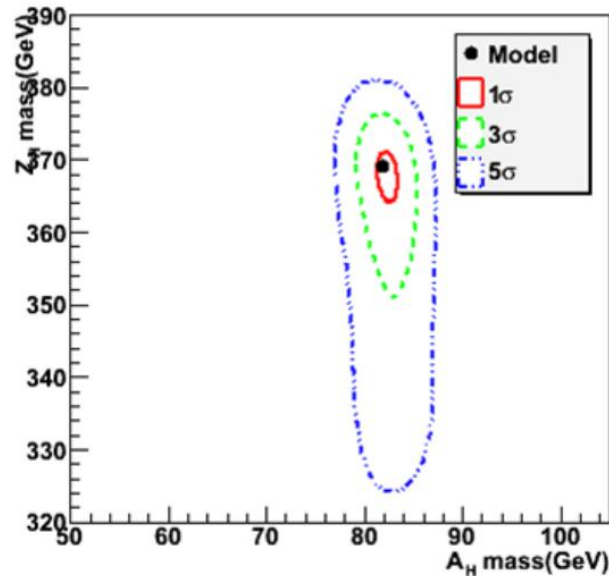
$$m_{\text{AH}} = 93.0 \pm 3.5 \text{ GeV,}$$

$$m_{\text{ZH}} = 335.4 \pm 4.7 \text{ GeV}$$

Preliminary

Process	obtained mass	solution 1 (GeV)	solution 2 (GeV)
$W_H W_H$	$(A_H, W_H)$	(81.9, 368.3)	(81.0, 218.0)
$Z_H Z_H$	$(A_H, Z_H)$	(82.7, 366.1)	(93.0, 335.4)

## Simultaneous fit using $W_H W_H$ & $Z_H Z_H$ signal



in order to obtain the true mass because true solution of  $m_{AH}$ .

Preliminary

Results

$A_H$	81.9 GeV	1.3%
$W_H$	369 GeV	0.20%
$Z_H$	368 GeV	0.56%

**High accuracy !**

➔  $\mathbf{f} = 580 \text{ GeV } 0.16 \%$

# Heavy lepton masses

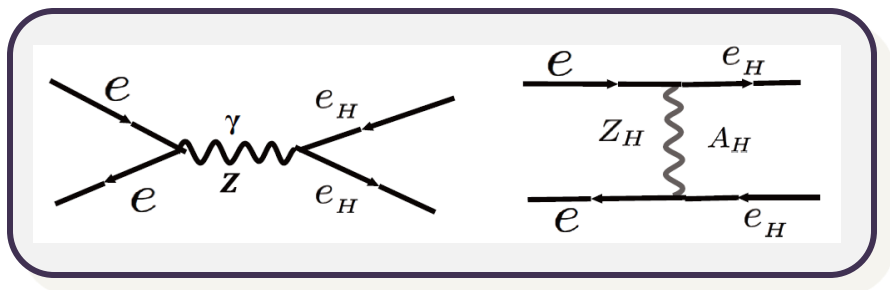
$$e^+e^- \rightarrow e_H e_H$$

$$e^+e^- \rightarrow \nu_H \nu_H$$

# Event

$$e^+e^- \rightarrow e_H e_H \rightarrow eeZ_H Z_H \rightarrow eeA_H A_H qqqq$$

Particle	Branch	Br
$e_H$	$eA_H$	30%
	$eZ_H$	25%
	$\nu W_H$	45%



$eeA_H A_H qqqq$  as the signal event

## Event

$$e^+e^- \rightarrow e_H e_H \rightarrow ee Z_H Z_H \rightarrow ee A_H A_H qqqq$$

