

Testing strong field QED with a laser on the electron beam at ILC/CLIC

Stefano Porto

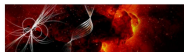
with Anthony Hartin and Gudrid Moortgat-Pick

ECFA Linear Collider Workshop 2013,

Hamburg, May 28th, 2013



Particles, Strings,
and the Early Universe
Collaborative Research Center SFB 676



- 1 Strong field QED
- 2 SLAC Experiment 144 and proposal at ILC/CLIC
- 3 Testing Zel'dovich quasi-levels and the Unruh effect
- 4 Conclusions and outlook

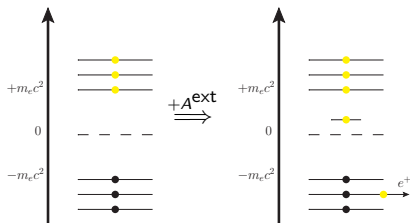
See also [Tony Hartin's talk at the working group RD9](#).

QED in strong external electromagnetic fields

Klein, Sauter, Schwinger:

$$E_{Cr} = \frac{m_e^2}{e} = 1.3 \cdot 10^{18} \text{ V/m}$$

$$B_{Cr} = \frac{m_e^2}{e} = 4.41 \cdot 10^9 \text{ T}$$



Schwinger critical field: **vacuum polarization**. Ex. laser $I \sim 10^{30} \text{ W/cm}^2$.

Can develop on magnetars, heavy ions collisions, and ... linear colliders.

- $\Upsilon \equiv \frac{e}{m_e^3} \sqrt{|(F_{\mu\nu} p^\nu)^2|} \sim \gamma \frac{E}{E_c} \text{ or } \gamma \frac{B}{B_c}$.

'Intensity' of the external field. Vacuum is **polarized** at $\Upsilon \sim 1$.

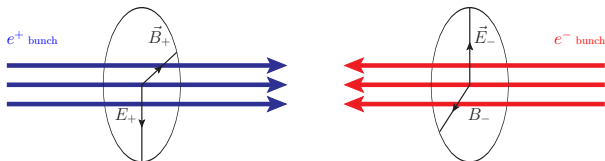
- $\eta = \frac{e}{m_e} \sqrt{|A_\mu A^\mu|}$.

Describes the multiplicity of photons from the external field in the initial states.

Strong fields at linear colliders

$$\text{High luminosity: } \mathcal{L} = f \frac{n_1 n_2}{4\pi\sigma_x\sigma_y}$$

needs **squeezed bunches** (small σ_x, σ_y) \implies large $\Upsilon \sim \frac{1}{\sigma_x + c\sigma_y}$



Intense charge bunches \longrightarrow strong field associated

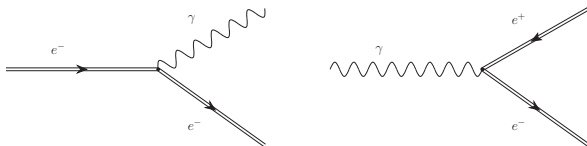
Machine	LEP II	SLC	ILC-1TeV	CLIC-3TeV
Energy (GeV)	94.5	46.6	500	1500
$\Upsilon_{average}$	0.00015	0.001	0.27	3.34

e^- sees 2 *almost* anticollinear **constant crossed fields**: $|\mathbf{E}| = |\mathbf{B}|$, $\mathbf{E} \cdot \mathbf{B} = 0$, **static**.

★ high density of photons + wavefunction overlap \implies external field \sim **classical field**.

Strong fields at linear colliders

Beamstrahlung; coherent pair production [Chen, Telnov '89]; incoherent pair production. Implemented in CAIN [Yokoya] and GUINEAPIG [Schulte].

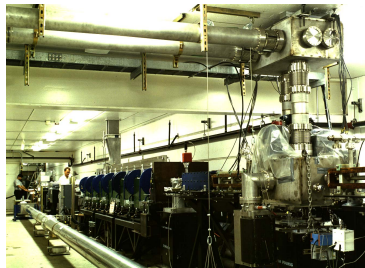
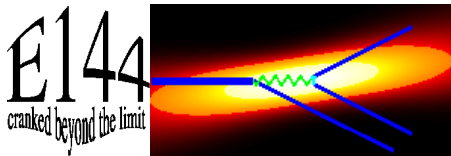


Need to consider all the processes in the strong field [SP, Hartin, Moortgat-Pick '13].

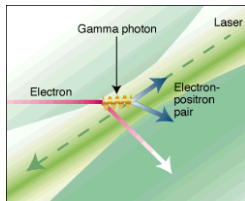
- Generalized Dirac equation: Furry picture [Furry '51]
- Classical electron motion: Quasi-classical operator approach [Baier & Katkov '51]

Strong Fields/Nonlinear QED to be considered in $\gamma\gamma$ interactions [Reiss '51], in particular:

- $\gamma\gamma$ -colliders [Ginzburg et al '81], [TESLA TDR].
- Laser physics ...



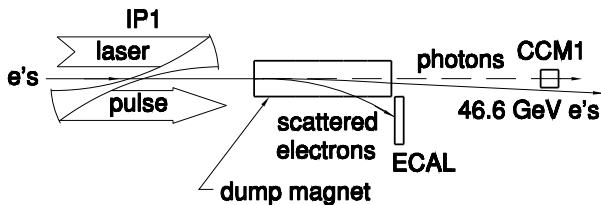
www.slac.stanford.edu/exp/e144/



- SLC Final Focus Test Beam (FFTB): 46.6 GeV
- Laser peak intensities: $\approx 0.5 \cdot 10^{18} \text{ W/cm}^2$
- Few weeks of operations

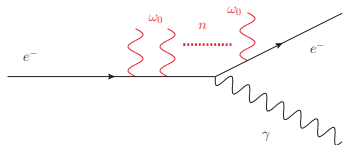
[Bamber *et al.* '99].

E144: set up for Nonlinear Compton scattering, $e^- + n\omega_0 \rightarrow e^- + \gamma$

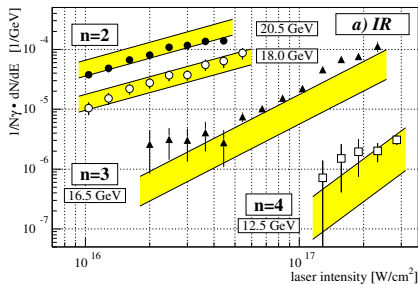
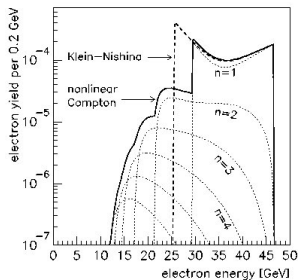


- $I = \frac{1}{377\Omega} E_{\text{lab}}^2$, $I_{\text{max}} \approx 0.5 \cdot 10^{18} \text{ W/cm}^2$
- $\eta^2 = 3.7 \cdot 10^{-19} I \lambda^2$, $\eta_{\text{max}} \sim 0.4$
- $\Upsilon_e = \eta \frac{p \cdot k}{m^2 c^4}$, $\Upsilon_{e, \text{max}} \sim 0.27$
- Terawatt pulses at $\lambda_{IR} = 1053 \text{ nm}$ and $\lambda_g = 527 \text{ nm}$, circular/linear polarized laser beams.

Nonlinear Compton scattering at E144, $e^- + n\omega_0 \rightarrow e^- + \gamma$



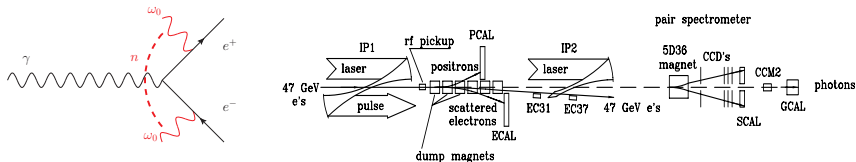
The rates are important, not the bad defined cross sections [Nikishov & Ritus '64].
Observed dependence on the number n of absorbed γ s.



$$\text{Rate at order } n \propto I^{n-1}$$

Adapted from [McDonald '98], [Bamber et al. '99].

Multiphoton Breit-Wheeler pair production at E144, $\gamma + n\omega_0 \rightarrow e^+e^-$



Theory: [Nikishov & Ritus '64-'65]. Dependences on the number n_γ and on Υ .

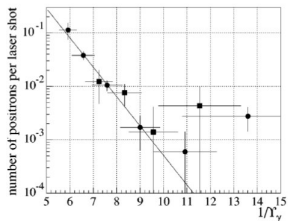
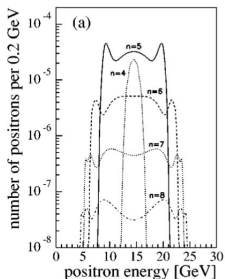


FIG. 49. Number of positrons per laser shot as a function of $1/Y_\gamma$. The circles are the 46.6 GeV data whereas the squares are the 49.1 GeV data. The solid line is a fit to the data.