## Signal

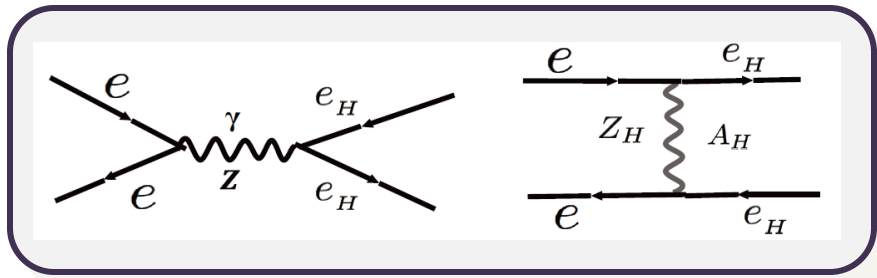
## Background

Process	cross section (fb)
$e_H e_H \rightarrow ee Z_H Z_H$ $\rightarrow ee A_H A_H qqqq$	3.91
$tt$	200.7
$ee WW$	1020
$\tau_H \tau_H$	3.32
$ttZ, ttH,$	114.5

We also consider

$evWZ, eeZZ, WWZ, ZZZ$

$ee A_H A_H qqqq$  as the signal event

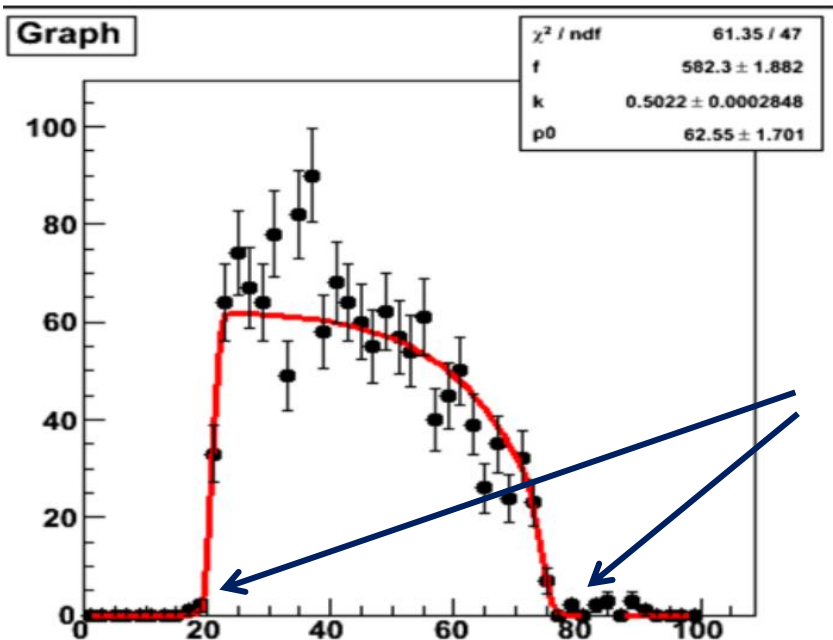


## Selection Cut

One isolated electrons & positron, and 4-jets

Two Higgs bosons ( $104 \text{ GeV} < m_h < 164 \text{ GeV}$ )

$P_{T, \text{miss}} > 50 \text{ GeV}/c$  : Missing transverse momentum



<electron energy distribution>

**The Energy distribution of reconstructed higgs boson. The distribution is fitted by a polynomial function.**

**Results**

$$m_{eH} = 412.8 \pm 1.7 \text{ GeV,}$$

$$m_{ZH} = 371.2 \pm 1.5 \text{ GeV}$$

**Preliminary**

**High accuracy !**

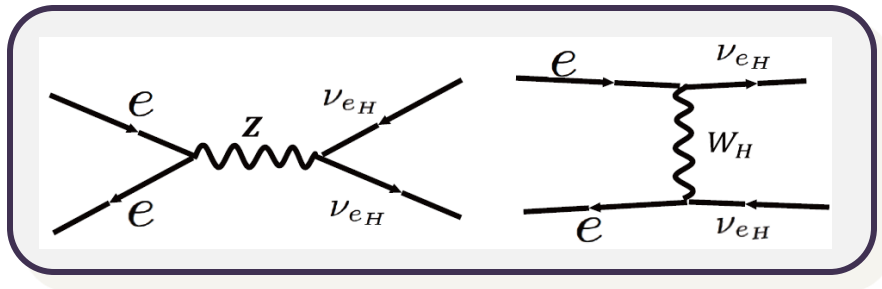
# Event

$$e^+e^- \rightarrow \nu_H \nu_H \rightarrow eeW_H W_H \rightarrow eeA_H A_H qqqq$$

Particle	Branch	Br
	$eW_H$	65%
$\nu_H$	$\nu Z_H$	11%
	$\nu A_H$	24%



$eeA_H A_H qqqq$  as the signal event



parameter	True value	measurement accuracy
$f$	580 GeV	0.16 %
$\kappa_\ell$	0.5	0.01 %

## Event

$$e^+e^- \rightarrow \nu_H \nu_H \rightarrow eeW_H W_H \rightarrow eeA_H A_H qqqq$$

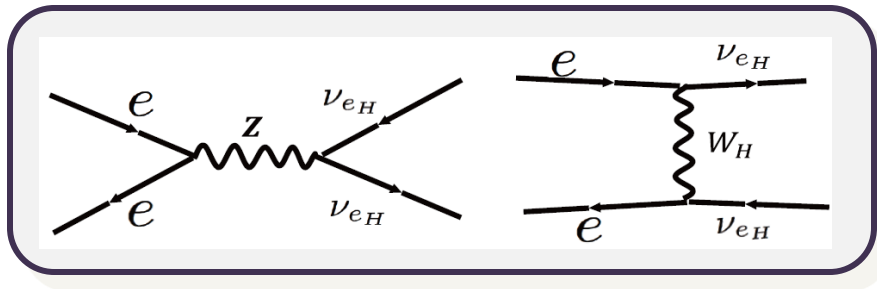
## Signal

## Background

Process	cross section (fb)
$\nu_H \nu_H \rightarrow eeW_H W_H$ $\rightarrow eeA_H A_H qqqq$	26.0
$\nu_{\tau H} \nu_{\tau H}, e_H e_H,$ $tt, ttZ, ttH,$	3.3 + 3.9 2482

We also consider

$evWZ, eeWW, WWZ, ZZZ$



$eeA_H A_H qqqq$  as the signal event

## Selection Cut

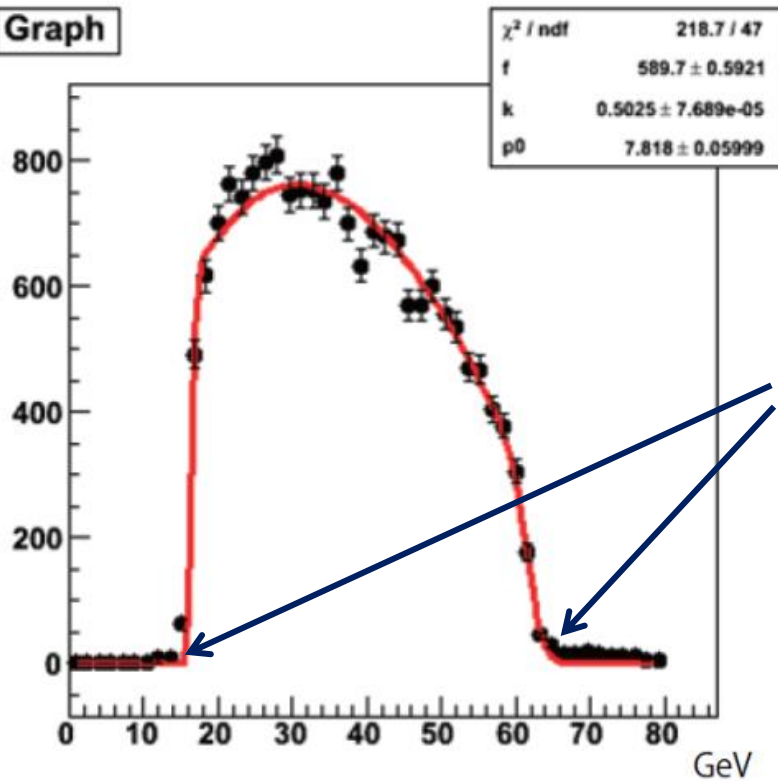
One isolated electrons & positron, and 4-jets

Two W bosons ( $60 \text{ GeV} < m_W < 100 \text{ GeV}$ )

$P_{T, \text{miss}} > 50 \text{ GeV}/c$  : Missing transverse momentum



Graph



<electron energy distribution>

The Energy distribution of reconstructed Higgs boson. The distribution is fitted by a polynomial function.