Adapted from [McDonald '98], [Bamber et al. '99].

Proposal to repeat E144 at ILC/CLIC

Today and next future:

- Daresbury, FZD: 10 MeV- e^- and high intensity ($\eta \sim 10$).
- Future very high intensities: XFEL 10^{27} W/cm², ELI 10^{26} W/cm².

No high energy physics.

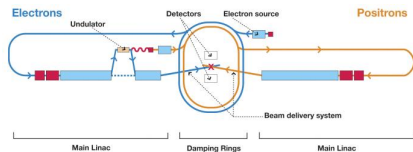
Proposal to repeat E144 at ILC/CLIC

Today and next future:

- Daresbury, FZD: 10 MeV- e^- and high intensity ($\eta \sim 10$).
- Future very high intensities: XFEL 10^{27} W/cm², ELI 10^{26} W/cm².

No high energy physics. Next linear collider?

Powerful electron beams 175-500 GeV: with similar laser as E144, $\Upsilon \gtrsim 1$:



Physical case for

- pair production
- vacuum polarization

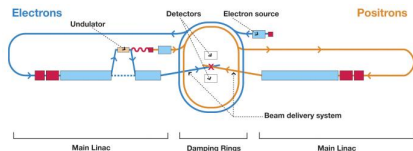
Proposal to repeat E144 at ILC/CLIC

Today and next future:

- Daresbury, FZD: 10 MeV- e^- and high intensity ($\eta \sim 10$).
- Future very high intensities: XFEL 10^{27} W/cm², ELI 10^{26} W/cm².

No high energy physics. Next linear collider?

Powerful electron beams 175-500 GeV: with similar laser as E144, $\Upsilon \gtrsim 1$:



Physical case for

- pair production
- vacuum polarization

Possibilities: **ADD-ON! no (compelling) DESIGN CHANGES**

- Exploiting the polarimeter laser or a specific laser on the beam dump.
- Exploiting Test-beam or parasitically the electron beam.
- Performing the same experiment on the positron beam (dump or not).

Side effect: electron mass-shift

The effective mass of an electron in an electromagnetic wave [Brown and Kibble '64], [Kibble '64]

$$\bar{m} = m\sqrt{1 + \eta^2}, \quad \eta = \frac{eE}{m\omega_0 c}$$

and the correspondent quasi-momentum $q_\mu = p_\mu + \frac{\eta^2 m^2}{2k \cdot p} k_\mu$ is such that $q^2 = \bar{m}^2$.

Side effect: electron mass-shift

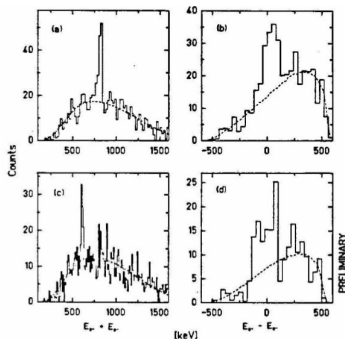
The effective mass of an electron in an electromagnetic wave [Brown and Kibble '64], [Kibble '64]

$$\bar{m} = m\sqrt{1 + \eta^2}, \quad \eta = \frac{eE}{m\omega_0 c}$$

and the correspondent quasi-momentum $q_\mu = p_\mu + \frac{\eta^2 m^2}{2k \cdot p} k_\mu$ is such that $q^2 = \bar{m}^2$.

Searching structures in the e^+e^- invariant-mass spectrum.

Darmstadt [Schweppe et al '83] and CERN-NA-046 [Bassompierre et al. '91] picks?



$$\bar{M}^2 = (q_1^2 + q_2^2) \propto \sqrt{n}$$

Complications because of \bar{m} .

$$1 + \frac{\Delta}{1 + \eta^2} \leq \frac{M^2}{4m^2} \leq 1 + \Delta$$

No decisive evidence for the mass shift at E144 [Bamber et al. '98].

Observed in Rochester [Meyerhofer et al. '95].

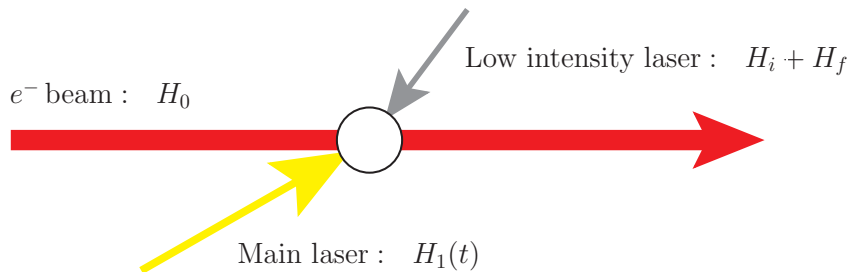
Radiative transitions between Zel'dovich quasi-levels in a laser

System H_0 interacting with em field $H_f + H_i$ under a T -periodic force $H_1(t)$: **NEW!**

$$H = H_0 + H_1(t) + H_f + H_i$$

\Rightarrow quasi-levels spectra with energy gaps $m \cdot \hbar\omega = \hbar 2\pi/T$ [Zel'dovich '67].

Transitions and radiations $m' \cdot \hbar\omega$ from **ground state** to **excited state** and viceversa due to H_1 .



Low I laser to the IP, frequency-tuned to stimulate quasi-levels transitions.

Outlook: probing the Unruh effect

A gravity accelerated observer out of a black hole sees Hawking radiation

[Hawking '74]:

$$T = \frac{\hbar g}{2\pi c k}$$

Equivalence



Principle

An accelerated observer in a region without gravity feels a radiation with a thermal bath [Fulling '72, Davies '75, Unruh '76]:

$$T = \frac{\hbar a}{2\pi c k}$$

Outlook: probing the Unruh effect

A gravity accelerated observer out of a black hole sees Hawking radiation

[Hawking '74]:

$$T = \frac{\hbar g}{2\pi c k}$$

Equivalence



Principle

An accelerated observer in a region without gravity feels a radiation with a thermal bath [Fulling '72, Davies '75, Unruh '76]:

$$T = \frac{\hbar a}{2\pi c k}$$

An electron observer Thomson-scatters with the Unruh thermal bath photons.

Unruh radiation:

$$\frac{dU_{\text{Unruh}}}{dt} = \frac{\hbar r_e^2 \left(\frac{eE}{m}\right)^4}{90\pi c^6} = \frac{\hbar r_e^2 a^4}{90\pi c^6} \propto T^4$$

Larmor radiation:

$$\frac{dU_{\text{Larmor}}}{dt} = \frac{\hbar 2e^2 a^2}{3c^3}$$

$\frac{dU_{\text{Unruh}}}{dt} \sim \frac{dU_{\text{Larmor}}}{dt}$ at $\sim 100 \cdot E_{\text{cr}}$. With 500 GeV electrons we would need:

$$E_{\text{lab}} \sim 10^{13} \text{ V/m}, \quad I_{\text{laser}} \sim 10^{19} - 10^{20} \text{ W/cm}^2$$

Large background. Can be observed at all?

Debate still open [Akhmedov, Singleton '07], [Martín-Martínez, Fuentes, Mann '11], [Ford, O'Connell '05].

- Strong external fields affect the motion and the states of particles (nonlinear QED).
- E144 studied the effect of a laser on a HE electron beam, in particular nonlinear Compton scattering and pair production.
- Next LC: opportunity to repeat E144, at **higher energy** and **statistics**. It would be important to discuss it **SOON**.
- Possibility of performing the same experiment on the positron beam.
- Studying mass-shift effects.
- Testing Zel'dovich quasi-levels theory.
- Check whether is possible to detect the Unruh radiation.

2nd International workshop on

Physics in Intense Fields

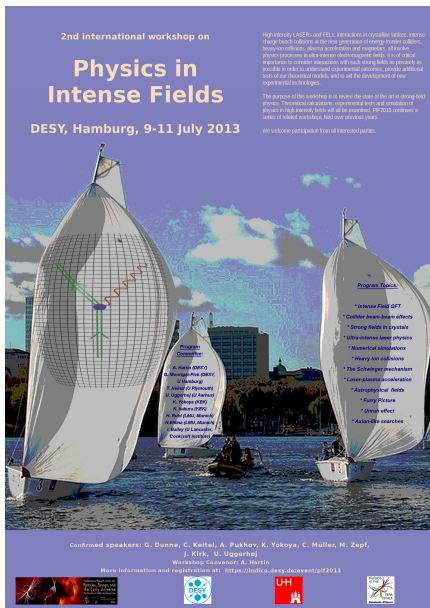
DESY, Hamburg, 9-11 July 2013

High intensity LASERS and FELs, interactions in crystalline lattices, intense charge bunches followed by the most generation of energy losses, electron-ion collisions, plasma acceleration and resonators, of exotic physics processes in ultra-short electromagnetic fields. It is of critical importance to consider interactions with such strong fields as presented as possible in order to understand experimental outcomes, provide additional fields of theoretical models, and/or aid the development of new experimental techniques.