Results

$$m_{\nu H} = 400.8 \pm 0.4 \text{ GeV,}$$

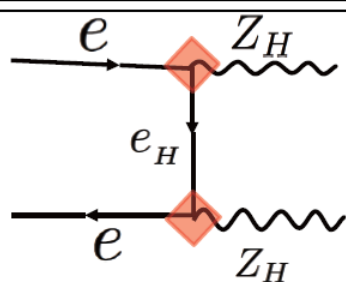
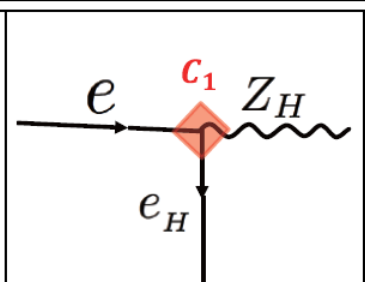
$$m_{WH} = 369.6 \pm 0.4 \text{ GeV}$$

Preliminary

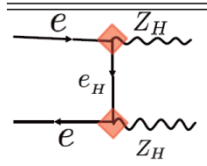
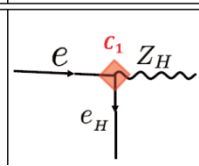
High accuracy !

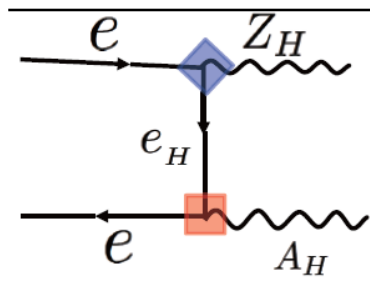
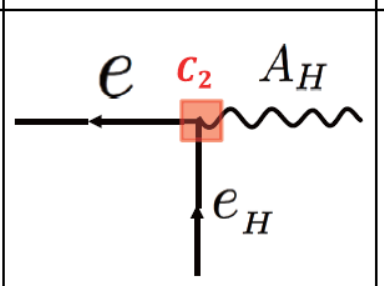
$$m_{eH} = \sqrt{2}\kappa_l f, \quad m_{\nu H} = \left(\frac{\sqrt{2} + \sqrt{1 + c_f}}{2}\right)\sqrt{2}\kappa_l f \quad \Rightarrow \quad \kappa_l = 0.5 \text{ GeV, } 0.01 \%$$