The purpose of this workshop is to review the state of the art in strong field physics. Theoretical calculations, experimental tests and simulation of physics in high intensity fields will also be covered. PIF2013 contains a series of related workshops held over previous years.


We welcome participation from all interested parties.

Confirmed speakers: G. Dunne, C. Keitel, A. Pukhov, K. Yokoya, C. Müller, M. Zepf, J. Kirk, U. Uggerhøj
Workshop Co-Chair: A. Martin
More information and registration at: <https://indico.desy.de/event/pif2013>



Confirmed speakers:
A. Martin (DESY)
G. Dunne (PIF DESY)
U. Uggerhøj
C. Keitel (F. Physik)
U. Uggerhøj (F. Physik)
A. Pukhov (DESY)
K. Yokoya (KIPAC)
M. Zepf (SLAC)
A. Pukhov (ILR, Mainz)
J. Kirk (Lund)
C. Müller (MPI)

Program Topics:
* Intense Field QFT
* Cluster beam-beam effects
* Strong fields in crystals
* Ultra-intense laser physics
* Nonlinear structures
* Heavy ion collisions
* The Schwinger mechanism
* Laser-plasma acceleration
* Astrophysical fields
* Flurry Picture
* Unruh effect
* Axion like particles



DESY, Hamburg 9th-11th July 2013

Topics: Strong field QED/QFT in

- LCs
- LASERS
- Crystals
- Heavy ion collisions etc.

Chair: Tony Martin

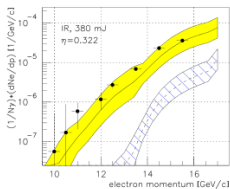
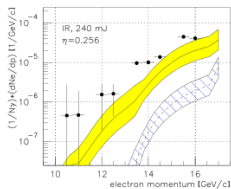
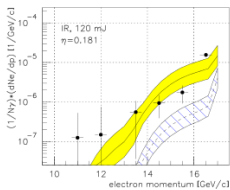
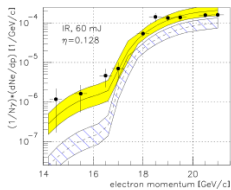
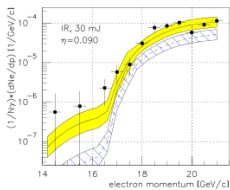
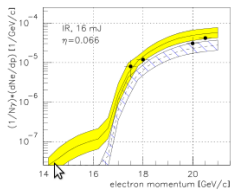
Registration deadline: 28th June

(Abstracts deadline: 31st May)

<https://indico.desy.de/event/pif2013>

Thanks!

Backup: Compton scattering at E144

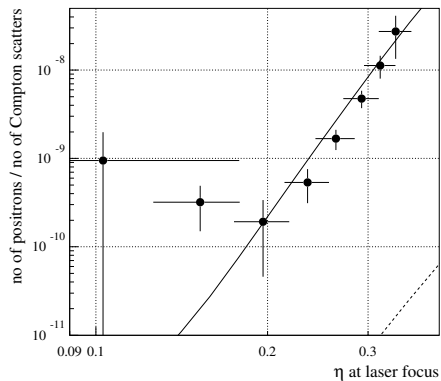


Scattering rates observed at E144 in terms of electron energy.

$n_\gamma = 2, 3$ -edges.

[McDonald '98]

Backup: Pair production at E144



[McDonald '98]

Backup: Observation of relativistic mass-shift effects

[Meyerhofer et al. '95]

Electrons (up to 175 keV) from the ionization of Kr by a laser:

- Laser: $\lambda = 1053 \text{ nm}$, 1 ps , $I_{\text{peak}} \simeq 10^{18} \text{ W/cm}^2$
- Relativistic mass-shift associated with the electrons quiver motion is apparent;
- Electron trajectory obtained without considering the mass-shift disagrees with the experimental observations.