# Coupling measurements

Process mode	Branch	derived vertex
		

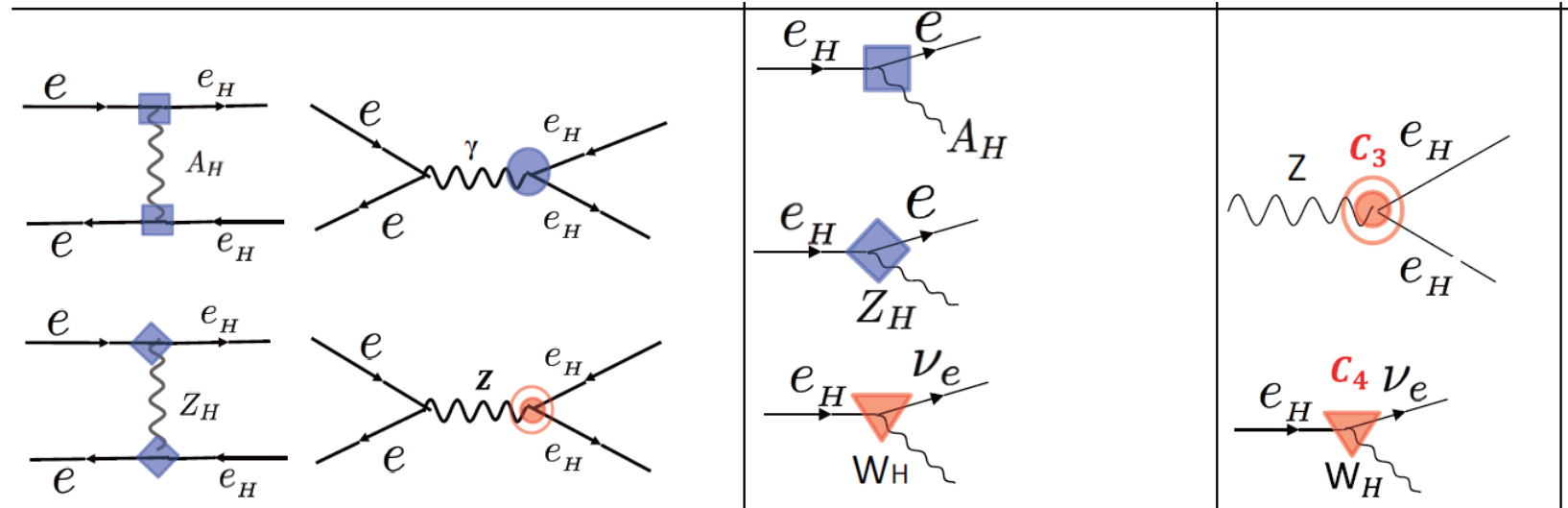
- Cross section

Process mode	Branch	derived vertex
		

		
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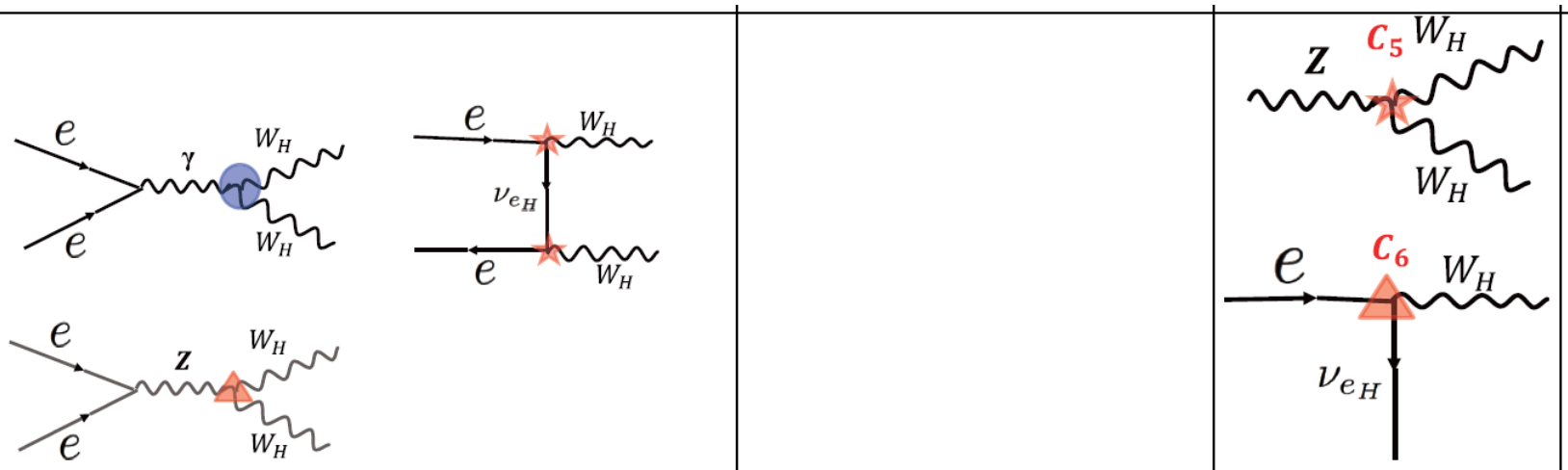
- Cross section

Process mode	Branch	derived vertex



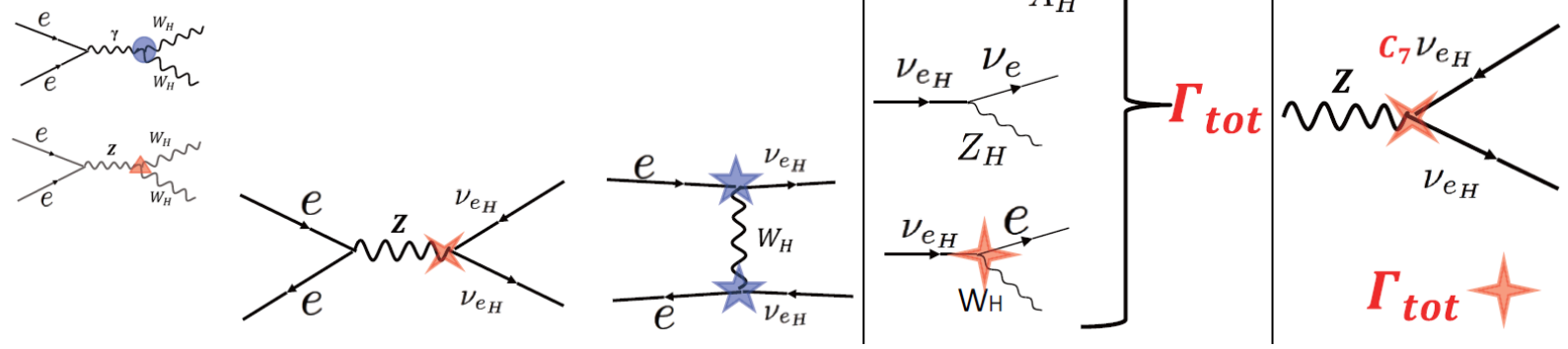
- Cross section
- polarization

Process mode	Branch	derived vertex



- Cross section
- polarization

Process mode	Branch	derived vertex



- Cross section
- polarization

Process mode	Branch	derived vertex

**0.21%**

**3.9%**

**15%**

**11%**

**0.87%**

**0.21%**

**0.28%**

**0.21%**

Preliminary



# Summary & Discussion

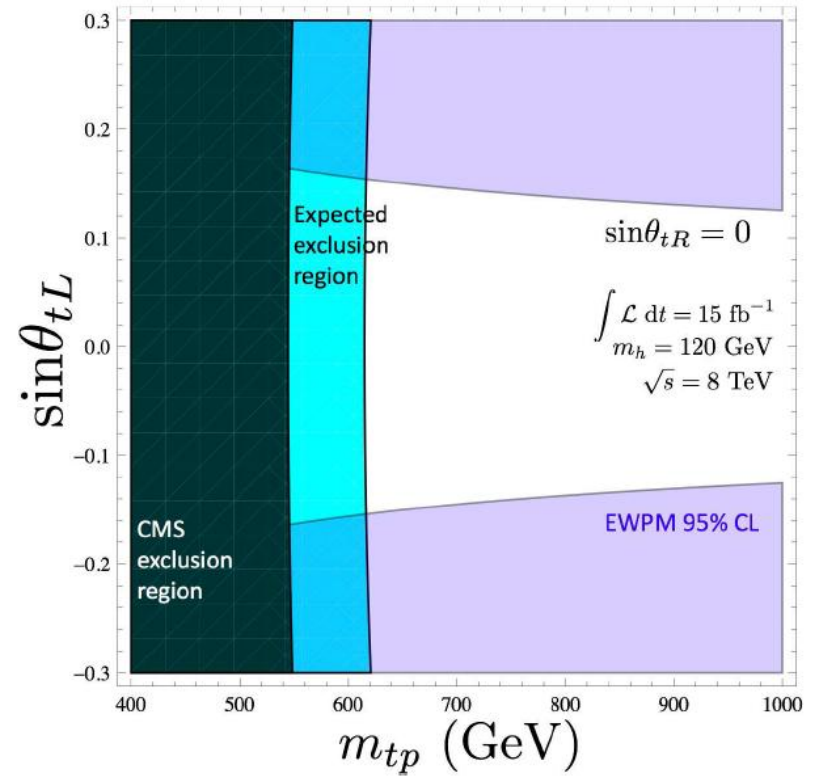
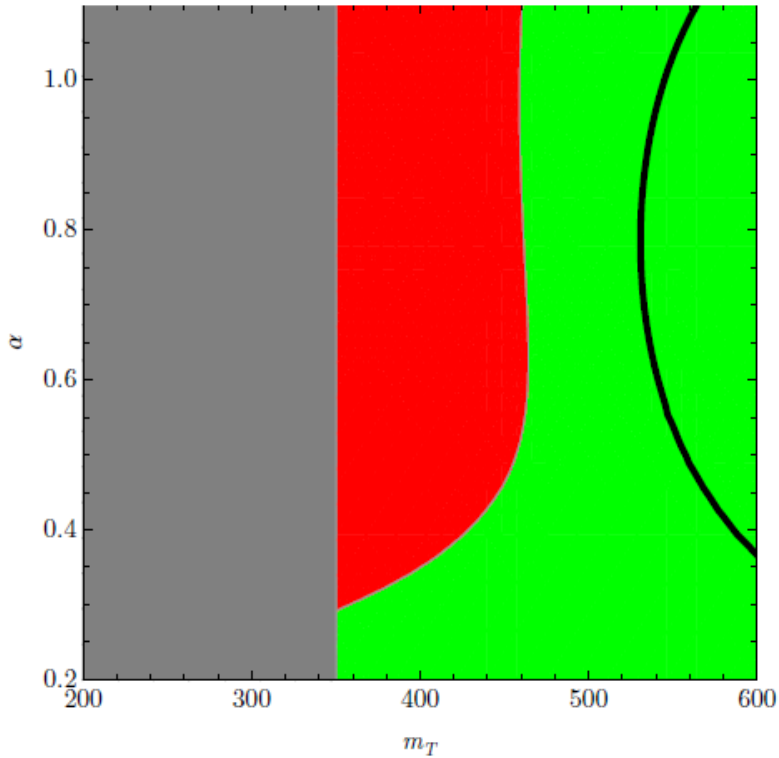
# Summary

- **Little Higgs** is one of the attractive models for physics beyond the SM.
- We study the phenomenology of non-colored new particles in **Litttest Higgs model with T-parity (LHT)** at the ILC.
- Heavy gauge boson and heavy lepton masses are determined with good accuracy  
& from these, we also obtain the VEV  $f$  & other model parameters.
- Furthermore, coupling of new particles can also be measured and it would help to confirm the TeV physics model.

Back up

Process	cross section (fb)	# of events	# of events after all cuts
$W_H W_H \rightarrow A_H A_H qqqq$	120.0	60,000	41,190
$WW \rightarrow qqqq$	1307	653,500	678
$eeWW \rightarrow eeqqqq$	490	245,000	46
$e\nu_e WZ \rightarrow e\nu_e qqqq$	24.5	12,250	3,797
$\nu\bar{\nu}WW \rightarrow \nu\bar{\nu}qqqq$	18.8	9,400	213
$ZWW \rightarrow \nu\bar{\nu}qqqq$	7.23	3,615	1,597
$Z_H Z_H \rightarrow A_H A_H qqqq$	5.61	2,805	1,533

Process	cross section (fb)	# of events	# of events after all cuts
$Z_H Z_H \rightarrow A_H A_H HH$	98.0	49760	18989
$\rightarrow A_H A_H qqqq$			
$WWZ$	63.86	31933	227
$\nu\nu WW$	14.67	7336	86
$WW$	3069	1947408	994
$tt$	192.9	96472	1749
$ZZ$	202.2	101094	114
$ZH$	17.98	8989	47
$W_H W_H$	108.6	54343	87



**Berger, Hubisz, Perelstein (1205.0013) Harigaya, Matsumoto, Nojiri, Tobioka (1204.2317)**

$$\chi^2 = \frac{(\text{rec } M_{SM1} - \text{tr } M_{SM})^2}{\sigma_{M_{SM}}^2} + \frac{(\text{rec } M_{SM2} - \text{tr } M_{SM})^2}{\sigma_{M_{SM}}^2} \quad (25)$$

